

NLO QCD corrections to Vector Boson Pair Production via Vector Boson Fusion

Giuseppe Bozzi

Institut für Theoretische Physik
Universität Karlsruhe

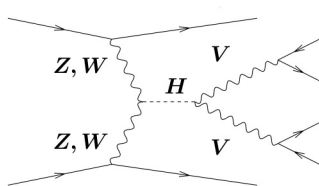
HEP 2007
Manchester, 19.07.2007

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

Higgs production via VBF



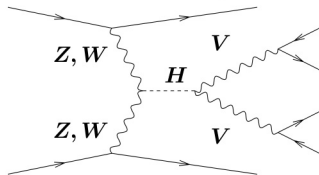
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



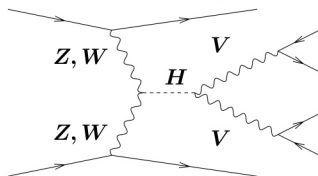
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC

- Clean experimental signature

- *two highly energetic outgoing jets*

- *large rapidity interval between jets*

- *no hadronic activity in the rapidity interval between jets*

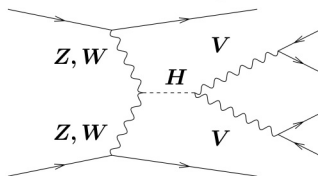
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



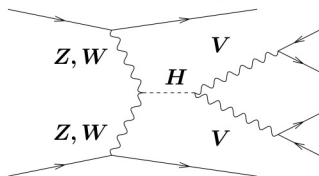
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



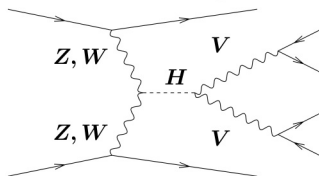
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



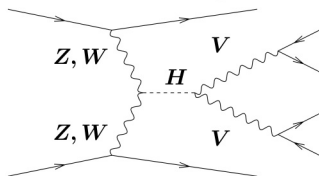
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



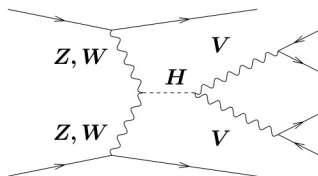
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

Higgs production via VBF



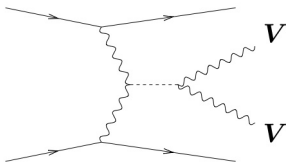
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

→ **Very promising channel at the LHC**

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

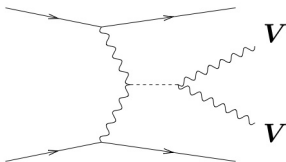
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



• Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

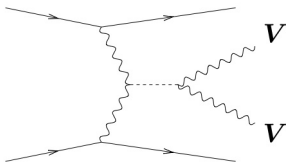
- similar features as H production \rightarrow *irreducible* background

• New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
- subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

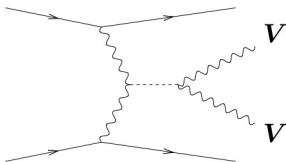
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

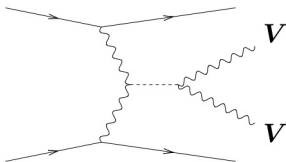
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

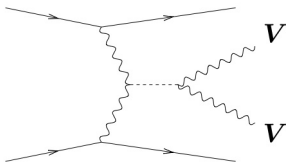
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

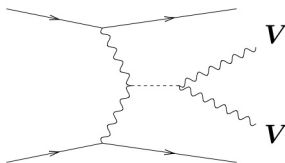
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

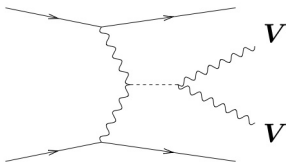
- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
- subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow Need accurate predictions for EW VVjj production!

VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

[Kauer, Plehn, Rainwater, Zeppenfeld (2001)]

- similar features as H production \rightarrow *irreducible* background

- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

\rightarrow **Need accurate predictions for EW VVjj production!**

Challenges

- Multi-parton process: huge number of Feynman diagrams

- $2 \rightarrow 4$ for $qq \rightarrow qqVV$

- $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$

→ how to speed up the evaluation?

- Suitable treatment of pentagon contributions

→ how to solve numerical instabilities?

→ Build a fully-flexible partonic Monte Carlo program allowing for

- computation of jet observables at NLO-QCD accuracy
- straightforward implementation of cuts

WW: [Jaeger, Oleari, Zeppenfeld (2006)]

ZZ: [Jaeger, Oleari, Zeppenfeld (2006)]

WZ: [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams

- $2 \rightarrow 4$ for $qq \rightarrow qqVV$

- $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$

→ how to speed up the evaluation?

- Suitable treatment of pentagon contributions

→ how to solve numerical instabilities?

→ Build a fully-flexible partonic Monte Carlo program allowing for

- computation of jet observables at NLO-QCD accuracy
- straightforward implementation of cuts

WW: [Jaeger, Oleari, Zeppenfeld (2006)]

ZZ: [Jaeger, Oleari, Zeppenfeld (2006)]

WZ: [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams

- $2 \rightarrow 4$ for $qq \rightarrow qqVV$

- $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$

→ how to speed up the evaluation?

- Suitable treatment of pentagon contributions

→ how to solve numerical instabilities?

→ Build a fully-flexible partonic Monte Carlo program allowing for

- computation of jet observables at NLO-QCD accuracy
- straightforward implementation of cuts

WW: [Jaeger, Oleari, Zeppenfeld (2006)]

ZZ: [Jaeger, Oleari, Zeppenfeld (2006)]

WZ: [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams
 - $2 \rightarrow 4$ for $qq \rightarrow qqVV$
 - $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$
 → how to speed up the evaluation?
- Suitable treatment of pentagon contributions
 - how to solve numerical instabilities?

- Build a fully-flexible partonic Monte Carlo program allowing for
- computation of jet observables at NLO-QCD accuracy
 - straightforward implementation of cuts

WW : [Jaeger, Oleari, Zeppenfeld (2006)]
 ZZ : [Jaeger, Oleari, Zeppenfeld (2006)]
 WZ : [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams
 - $2 \rightarrow 4$ for $qq \rightarrow qqVV$
 - $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$
 → how to speed up the evaluation?
- Suitable treatment of pentagon contributions
 - how to solve numerical instabilities?

- Build a fully-flexible partonic Monte Carlo program allowing for
- computation of jet observables at NLO-QCD accuracy
 - straightforward implementation of cuts

WW : [Jaeger, Oleari, Zeppenfeld (2006)]
 ZZ : [Jaeger, Oleari, Zeppenfeld (2006)]
 WZ : [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams
 - $2 \rightarrow 4$ for $qq \rightarrow qqVV$
 - $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$
 - how to speed up the evaluation?
 - Suitable treatment of pentagon contributions
 - how to solve numerical instabilities?
- Build a fully-flexible partonic Monte Carlo program allowing for
- computation of jet observables at NLO-QCD accuracy
 - straightforward implementation of cuts

WW: [Jaeger, Oleari, Zeppenfeld (2006)]
 ZZ: [Jaeger, Oleari, Zeppenfeld (2006)]
 WZ: [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Challenges

- Multi-parton process: huge number of Feynman diagrams
 - $2 \rightarrow 4$ for $qq \rightarrow qqVV$
 - $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$
 → how to speed up the evaluation?
 - Suitable treatment of pentagon contributions
 - how to solve numerical instabilities?
- Build a fully-flexible partonic Monte Carlo program allowing for
- computation of jet observables at NLO-QCD accuracy
 - straightforward implementation of cuts

WW: [Jaeger, Oleari, Zeppenfeld (2006)]

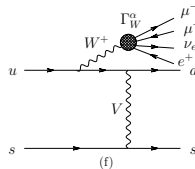
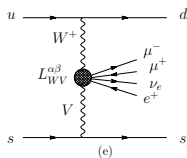
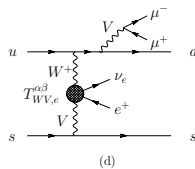
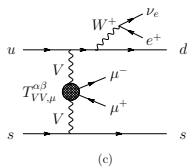
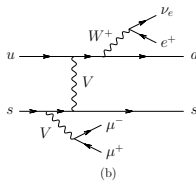
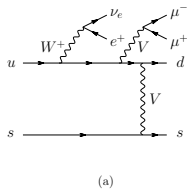
ZZ: [Jaeger, Oleari, Zeppenfeld (2006)]

WZ: [gb, Jaeger, Oleari, Zeppenfeld (2007)]

Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - **Tree-level features**
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

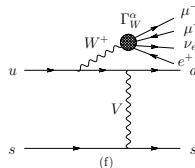
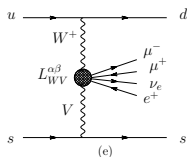
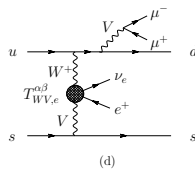
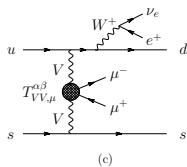
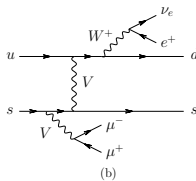
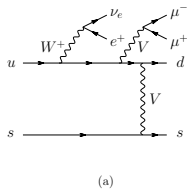
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

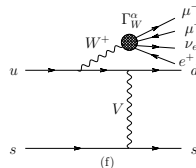
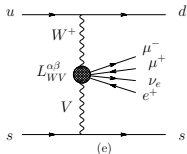
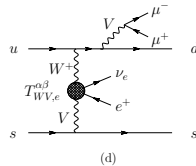
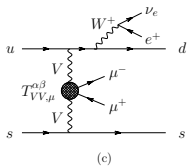
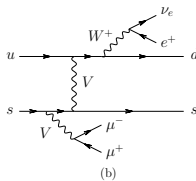
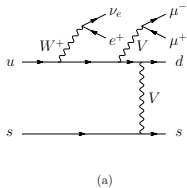
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

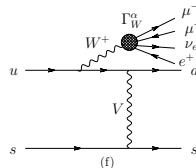
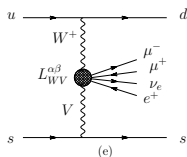
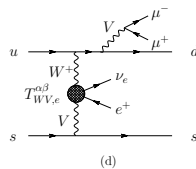
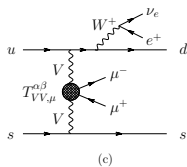
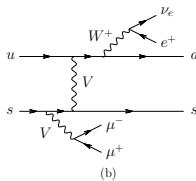
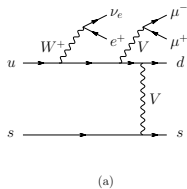
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

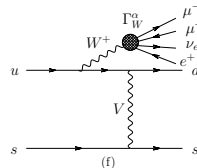
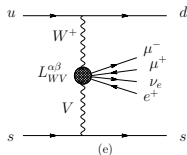
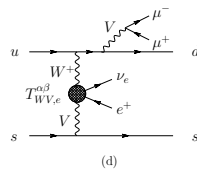
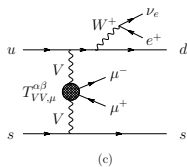
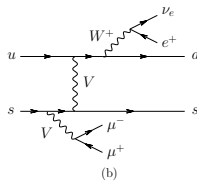
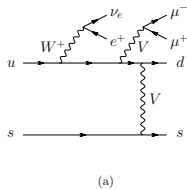
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

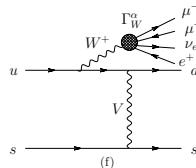
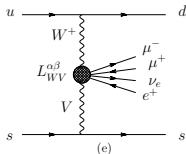
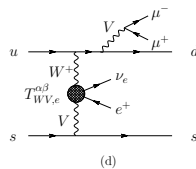
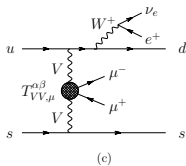
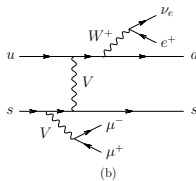
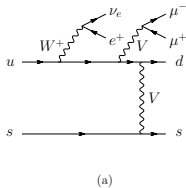
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

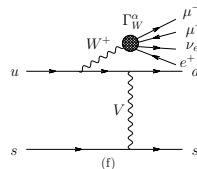
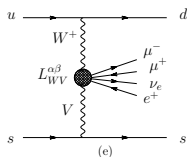
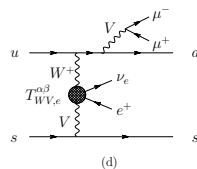
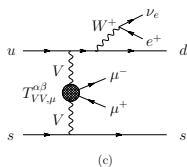
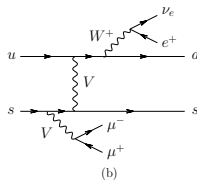
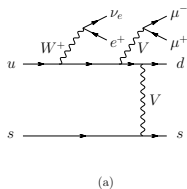
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

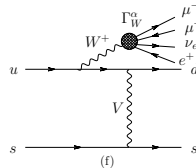
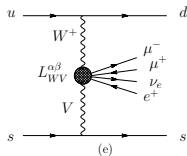
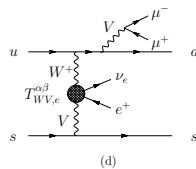
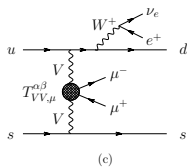
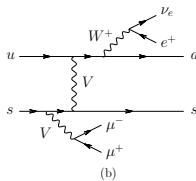
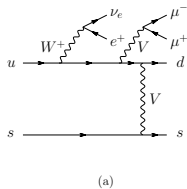
(c) quark line +
leptonic tensor $T_{VV,\mu}$

(d) quark line +
leptonic tensor $T_{WV,e}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

(c) quark line +
leptonic tensor $T_{VV, \mu}^{\alpha\beta}$

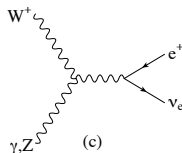
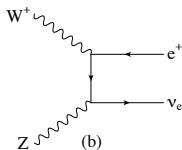
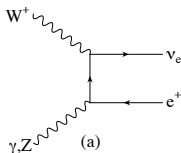
(d) quark line +
leptonic tensor $T_{WV, e}^{\alpha\beta}$

(e) leptonic tensor L_{WV}
VBF topology

(f) leptonic tensor Γ_W

Leptonic tensors

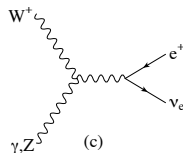
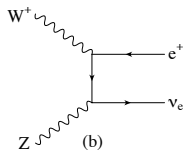
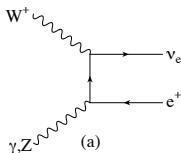
Example: T_{WV} built up from



- Sum of sub-amplitudes involving EW bosons and leptons only
- Use them for diagrams with *same* topology but *differences* in quark propagators
- Develop **modular structure** to speed up the calculation:
 - compute common building blocks of several diagrams only **once** per phase-space point
 - straightforward (future) implementation of **new-physics** effects in the bosonic/leptonic sector

Leptonic tensors

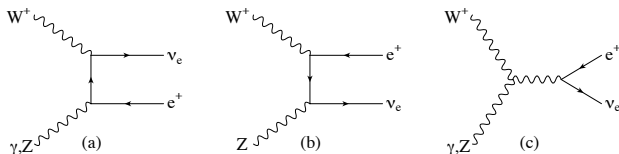
Example: T_{WV} built up from



- Sum of sub-amplitudes involving EW bosons and leptons only
- Use them for diagrams with *same* topology but *differences* in quark propagators
- Develop **modular structure** to speed up the calculation:
 - compute common building blocks of several diagrams only **once** per phase-space point
 - straightforward (future) implementation of **new-physics** effects in the bosonic/leptonic sector

Leptonic tensors

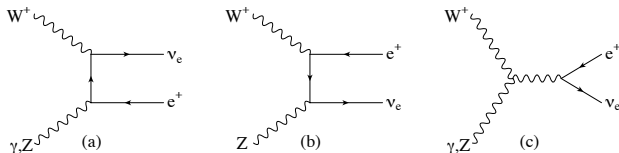
Example: T_{WV} built up from



- Sum of sub-amplitudes involving EW bosons and leptons only
- Use them for diagrams with *same* topology but *differences* in quark propagators
- Develop **modular structure** to speed up the calculation:
 - compute common building blocks of several diagrams only **once** per phase-space point
 - straightforward (future) implementation of **new-physics** effects in the bosonic/leptonic sector

Leptonic tensors

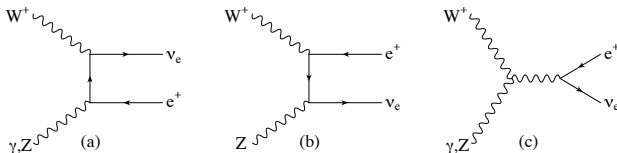
Example: T_{WV} built up from



- Sum of sub-amplitudes involving EW bosons and leptons only
- Use them for diagrams with *same* topology but *differences* in quark propagators
- Develop **modular structure** to speed up the calculation:
 - compute common building blocks of several diagrams only **once** per phase-space point
 - straightforward (future) implementation of **new-physics** effects in the bosonic/leptonic sector

Leptonic tensors

Example: T_{WV} built up from



- Sum of sub-amplitudes involving EW bosons and leptons only
- Use them for diagrams with *same* topology but *differences* in quark propagators
- Develop **modular structure** to speed up the calculation:
 - compute common building blocks of several diagrams only **once** per phase-space point
 - straightforward (future) implementation of **new-physics** effects in the bosonic/leptonic sector

Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - **NLO: real contributions**
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

Real corrections

- Attach a gluon to the quark lines in all possible ways
 - Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
 - **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction
- [Catani, Seymour (1997)]
- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1+\epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Real corrections

- Attach a gluon to the quark lines in all possible ways
- Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
- **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction

[Catani, Seymour (1997)]

- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1+\epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Real corrections

- Attach a gluon to the quark lines in all possible ways
- Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
- **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction

[Catani, Seymour (1997)]

- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1+\epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Real corrections

- Attach a gluon to the quark lines in all possible ways
- Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
- **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction

[Catani, Seymour (1997)]

- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1+\epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Real corrections

- Attach a gluon to the quark lines in all possible ways
- Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
- **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction

[Catani, Seymour (1997)]

- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1+\epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Real corrections

- Attach a gluon to the quark lines in all possible ways
- Crossing diagrams: initial gluon splitting in a $q\bar{q}$ pair
- **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction

[Catani, Seymour (1997)]

- Divergences only depend on the colour structure of the external partons
 - subtraction terms *identical* to Higgs production via VBF

$$\langle I(\epsilon) \rangle = |M_B|^2 \frac{\alpha_s(\mu_R)}{2\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + 9 - \frac{4}{3}\pi^2 \right]$$

[Figy, Oleari, Zeppenfeld (2003)]

(Q =momentum transfer between initial and final state quark)

Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - **NLO: virtual contributions**
- 3 Selected results
 - Differential distributions at the LHC

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → **up to pentagons!**

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → up to pentagons!

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → up to pentagons!

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → up to pentagons!

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → up to pentagons!

Virtual corrections

- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → up to pentagons!

Virtual corrections

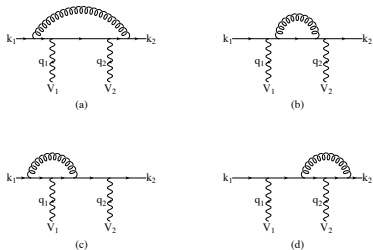
- Add interference between Born and virtual amplitudes
- EW bosons exchanged in the t -channel are colour-singlet
 - no contributions from gluons attached both to upper and lower quark lines!
 - consider radiative corrections to single quark line only
- Regularization performed in the Dimensional Reduction scheme
 - PV reduction performed in $d = 4 - 2\epsilon$ dimensions
 - algebra of γ , p , ϵ performed in $d = 4$ dimensions

[Siegel (1979)]

- Three classes of contributions: virtual corrections along a quark line with 1,2,3 vector boson(s) attached → **up to pentagons!**

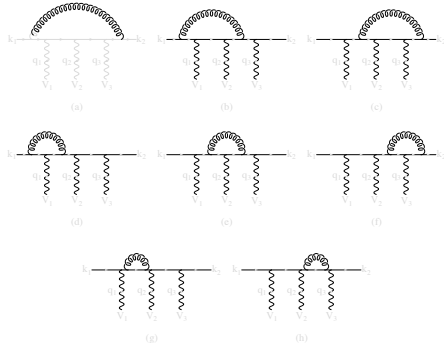
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

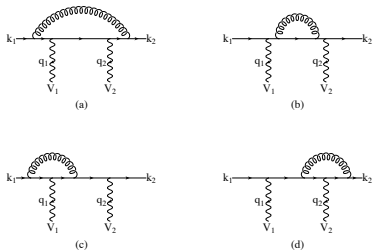
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

Topologies

quark line with 2 bosons attached

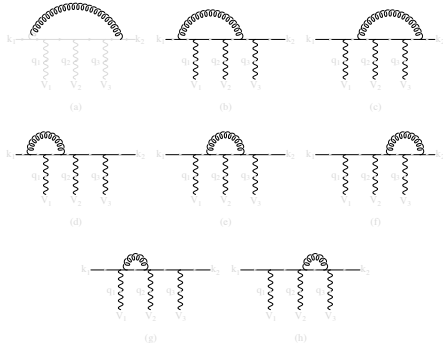


- self-energies

- triangles

- boxes

quark line with 3 bosons attached



- self-energies

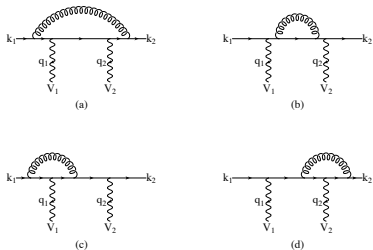
- triangles

- boxes

- pentagons

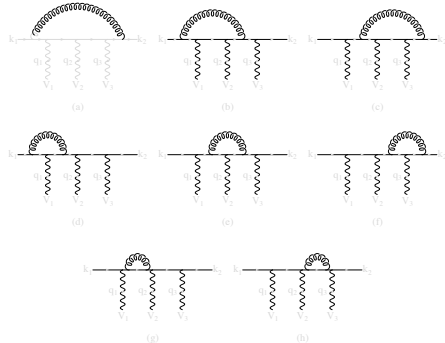
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

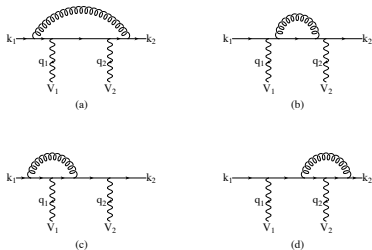
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

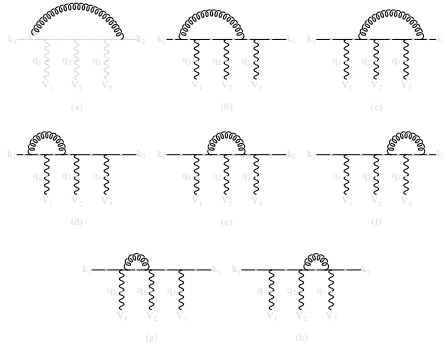
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

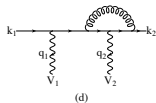
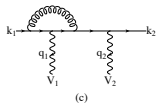
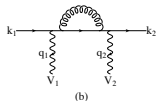
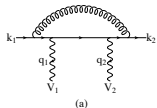
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

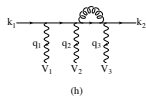
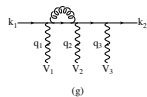
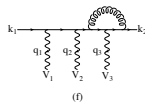
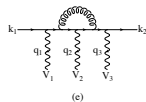
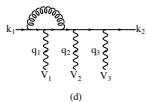
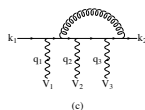
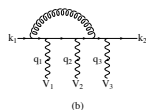
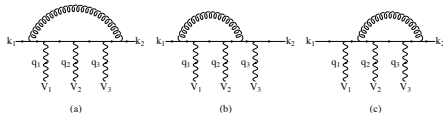
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

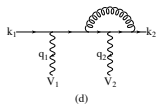
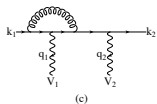
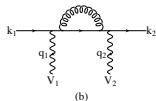
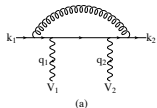
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

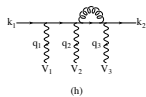
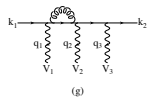
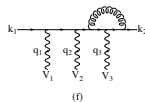
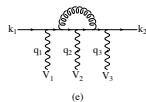
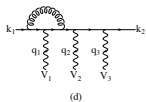
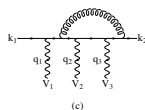
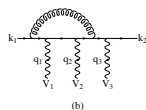
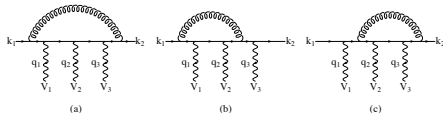
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

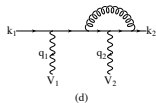
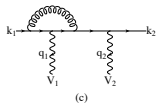
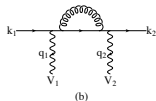
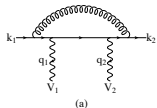
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

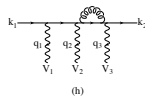
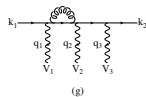
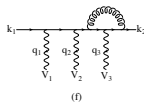
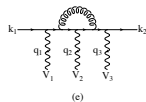
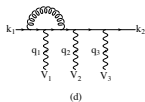
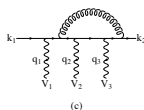
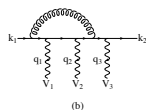
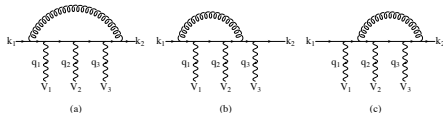
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

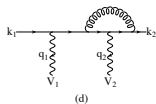
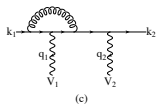
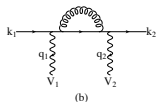
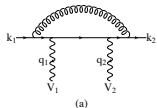
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

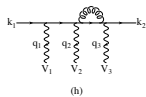
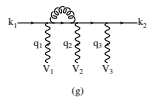
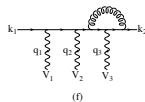
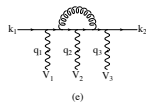
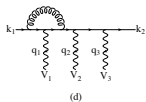
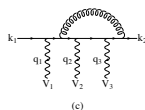
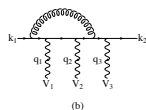
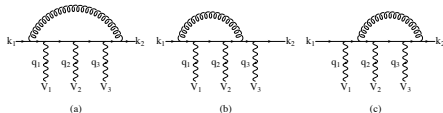
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

quark line with 3 bosons attached

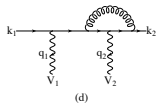
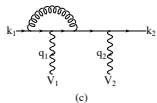
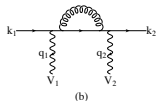
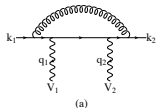


- self-energies
- triangles
- boxes

• pentagons

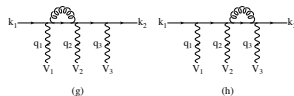
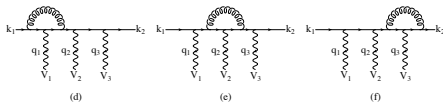
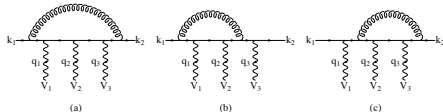
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- **pentagons**

Finite contributions

Summing up:

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{4\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + c_{\text{virt}} \right] + \widetilde{\mathcal{M}}_V$$

- divergent part proportional to Born amplitude
 - exactly cancels the phase-space integral of the dipole terms
- finite term proportional to Born amplitude
- finite *non-universal* term $\widetilde{\mathcal{M}}_V$
 - can be computed in $d = 4$ dimensions
 - given in terms of the **finite parts** of the Passarino-Veltman $B_{ij}, C_{ij}, D_{ij}, E_{ij}$ coefficient functions

Finite contributions

Summing up:

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{4\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + c_{\text{virt}} \right] + \widetilde{\mathcal{M}}_V$$

- divergent part proportional to Born amplitude
 - exactly cancels the phase-space integral of the dipole terms
- finite term proportional to Born amplitude
- finite *non-universal* term $\widetilde{\mathcal{M}}_V$
 - can be computed in $d = 4$ dimensions
 - given in terms of the **finite parts** of the Passarino-Veltman $B_{ij}, C_{ij}, D_{ij}, E_{ij}$ coefficient functions

Finite contributions

Summing up:

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{4\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + c_{\text{virt}} \right] + \widetilde{\mathcal{M}}_V$$

- divergent part proportional to Born amplitude
 - exactly cancels the phase-space integral of the dipole terms
- finite term proportional to Born amplitude
- finite *non-universal* term $\widetilde{\mathcal{M}}_V$
 - can be computed in $d = 4$ dimensions
 - given in terms of the **finite parts** of the Passarino-Veltman $B_{ij}, C_{ij}, D_{ij}, E_{ij}$ coefficient functions

Finite contributions

Summing up:

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{4\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + c_{\text{virt}} \right] + \widetilde{\mathcal{M}}_V$$

- divergent part proportional to Born amplitude
 - exactly cancels the phase-space integral of the dipole terms
- finite term proportional to Born amplitude
- finite *non-universal* term $\widetilde{\mathcal{M}}_V$
 - can be computed in $d = 4$ dimensions
 - given in terms of the **finite parts** of the Passarino-Veltman $B_{ij}, C_{ij}, D_{ij}, E_{ij}$ coefficient functions

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator

→ keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

→ computed through Passarino-Veltman reduction procedure

→ numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

→ numerical instabilities if kinematical invariants (Gram determinant) become small

→ use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator
 - keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals
 - computed through Passarino-Veltman reduction procedure
 - numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals
 - numerical instabilities if kinematical invariants (Gram determinant) become small
 - use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator
 - keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

- computed through Passarino-Veltman reduction procedure
- numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

- numerical instabilities if kinematical invariants (Gram determinant) become small
- use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator
 - keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

- computed through Passarino-Veltman reduction procedure
- numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

- numerical instabilities if kinematical invariants (Gram determinant) become small
- use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator

→ keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

→ computed through Passarino-Veltman reduction procedure

→ numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

→ numerical instabilities if kinematical invariants (Gram determinant) become small

→ use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator

→ keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

→ computed through Passarino-Veltman reduction procedure
 → numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

→ numerical instabilities if kinematical invariants (Gram determinant) become small
 → use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator
 - keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals
 - computed through Passarino-Veltman reduction procedure
 - numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals
 - numerical instabilities if kinematical invariants (Gram determinant) become small
 - use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator
 - keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals
 - computed through Passarino-Veltman reduction procedure
 - numerically stable in phase-space regions relevant for VBF

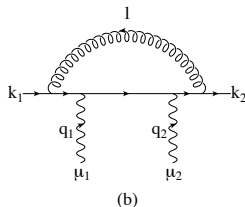
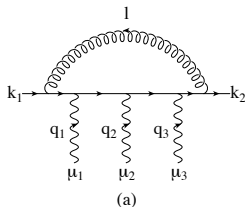
[Passarino, Veltman (1979)]

- Five-point tensor integrals
 - numerical instabilities if kinematical invariants (Gram determinant) become small
 - use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l + k_1 + \not{q}_{123}} \gamma_{\mu_3} \frac{1}{l + k_1 + \not{q}_{12}} \gamma_{\mu_2} \frac{1}{l + k_1 + \not{q}_1} \gamma_{\mu_1} \frac{1}{l + k_1} \gamma_\alpha \frac{1}{\beta}$$

$$q_1^{\mu_1} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_2 \mu_3}(k_1, q_1 + q_2, q_3) - \mathcal{D}_{\mu_2 \mu_3}(k_1 + q_1, q_2, q_3)$$

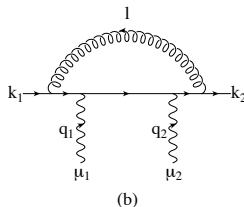
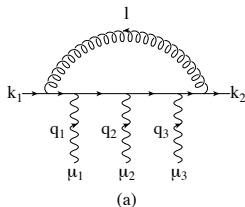
$$q_2^{\mu_2} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1, q_2 + q_3) - \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1 + q_2, q_3)$$

$$q_3^{\mu_3} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2) - \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2 + q_3)$$

Express $\mathcal{E}_{\mu_1 \mu_2 \mu_3}$ ($\mathcal{D}_{\mu_1 \mu_2}$) as a sum of coefficients up to E_{ij} (D_{ij}) and verify the Ward identities \rightarrow **strong check** of the code!

Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l + k_1 + \not{q}_{123}} \gamma_{\mu_3} \frac{1}{l + k_1 + \not{q}_{12}} \gamma_{\mu_2} \frac{1}{l + k_1 + \not{q}_1} \gamma_{\mu_1} \frac{1}{l + k_1} \gamma_\alpha \frac{1}{l^2}$$

$$q_1^{\mu_1} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_2 \mu_3}(k_1, q_1 + q_2, q_3) - \mathcal{D}_{\mu_2 \mu_3}(k_1 + q_1, q_2, q_3)$$

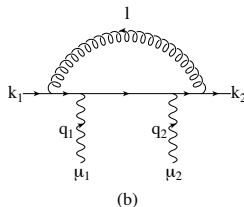
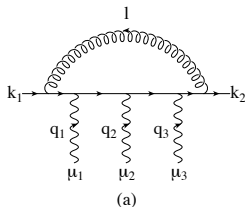
$$q_2^{\mu_2} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1, q_2 + q_3) - \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1 + q_2, q_3)$$

$$q_3^{\mu_3} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2) - \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2 + q_3)$$

Express $\mathcal{E}_{\mu_1 \mu_2 \mu_3}$ ($\mathcal{D}_{\mu_1 \mu_2}$) as a sum of coefficients up to E_{ij} (D_{ij}) and verify the Ward identities \rightarrow **strong check** of the code!

Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l+k_1+\not{q}_{123}} \gamma_{\mu_3} \frac{1}{l+k_1+\not{q}_{12}} \gamma_{\mu_2} \frac{1}{l+k_1+\not{q}_1} \gamma_{\mu_1} \frac{1}{l+k_1} \gamma_\alpha \frac{1}{l^2}$$

$$q_1^{\mu_1} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_2\mu_3}(k_1, q_1 + q_2, q_3) - \mathcal{D}_{\mu_2\mu_3}(k_1 + q_1, q_2, q_3)$$

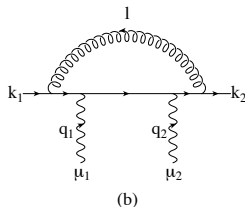
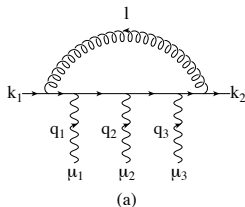
$$q_2^{\mu_2} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1\mu_3}(k_1, q_1, q_2 + q_3) - \mathcal{D}_{\mu_1\mu_3}(k_1, q_1 + q_2, q_3)$$

$$q_3^{\mu_3} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1\mu_2}(k_1, q_1, q_2) - \mathcal{D}_{\mu_1\mu_2}(k_1, q_1, q_2 + q_3)$$

Express $\mathcal{E}_{\mu_1\mu_2\mu_3}$ ($\mathcal{D}_{\mu_1\mu_2}$) as a sum of coefficients up to E_{ij} (D_{ij}) and verify the Ward identities \rightarrow strong check of the code!

Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l+k_1+\not{q}_{123}} \gamma_{\mu_3} \frac{1}{l+k_1+\not{q}_{12}} \gamma_{\mu_2} \frac{1}{l+k_1+\not{q}_1} \gamma_{\mu_1} \frac{1}{l+k_1} \gamma_\alpha \frac{1}{l^2}$$

$$q_1^{\mu_1} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_2\mu_3}(k_1, q_1 + q_2, q_3) - \mathcal{D}_{\mu_2\mu_3}(k_1 + q_1, q_2, q_3)$$

$$q_2^{\mu_2} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1\mu_3}(k_1, q_1, q_2 + q_3) - \mathcal{D}_{\mu_1\mu_3}(k_1, q_1 + q_2, q_3)$$

$$q_3^{\mu_3} \mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1\mu_2}(k_1, q_1, q_2) - \mathcal{D}_{\mu_1\mu_2}(k_1, q_1, q_2 + q_3)$$

Express $\mathcal{E}_{\mu_1\mu_2\mu_3}(\mathcal{D}_{\mu_1\mu_2})$ as a sum of coefficients up to E_{ij} (D_{ij}) and verify the Ward identities → **strong check** of the code!

“True” pentagons

- Loop amplitudes eventually contracted with leptonic currents

- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

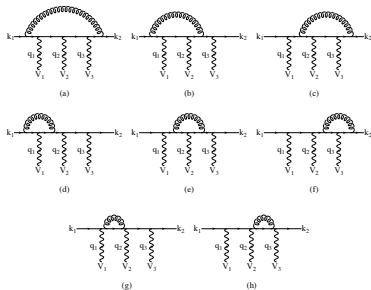
- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1 \mu_2 \mu_3}(k_1, q_+, q_-)$

- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

→ $M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1 \mu_2 \mu_3} + \text{boxes}$

- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



“True” pentagons

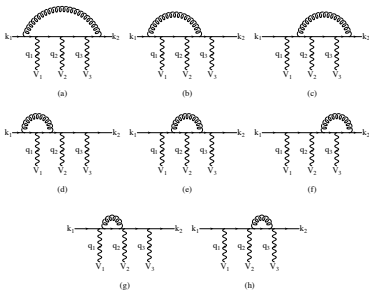
- Loop amplitudes eventually contracted with leptonic currents
- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1\mu_2\mu_3}(k_1, q_+, q_-, q_0)$
- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

$$\rightarrow M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1\mu_2\mu_3} + \text{boxes}$$

- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



“True” pentagons

- Loop amplitudes eventually contracted with leptonic currents
- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

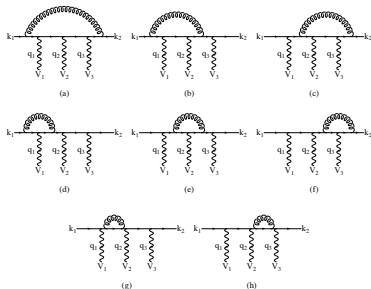
- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1 \mu_2 \mu_3}(k_1, q_+, q_-, q_0)$

- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

→ $M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1 \mu_2 \mu_3} + \text{boxes}$

- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



“True” pentagons

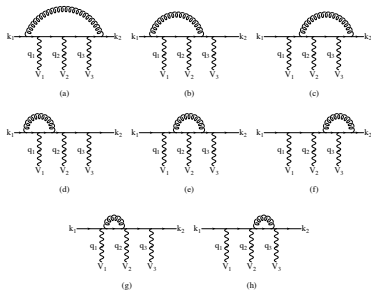
- Loop amplitudes eventually contracted with leptonic currents
- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1\mu_2\mu_3}(k_1, q_+, q_-, q_0)$
- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

$$\rightarrow M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1\mu_2\mu_3} + \text{boxes}$$

\rightarrow Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



“True” pentagons

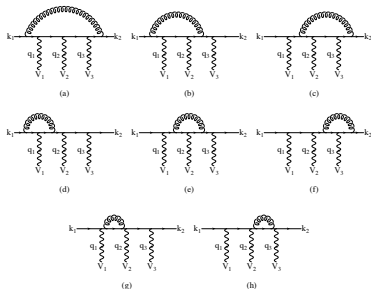
- Loop amplitudes eventually contracted with leptonic currents
- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1 \mu_2 \mu_3}(k_1, q_+, q_-, q_0)$
- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

$$\rightarrow M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1 \mu_2 \mu_3} + \text{boxes}$$

- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



“True” pentagons

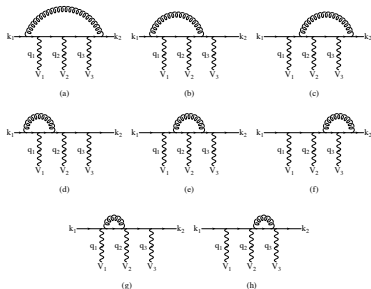
- Loop amplitudes eventually contracted with leptonic currents
- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \mathcal{P}_{\mu_1 \mu_2 \mu_3}(k_1, q_+, q_-, q_0)$

- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

→ $M_5 = r_+^{\mu_1} r_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1 \mu_2 \mu_3} + \text{boxes}$

- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon contribution to cross section



Numerical stability of pentagon contributions

- Gauge-check procedure

- Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
- Discard these points for the calculation of the finite parts
- Correct the remaining pentagon contributions by a factor $1/(1 - f)$
- Error induced by this approximation: far below numerical accuracy of Monte Carlo program

- PV formalism: $f \sim 15\%$

- DD formalism: $f \sim 0.1\%$

→ Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure

- Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
- Discard these points for the calculation of the finite parts
- Correct the remaining pentagon contributions by a factor $1/(1 - f)$
- Error induced by this approximation: far below numerical accuracy of Monte Carlo program

- PV formalism: $f \sim 15\%$

- DD formalism: $f \sim 0.1\%$

→ Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
 - PV formalism: $f \sim 15\%$
 - DD formalism: $f \sim 0.1\%$
- Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
 - PV formalism: $f \sim 15\%$
 - DD formalism: $f \sim 0.1\%$
- Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
 - PV formalism: $f \sim 15\%$
 - DD formalism: $f \sim 0.1\%$
- Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
 - PV formalism: $f \sim 15\%$
 - DD formalism: $f \sim 0.1\%$
- Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
- PV formalism: $f \sim 15\%$
- DD formalism: $f \sim 0.1\%$

→ Pentagons under control using DD formalism!

Numerical stability of pentagon contributions

- Gauge-check procedure
 - Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
 - Discard these points for the calculation of the finite parts
 - Correct the remaining pentagon contributions by a factor $1/(1 - f)$
 - Error induced by this approximation: far below numerical accuracy of Monte Carlo program
 - PV formalism: $f \sim 15\%$
 - DD formalism: $f \sim 0.1\%$
- Pentagons under control using DD formalism!

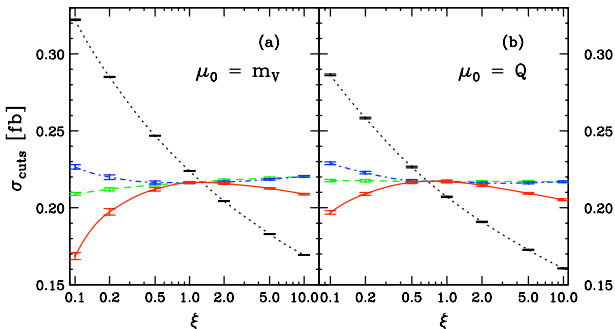
Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

VBF cuts

Tagging Jets	$p_{Tj} \geq 20 \text{ GeV}, \quad y_j \leq 4.5$ $\Delta y_{jj} = y_{j_1} - y_{j_2} > 4,$ $y_{j_1} \cdot y_{j_2} < 0$
Charged Leptons	$p_{Tl} > 20 \text{ GeV}, \quad \eta_l \leq 2.5$ $y_{j,min} < \eta_l < y_{j,max}$ $\Delta R_{jl} \geq 0.4$
Higgs on/off	$M_{VV} > M_H + 10 \text{ GeV}$ (WW,ZZ continuum only)

Scale dependence - total σ (WZ case)



- Two possible scales

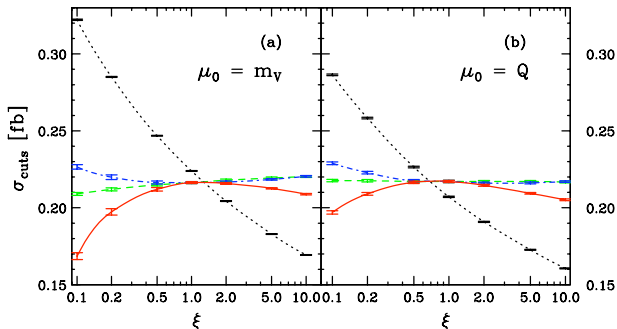
- $m_V = (m_Z + m_W)/2$
- $Q = \text{momentum-transfer of exchanged vector boson in VBF graphs}$

- K-factor ~ 1 (\pm few percent) in both cases

- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)

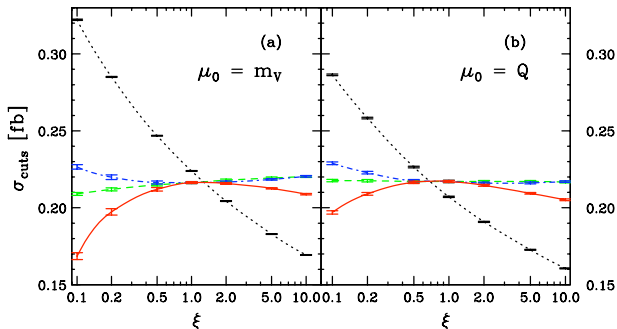
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

Scale dependence - total σ (WZ case)



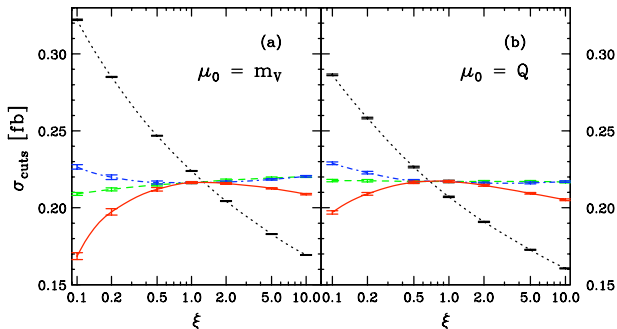
- Two possible scales
 - $m_V = (m_Z + m_W)/2$
 - $Q = \text{momentum-transfer of exchanged vector boson in VBF graphs}$
- K-factor ~ 1 (\pm few percent) in both cases
- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

Scale dependence - total σ (WZ case)



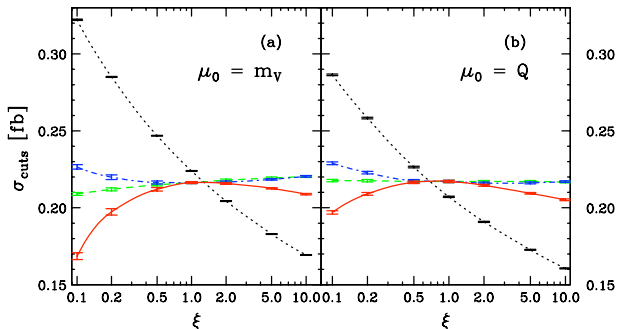
- Two possible scales
 - $m_V = (m_Z + m_W)/2$
 - Q =momentum-transfer of exchanged vector boson in VBF graphs
- K-factor ~ 1 (\pm few percent) in both cases
- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

Scale dependence - total σ (WZ case)



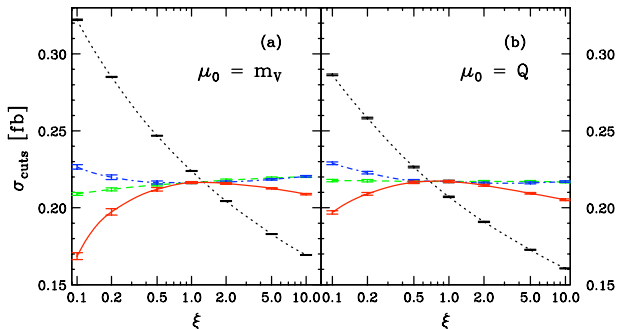
- Two possible scales
 - $m_V = (m_Z + m_W)/2$
 - Q =momentum-transfer of exchanged vector boson in VBF graphs
- K-factor ~ 1 (\pm few percent) in both cases
- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

Scale dependence - total σ (WZ case)



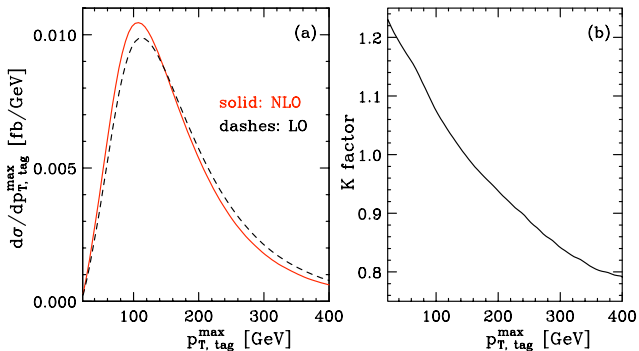
- Two possible scales
 - $m_V = (m_Z + m_W)/2$
 - Q =momentum-transfer of exchanged vector boson in VBF graphs
- K-factor ~ 1 (\pm few percent) in both cases
- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

Scale dependence - total σ (WZ case)



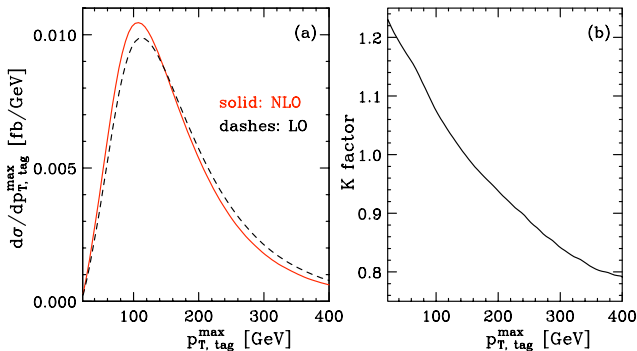
- Two possible scales
 - $m_V = (m_Z + m_W)/2$
 - $Q = \text{momentum-transfer of exchanged vector boson in VBF graphs}$
- K-factor ~ 1 (\pm few percent) in both cases
- LO depends on μ_F only $\rightarrow \sim 10\%$ dependence ($0.5 < \xi < 2$)
- NLO improvement $\rightarrow \sim 2\%$ dependence ($0.5 < \xi < 2$)

K-factor - $p_{T,jet}^{max}$ (WW case)



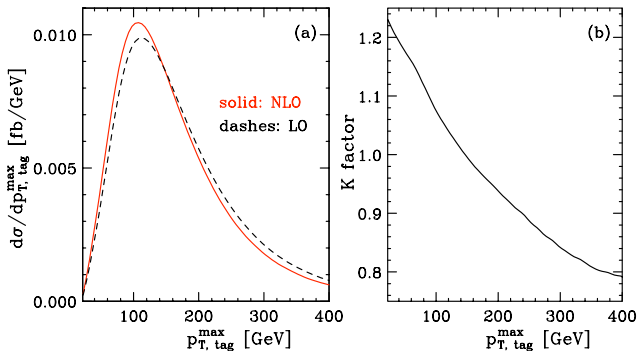
- Strong change in shape \rightarrow shift to smaller p_T at NLO
- Mainly due to extra parton from real emission
- K-factor varying between 1.2 and 0.8 ($20 \text{ GeV} < p_T < 400 \text{ GeV}$)

K-factor - $p_{T,jet}^{max}$ (WW case)



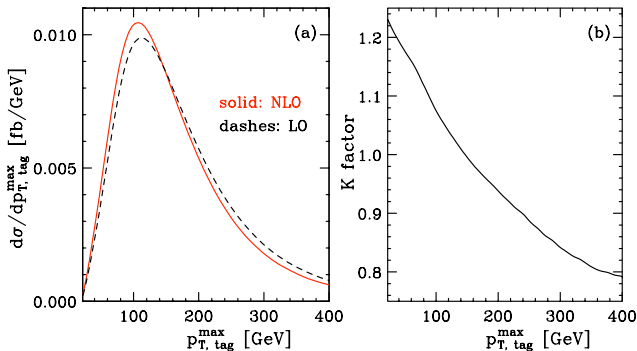
- Strong change in shape → **shift to smaller p_T at NLO**
- Mainly due to extra parton from real emission
- **K-factor** varying between **1.2** and **0.8** ($20 \text{ GeV} < p_T < 400 \text{ GeV}$)

K-factor - $p_{T,jet}^{max}$ (WW case)



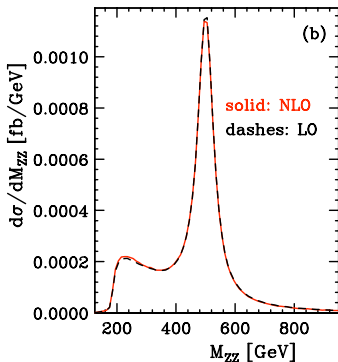
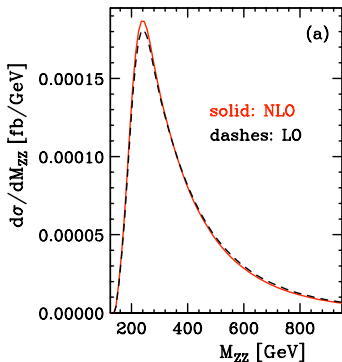
- Strong change in shape → **shift to smaller p_T at NLO**
- Mainly due to extra parton from real emission
- **K-factor** varying between **1.2** and **0.8** ($20 \text{ GeV} < p_T < 400 \text{ GeV}$)

K-factor - $p_{T,jet}^{max}$ (WW case)



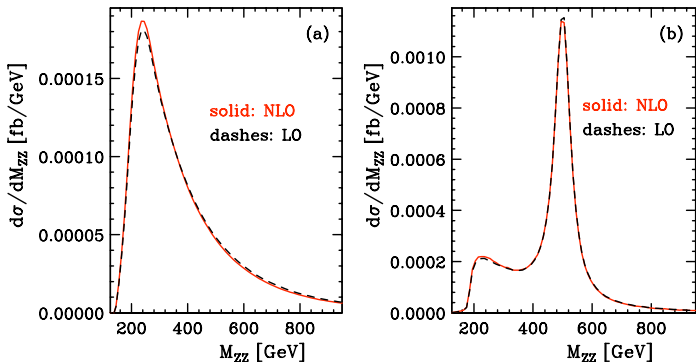
- Strong change in shape \rightarrow **shift to smaller p_T at NLO**
- Mainly due to extra parton from real emission
- **K-factor** varying between **1.2** and **0.8** ($20 \text{ GeV} < p_T < 400 \text{ GeV}$)

Invariant mass - lepton pairs (ZZ case)



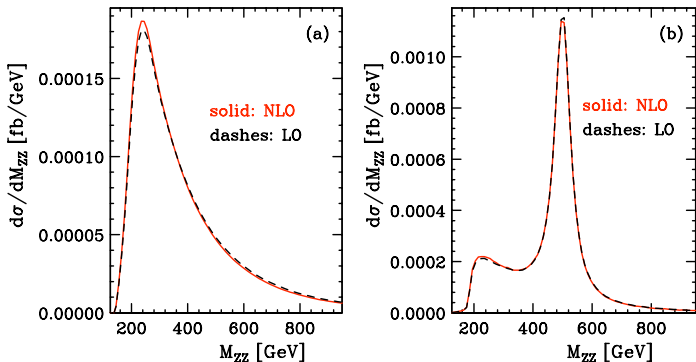
- Continuum ZZ (left) vs. Higgs contribution (right): $\mu_0 = M_Z$
- Pronounced resonance behaviour for $M_H < 800$ GeV
- LO and NLO virtually indistinguishable → excellent stability!

Invariant mass - lepton pairs (ZZ case)



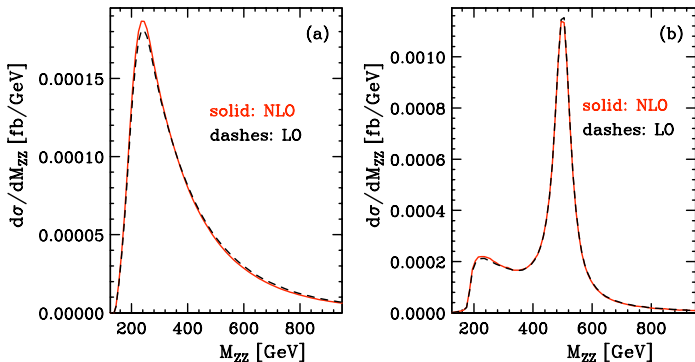
- Continuum ZZ (left) vs. Higgs contribution (right): $\mu_0 = M_Z$
- Pronounced resonance behaviour for $M_H < 800$ GeV
- LO and NLO virtually indistinguishable \rightarrow excellent stability!

Invariant mass - lepton pairs (ZZ case)



- Continuum ZZ (left) vs. Higgs contribution (right): $\mu_0 = M_Z$
- **Pronounced resonance behaviour** for $M_H < 800$ GeV
- LO and NLO virtually indistinguishable \rightarrow **excellent stability!**

Invariant mass - lepton pairs (ZZ case)



- Continuum ZZ (left) vs. Higgs contribution (right): $\mu_0 = M_Z$
- **Pronounced resonance behaviour** for $M_H < 800$ GeV
- LO and NLO virtually indistinguishable \rightarrow **excellent stability!**

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

- Modular structure: separate (one-time) computation of leptonic tensors and decay width
- Fast evaluation and good numerical stability
- NLO corrections under excellent control (modest K-factors and scale dependence)
- *Outlook: VBFNLO*
Monte Carlo program of relevant VBF processes at NLO QCD

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

- Modular structure: separate (one-time) computation of leptonic tensors and decay width
- Fast evaluation and good numerical stability
- NLO corrections under excellent control (modest K-factors and scale dependence)
- *Outlook: VBFNLO*
Monte Carlo program of relevant VBF processes at NLO QCD

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

- Modular structure: separate (one-time) computation of leptonic tensors and decay width
- Fast evaluation and good numerical stability
- NLO corrections under excellent control (modest K-factors and scale dependence)
- *Outlook: VBFNLO*
Monte Carlo program of relevant VBF processes at NLO QCD

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

- Modular structure: separate (one-time) computation of leptonic tensors and decay width
- Fast evaluation and good numerical stability
- NLO corrections under excellent control (modest K-factors and scale dependence)

- Outlook: VBFNLO*

Monte Carlo program of relevant VBF processes at NLO QCD

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

- Modular structure: separate (one-time) computation of leptonic tensors and decay width
- Fast evaluation and good numerical stability
- NLO corrections under excellent control (modest K-factors and scale dependence)
- Outlook: **VBFNLO**
Monte Carlo program of relevant VBF processes at NLO QCD*