



# Lepton Flavour Violation in the Littlest Higgs Model with T-Parity (LHT)

1. Some Notions about Little Higgs Models
2. Lepton Flavour Violating Decays in LHT

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# A brief theoretical introduction...

## The Little Hierarchy Problem

“New Physics (NP) at 1 TeV is expected but its effects are not observed”

From the **instability**  
of the **Higgs mass**

Parameterizing NP by higher-dimensional  
operators suppressed by the NP scale  $\Lambda$ :  
 $(h^\dagger D_\mu h)^2/\Lambda^2, (D^2 h^\dagger D^2 h)/\Lambda^2, \dots$

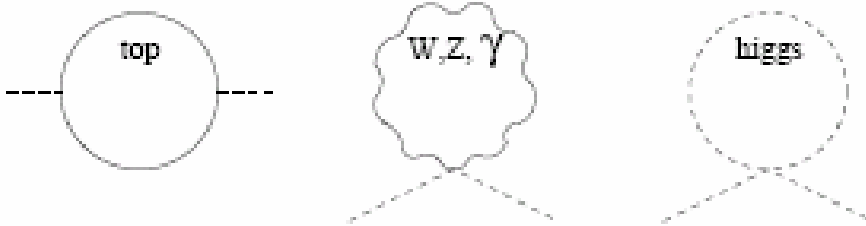
Ew precision tests yield  $\Lambda \geq 5-10\text{TeV}$

**Little Higgs Models can stabilize the Higgs mass  
without violating this bound!**

# SUSY vs Little Higgs



Problematic quadratic divergences in  $m_H^2$



	SUSY	Little Higgs
Quadratic divergences canceled by:	(different statistics) super-partners	(same statistics) heavy partners
Coupling relationships due to:	boson-fermion symmetry	global symmetry

## More formally, in Little Higgs Models:

[N. Arkani-Hamed, A.G. Cohen, H. Georgi (2001)]

1. The **Higgs** is **light** as it is the **Goldstone boson** of a spontaneously broken **global symmetry (G)**
2. **Gauge and Yukawa couplings** of the Higgs are introduced by **gauging a subgroup of G**
3. ``**Dangerous**`` **quadratic corrections are avoided at one-loop** through **Collective Symmetry Breaking**  
(the Higgs becomes massive only when two couplings are non-vanishing)

# The most economical in matter content: Littlest Higgs (LH)

[N. Arkani-Hamed, A.G. Cohen, E. Katz, A.E. Nelson (2002)]

**Global Spontaneous SB:**  $SU(5) \xrightarrow{f \approx O(1\text{TeV})} SO(5)$

**Gauging:**  $[SU(2) \otimes U(1)]_1 \otimes [SU(2) \otimes U(1)]_2 \xrightarrow{f} SU(2)_L \otimes U(1)_Y$   
 $(g_1) \quad (g'_1) \quad (g_2) \quad (g'_2)$

**Collective SB:**  $\delta m_H^2 \propto g_1^{(\prime)2} g_2^{(\prime)2}$

UV-cutoff  $\Lambda = (4\pi f)$

**New Heavy Particles** (with  $O(f)$  masses)

**Gauge Bosons:**  $W^\pm_H, Z^0_H, A^0_H$

**Fermions:** T

**Scalars:**  $\Phi(\text{triplet})$

## Electroweak (ew) precision tests

Tree-level heavy gauge boson contributions and the triplet  $\Phi$  vev make ew precision tests highly constraining

[Han, Logan, McElrath, Wang (2003)]  
[Csaki, Hubisz, Kribs, Meade, Terning (2003)]

$$f \geq 2-3 \text{ TeV}$$

The little hierarchy problem is back!

The solution comes from a discrete symmetry:

T-Parity [H.C. Cheng, I. Low (2003)]

Symmetry under  $[\text{SU}(2) \otimes \text{U}(1)]_1 \longleftrightarrow [\text{SU}(2) \otimes \text{U}(1)]_2$   
 $\longleftrightarrow$   $g_1 = g_2$   $g'_1 = g'_2$

T-parity forbids the unwanted contributions:

- SM particles are T-even,
- new particles are T-odd  
(similarly to R-parity in SUSY)

smaller  $f$  allowed by ew tests  
[Hubisz, Meade, Noble, Perelstein (2005)]

$$f \geq 500 \text{ GeV}$$

# The Littlest Higgs Model with T-Parity (LHT)

**T-even Sector:**  
SM Particles +  $T_+$

**T-odd Sector:**  
Gauge Bosons:  $W_H^\pm, Z_H^0, A_H^0$   
Fermions:  $T_-,$  **Mirror Fermions ( $f_H$ )**  
Scalars:  $\Phi$

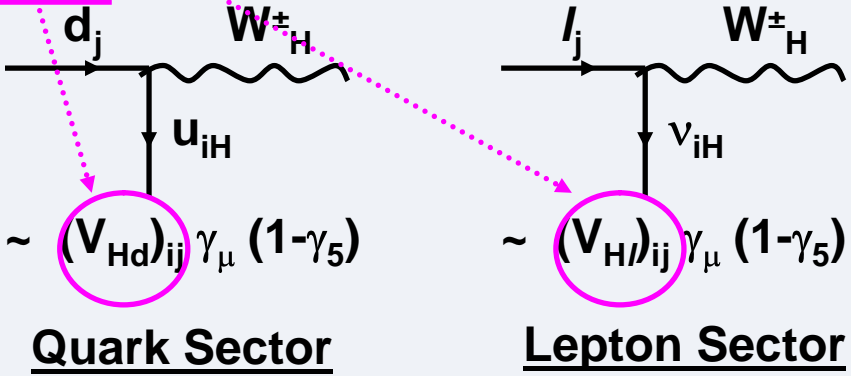
Dark Matter candidate

**New mixing matrices**  
in addition to  
 $V_{CKM}$  and  $V_{PMNS}$

**with NEW flavour interactions**

**New parameters of the LHT model:**

- Global symmetry breaking scale:  $f \sim 1\text{TeV}$
- 1 parameter describing the  $T_+$ :  $x_L$
- $V_{Hd}$  and  $V_{Hl}$ : 3 angles and 3 phases each
- 3 generations of mirror quarks and leptons (6 masses)



**LHT goes beyond**  
**Minimal Flavour Violation (MFV)**  
“visible effects in flavour physics are possible”

**Flavour Physics  
Analyses in LHT**

J.Hubisz et al., hep-ph/0512169

← Quark sector

A.Goyal, hep-ph/0609095

← Lepton sector

S.R. Choudhury et al., hep-ph/0612327

**LHT Team at the Technical University Munich**



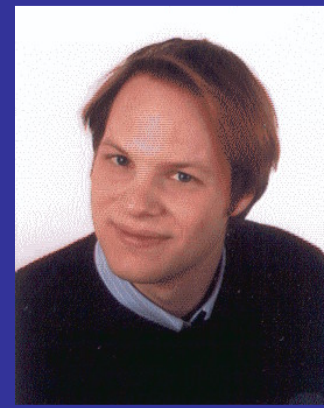
M.Blanke



A.J Buras



B.Duling



A.Poschenrieder

Quark sector

hep-ph/0605214

hep-ph/0609284

hep-ph/0610298

hep-ph/0703254

0704.3329[hep-ph]

Lepton sector

hep-ph/0702136



S.Recksiegel



CT



S.Uhlig



A.Weiler



# Lepton Flavour Violating (LFV) decays

LFV decays are **strongly suppressed** in the **SM**, due to **tiny neutrino masses**

E.g.

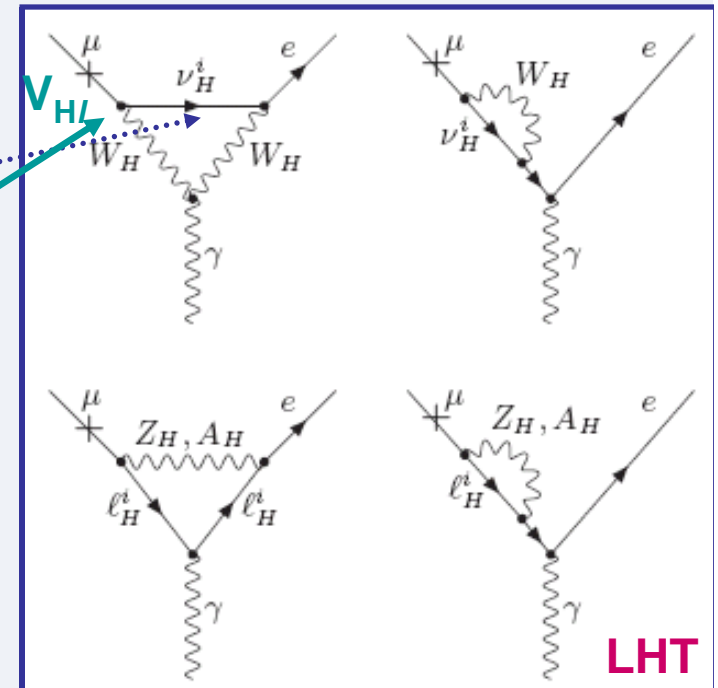
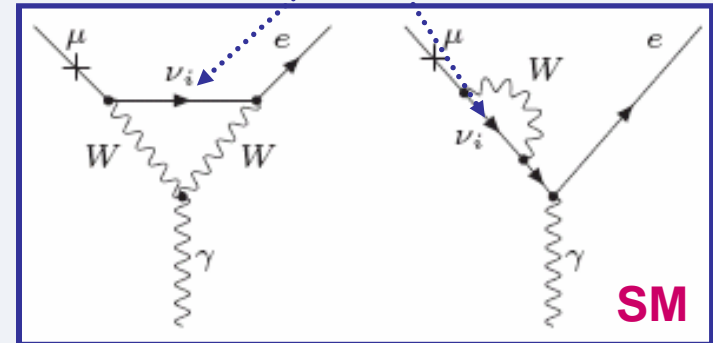
$$Br(\mu \rightarrow e \gamma)_{SM} < 10^{-54}$$

$$10^{-11} \overset{\wedge}{-} 10^{-13}$$

present (MEGA) – near future (MEG)  
experimental upper bounds

In the **LHT** model:

- Mirror leptons have masses of **O(1 TeV)**
- New flavour violating interactions appear



**Spectacular departures** from the SM are possible  
(e.g. of **45 orders of magnitude** in  $\mu \rightarrow e \gamma$ )

In hep-ph/0702136, we have calculated in LHT the LFV decays:

$$\begin{aligned} \mu &\rightarrow e\gamma \\ \tau &\rightarrow \mu\gamma \\ \tau &\rightarrow e\gamma \end{aligned}$$

$$\begin{aligned} \mu^- &\rightarrow e^- e^+ e^- \\ \tau^- &\rightarrow \mu^- \mu^+ \mu^- \\ \tau^- &\rightarrow e^- e^+ e^- \end{aligned}$$

$$\begin{aligned} K_{L,S} &\rightarrow \mu e & \Delta L=1 \\ & & \Delta S=1 \\ B_{d,s} &\rightarrow \mu e & (\Delta B=1) \\ B_{d,s} &\rightarrow \tau e \\ B_{d,s} &\rightarrow \tau \mu \end{aligned}$$

$$K_{L,S} \rightarrow \pi^0 \mu e$$

$$\begin{aligned} \tau &\rightarrow \mu\pi \\ \tau &\rightarrow e\pi \\ \tau &\rightarrow \mu\eta \\ \tau &\rightarrow e\eta \\ \tau &\rightarrow \mu\eta' \\ \tau &\rightarrow e\eta' \end{aligned}$$

$\Delta L=2$

$$\begin{aligned} \tau^- &\rightarrow e^- \mu^+ e^- \\ \tau^- &\rightarrow \mu^- e^+ \mu^- \end{aligned}$$

$(\Delta L=1, \Delta L=2)$

$$\begin{aligned} \tau^- &\rightarrow \mu^- e^+ e^- \\ \tau^- &\rightarrow e^- \mu^+ \mu^- \end{aligned}$$

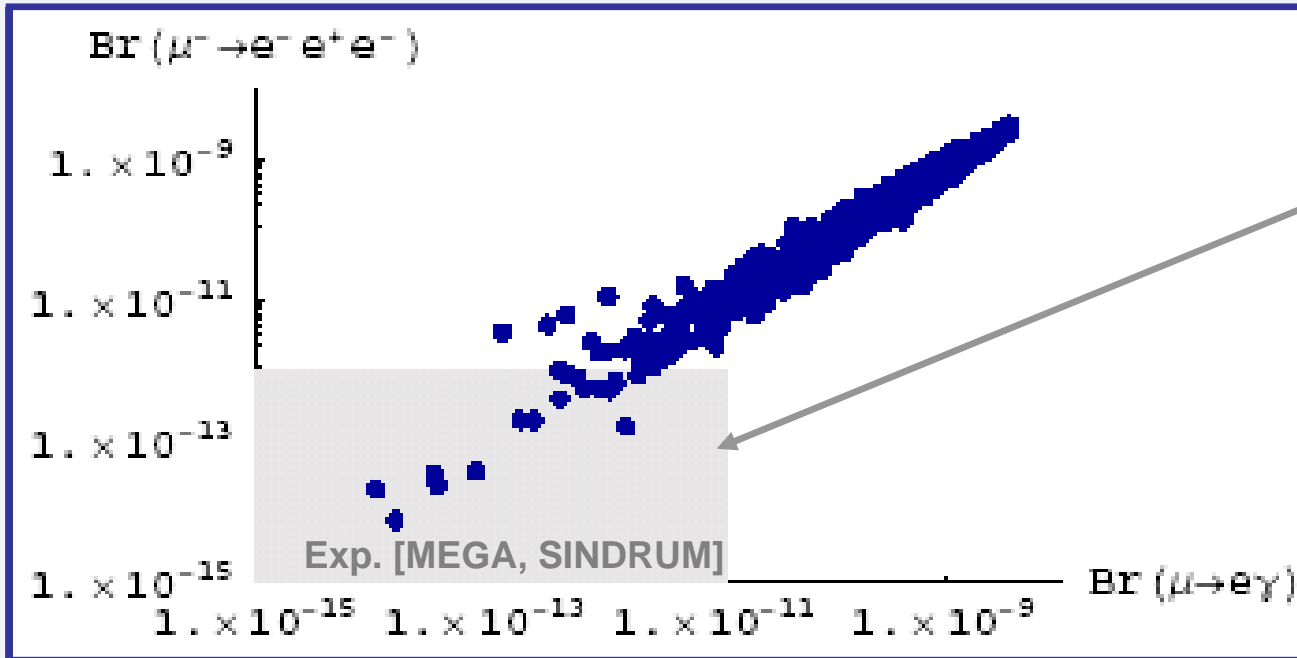
$$\mu \text{Ti} \rightarrow e \text{Ti}$$

AND  $(g-2)_\mu$

(LHT effects are found to be a factor 5 below the experimental uncertainty)

$\mu^- \rightarrow e^- e^+ e^-$  vs.  $\mu^- \rightarrow e^- \gamma$

- $f = 1 \text{ TeV}$  (or  $500 \text{ GeV}$ )
- $300 \text{ GeV} \leq m'_{H_i} \leq 1.5 \text{ TeV}$
- general scan over  $V_{H_i}$



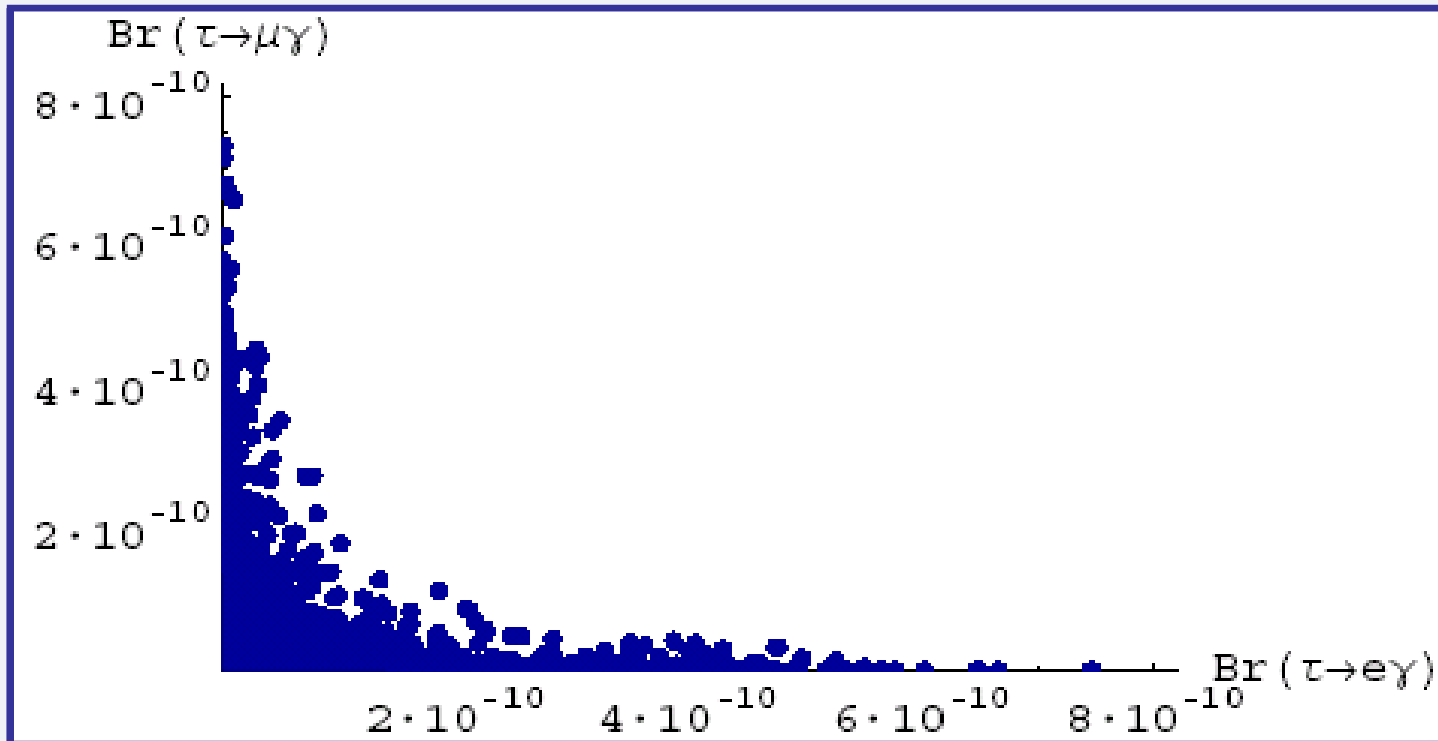
The **exp. constraints** rule out a large region of parameter space

We impose them in the rest of the analysis

To satisfy the experimental constraints:

- $V_{H_i}$  must have a **strong hierarchy** (different from  $V_{\text{PMNS}}^\dagger$  and  $V_{\text{CKM}}$ )  
or
- **Mirror leptons** have to be **quasi-degenerate**

$\tau \rightarrow \mu \gamma$  vs.  $\tau \rightarrow e \gamma$











- $\tau \rightarrow \mu \gamma$  and  $\tau \rightarrow e \gamma$  can be individually as high as  $8 \cdot 10^{-10}$  (close to the planned sensitivity at a Super Flavour Factory)
- Larger values of  $\tau \rightarrow \mu \gamma$  correspond to smaller values of  $\tau \rightarrow e \gamma$  (from the  $V_{Hl}$  structure)


A similar pattern is found in  $\tau \rightarrow \mu \pi(\eta)$  vs.  $\tau \rightarrow e \pi(\eta)$ , that can be individually as high as  $2 \cdot 10^{-9}$  ( $6 \cdot 10^{-10}$ ) [one (two) order of magnitude below present exp. bounds: Belle+BaBar]

S. Banerjee,  
hep-ex/0702017

Belle+BaBar

**LHT vs. Exp.  
upper bounds**

decay	$f = 1000 \text{ GeV}$	$f = 500 \text{ GeV}$	exp. upper bound
 $\mu \rightarrow e\gamma$	$1.2 \cdot 10^{-11}$ ( $1 \cdot 10^{-11}$ )	$1.2 \cdot 10^{-11}$ ( $1 \cdot 10^{-11}$ )	$1.2 \cdot 10^{-11}$ [17]
 $\mu^- \rightarrow e^- e^+ e^-$	$1.0 \cdot 10^{-12}$ ( $1 \cdot 10^{-12}$ )	$1.0 \cdot 10^{-12}$ ( $1 \cdot 10^{-12}$ )	$1.0 \cdot 10^{-12}$ [42]
 $\mu\text{Ti} \rightarrow e\text{Ti}$	$2 \cdot 10^{-10}$ ( $5 \cdot 10^{-12}$ )	$4 \cdot 10^{-11}$ ( $5 \cdot 10^{-12}$ )	$4.3 \cdot 10^{-12}$ [29]
 $\tau \rightarrow e\gamma$	$8 \cdot 10^{-10}$ ( $7 \cdot 10^{-10}$ )	$1 \cdot 10^{-8}$ ( $1 \cdot 10^{-8}$ )	$9.4 \cdot 10^{-8}$ [33]
 $\tau \rightarrow \mu\gamma$	$8 \cdot 10^{-10}$ ( $8 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $1 \cdot 10^{-8}$ )	$1.6 \cdot 10^{-8}$ [33]
$\tau^- \rightarrow e^- e^+ e^-$	$7 \cdot 10^{-10}$ ( $6 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $2 \cdot 10^{-8}$ )	$2.0 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow \mu^- \mu^+ \mu^-$	$7 \cdot 10^{-10}$ ( $6 \cdot 10^{-10}$ )	$3 \cdot 10^{-8}$ ( $3 \cdot 10^{-8}$ )	$1.9 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow e^- \mu^+ \mu^-$	$5 \cdot 10^{-10}$ ( $5 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $2 \cdot 10^{-8}$ )	$2.0 \cdot 10^{-7}$ [72]
$\tau^- \rightarrow \mu^- e^+ e^-$	$5 \cdot 10^{-10}$ ( $5 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $2 \cdot 10^{-8}$ )	$1.9 \cdot 10^{-7}$ [72]
$\tau^- \rightarrow \mu^- e^+ \mu^-$	$5 \cdot 10^{-14}$ ( $3 \cdot 10^{-14}$ )	$2 \cdot 10^{-14}$ ( $2 \cdot 10^{-14}$ )	$1.3 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow e^- \mu^+ e^-$	$5 \cdot 10^{-14}$ ( $3 \cdot 10^{-14}$ )	$2 \cdot 10^{-14}$ ( $2 \cdot 10^{-14}$ )	$1.1 \cdot 10^{-7}$ [71]
 $\tau \rightarrow \mu\pi$	$2 \cdot 10^{-9}$ ( $2 \cdot 10^{-9}$ )	$5.8 \cdot 10^{-8}$ ( $5.8 \cdot 10^{-8}$ )	$5.8 \cdot 10^{-8}$ [33]
 $\tau \rightarrow e\pi$	$2 \cdot 10^{-9}$ ( $2 \cdot 10^{-9}$ )	$4.4 \cdot 10^{-8}$ ( $4.4 \cdot 10^{-8}$ )	$4.4 \cdot 10^{-8}$ [33]
$\tau \rightarrow \mu\eta$	$6 \cdot 10^{-10}$ ( $6 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $2 \cdot 10^{-8}$ )	$5.1 \cdot 10^{-8}$ [33]
$\tau \rightarrow e\eta$	$6 \cdot 10^{-10}$ ( $6 \cdot 10^{-10}$ )	$2 \cdot 10^{-8}$ ( $2 \cdot 10^{-8}$ )	$4.5 \cdot 10^{-8}$ [33]
$\tau \rightarrow \mu\eta'$	$7 \cdot 10^{-10}$ ( $7 \cdot 10^{-10}$ )	$3 \cdot 10^{-8}$ ( $3 \cdot 10^{-8}$ )	$5.3 \cdot 10^{-8}$ [33]
$\tau \rightarrow e\eta'$	$7 \cdot 10^{-10}$ ( $7 \cdot 10^{-10}$ )	$3 \cdot 10^{-8}$ ( $3 \cdot 10^{-8}$ )	$9.0 \cdot 10^{-8}$ [33]
 $K_L \rightarrow \mu e$	$4 \cdot 10^{-13}$ ( $2 \cdot 10^{-13}$ )	$3 \cdot 10^{-14}$ ( $3 \cdot 10^{-14}$ )	$4.7 \cdot 10^{-12}$ [50]
$K_L \rightarrow \pi^0 \mu e$	$4 \cdot 10^{-15}$ ( $2 \cdot 10^{-15}$ )	$5 \cdot 10^{-16}$ ( $5 \cdot 10^{-16}$ )	$6.2 \cdot 10^{-9}$ [73]
$B_d \rightarrow \mu e$	$5 \cdot 10^{-16}$ ( $2 \cdot 10^{-16}$ )	$9 \cdot 10^{-17}$ ( $9 \cdot 10^{-17}$ )	$1.7 \cdot 10^{-7}$ [74]
$B_s \rightarrow \mu e$	$5 \cdot 10^{-15}$ ( $2 \cdot 10^{-15}$ )	$9 \cdot 10^{-16}$ ( $9 \cdot 10^{-16}$ )	$6.1 \cdot 10^{-6}$ [75]
$B_d \rightarrow \tau e$	$3 \cdot 10^{-11}$ ( $2 \cdot 10^{-11}$ )	$2 \cdot 10^{-10}$ ( $2 \cdot 10^{-10}$ )	$1.1 \cdot 10^{-4}$ [76]
$B_s \rightarrow \tau e$	$2 \cdot 10^{-10}$ ( $2 \cdot 10^{-10}$ )	$2 \cdot 10^{-9}$ ( $2 \cdot 10^{-9}$ )	—
$B_d \rightarrow \tau \mu$	$3 \cdot 10^{-11}$ ( $3 \cdot 10^{-11}$ )	$3 \cdot 10^{-10}$ ( $3 \cdot 10^{-10}$ )	$3.8 \cdot 10^{-5}$ [76]
$B_s \rightarrow \tau \mu$	$2 \cdot 10^{-10}$ ( $2 \cdot 10^{-10}$ )	$3 \cdot 10^{-9}$ ( $3 \cdot 10^{-9}$ )	—

 = LHT effects  
could be seen  
in the near future

# Distinction between LHT and MSSM from LFV correlations

☺ =  
LHT and MSSM  
can be clearly  
distinguished



Correlations between BR's  
• are less parameter-dependent  
• can provide a clear model signal

J.R. Ellis et al., hep-ph/0206110  
A. Brignole and A. Rossi, hep-ph/0404211  
E. Arganda and M.J. Herrero, hep-ph/0510405  
P. Paradisi, hep-ph/0505046,0508054,0601110

ratio	LHT	MSSM (dipole)	MSSM (Higgs)
$\frac{Br(\mu^- \rightarrow e^- e^+ e^-)}{Br(\mu \rightarrow e \gamma)}$ ☺	0.4... 2.5	$\sim 6 \cdot 10^{-3}$	$\sim 6 \cdot 10^{-3}$
$\frac{Br(\tau^- \rightarrow e^- e^+ e^-)}{Br(\tau \rightarrow e \gamma)}$ ☺	0.4... 2.3	$\sim 1 \cdot 10^{-2}$	$\sim 1 \cdot 10^{-2}$
$\frac{Br(\tau^- \rightarrow \mu^- \mu^+ \mu^-)}{Br(\tau \rightarrow \mu \gamma)}$ ☺	0.4... 2.3	$\sim 2 \cdot 10^{-3}$	0.06... 0.1
$\frac{Br(\tau^- \rightarrow e^- \mu^+ \mu^-)}{Br(\tau \rightarrow e \gamma)}$ ☺	0.3... 1.6	$\sim 2 \cdot 10^{-3}$	0.02... 0.04
$\frac{Br(\tau^- \rightarrow \mu^- e^+ e^-)}{Br(\tau \rightarrow \mu \gamma)}$ ☺	0.3... 1.6	$\sim 1 \cdot 10^{-2}$	$\sim 1 \cdot 10^{-2}$
$\frac{Br(\tau^- \rightarrow e^- e^+ e^-)}{Br(\tau^- \rightarrow e^- \mu^+ \mu^-)}$	1.3... 1.7	$\sim 5$	0.3... 0.5
$\frac{Br(\tau^- \rightarrow \mu^- \mu^+ \mu^-)}{Br(\tau^- \rightarrow \mu^- e^+ e^-)}$	1.2... 1.6	$\sim 0.2$	5... 10
$\frac{R(\mu Ti \rightarrow e Ti)}{Br(\mu \rightarrow e \gamma)}$	$10^{-2} \dots 10^2$	$\sim 5 \cdot 10^{-3}$	0.08... 0.15

The MSSM is dominated by the dipole operator, while the LHT by box- and Z-penguin diagrams

☺ & the double ratios ( $\mu \leftrightarrow e$ ):

$$R_1 = \frac{Br(\tau^- \rightarrow e^- e^+ e^-) Br(\tau^- \rightarrow \mu^- e^+ e^-)}{Br(\tau^- \rightarrow \mu^- \mu^+ \mu^-) Br(\tau^- \rightarrow e^- \mu^+ \mu^-)},$$

$$R_2 = \frac{Br(\tau^- \rightarrow e^- e^+ e^-) Br(\tau \rightarrow \mu \gamma)}{Br(\tau^- \rightarrow \mu^- \mu^+ \mu^-) Br(\tau \rightarrow e \gamma)},$$

$$R_3 = \frac{Br(\tau^- \rightarrow e^- \mu^+ \mu^-) Br(\tau \rightarrow \mu \gamma)}{Br(\tau^- \rightarrow \mu^- e^+ e^-) Br(\tau \rightarrow e \gamma)}.$$

$0.8 \leq R_{1,2,3} \leq 1.3$  [LHT]  
 $R_1 \approx 20, R_2 \approx 5, R_3 \approx 0.2$  [MSSM]

## Conclusions

### The Littlest Higgs Model with T-parity

- solves the **little hierarchy problem**
- is compatible with **ew precision tests**
- introduces **new flavour violating interactions**
- can yield **large effects in Flavour Physics**
- in particular in **Lepton Flavour Violating decays**  
(strongly suppressed within the SM by tiny  $\nu$  masses)

• Many **LHT upper bounds** for **LFV decays** are **close to present and near future exp. upper bounds !!**

• **Correlations** of  $Br$ 's could provide a **clear distinction** between **LHT** and **MSSM !!**

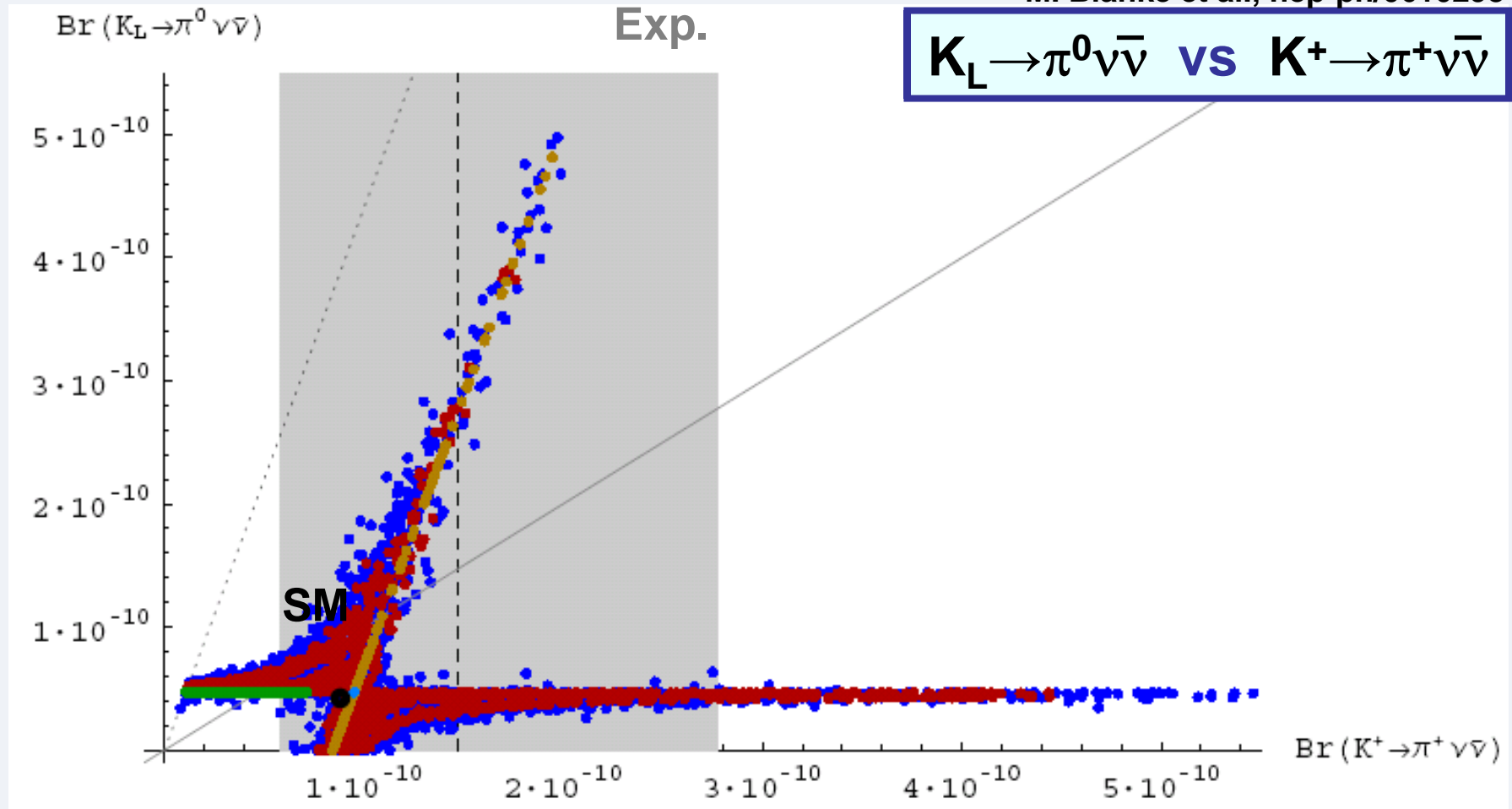
**BACKUP**



LHT effects in  
the **quark sector**  
in 1 slide

The **largest effects** are found in **rare K decays**,  
where the SM contribution is CKM suppressed by  $(V_{ts}^* V_{td})$

M. Blanke et al., hep-ph/0610298



Two distinguished branches appear!  
~10 times enhancement in  $K_L \rightarrow \pi^0 \nu \bar{\nu}$   
~5 times enhancement in  $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

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