

# $V_{ub}$ determination using $B \rightarrow \pi$ form factor from Light-Cone Sum Rules

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## Outline

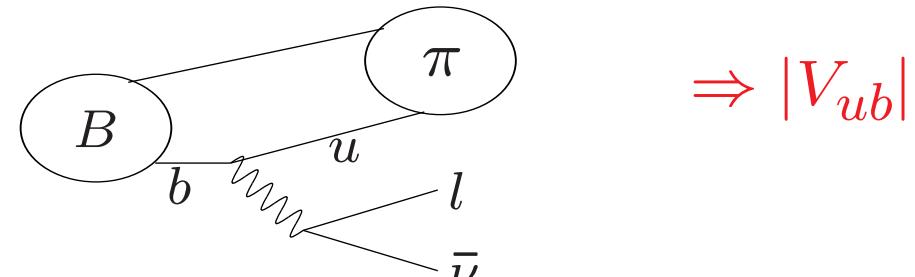
- minireview of exclusive  $|V_{ub}|$  determinations, mainly  $B \rightarrow \pi$
- new (preliminary) results on  $f_{B\pi}^+(q^2)$  from LCSR

[G. Duplancić, A.K., Th. Mannel, B. Melić and N. Offen]

## Exclusive $b \rightarrow u$ transitions sensitive to $|V_{ub}|$

Channel	BR $\times 10^4$	hadronic input	theory
$B^- \rightarrow \tau^- \bar{\nu}_\tau$	$1.79^{+0.56+0.39}_{-0.49-0.46}$ [Belle] $< 1.7$ (90% CL) [BaBar]	$f_B$	Lattice ( $2 \oplus 1$ ) QCD SR
$\bar{B}^0 \rightarrow \pi^+ l^- \bar{\nu}_l$	$1.39 \pm 0.06 \pm 0.06$ [HFAG] $q^2$ -shape [BaBar]	FF $f_{B\pi}^+(q^2)$	Lattice ( $2 \oplus 1$ ), SCET, LCSR
$\bar{B}^0 \rightarrow \rho^+ l^- \bar{\nu}_l$	$2.14 \pm 0.21 \pm 0.51 \pm 0.28$ [BaBar] $2.17 \pm 0.54 \pm 0.31 \pm 0.08$ [Belle]	three FF's	Lattice (quench.) LCSR("quench.")
$B^- \rightarrow l^- \bar{\nu}_l \gamma$	$< 0.05$ (90% CL) [BaBar]	two FF's	QCDF,LCSR
$B^- \rightarrow \pi^- \pi^0$	$0.057 \pm 0.004$ [HFAG]	hadr. ampl.	input: $B \rightarrow \pi$ FF QCDF, SCET

## $B \rightarrow \pi$ form factor (FF)



- Hadronic matrix element:

$$\langle \pi^+(p) | \bar{u} \gamma_\mu b | \bar{B}^0(p+q) \rangle = f_{B\pi}^+(q^2)(2p_\mu + q_\mu) + \dots$$

.... - the second FF  $f^-$  ( $f^0$ ) only in  $B \rightarrow \pi \tau \nu_\tau$

- $B \rightarrow \pi \mu \nu_\mu, \pi e \nu_e$  region:  $0 < q^2 < (m_B - m_\pi)^2 = q_{max}^2 \simeq 26.4 \text{ GeV}^2$
- differential decay distribution (neglecting  $m_l$ )

$$\frac{d\Gamma(\bar{B}^0 \rightarrow \pi^+ l^- \bar{\nu})}{dq^2} = \frac{G_F^2 |V_{ub}|^2}{24\pi^3} (E_\pi^2 - m_\pi^2)^{3/2} |f_{B\pi}^+(q^2)|^2$$

$E_\pi = (m_B^2 + m_\pi^2 - \textcolor{red}{q}^2)/(2m_B)$  - pion energy in the  $B$  rest frame

## Employing the analyticity of $f_{B\pi}^+(q^2)$

- $q^2 \Rightarrow$  complex variable,  $f_{B\pi}^+(q^2) \Rightarrow$  analytic function, singularities (poles, branch points) given by unitarity relation (consult your favorite Quantum Field Theory textbook)
- $f_{B\pi}^+(q^2)$  real at  $q^2 < m_{B^*}^2$ ,  $m_{B^*}^2 > q_{max}^2$ , pole at  $q^2 = m_{B^*}^2$ , branch points (and poles) at  $q^2 > (m_B + m_\pi)^2$  (radial excitations of  $B^*$ )
- dispersion relation (derived from Cauchy theorem in  $q^2$ -plane)

$$f_{B\pi}^+(q^2) = \frac{f_{B^*} g_{B^* B\pi}}{m_{B^*}^2 - q^2} + \frac{1}{\pi} \int_{(m_B+m_\pi)^2}^{\infty} ds \frac{\text{Im} f_{B\pi}^+(s)}{s - q^2}$$

## Employing the analyticity of $f_{B\pi}^+(q^2)$

- analyticity  $\oplus$  perturbative QCD bounds  
(unitarity for correlation function of heavy-light vector currents)  
 $\Rightarrow$  knowledge of  $f_{B\pi}^+(q^2)$  at a few points tightly bounds ( $\simeq$  reproduces) the FF at all  $q^2 < q_{max}^2$
- mathematical framework: [N. Meiman (1963)];...,  
applications to  $B \rightarrow \pi$  FF: [C.G.Boyd, B.Grinstein, R.Lebed (1995)],  
L.Lellouch (1996)],..., [T.Becher, R.Hill (2006)],..
- Use of the relation between  $f_{B\pi}(q^2)$  and the elastic phase of  $\pi B \rightarrow \pi B$  strong scattering amplitude:  
[Muskhelishvili (1950)], [ Omnes (1950's) ], ...  
applications to  $B \rightarrow \pi$  FF: [J. Flynn, J. Nieves (2001)]
- Practical use: FF parameterizations

## Simplest parameterizations

- dispersion relation as a starting point, replace the integral by an effective pole:

$$[f_{B\pi}^+(q^2)]_{disp.\,rel} \Rightarrow \frac{r_1}{1 - q^2/m_{B^*}^2} + \frac{r_2}{1 - q^2/m_{B_{fit}}^2}$$

- BK parameterization [D.Becirevic, A.Kaidalov, 2000]  
(3 parameters  $\rightarrow$  2 , motivated by HQ limit)

$$[f_{B\pi}^+(q^2)]_{BK} = \frac{f_{B\pi}(0)}{(1 - q^2/m_{B^*}^2)(1 - \alpha q^2/m_{B^*}^2)}$$

- a generic 3-parameter two-pole parameterization: [P.Ball, R.Zwicky, 2004]

$$[f_{B\pi}^+(q^2)]_{BZ} = f_{B\pi}(0) \left( \frac{1}{1 - q^2/m_{B^*}^2} + \frac{rq^2/m_{B^*}^2}{(1 - q^2/m_{B^*}^2)(1 - \alpha_{BZ}q^2/m_{B^*}^2)} \right)$$

reduced to BK at  $\alpha_{BZ} = r = \alpha_{BK}$

## Advanced parameterizations

- BGL-parameterization [Boyd, Grinstein, Lebed]:  
a power expansion with  $\sim 4, 5$  parameters

$$[f_{B\pi}^+(q^2)]_{BGL} = \frac{\sum_k a_k [z(q^2, q_0^2)]^k}{P(q^2, m_{B^*}^2) \phi(q^2, q_0^2)}$$

$z < 1$ ,  $P$ ,  $\phi$  -known functions,  $\sum_k a_k^2 \leq 1$  (unitarity), truncated at  $k = 2$

- AFHNV parameterization, based on Omnes-representation, with 4 shape parameters: [Albertus, Flynn, Hernandez, Nieves, Verde-Velasco ( 2005)]

$$f_{B\pi}^+(q_i^2)/f_{B\pi}^+(0) , 0.25q_{max}^2 < q_i^2 < q_{max}^2 \oplus \text{overall normaliz.}$$

Give me four parameters, and I can fit an elephant.  
Give me five, and I can wiggle its trunk”  
(J. von Neumann)

## Fixing $|V_{ub}f_{B\pi}^+(0)|$ from data

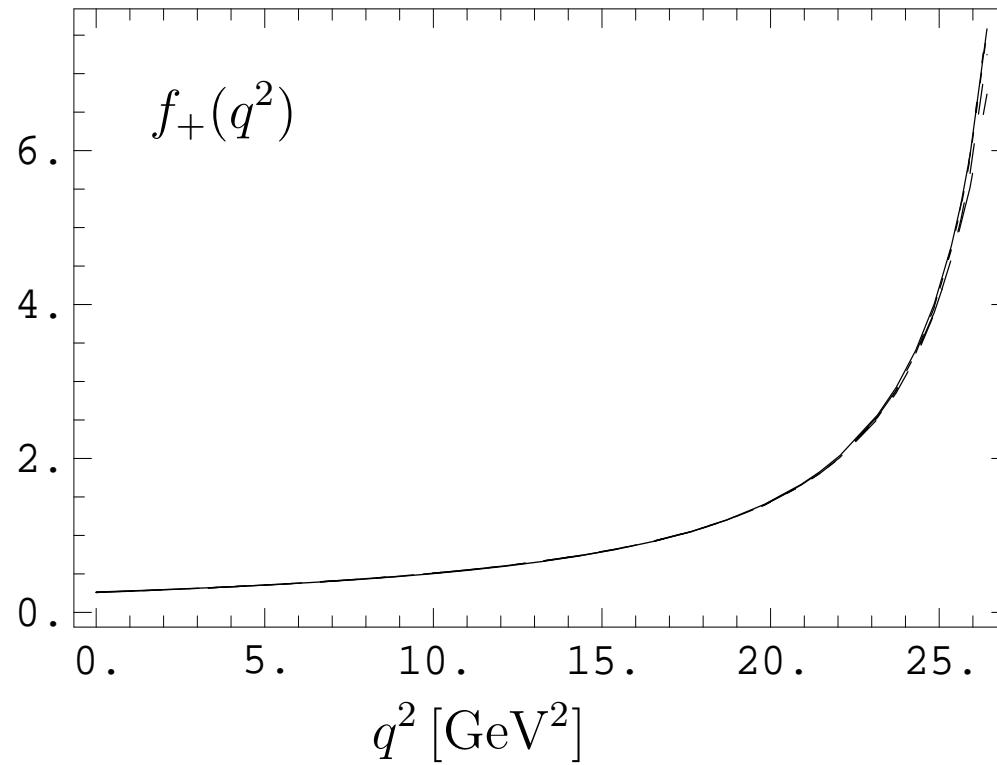
- fit various parameterizations to the exp. differential shape and  $BR_{tot}(B \rightarrow \pi l \nu_l)$  [HFAG]
- using combined CLEO,BaBar,Belle shape measurement and BGL form [T. Becher, R.Hill , hep-ph/0509090]

$$|V_{ub}f_{B\pi}^+(0)| = (0.92 \pm [0.11]_{shape} \pm [0.03]_{BR}) \times 10^{-3}$$

- using recent BR  $q^2$ -bins [BaBar], fitted five diff. parameterizations (BK, BZ, BGLa, BGLb, AFHNV) [P. Ball , hep-ph/0611108]

$$|V_{ub}f_{B\pi}^+(0)| = (0.91 \pm [0.06]_{shape} \pm [0.03]_{BR}) \times 10^{-3}$$

- it remains to calculate  $|f_{B\pi}^+(0)|$  in QCD (or  $f_{B\pi}^+(q^2)$  at any single point  $q^2$ )  
 $\Rightarrow$  extract  $|V_{ub}|$  .



best-fit form factors of 5 different parameterizations,  
normalized with  $|V_{ub}| = (3.50 \pm 0.18) \times 10^{-3}$  (UT fits)  
from P. Ball, talk at FPCP 2007, Bled [hep-ph 07052290]

## Lattice QCD

- $b$ -quark too heavy, pion too light for the lattice ,  
only small  $E_\pi$  - large  $q^2 > 16 \text{ GeV}^2$  accessible
- recent progress: calculation of heavy-light hadronic matrix elements with  
 $n_f = 2 \oplus 1$  dynamical quark flavours :
- Fermilab-MILC , “Fermilab treatment” of heavy quarks,  
[M.Okamoto, hep-lat/050113]
- HPQCD [E.Gulez et. al, hep-lat/0601021 (+errata)], use NRQCD,  

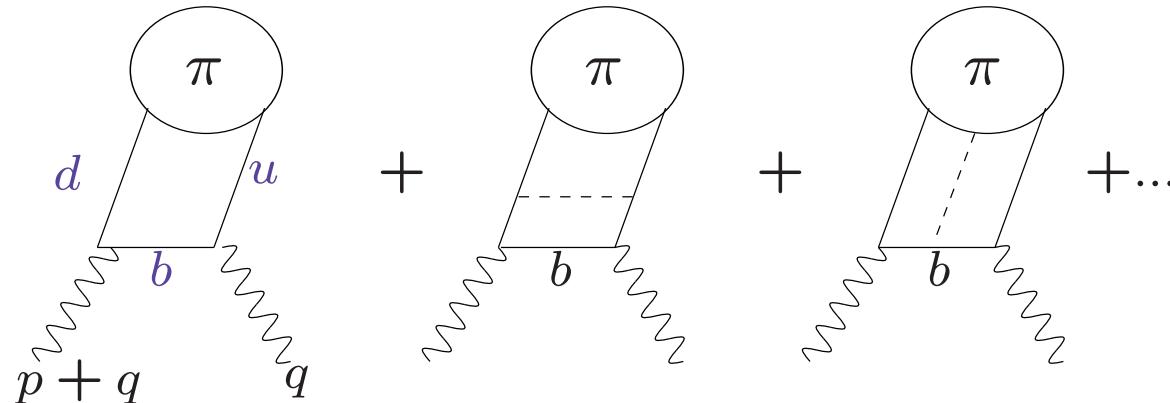
$q^2 [\text{GeV}^2]$	$f_{B\pi}^+(q^2)$
17.34	$1.101 \pm 0.053$
18.39	$1.273 \pm 0.099$
19.45	$1.458 \pm 0.142$
20.51	$1.627 \pm 0.185$
21.56	$1.816 \pm 0.126$

fit to BZ param.

$$\frac{1}{|V_{ub}|^2} \int_{16 \text{ GeV}^2}^{q_{max}^2} dq^2 d\Gamma/dq^2 = 2.07 \pm 0.57 \text{ ps}^{-1}$$

(dominant error from chiral extrapolation)

# Light-Cone Sum Rules



- The correlation function:  $q^2, (p+q)^2 \ll m_b^2$ ,  **$b$ -quark highly virtual**  

$$F_\lambda(q, p) = i \int d^4x e^{iqx} \langle \pi(p) | T\{\bar{u}(x)\gamma_\lambda b(x), \bar{b}(0)i\gamma_5 d(0)\} | 0 \rangle$$
- operator-product-expansion (OPE) **near the light-cone**,  $x^2 \sim 0$
- universal input:  $\langle \pi(p) | \bar{u}(x) \dots d(0) | 0 \rangle \sim \varphi_\pi^{(t)}(u)$   
 -pion distribution amplitudes (DA) of twist  $t = 2, 3, 4..$   
 [P. Ball, R. Zwicky (2001), 2004].  
 [G. Duplancić, A.K., Th. Mannel, B. Melić and N. Offen, paper in preparation]

## Derivation of LCSR

- Hadronic dispersion relation in  $(p+q)^2$ : ( $q^2 \ll m_b^2$  fixed)

$$F(q^2, (p+q)^2) = \begin{array}{c} \text{Diagram: } p+q \rightarrow B \rightarrow \pi + b \\ \text{Diagram: } p+q \rightarrow B_h \rightarrow \pi + b \end{array} + \sum_h \begin{array}{c} \text{Diagram: } p+q \rightarrow B_h \rightarrow \pi + b \\ \text{Diagram: } p+q \rightarrow B_h \rightarrow \pi + b \end{array}$$

$$f_B f_{B\pi}^+(q^2) \quad \sum_{B_h} \rightarrow duality (s_0^B)$$

- matching at  $\langle -(p+q)^2 \rangle \sim \mu^2 \sim m_b \Lambda$  and using duality

$$[F((p+q)^2, q^2)]_{OPE} = \frac{m_B^2 f_B f_{B\pi}^+(q^2)}{m_B^2 - (p+q)^2} + \int_{s_0^B}^{\infty} ds \frac{[\text{Im}F(s, q^2)]_{OPE}}{s - (p+q)^2}$$

- inputs:

- \*  $m_b$ ,  $\alpha_s$ , (full QCD), we use  $\overline{MS}$  mass :  $m_b(m_b) = 4.164 \pm 0.025$  GeV  
[J. Kühn, M. Steinhauser, C. Sturm]

- \*  $\varphi_\pi^{(t)}(u)$ ,  $t = 2, 3, 4$ ,

main sensitivity to  $a_2, a_4$  -parameters of twist 2 DA (Gegenbauer moments)

- \*  $s_0^B$ ,  $M^2$  fixed by calculating  $m_B$ ,

- \* we fix  $a_2, a_4, s_0^B$  by fitting  $q^2$  dep. to the measured slope of FF

- \*  $f_B$  - Determined from two-point sum rule;

$$f_B = 210 \pm 19 \text{ MeV}$$

[M. Jamin, B. Lange (2001)]

- uncertainties:

- \* induced by variation of (universal) input parameters, errors of exp.inputs

- \* parton-hadron duality (suppressed with Borel transformation, constrained by  $q^2$  shape)

- recent LCSR results: (with one-loop pole mass  $m_b$ )

$$f_{B\pi}^+(0) = 0.258 \pm 0.031 \quad [\text{P.Ball, R.Zwicky(2004)}]$$

$$f_{B\pi}^+(0) = 0.26 \pm 0.02_{[a_{2,4}^\pi]} \pm 0.03_{[\text{param}]} \quad (\text{without twist-3 NLO})$$

[A. K., T. Mannel, M. Melcher and B. Melic, PRD (2005), hep-ph/0509049]

- preliminary result of our new LCSR calculation  
with  $\overline{MS}$  mass and tw.-2,3 NLO:

$$f_{B\pi}^+(0) = 0.25 \pm 0.04$$

(combining all individual uncertainties in quadrature)

## Recent $|V_{ub}|$ determinations using $f_{B\pi}^+$

Method	$ V_{ub}  \times 10^3$	use of	ref.
Lattice $n_f = 3$	$3.78 \pm 0.25 \pm 0.52$	BK	Fermilab/MILC '05
Lattice $n_f = 3$	$3.35 \pm 0.25 \pm 0.50$	BZ-par.	HPQCD '07
comb.	$3.54 \pm 0.17 \pm 0.44$	BGL $\oplus$ Latt. $\oplus$ SCET $B \rightarrow \pi\pi$	Arnesen et al. '05
comb.	$3.7 \pm 0.2 \pm 0.1$	BGL $\oplus$ Latt.	Becher,Hill '06
comb.	$3.47 \pm 0.29 \pm 0.03$	Omnes repr. $\oplus$ Latt $\oplus$ LCSR (BZ)	Flynn, Nieves '07
LCSR	$3.5 \pm 0.4 \pm 0.1$	BGL	Ball '06

- using  $|V_{ub}f^+(0)|$  from P. Ball analysis, our *preliminary* result:  
 $|V_{ub}| = (3.64 \pm [0.3]_{exp} \pm [0.58]_{th}) \times 10^{-3}$
- UT fits (only CP- observables):  $|V_{ub}| = (3.50 \pm 0.18) \times 10^{-3}$  [Utfit, CKMfitter]
- $B \rightarrow X_u l \nu_l$  (“selected” analyses):  $|V_{ub}| = 4.10 \pm 0.30(exp) \pm 0.29(th) \times 10^{-3}$   
[M.Neubert, talk at FPCP-2007, Bled]