

Influence of graviton on top-antitop production at the LHC

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Content

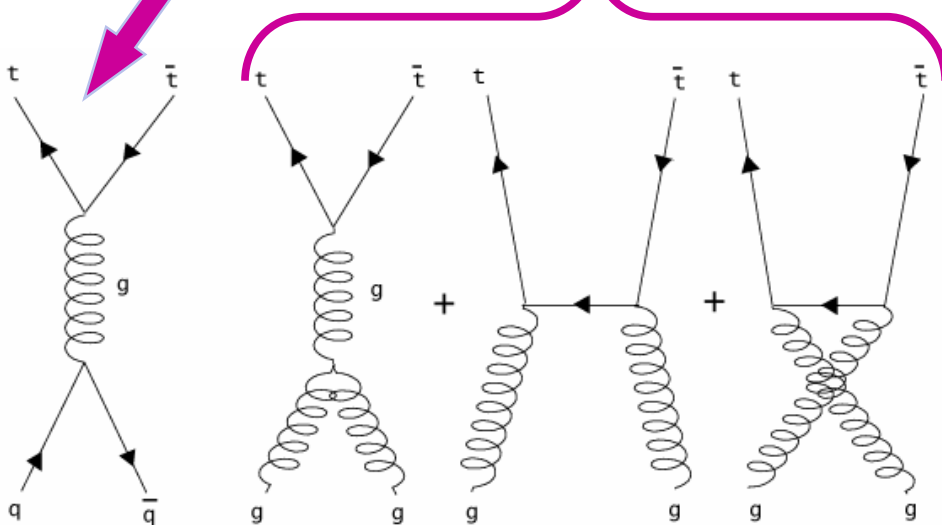
- Basic facts about the top quark
- Spin properties of the top quark (polarization, spin correlation)
- Top production in theories with extra dimensions
 - ✓ ADD model
 - ✓ Randall-Sundrum model
- Top production in the Randall-Sundrum scenario – predictions for the LHC
- Conclusion

Top quark

- The heaviest quark of the Standard Model.
- High mass: 175 GeV (as the atom of gold).

LHC: Proton-proton interactions with CMS energy 14 TeV
 The most of the top quarks produced in top-antitop pairs

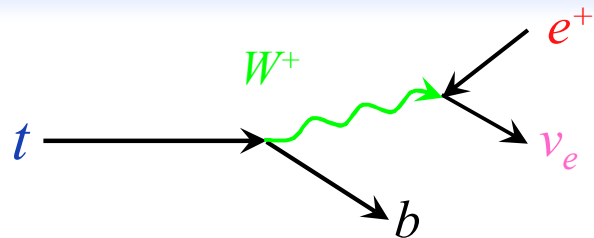
- gluon-gluon fusion (87 %)
 - quark-antiquark annihilation (13 %)
- } $\sim 10 \cdot 10^6$ pairs/year



LEPTONS		
Electron neutrino Mass: 0?	Muon neutrino 0?	Tau neutrino 0?
Electron 511	Muon 105.7	Tau 1,777
QUARKS		
Up Mass: 5	Charm 1,500	Top ~180,000
Down 8	Strange 160	Bottom 4,250

- Observed in 1995 in Fermilab (produced ~ 200 pairs).
- Lifetime: 10^{-24} s -> **does not hadronise**, the angular distribution of decay products is influenced by the spin properties of t quark.
- The only one quark, where **we can study its spin properties**.
- Spin properties of t quarks sensitive to some effects beyond the Standard Model.

Decay of top quark



$$t \rightarrow b + W^+ \quad - 98.8 \%$$

$$W^+ \rightarrow \begin{cases} \bar{d} + u \\ \bar{s} + c \\ \dots \end{cases} \quad \left. \vphantom{W^+} \right\} 67.6 \%$$

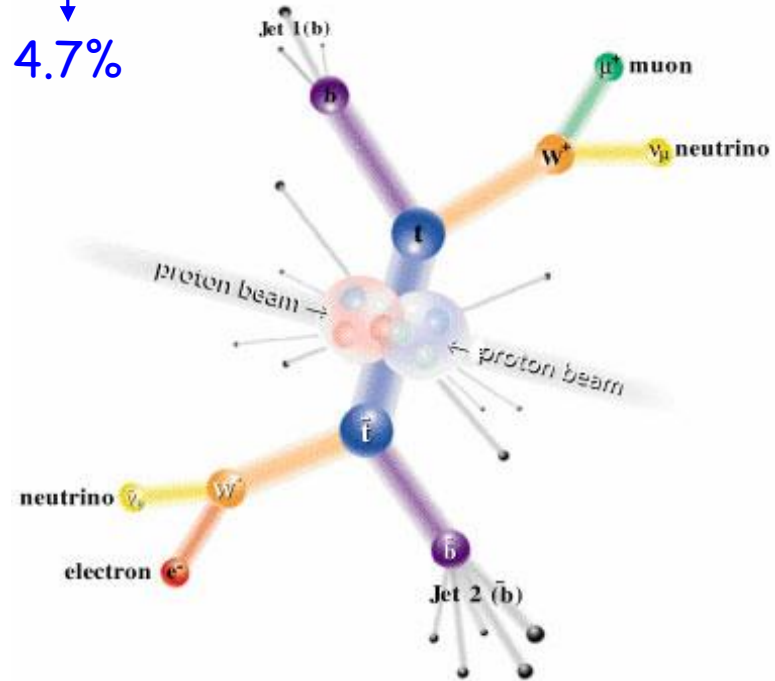
$$\left(\begin{matrix} e^+ + \nu_e \\ \mu^+ + \nu_\mu \end{matrix} \right) \quad \left. \vphantom{\begin{matrix} e^+ + \nu_e \\ \mu^+ + \nu_\mu \end{matrix}} \right\} 21.6 \%$$

$$\tau^+ + \nu_\tau \quad - 10.8 \%$$

$$\left(\begin{matrix} e^+ + \nu_e + \nu_\tau \\ \mu^+ + \nu_\mu + \nu_\tau \end{matrix} \right) \quad \left. \vphantom{\begin{matrix} e^+ + \nu_e + \nu_\tau \\ \mu^+ + \nu_\mu + \nu_\tau \end{matrix}} \right\} 35.2 \%$$

$$\begin{aligned} t &\rightarrow b + W^+ \rightarrow b + e^+ + \nu \\ \bar{t} &\rightarrow \bar{b} + W^- \rightarrow \bar{b} + e^- + \bar{\nu} \end{aligned}$$

4.7%



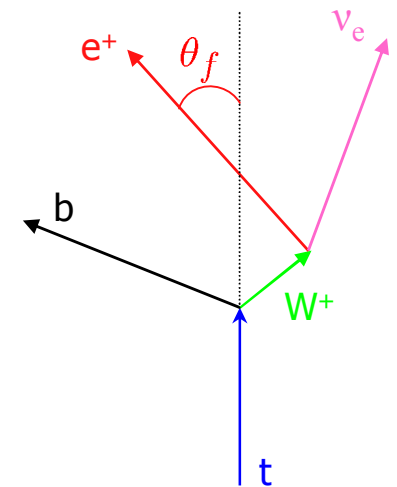
Polarization of the top quark

- At the LHC, the top (antitop) quarks are produced (in a good approximation) as the **helicity eigen-states**.
- The top and antitop quarks are produced as **unpolarized** – the same number of left- and right-handed top quarks.
- It is possible to **study the polarization** of the top quark **using the decay products**:

$$\frac{1}{N} \frac{dN}{d\cos\theta_f} = \frac{1}{2}(1 + \alpha_f \cos\theta_f)$$

The angle between the direction of movement of particle f in the top rest frame and the direction of top quark spin.

-0.41 for b
 0.41 for W^+
 0.35 for jet
1.0 for e^+, μ^+



Spin correlation of top-antitop pairs

The number of top-antitop pairs with the same and opposite helicity is not the same.

$$\begin{aligned} A &= 4 \langle (\hat{\mathbf{p}}_t \cdot \mathbf{S}_t)(\hat{\mathbf{p}}_{\bar{t}} \cdot \mathbf{S}_{\bar{t}}) \rangle \\ &= \frac{\sigma(t_{\uparrow}\bar{t}_{\uparrow}) + \sigma(t_{\downarrow}\bar{t}_{\downarrow}) - \sigma(t_{\uparrow}\bar{t}_{\downarrow}) - \sigma(t_{\downarrow}\bar{t}_{\uparrow})}{\sigma(t_{\uparrow}\bar{t}_{\uparrow}) + \sigma(t_{\downarrow}\bar{t}_{\downarrow}) + \sigma(t_{\uparrow}\bar{t}_{\downarrow}) + \sigma(t_{\downarrow}\bar{t}_{\uparrow})} \\ &= 1 - 2 \frac{\sigma(t_{\uparrow}\bar{t}_{\downarrow}) + \sigma(t_{\downarrow}\bar{t}_{\uparrow})}{\underbrace{\sigma(t_{\uparrow}\bar{t}_{\uparrow}) + \sigma(t_{\downarrow}\bar{t}_{\downarrow}) + \sigma(t_{\uparrow}\bar{t}_{\downarrow}) + \sigma(t_{\downarrow}\bar{t}_{\uparrow})}_{\text{Fraction of top-antitop pairs with the opposite helicities}}} \neq 0 \end{aligned}$$

Fraction of top-antitop pairs with the opposite helicities

SM prediction:

$$A = 0.319$$

- If the top quark is coupled to a new **physics beyond the SM**, the top-antitop **spin correlation could be altered**.

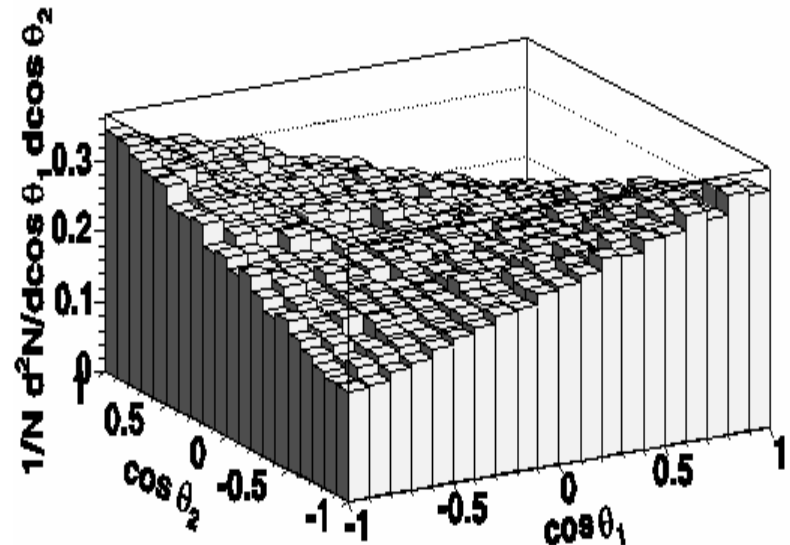
How to measure spin correlation

The **double differential angular distribution** of top and anti-top decay products:

$$\frac{1}{N} \frac{d^2 N}{d\cos\theta_f d\cos\theta_{\bar{f}}} = \frac{1}{4} (1 - A \underbrace{|\alpha_f \alpha_{\bar{f}}|}_{=1 \text{ for double-lepton channel}} \cos\theta_f \cos\theta_{\bar{f}})$$

The best statistical unbiased estimator:

$$A = -9 \langle \cos\theta_f \cos\theta_{\bar{f}} \rangle$$



$$A = 0.319$$

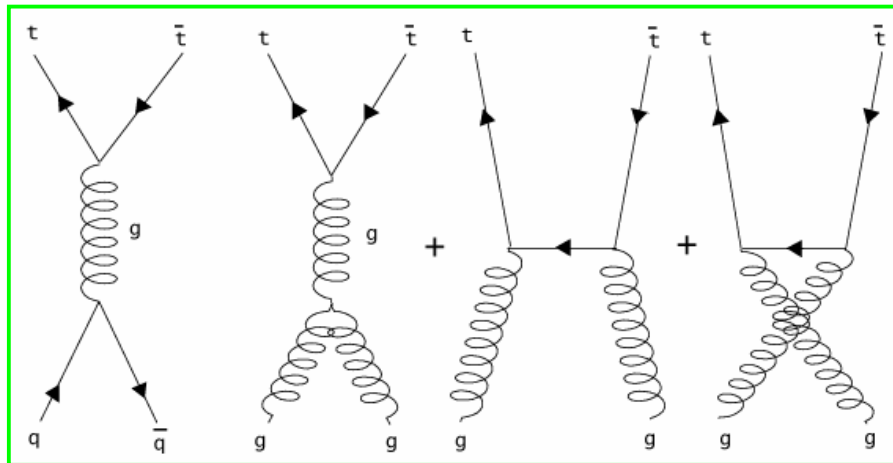
Measurement of top spin correlation

In the ATLAS experiment:

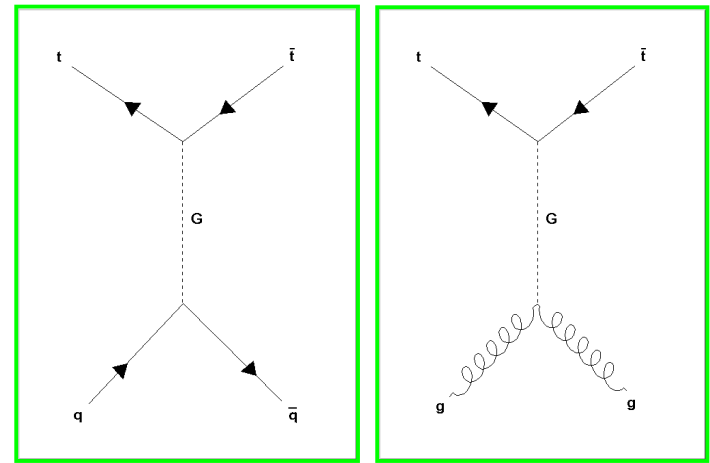
- F. Hubaut, E. Monnier, P. Pralavorio, K. Smolek, V. Šimák: *ATLAS sensitivity to top quark and W boson polarization in $t\bar{t}$ events*, Eur.Phys.J. C44 (2005) 13-33.
- Semileptonic and dileptonic top-antitop channel.
- Combining the results of both channels allows to measure the SM spin correlation A with a **3% precision** for 10 fb^{-1} .

Top quarks in theories with extra dimensions

- We studied two brane world scenarios:
 - ✓ **ADD** (Arkani-Hamed, Dimopoulos, Dvali)
 - ✓ **RS I** (Randal, Sundrum)
- Kaluza-Klein states of gravitons can contribute to the top-antitop production.



SM contribution



KK states contribution

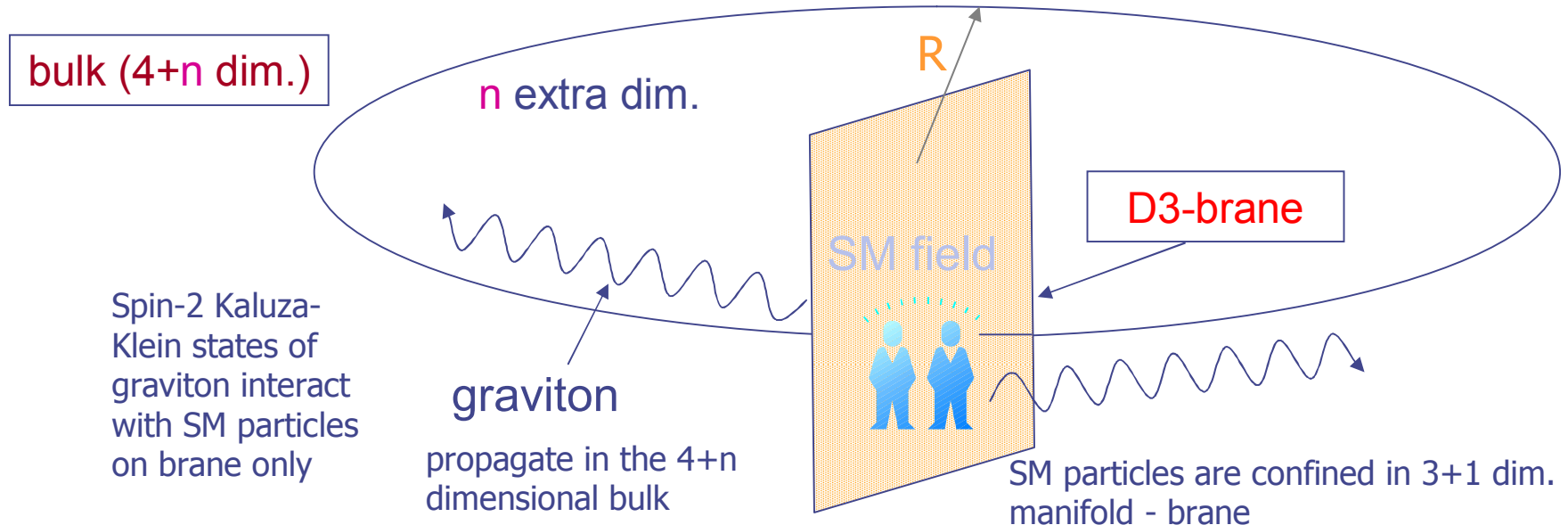
- KK gravitons can give rise to characteristic angular distributions and spin configurations of outgoing particles, which reflect the spin-2 nature of KK gravitons.

ADD model with large extra dimensions

- Theory with n extra-dimensions compactified with large radii.
N. Arkani-Hamed, et al, PLB429 (1998) 263, hep-ph/9803315
I. Antoniadis, et al, PLB436 (1998) 257, hep-ph/9804398
- n -extra dimensions are compactified on n -torus with common radius R
- **D3-brane** is embedded in $4+n$ dimensional bulk

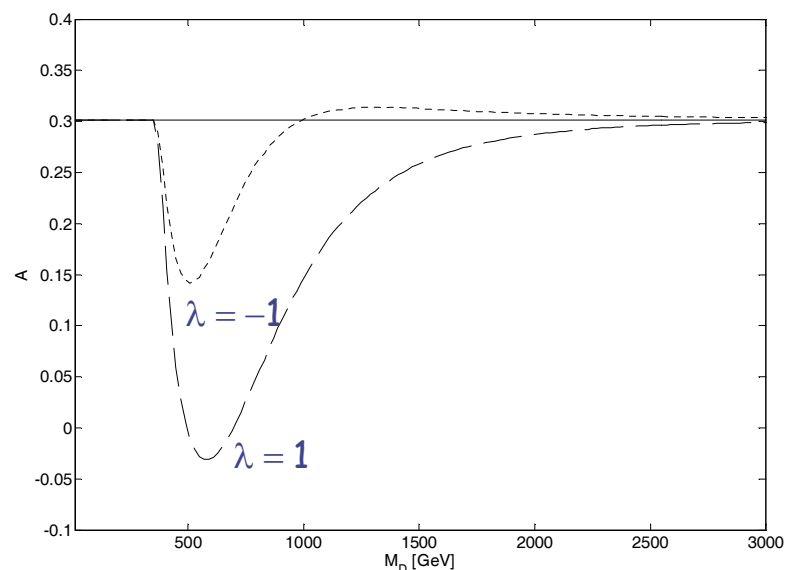
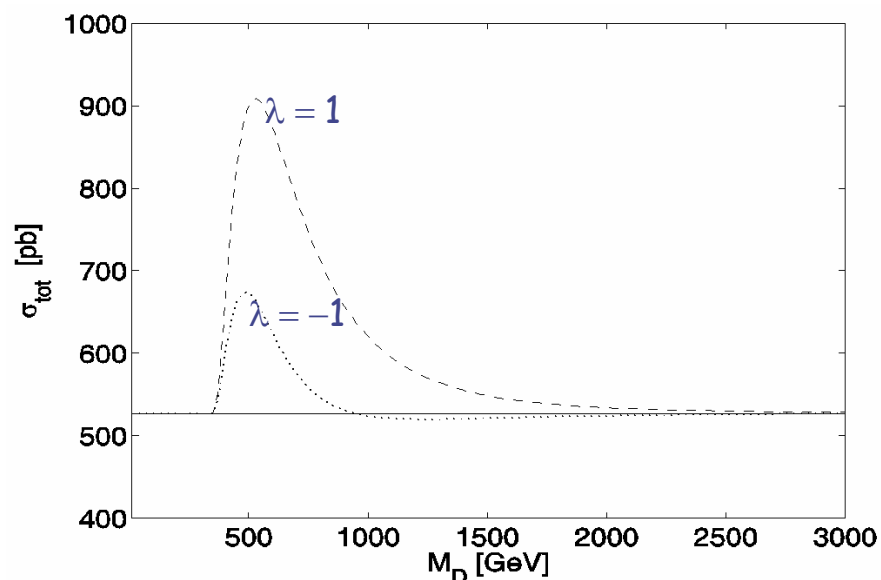
$$M_{\text{PL}}^2 = M_{\text{D}}^{n+2} R^n$$

M_{D} = low-energy effective string scale (~ 1 TeV for $R \sim 1$ mm, $n = 2$)



Top production in ADD model

- We computed full density matrix for top-antitop production.
- We studied spin correlation of top-antitop in ADD model.
- M. Arai, N. Okada, K. Smolek, V. Šimák: Phys.Rev. D70 (2004) 115015



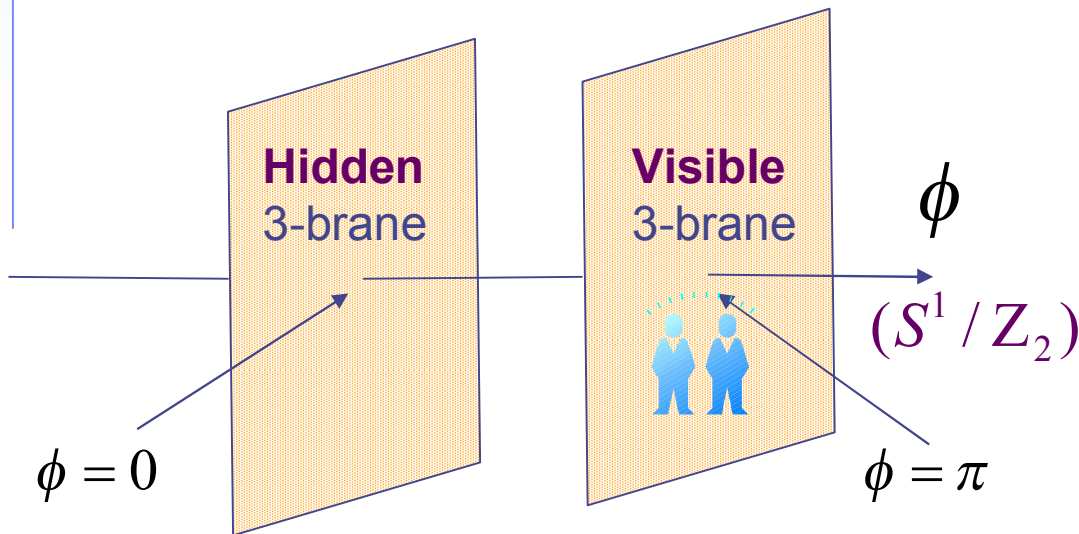
- $\lambda = \pm 1$ – connected to the regularization procedure for the contributions from the infinite number of KK gravitons. λ represents the sign of the interference term between SM and ADD contribution in the $gg \rightarrow t\bar{t}$ process.
- A sizable deviation of the top spin correlations from the SM one can be visible for the scale M_D below 2 TeV.

Randall-Sundrum scenario

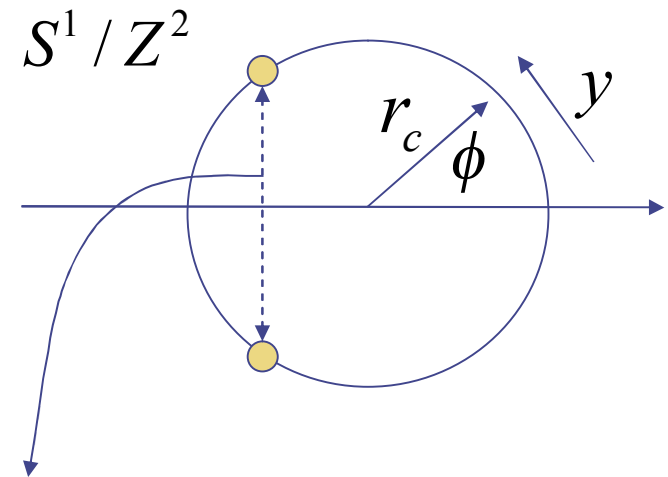
- In ADD scenario, the energy density of brane (gravitational field that brane produces) is ignored.

RS scenario (Randall, Sundrum, PRL83 (1999) 3370; 4690):

- 5 dimensional theory.
- Warped extra dimension. 5th dimension is compactified with orbifold symmetry.
- **M. Arai, N. Okada, K. Smolek, V. Šimák: Phys.Rev. D75 (2007) 095008**



r_c - compactification radius



Two points are identified.

$$y \leftrightarrow -y$$

$$y \rightarrow y + 2\pi r_c$$

Randall-Sundrum scenario

- The effective interaction Lagrangian:

$$\mathcal{L}_{\text{int}} = -\frac{1}{\bar{M}_{\text{pl}}} T^{\mu\nu}(x) h_{\mu\nu}^{(0)}(x) - \frac{1}{\Lambda_\pi} T^{\mu\nu}(x) \sum_{n=1}^{\infty} h_{\mu\nu}^{(n)}(x),$$

$h_{\mu\nu}^{(n)}$ - n -th graviton KK mode $T^{\mu\nu}$ - energy-momentum tensor of SM fields on the visible brane
 \bar{M}_{pl} - reduced Plack mass

- Sum of all intermediate KK gravitons gives a finite value.
- The graviton zero mode couples with the usual strength -> negligible effect.
- Each KK graviton strongly couples to SM fields with Λ_π suppressed couplings.

$$\Lambda_\pi = e^{-\kappa r_c \pi} \bar{M}_{\text{pl}} = \frac{m_1}{x_1} \left(\frac{\bar{M}_{\text{pl}}}{\kappa} \right) \sim \text{TeV}$$

κ - 5-dimensional curvature

- For $\kappa r_c \simeq 12$, $\Lambda_\pi = \mathcal{O}(1 \text{ TeV})$ and give a natural solution to the gauge hierarchy problem.

Randal-Sundrum scenario

- Mass spectrum of gravitons

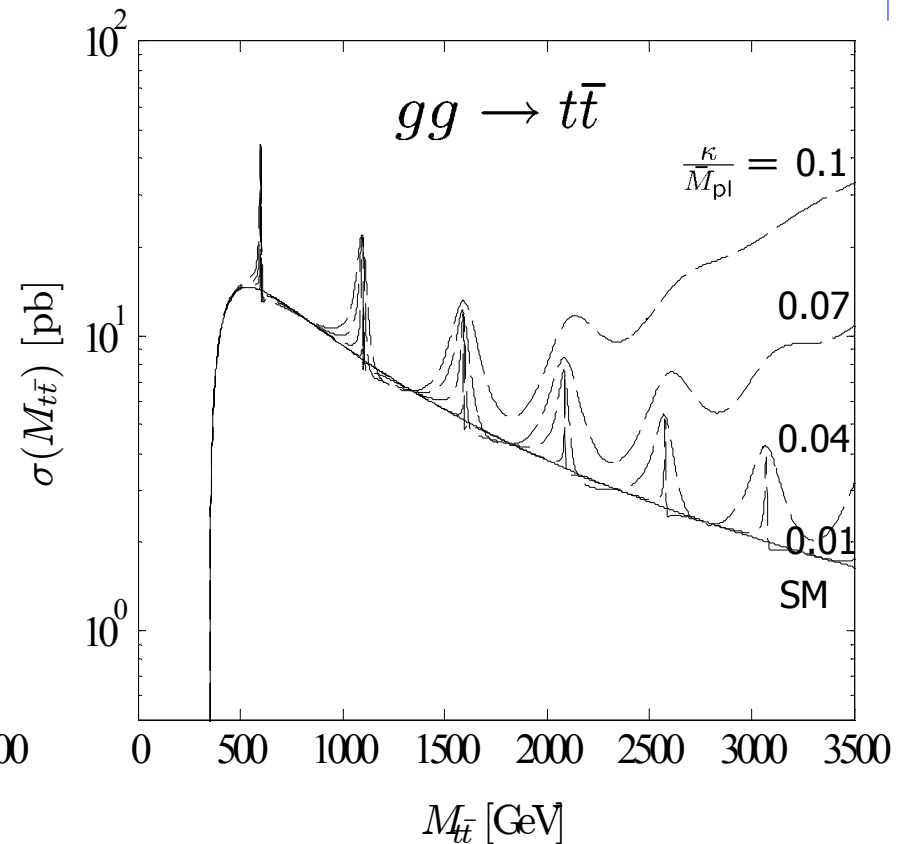
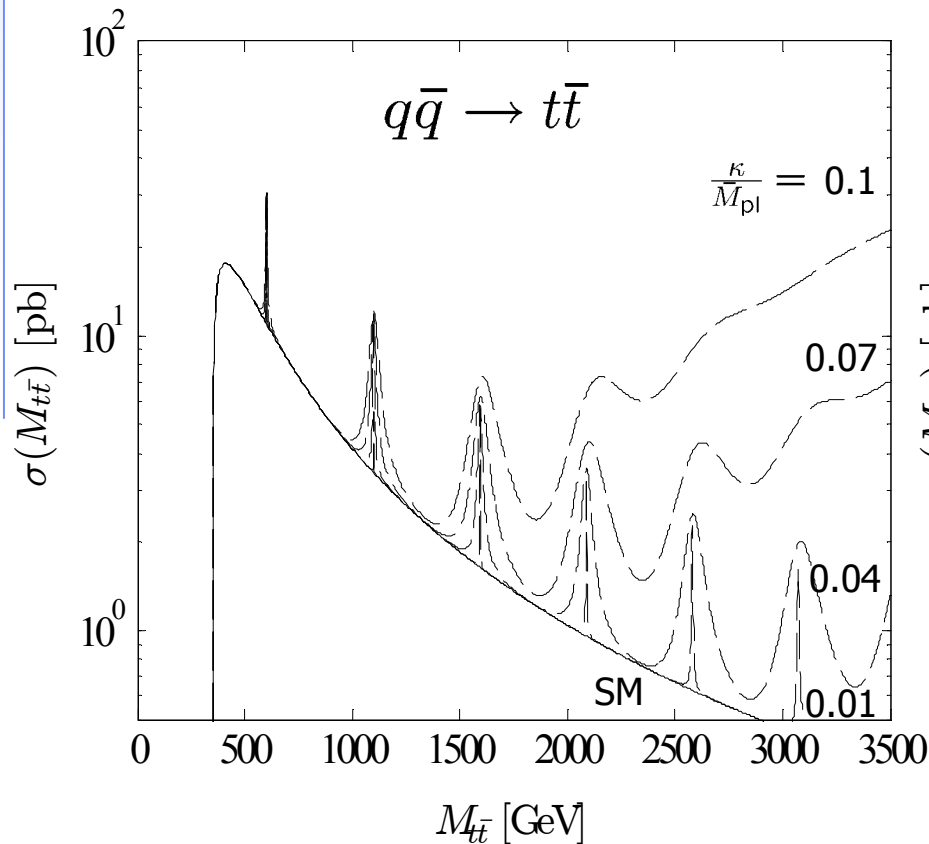
$$m_n = \kappa x_n e^{\kappa r c \pi} = m_1 \frac{x_n}{x_1}$$

x_n - roots of the Bessel function of the first order ($x_1 = 3.83$, $x_2 = 7.02, \dots$)

- We can expect a resonant production of KK gravitons at colliders.
- The resonance gives rise to an enhancement of production of the top-antitop pairs and provide a big statistical advantage for studying the top spin correlations around the resonance pole.
- In our analysis we used:
 - ✓ $m_1 = 600 \text{ GeV}/c^2$ - $m_1 \geq 600 \text{ GeV}/c^2$ from D0 experiment
 - ✓ $\frac{\kappa}{M_{\text{pl}}} = 0.01, 0.04, 0.07, 0.1$ - guarantees the perturbation of the graviton
 - ✓ $m_t = 175 \text{ GeV}/c^2$
 - ✓ PDF CTEQ6L

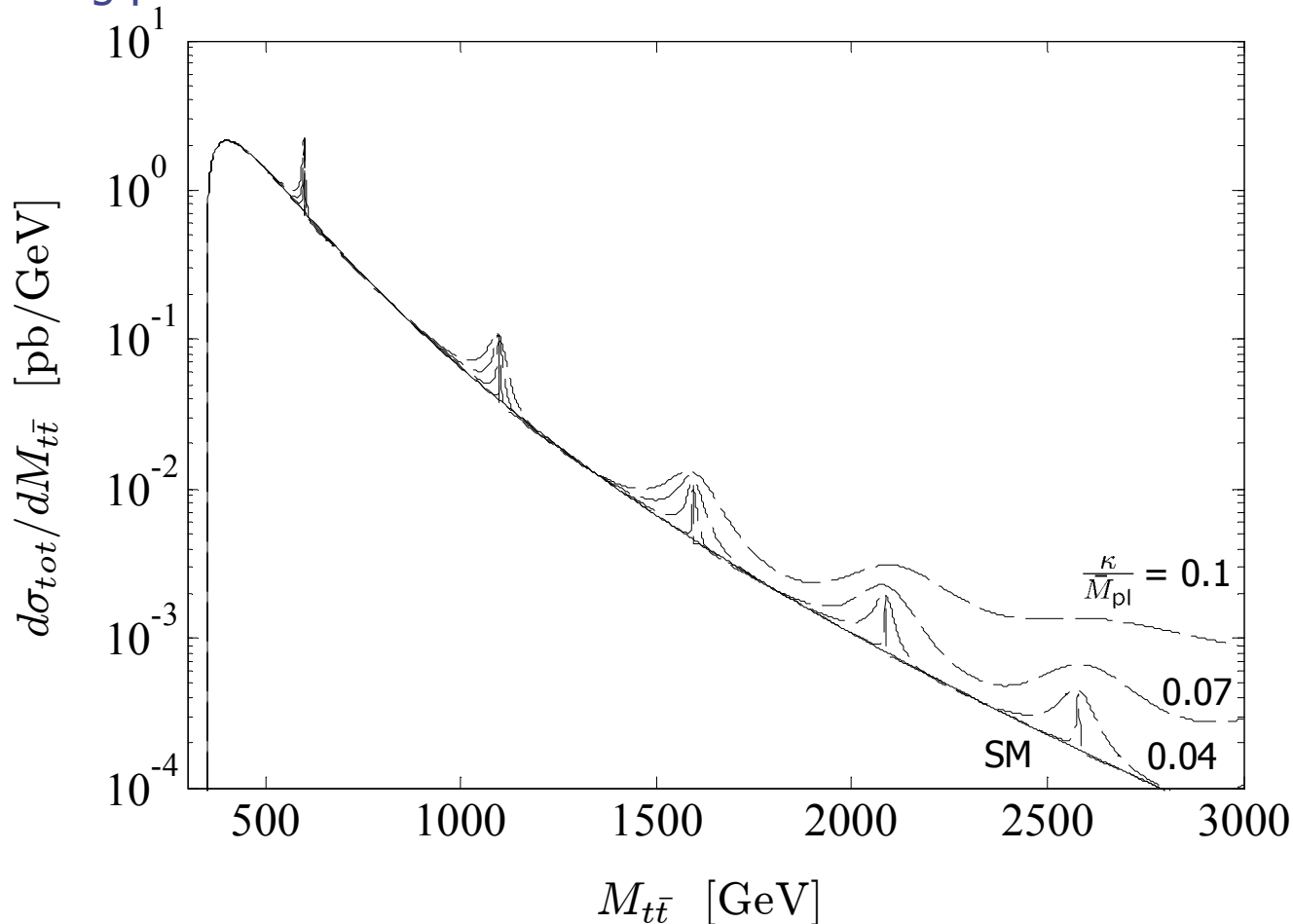
Top quark in Randal-Sundrum scenario - results

The dependence of the cross section of the top-antitop quark pair production by quark annihilation and gluon fusion on the CMS energy of colliding partons.



Top quark in Randal-Sundrum scenario - results

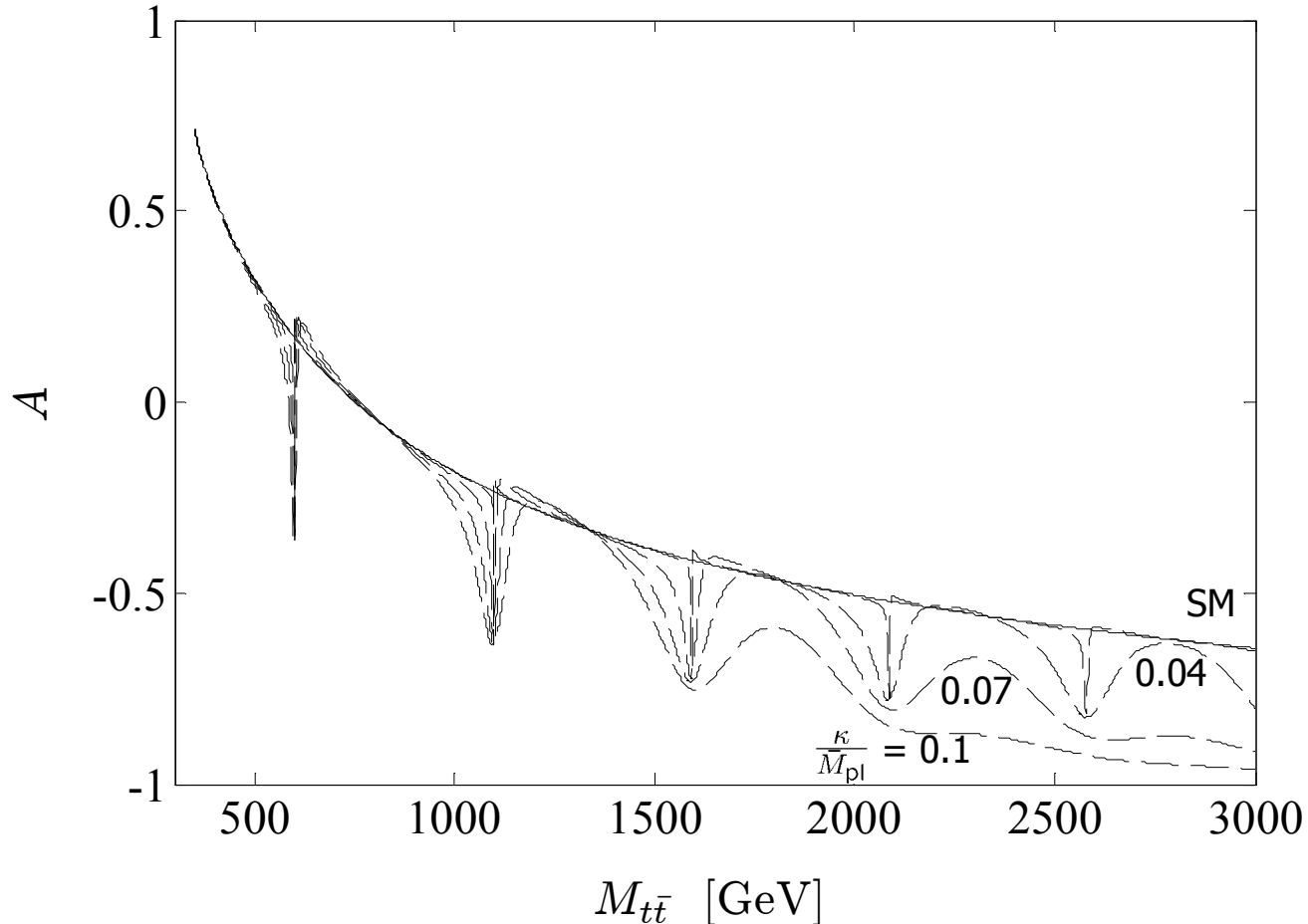
Total differential cross section $\frac{d\sigma_{tot}(pp \rightarrow t\bar{t})}{d\sqrt{M_{t\bar{t}}}}$ as a function of the CMS energy of colliding partons.



- Resonant production of the KK gravitons give rise to an enhancement of the deviations from the SM.

Top quark in Randal-Sundrum scenario - results

Spin asymmetry A as a function of the CMS energy of colliding partons.

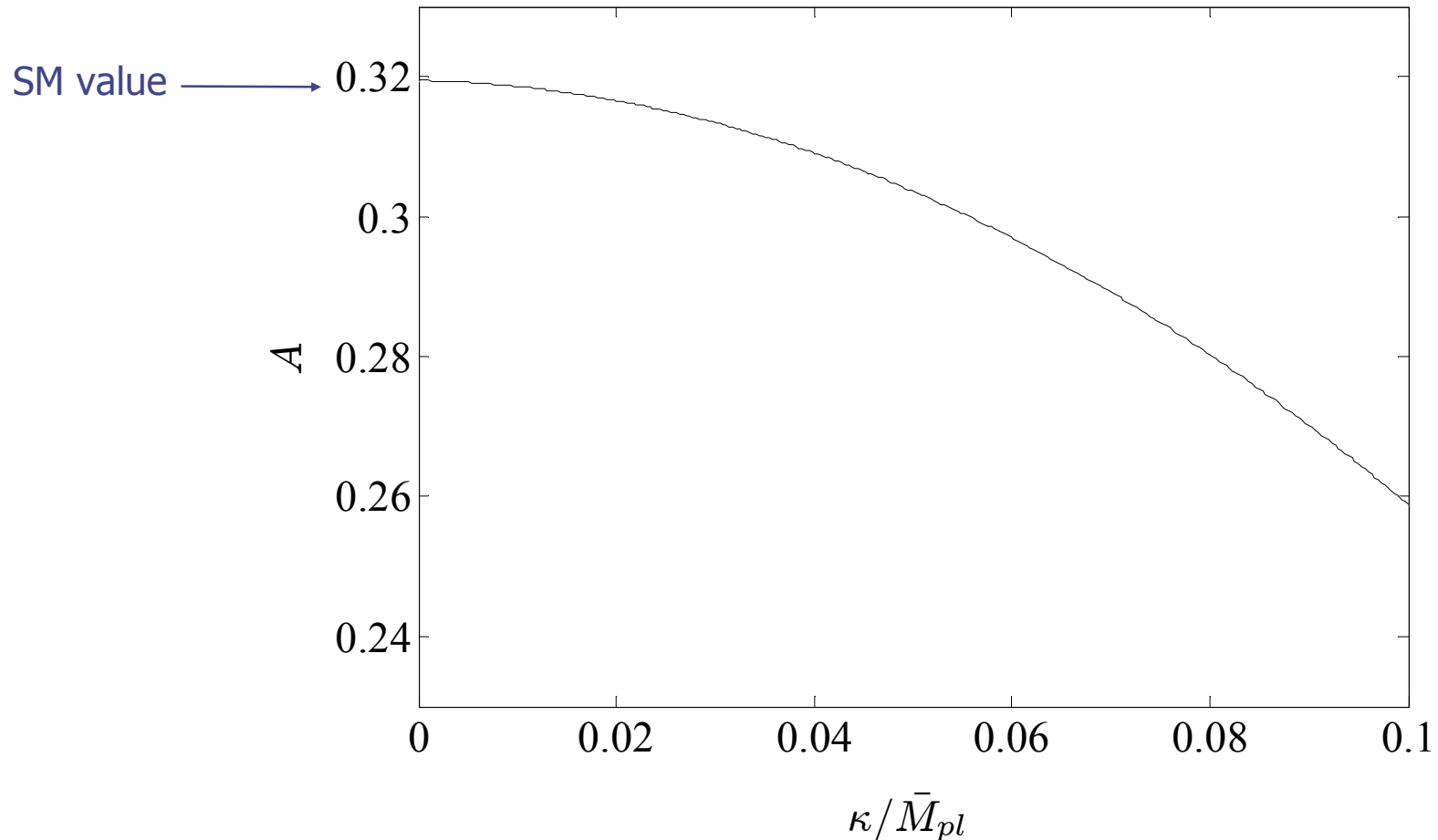


- Resonant production of the KK gravitons give rise to an enhancement of the deviations from the SM.

Top quark in Randal-Sundrum scenario - results

Spin asymmetry A as a function of $\frac{\kappa}{\bar{M}_{pl}}$.

$$m_1 = 600 \text{ GeV}/c^2$$



Conclusions

- Because of its high mass, the top quark is an ideal place to search for physics beyond the SM.
- The ADD model with large extra dimensions or RS model is an example of such physics.
- In addition to cross section and various kinematical distributions, the spin correlation is sensitive to the existence extra dimensions.
- We studied in detail the production of top-antitop quarks at LHC for the RS scenario.
- The influence of gravitons in the RS model on the spin correlation of top-antitop quarks could be visible at the LHC.
- Resonant production of the KK gravitons give rise to a remarkable enhancement of the deviations from the SM. This is a crucial difference from the case in the ADD model.
- M. Arai, N. Okada, K. Smolek, V. Simak: Phys.Rev. D70 (2004) 115015
- M. Arai, N. Okada, K. Smolek, V. Simak: Phys.Rev. D75 (2007) 095008.