



Medium-Modified Fragmentation Functions

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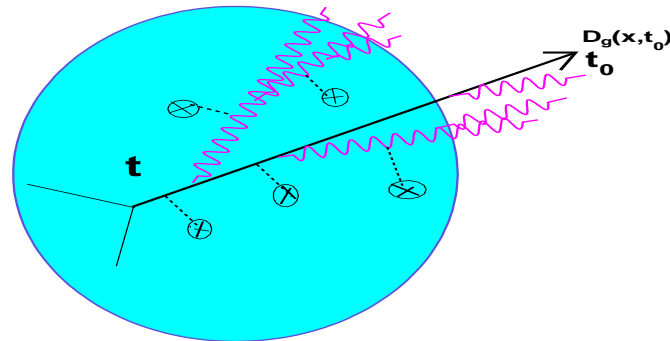
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Outline

- Introduction: Medium Induced Gluon Radiation
- Medium-Modified Splitting Functions
- Sudakov Form Factors
- Medium Modified Fragmentation Functions
- Particle Spectra $\longrightarrow R_{AA} \longrightarrow \hat{q}$
- Summary

Introduction: Medium-induced gluon radiation

- Medium-Induced Radiation: A hard parton produced in a Heavy Ion Collision travels through the hot medium. The scattering centers induce successive gluon radiations.



- The hard parton loses virtuality from the initial scale t to the final hadronization one $t_0 \simeq \lambda_{QCD}^2$. Hadronization happens outside the medium.
- We end up with a parent parton degraded in momentum and a surrounding cone of soft partons.
- **Medium-Induced gluon radiation is the standard explanation for Jet Quenching observed at RHIC.**

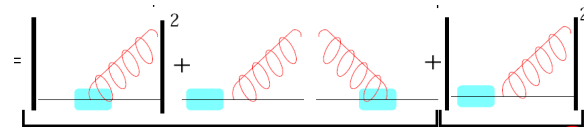
Introduction: Medium-induced gluon radiation

- The single inclusive distribution of medium induced gluons with energy ω and k_t from a parent parton of energy E :

$$\omega \frac{dI}{d\omega dk_T} = \frac{\alpha_S C_R}{(2\pi)^2 \omega^2} 2 \text{Re} \int_0^\infty dy_l \int_{y_l}^\infty d\bar{y}_l \int du e^{-ik_T u} e^{-\frac{1}{2} \int_{y_l}^\infty d\zeta n(\zeta) \sigma(u)}$$

$$\frac{\partial}{\partial y} \frac{\partial}{\partial u} \int_{y=0=r(y_l)}^{u=r(\bar{y}_l)} \mathcal{D}r e^{i \int_{y_l}^{\bar{y}_l} d\zeta \frac{\omega}{2} (r^2 - \frac{n(\zeta) \sigma(r)}{i\omega})}$$

- Two approaches:
 - opacity expansion (in powers of $n(\zeta) \sigma(\zeta)$)
 - BDMPS approximation: $n(\zeta) \sigma(r) \simeq \frac{1}{2} \hat{q}(\zeta) r^2$ The path integral is the one of an harmonic oscillator.
 - $\hat{q}(\zeta)$ is the transport coefficient $\longrightarrow \langle q_T^2 \rangle / \lambda$.



$$\diamond \omega \frac{dI}{d\omega dk_T} \longrightarrow F\left(\frac{\omega}{\omega_c}, \kappa^2\right), \quad \omega_c = \frac{1}{2} \hat{q} L^2, \quad \kappa^2 = \frac{k_T^2}{\hat{q} L}$$

Previous MMFF Calculations

- A **Poissonian** distribution of **independent** radiations was **assumed** (**BMDS**)

$$P_E(\epsilon) = \sum_{n=0}^{\infty} \frac{1}{n!} [\prod_{i=1}^n \int d\omega_i \frac{d((\omega_i))}{d\omega}] \delta(\epsilon - \sum_{i=1}^n \frac{\omega_i}{E}) e^{-\int d\omega \frac{dI}{d\omega}}$$

- The MMFF were calculated shifting the vacuum ones

$$D_{kh}^{(med)}(z, Q^2) = \int_0^1 d\epsilon P_E(\epsilon) \frac{1}{1-\epsilon} D_k(\frac{z}{1-\epsilon}, Q^2)$$

- Limitations:
 - Energy and momentum are not conserved (**only a posteriori**)
 - There's no evolution in virtuality
 - Medium an vacuum are treated differently

Now: Medium-modified splitting functions

- $\frac{dI^{vac}}{dzdk_T^2} = \frac{\alpha_s P(z)_{z \rightarrow 1}^{vac}}{2\pi k_T^2}$, $P(z)_{z \rightarrow 1}^{vac} \simeq \frac{2C_R}{1-z}$, $z = 1 - x$

- The ansatz is an extension of the former eq. to medium: **Salgado and Polosa (hep-ph/0607295)**

$$\frac{dI^{MED}}{dzdk_T^2} = \frac{\alpha_s P(z)_{z \rightarrow 1}^{MED}}{2\pi k_T^2} , P(z)_{z \rightarrow 1}^{med} = \frac{2\pi z t}{\hat{q} L} F\left(\frac{\omega}{\omega_c}, \kappa^2\right)$$

- The formalism of medium induced radiation relies on high energy approximations where previous eqs. are valid.
- The total splitting distribution is assumed to be the vacuum + medium ones:

$$P^{TOTAL}(z) = P^{VACUUM}(z) + P^{MEDIUM}(z, t, \hat{q}, L)$$

- **Borghini, Wiedemann** proposal: medium multiplicative factor

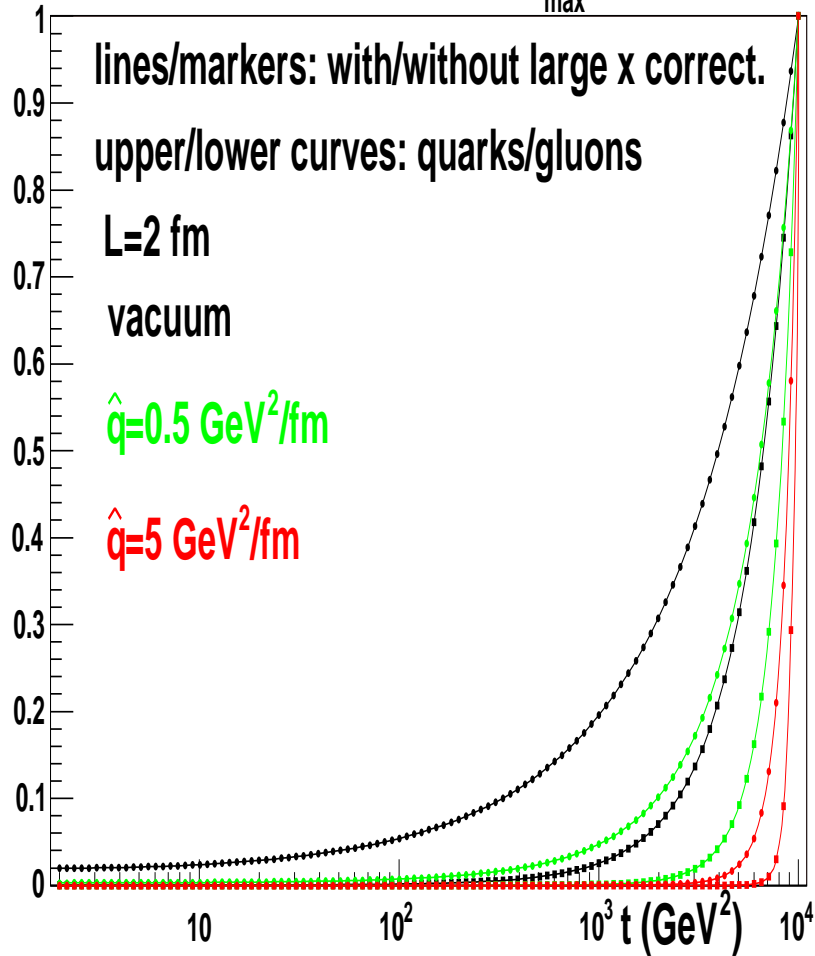
Sudakov Form Factors

- $\Delta_a(t, t_0^a) = e^{-\sum_{a-cc'} \int_{t_0^a}^t \frac{dt'}{t'} \int_{z_{min}(t')}^{1-z_{min}(t')} dz \frac{\alpha_S(t', z)}{2\pi} P_{ca}(z)}$
- $\Delta_a(t, t_0^a)$ means the probability for parton a not to branch(resolved) while evolving from an initial virtuality t_0 to a final scale t .
- $\frac{\Delta(t_0, t)}{\Delta(t_0, t')}$ stands for the probability of evolving from the intermediate scale t' to t without branching.
- independent radiation: branching at a scale t does not affect branching at a scale t' .
- We modify the Sudakov:

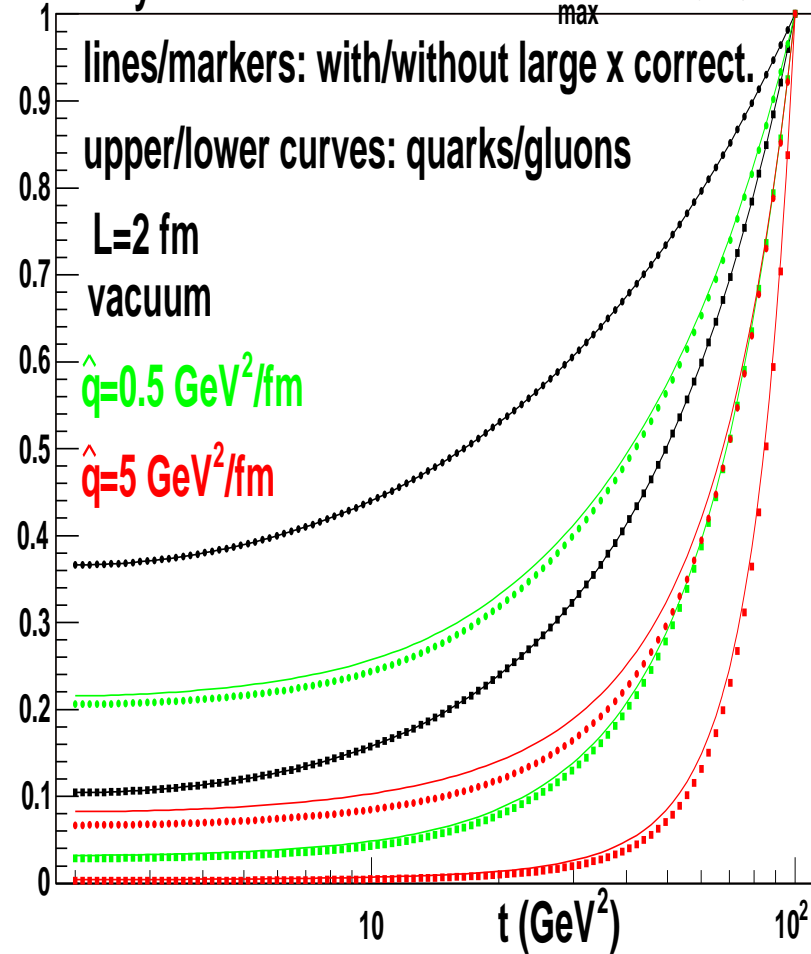
$$\Delta_a(t, t_0^a) = e^{-\sum_{a-cc'} \int_{t_0^a}^t \frac{dt'}{t'} \int_{z_{min}(t')}^{1-z_{min}(t')} dz \frac{\alpha_S(t', z)}{2\pi} (P_{ca}(z)^{vac} + P_{ca}(z)^{med})}$$

Medium Modified Sudakov Factors

Probability of no emission between $t = E^2 = 10^4 \text{ GeV}^2$ and t_{max}



Probability of no emission between $t = E^2 = 10^2 \text{ GeV}^2$ and t_{max}



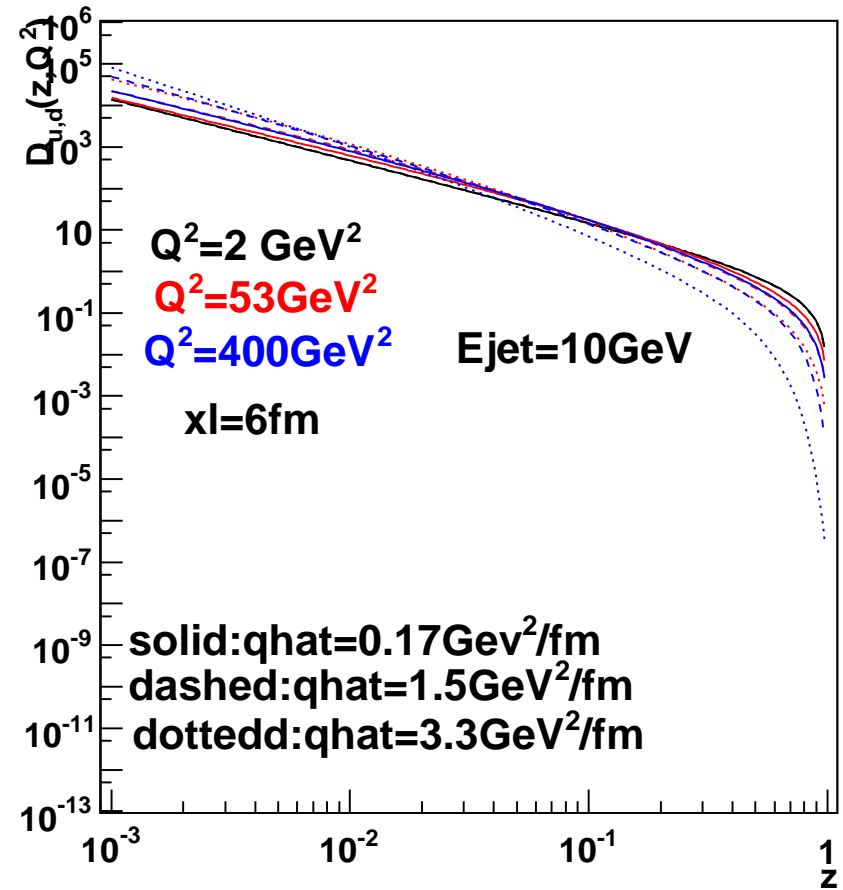
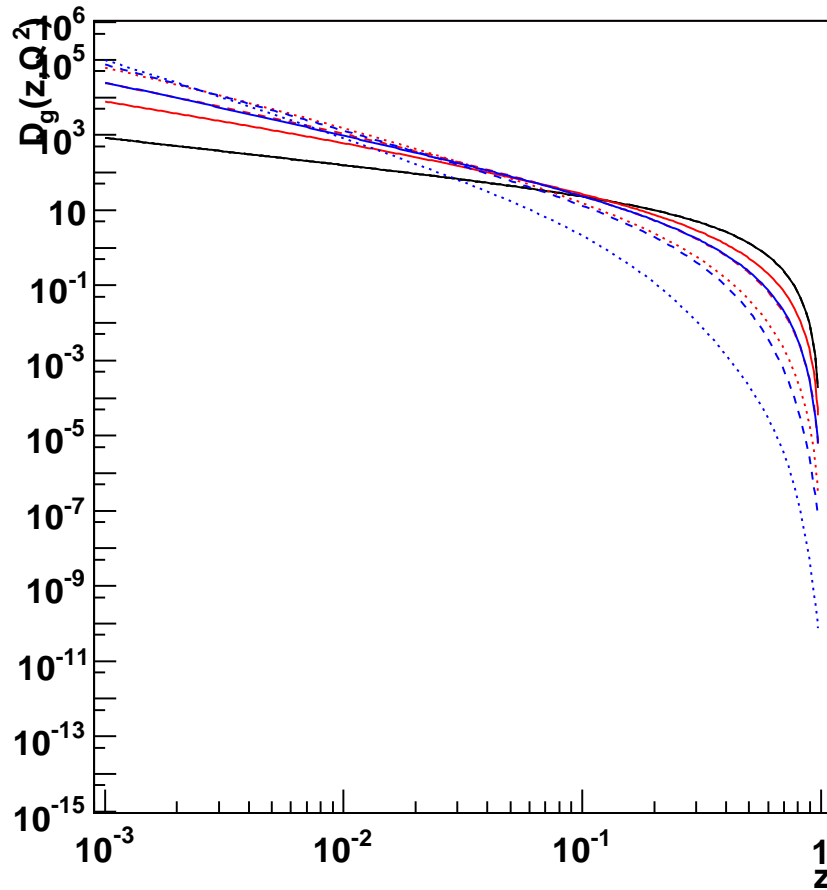
Medium Modified Fragmentation Functions

- DGLAP can be written in terms of the Sudakovs:

$$t \frac{\partial}{\partial t} \left(\frac{D_a^h(x,t)}{\Delta_a(t,t_0^a)} \right) = \int_x^{1-z_{min}(t)} \frac{dz}{z} \frac{\alpha_S(k_T^2, z)}{2\pi} P_{ba}(z) \frac{D_b^h(\frac{x}{z}, t)}{\Delta_a(t,t_0^a)}$$

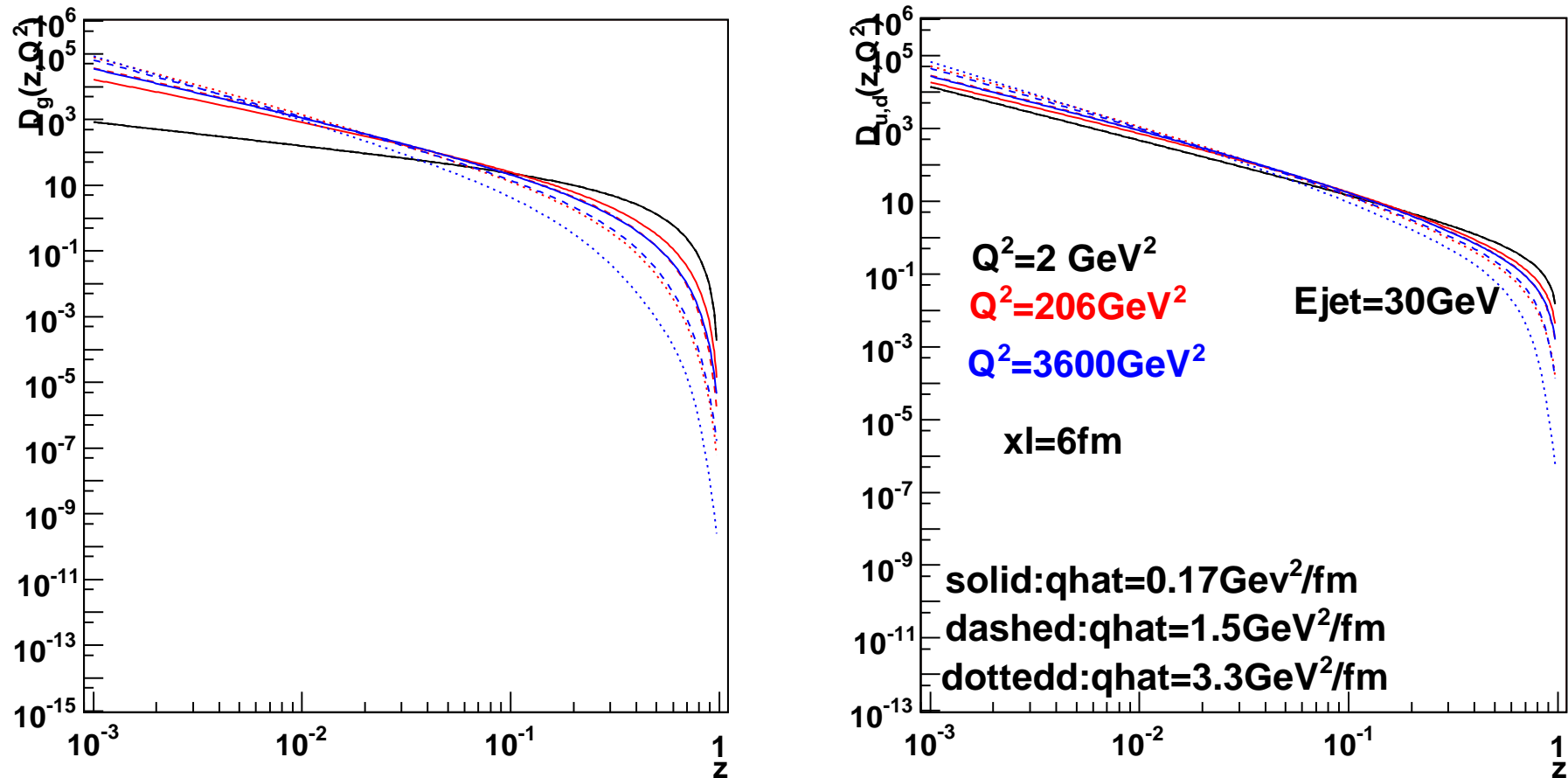
- We consider only 3 flavors(u,d,s).
- The renormalization scale is $t(1-z)z = k_T^2$.
- As initial values for the FF we take the KKP ones at virtuality $t_0 = 2GeV^2$
- For each parton energy, $t_0 < t < 4E^2$ and $t_0/t < z(t) < 1 - t_0/t$
- Our evolution depends on the initial parton energy through the scale range in the Sudakovs.
- The induced gluon radiation accelerates the evolution.

Medium Modified Fragmentation Functions



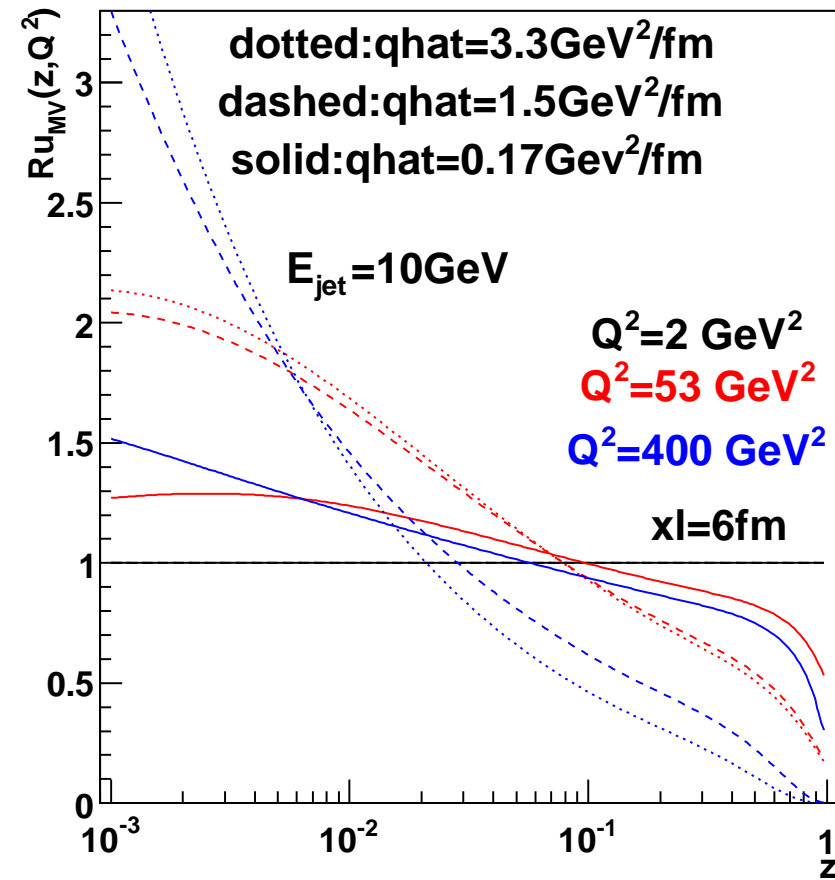
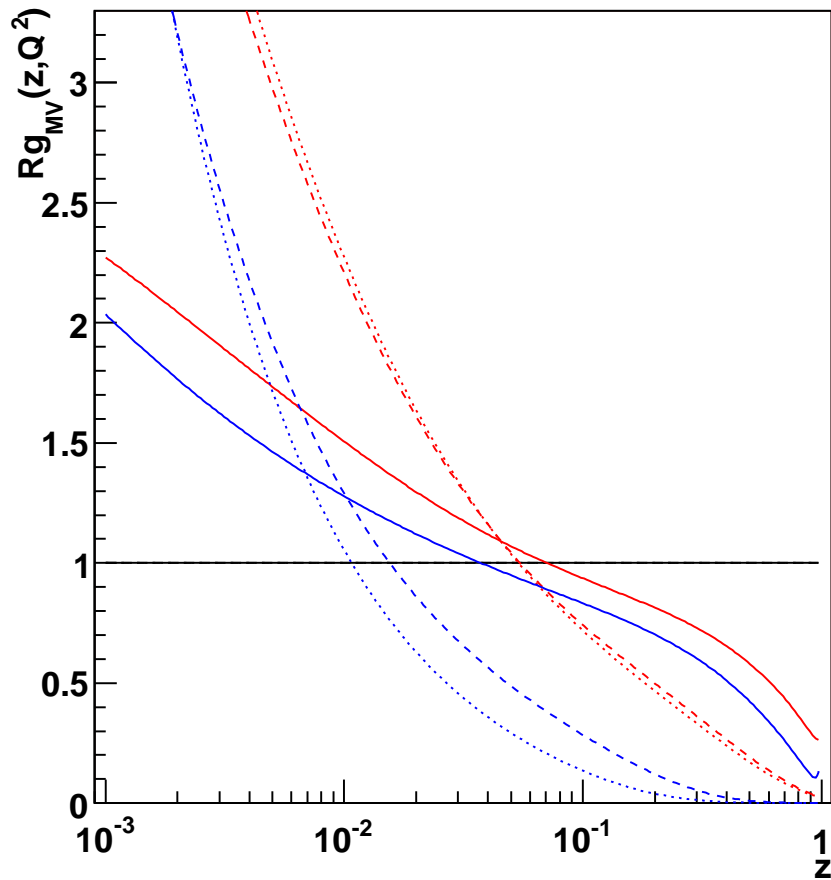
Fragmentation Functions for different medium densities for $E_{jet} = 10 \text{ GeV}$

Medium Modified Fragmentation Functions



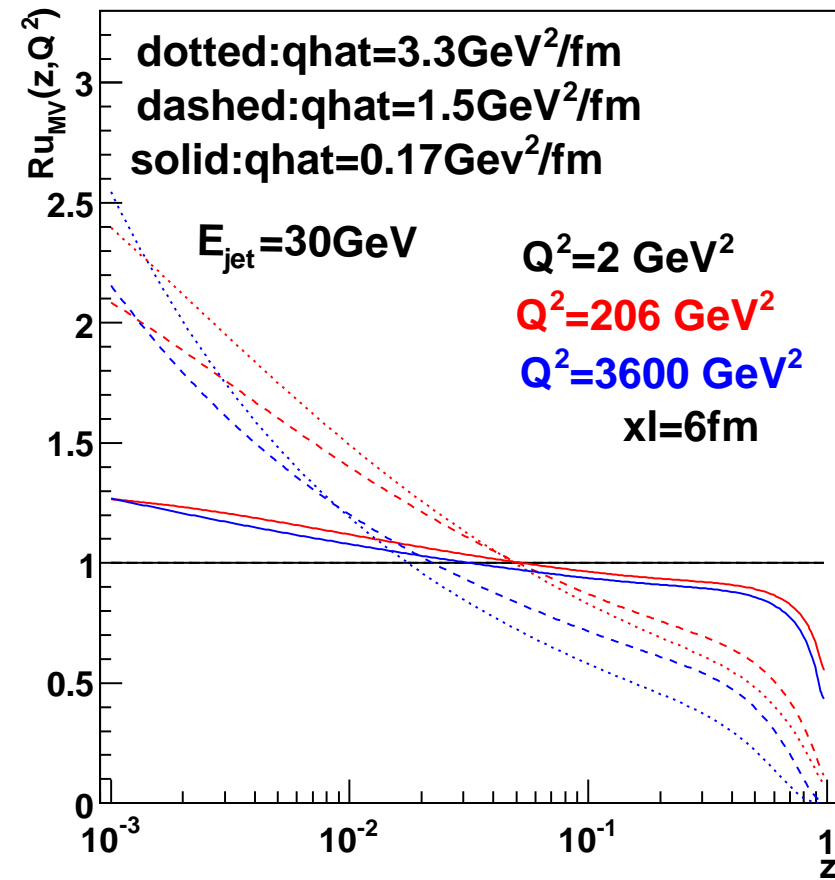
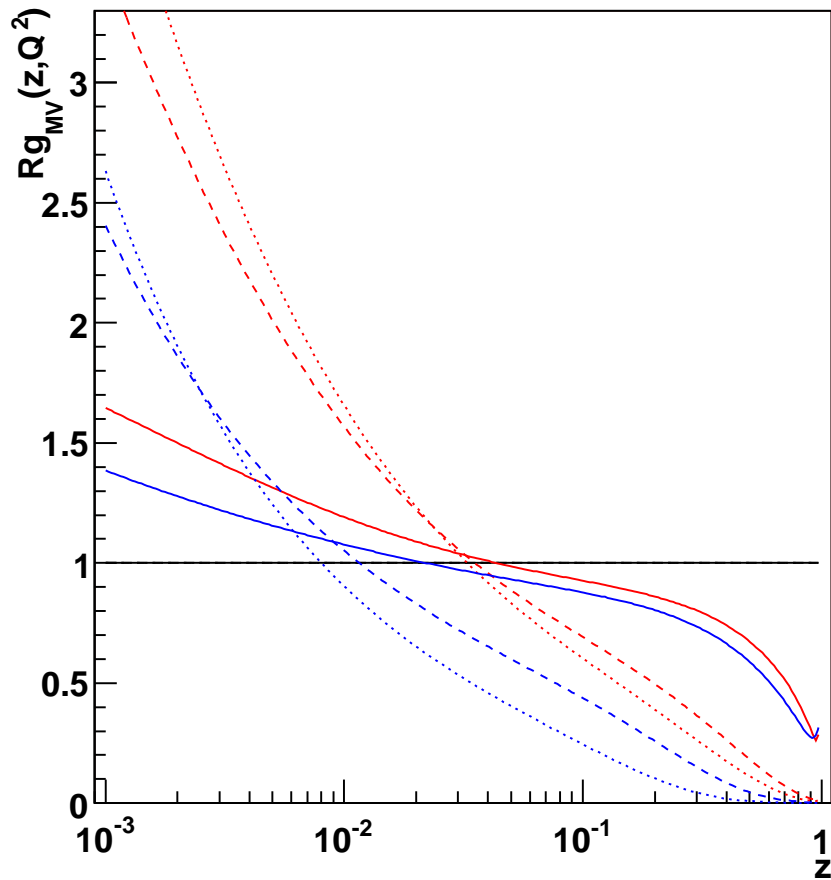
Fragmentation Functions for different medium densities for $E_{jet} = 30 \text{ GeV}$

Medium Modified Fragmentation Functions



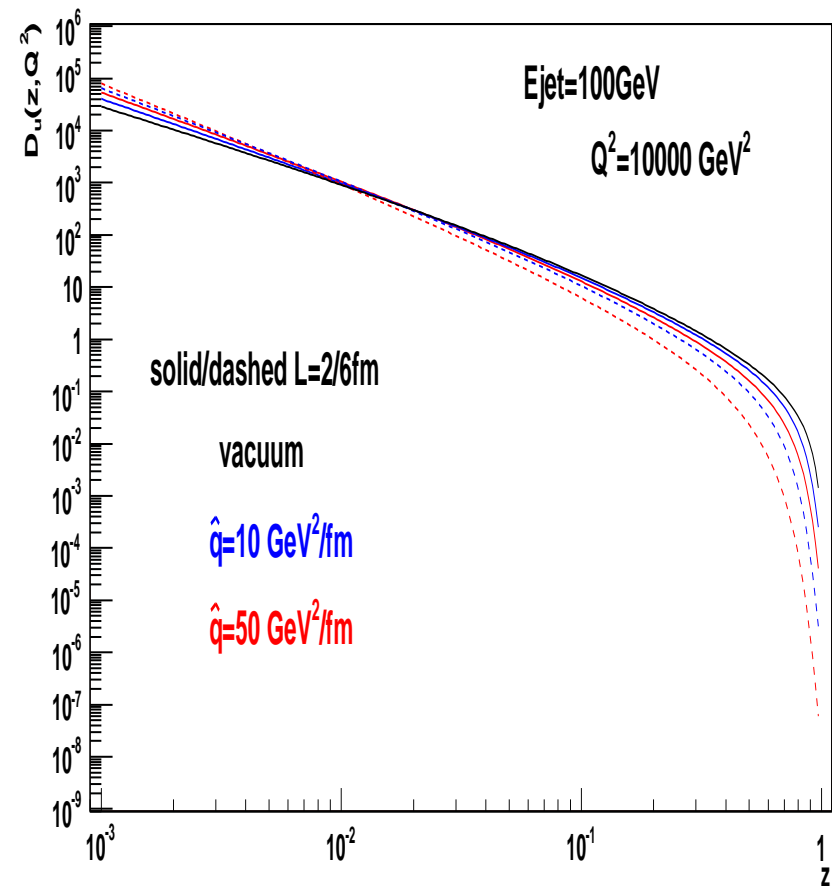
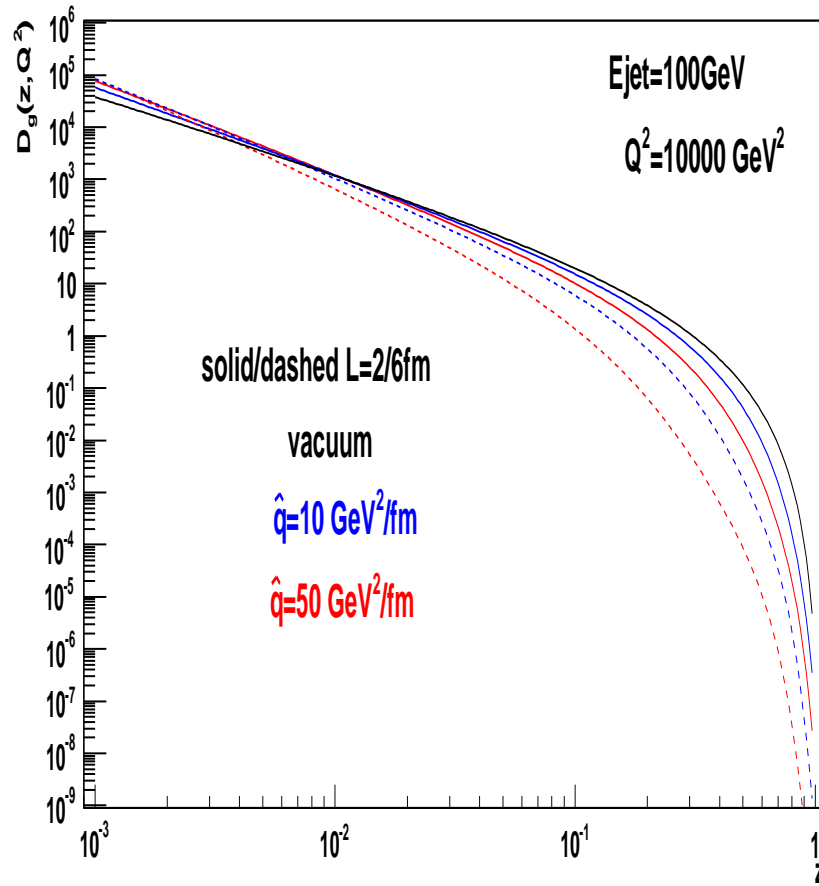
Ratio Medium/Vacuum of the Fragmentation Functions of $E_{jet} = 10 GeV$.

Medium Modified Fragmentation Functions

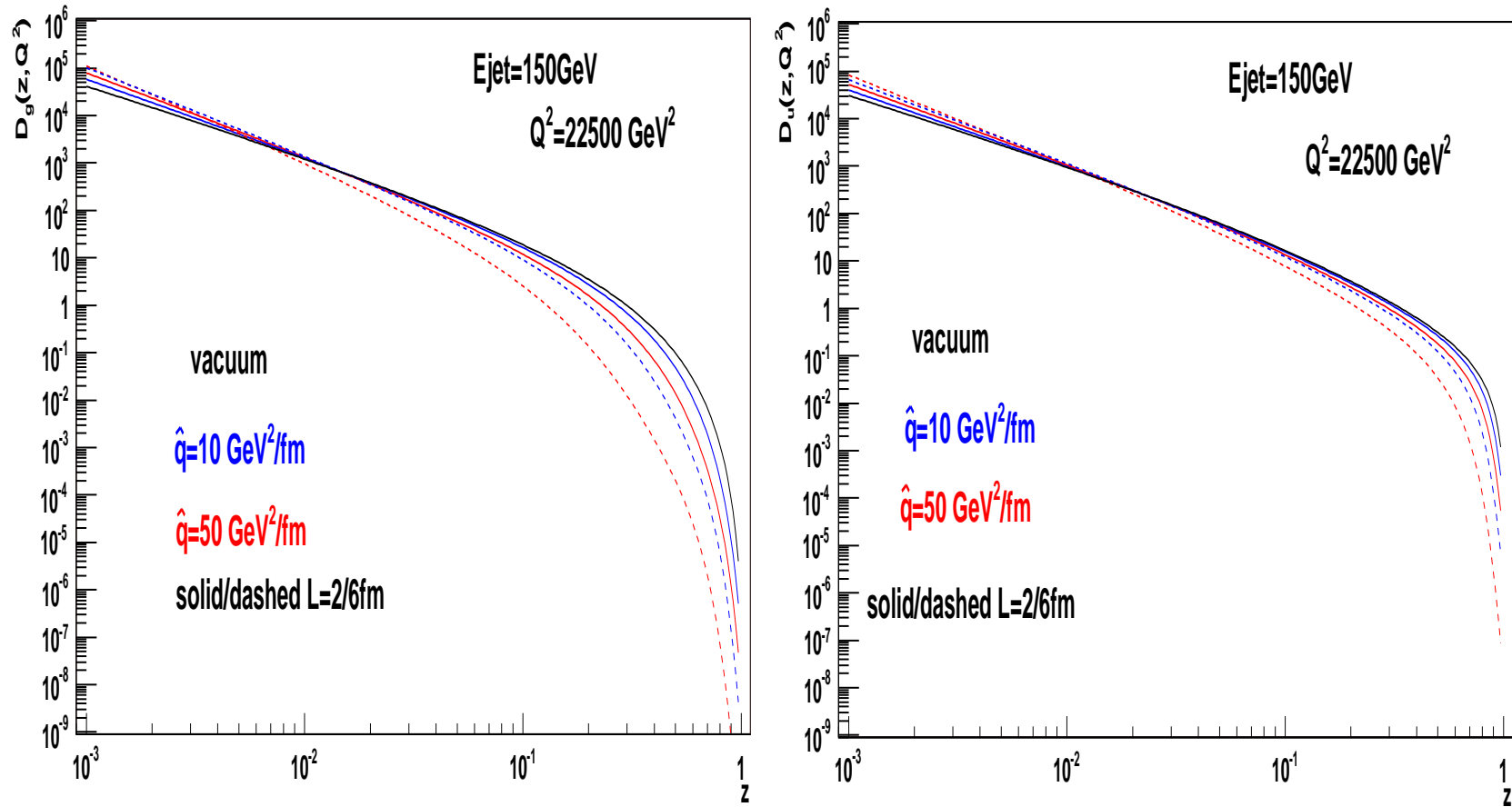


Ratio Medium/Vacuum of the Fragmentation Functions of $E_{jet} = 30 \text{ GeV}$.

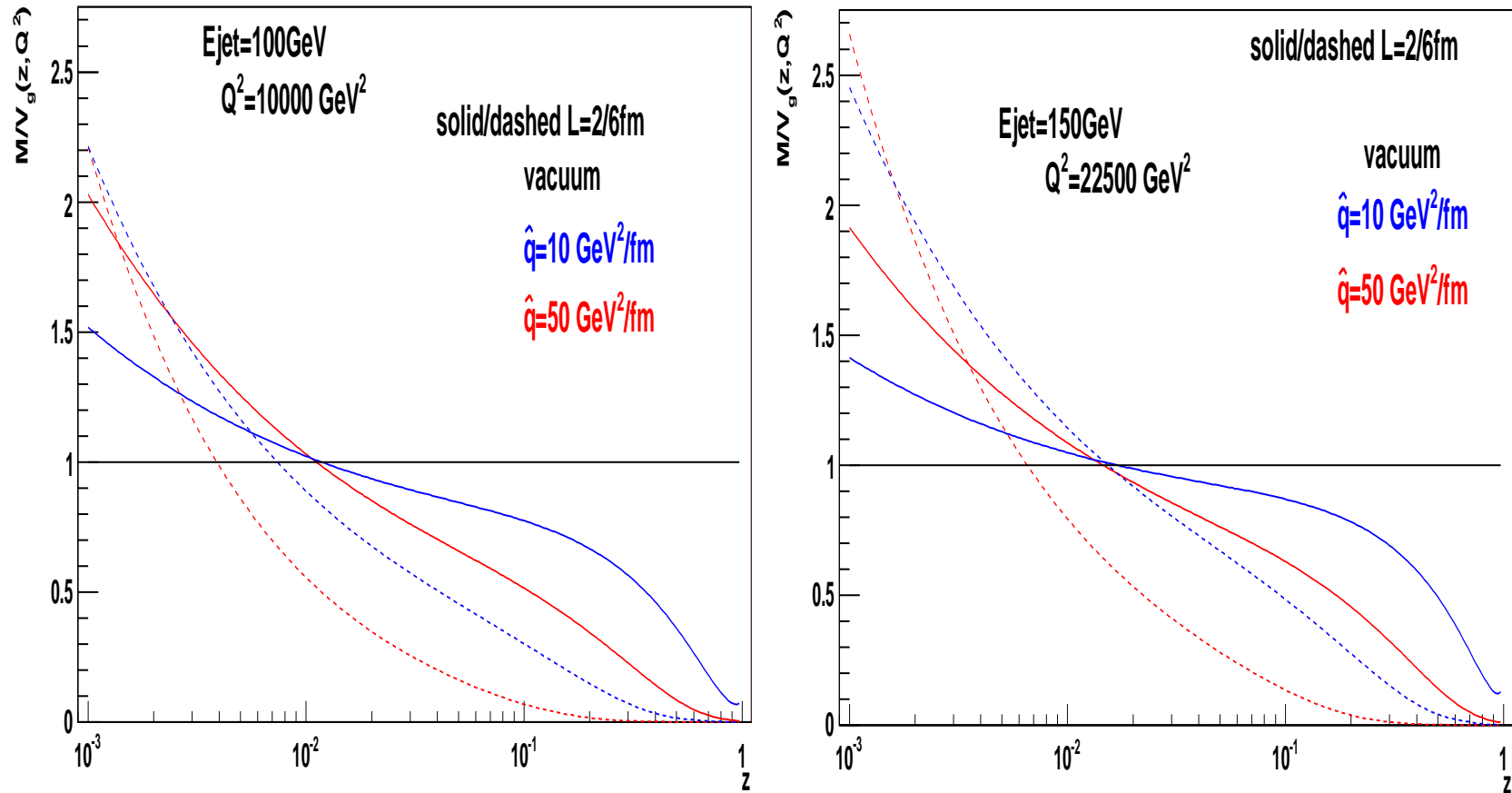
Medium Modified Fragmentation Functions:LHC



Medium Modified Fragmentation Functions:LHC



Medium Modified Fragmentation Functions:LHC

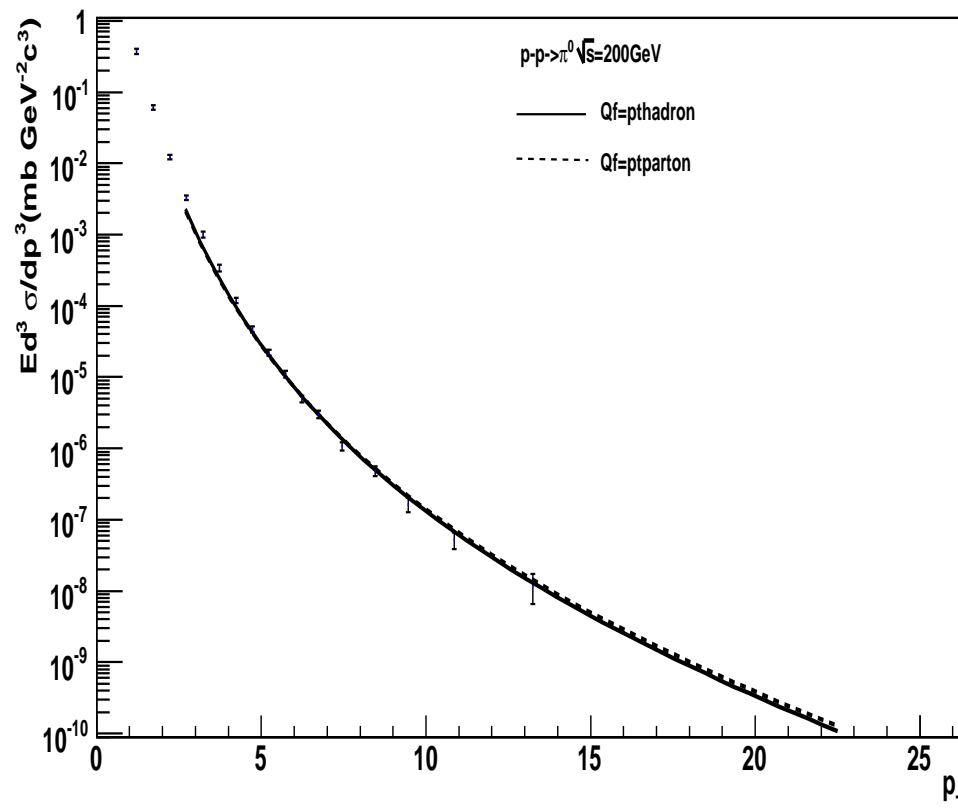


Particle Spectra: formalism

- A typical hard cross section can be written in the form:
- $\sigma^{ABh} = f_A(x_1, Q^2) f_B(x_2, Q^2) \otimes \sigma(x_1, x_2, Q^2) \otimes D_{i \rightarrow h}(z, Q^2)$
 - $D_{i \rightarrow h}(z, Q^2)$ long distance non perturbative object \rightarrow we modify its perturbative evolution.
- We define the nuclear modification factor as:

$$R_{AA} = \frac{\frac{d\sigma}{dydq_T^2}(pdf + EKS + MMFF)}{\frac{d\sigma}{dydq_T^2}(pdf + VACFF)}$$

Particle Spectra: pp Reference



- Vacuum pp spectra as a reference.

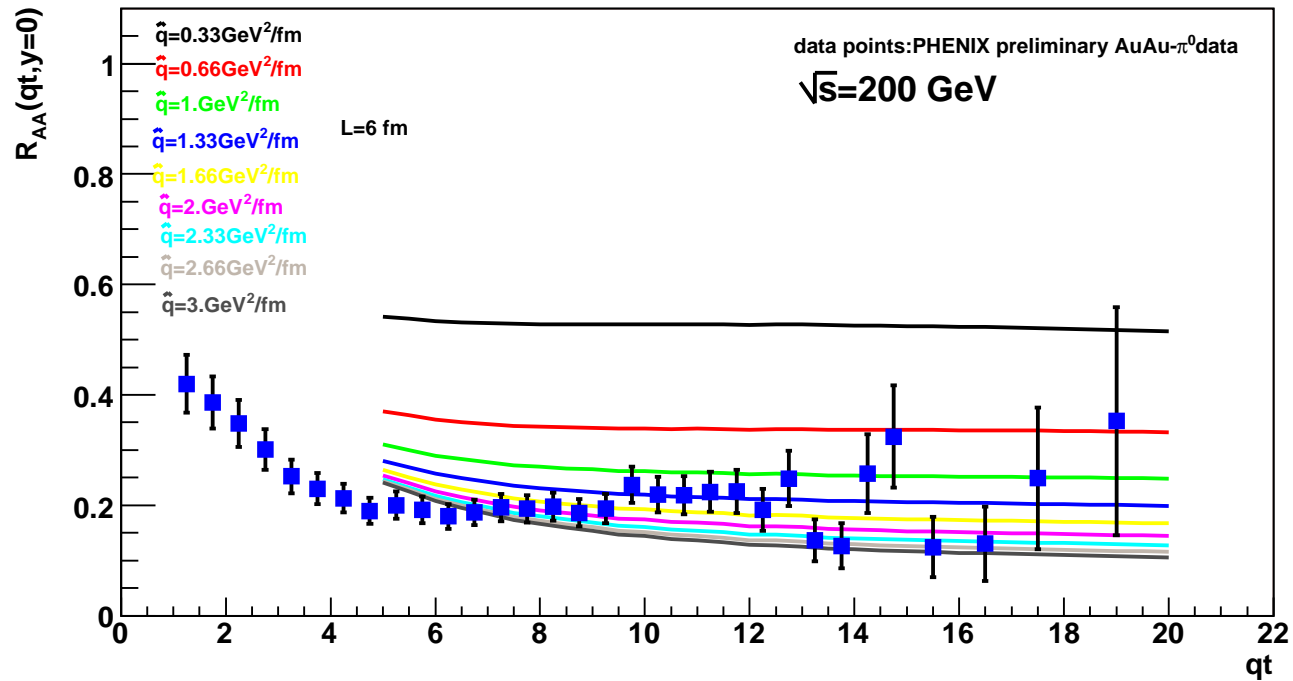
- CTEQ 4L pdf's, LO.

$$Q_{\text{fac}} = Q_{\text{renor}} = k_{T_{\text{parton}}}$$

- Fragmentation scale:

- k_T hadron/parton
 → different $K(\sqrt{s})$
 -see **Eskola et al(Nucl.Phys.A713(2003))**

Nuclear Modification Factor



The fragmentation scale is the internal parton momentum

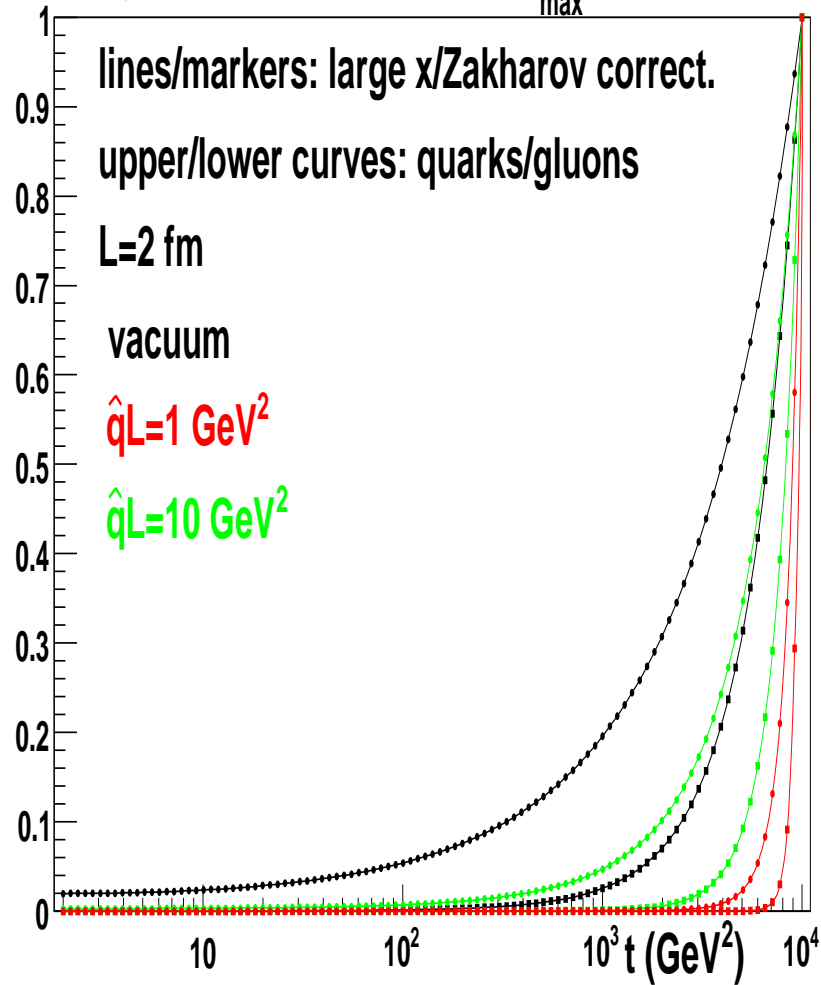
Summary

- Vacuum splitting functions \longrightarrow medium splitting functions
- Sudakov Form Factors in medium
- Medium modified DGLAP evolution via Sudakov Factors
- Medium modified Fragmentation Functions \longrightarrow A code will be soon publicly available!
- Some phenomenological applications:
 - Perturbative convolution \longrightarrow Particle distributions
 - Comparison to experimental data: R_{AA}
 - we determine the value of the transport coefficient \hat{q} for a fixed $L=6\text{fm} \longrightarrow \hat{q} \simeq 1\text{GeV}^2/\text{fm}$
- Step towards a medium-modified parton branching

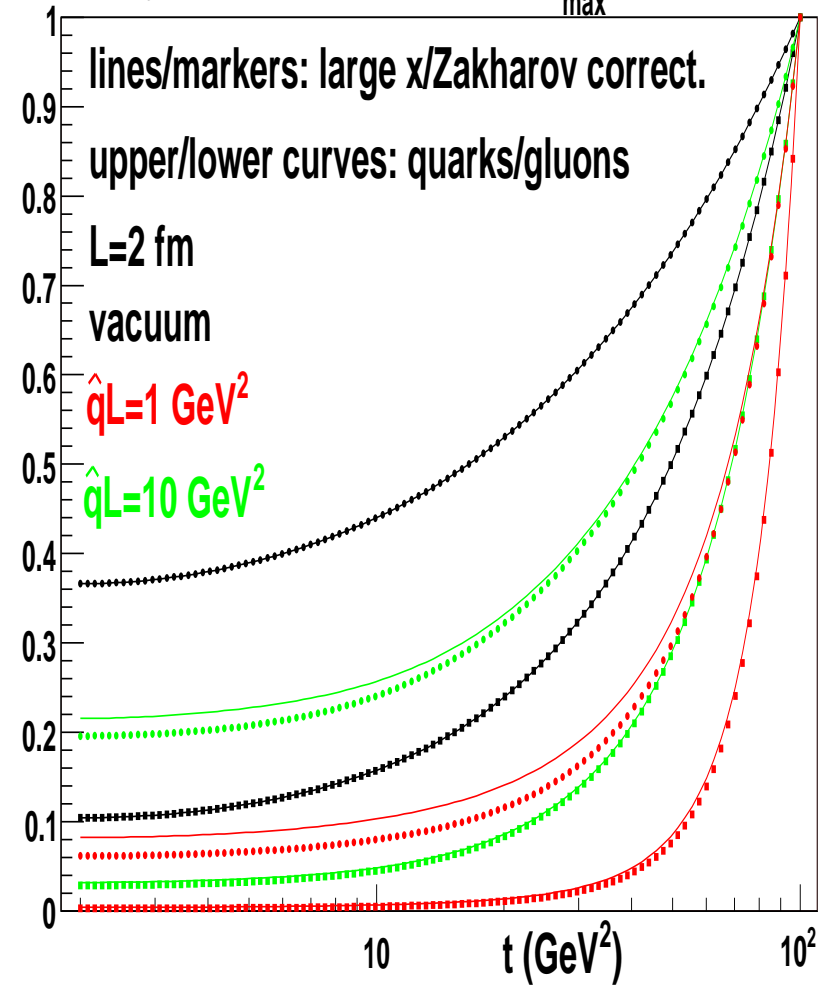
EXTRAS

Medium Modified Sudakov Factors

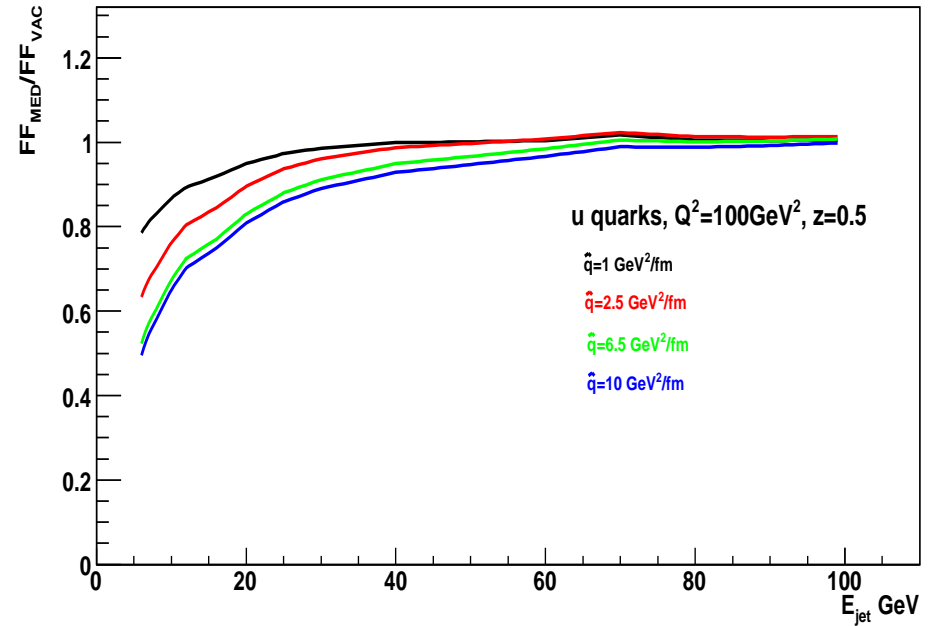
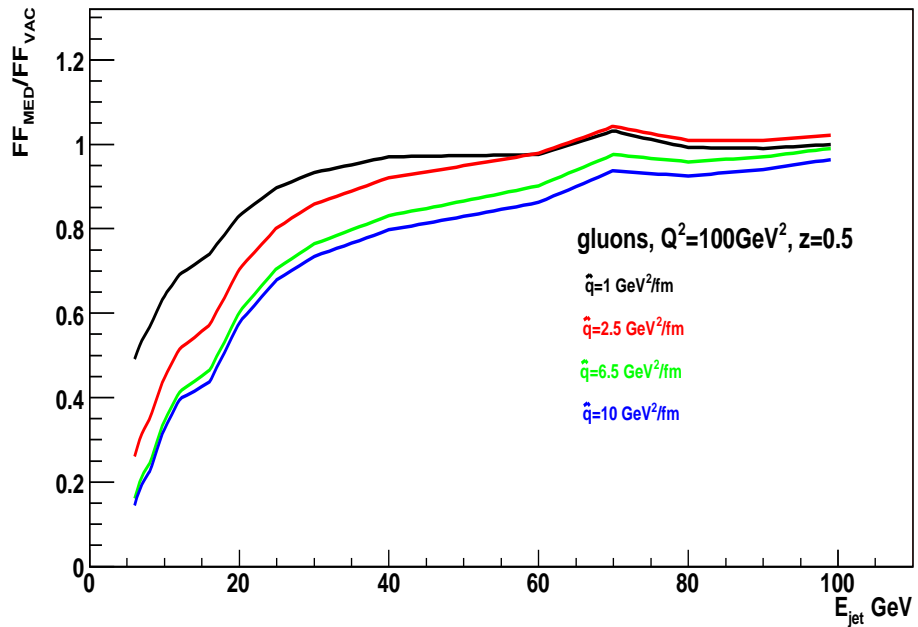
Probability of no emission between $t = E^2 = 10^4 \text{ GeV}^2$ and t_{max}



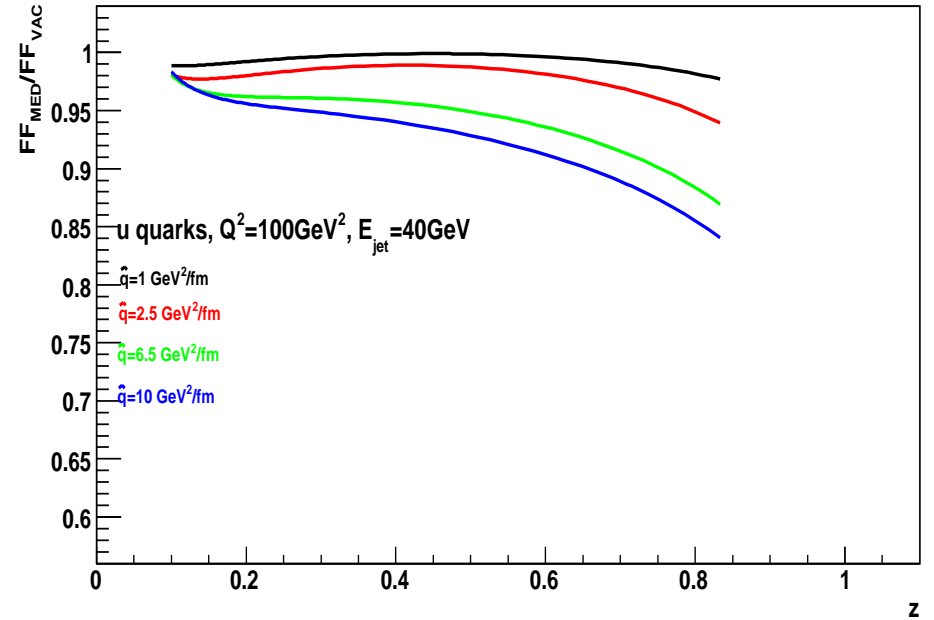
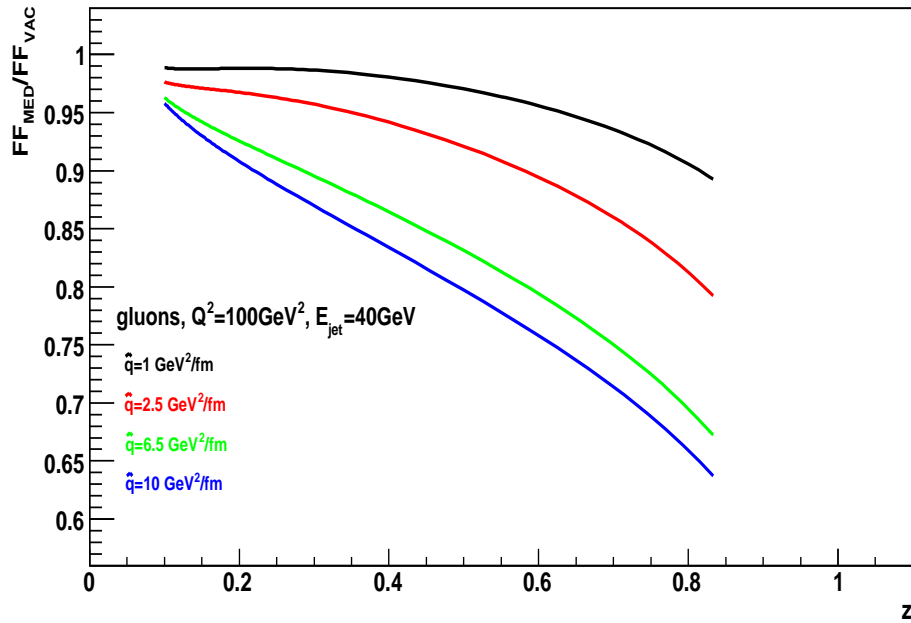
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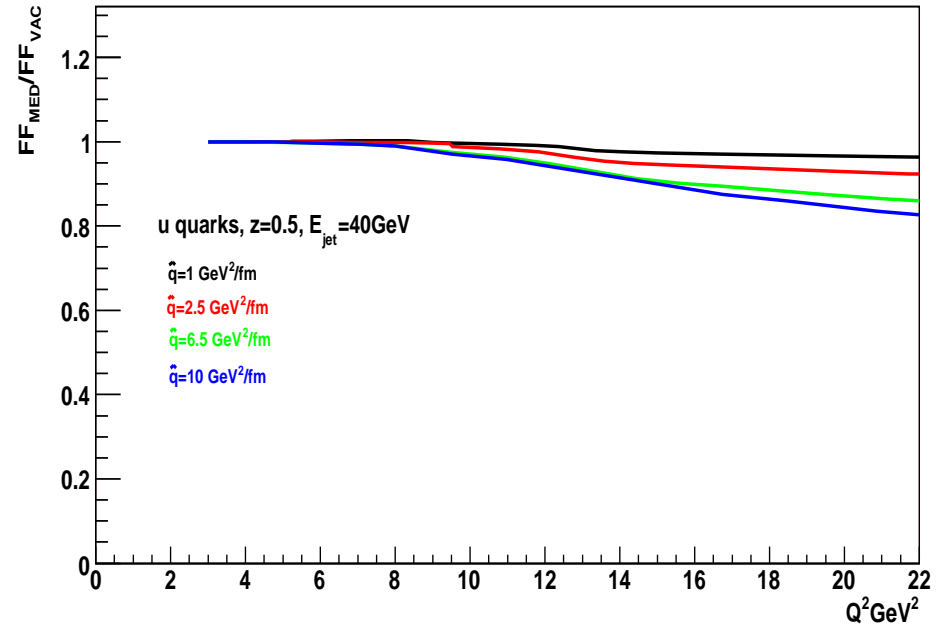
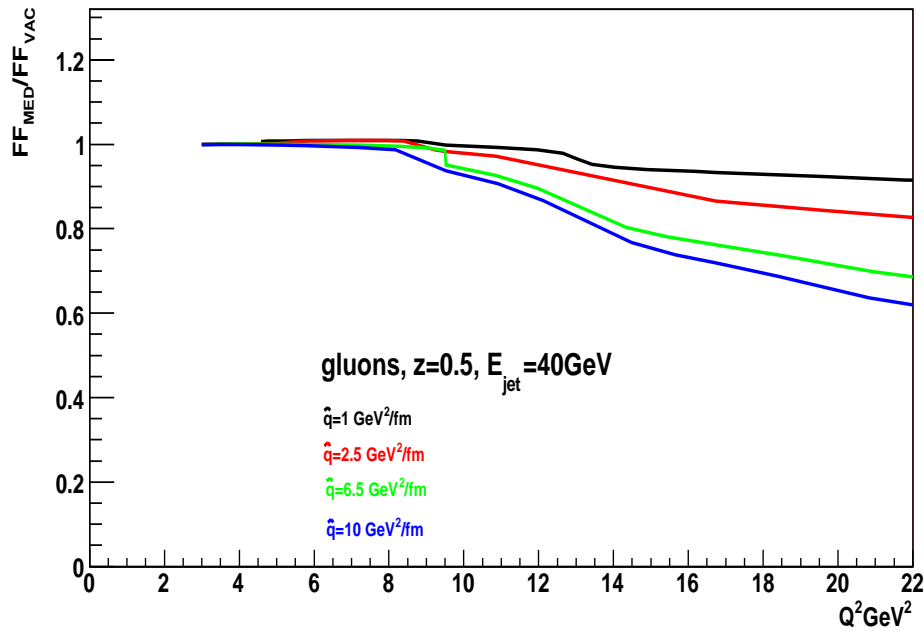
Medium Modified Fragmentation Functions



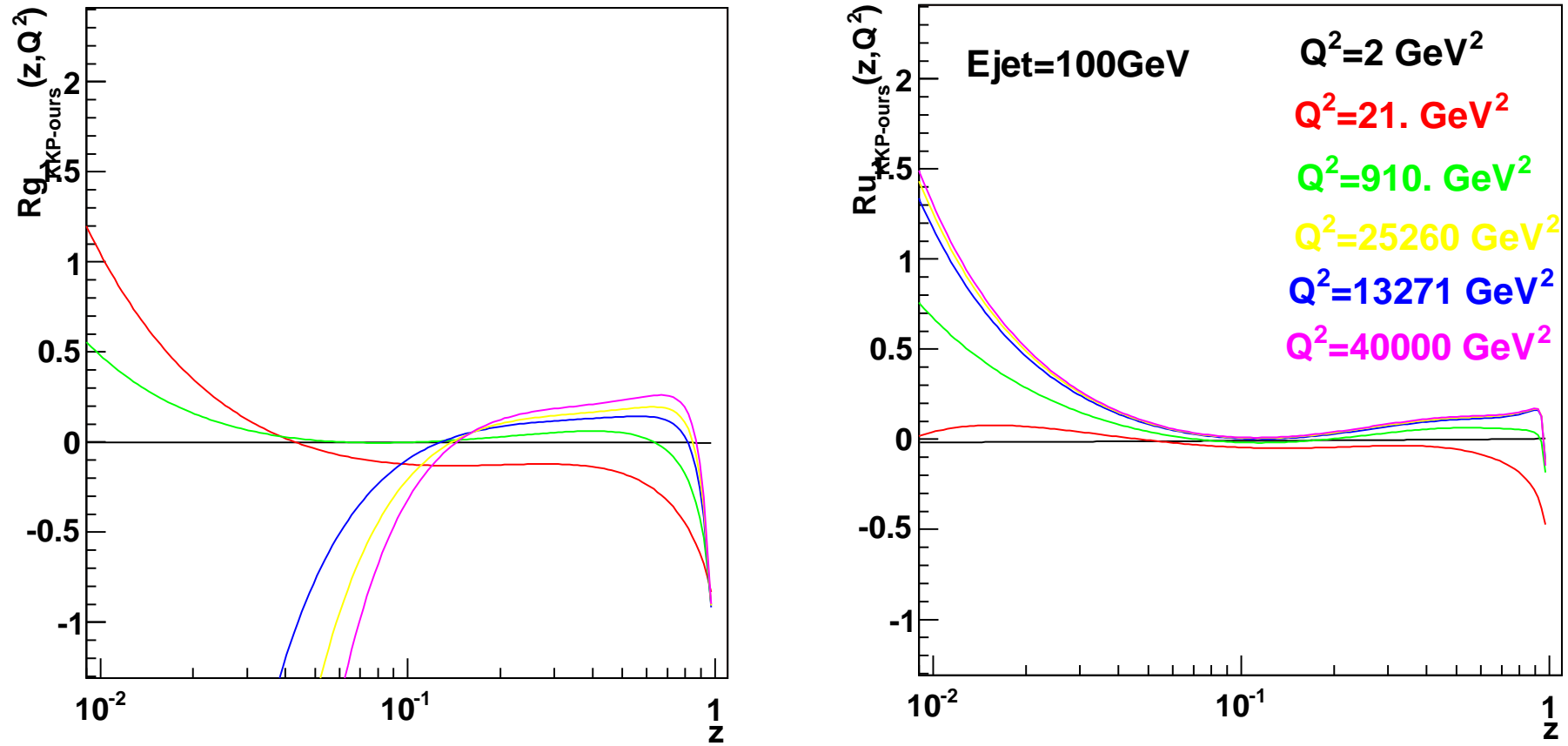
Medium Modified Fragmentation Functions



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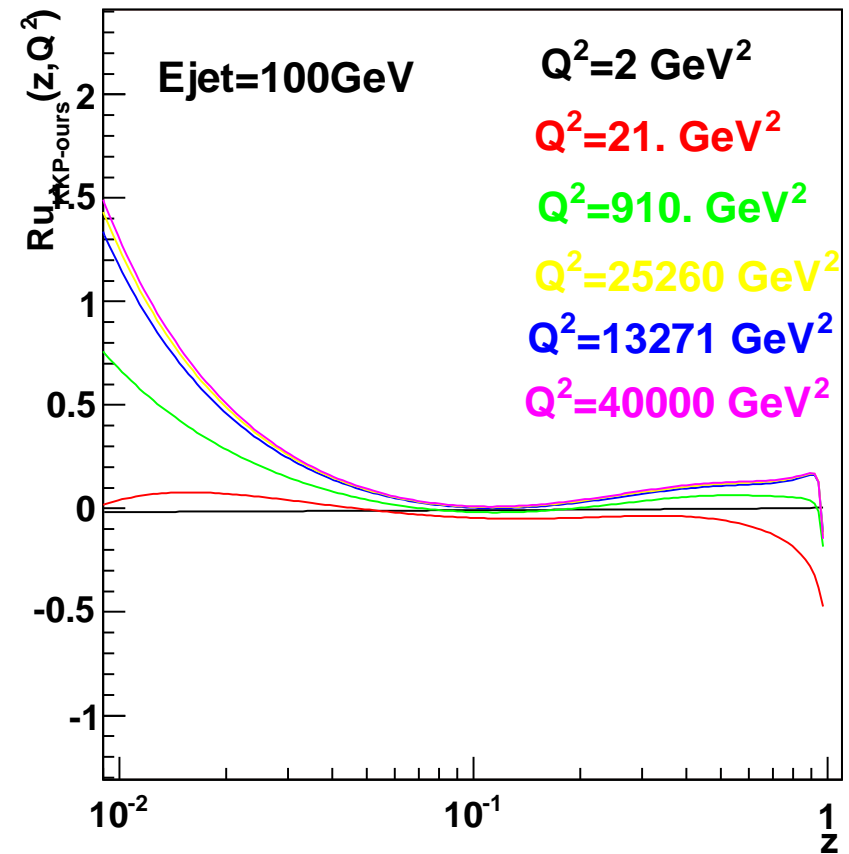
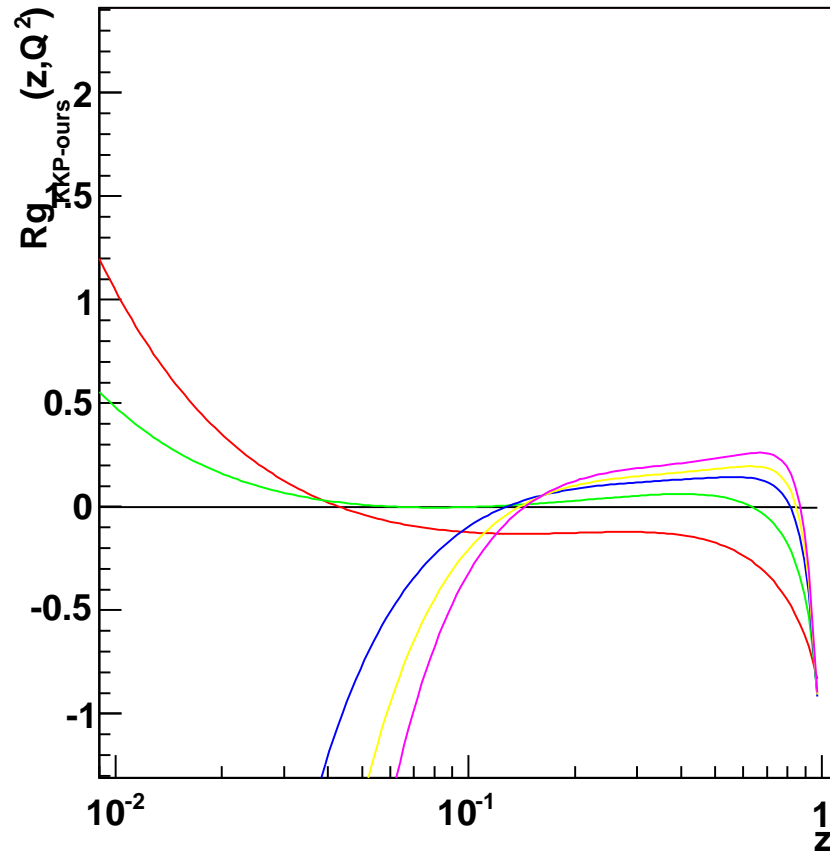


Medium Modified Fragmentation Functions



The same but for an initial parton of energy $E_{jet} = 100 \text{ GeV}$.

Medium Modified Fragmentation Functions



The same but for an initial parton of energy $E_{jet} = 100 \text{ GeV}$.