

Double Beta and TeV Scale Physics

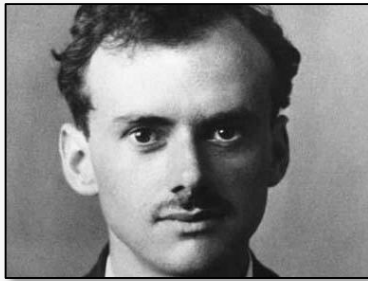
Frank Deppisch

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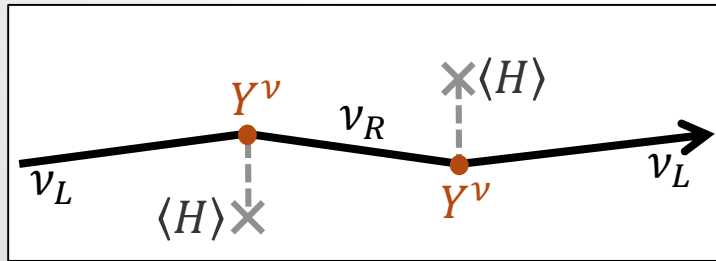
University College London

Dirac vs Majorana

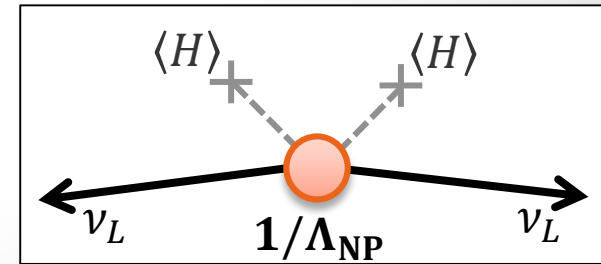
- ▶ Origin of neutrino masses beyond the Standard Model
- ▶ Two possibilities to define neutrino mass



Dirac mass analogous to other fermions but with $m_\nu / \Lambda_{EW} \approx 10^{-12}$ couplings to Higgs



Majorana mass, using only a left-handed neutrino
 → Lepton Number Violation



Beta Decays and ν Mass

- ▶ Single beta decay

$$(A, Z) \rightarrow (A, Z + 1) + e^- + \bar{\nu}_e$$

- Tritium decay, KATRIN: $m_\beta \approx 0.2 \text{ eV}$
- Project 8: Atomic Tritium + Cyclotron Radiation Spectroscopy: $m_\beta \approx 0.05 \text{ eV}$
- HOLMES: e^- capture in ^{163}Ho : $m_\beta \approx 0.1 \text{ eV?}$

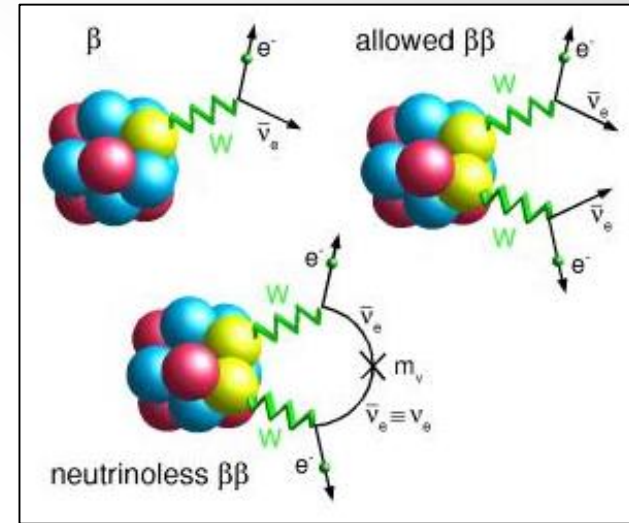
- ▶ Allowed double beta ($2\nu\beta\beta$) decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}_e$$

- ▶ Neutrinoless double beta ($0\nu\beta\beta$) decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

- Violation of lepton number
- Mediated by Majorana neutrinos

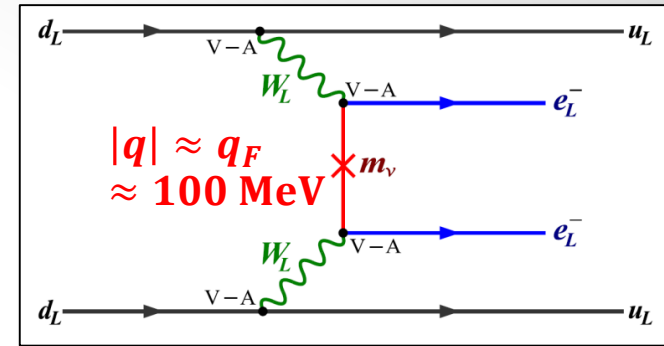


Three Active Neutrinos

▶ Half-life

$$T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$$

▶ Particle Physics



$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 U_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + m_{\nu_i}}{q^2 - m_{\nu_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4q^2} \sum_{i=1}^3 U_{ei}^2 m_{\nu_i} \rightarrow m_{\beta\beta}$$

▶ Atomic Physics

- Leptonic phase space $G^{0\nu} \propto Q^5$

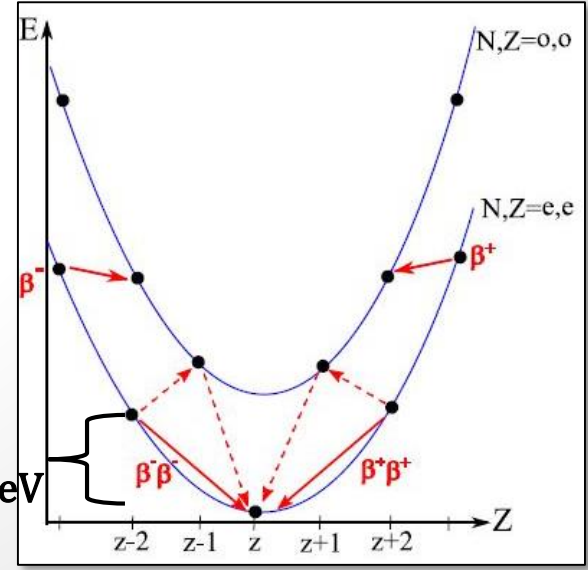
▶ Nuclear Physics

- Nuclear transition matrix element $M^{0\nu} \approx 1$

$$T_{1/2}^{-1} \propto |m_{\beta\beta}|^2 q_F^2 G_F^4 Q^5$$

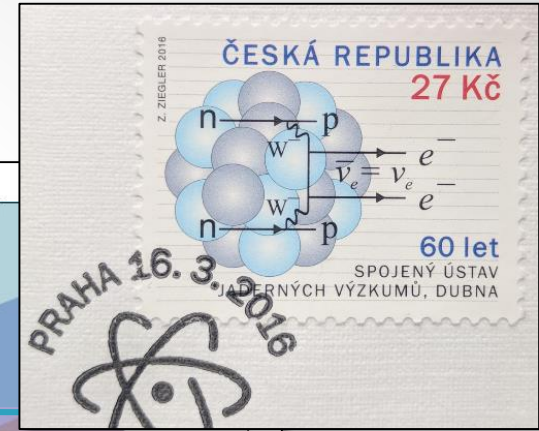
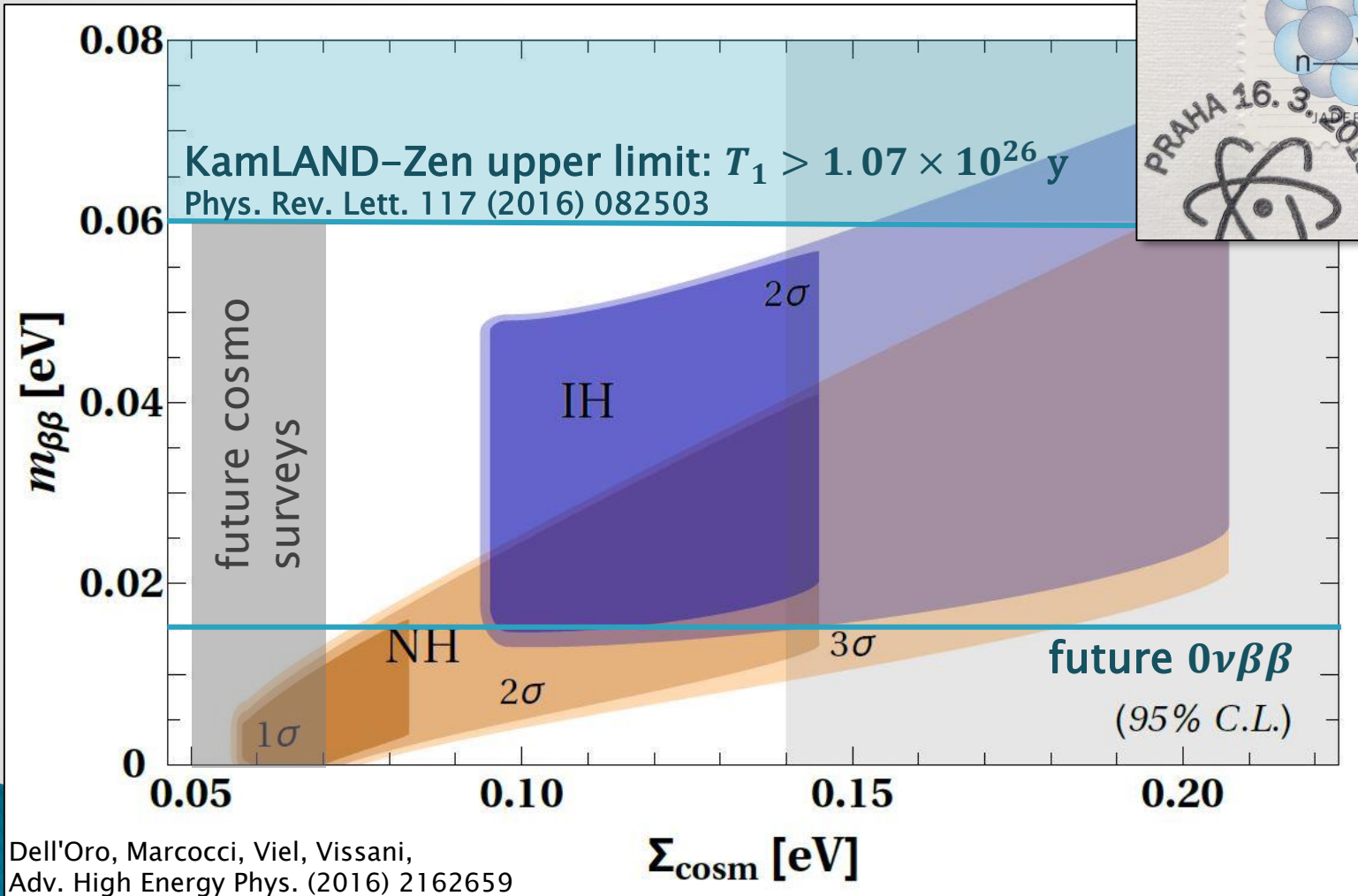
$$\frac{10^{25} y}{T_{1/2}} \approx \left(\frac{|m_{\beta\beta}|}{\text{eV}} \right)^2$$

$$Q + 2m_e \approx 3-5 \text{ MeV}$$



Three Active Neutrinos

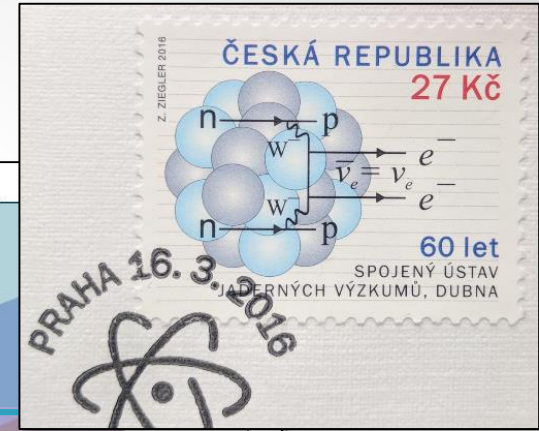
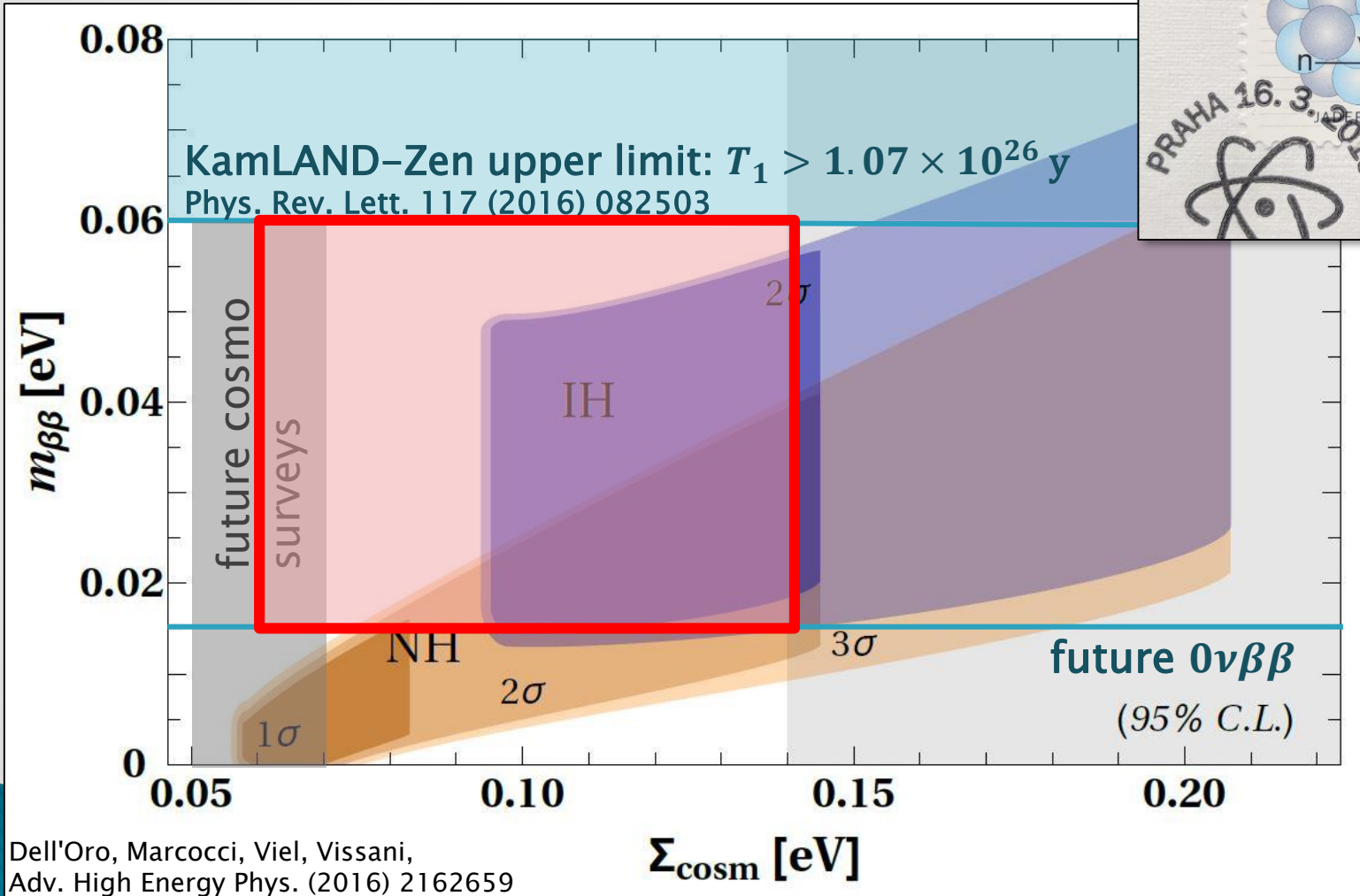
▶ Effective $0\nu\beta\beta$ Mass



Dell'Oro, Marcocci, Viel, Vissani,
Adv. High Energy Phys. (2016) 2162659

Three Active Neutrinos

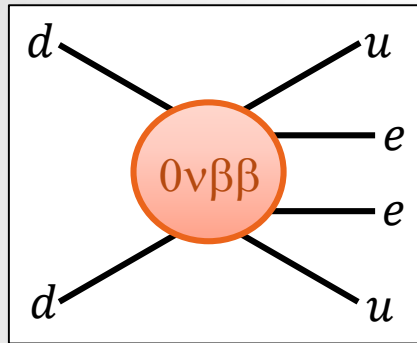
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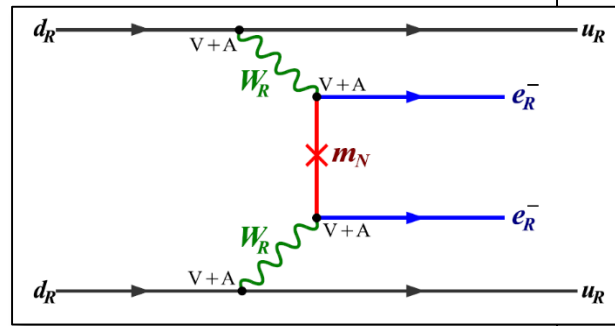
Dell'Oro, Marcocci, Viel, Vissani,
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New Physics and $0\nu\beta\beta$

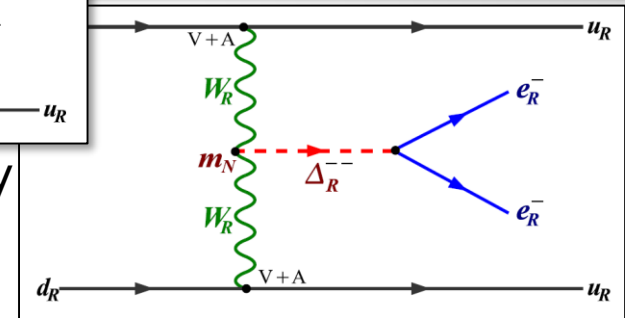
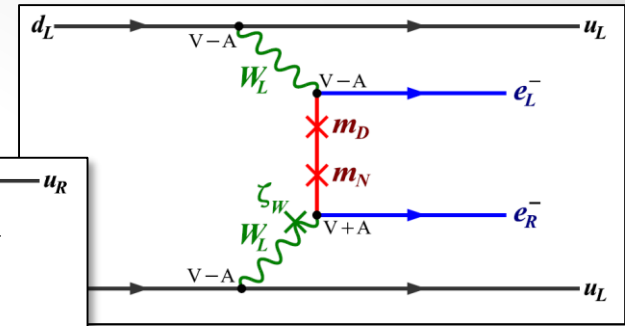
► Plethora of New Physics scenarios



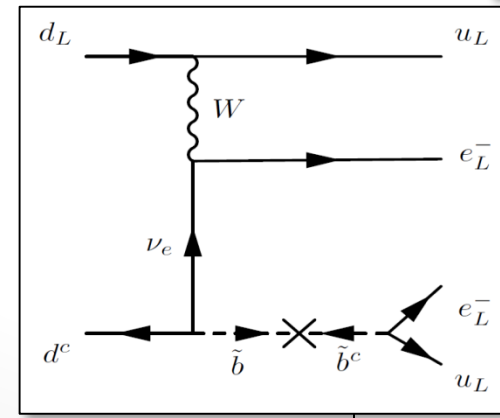
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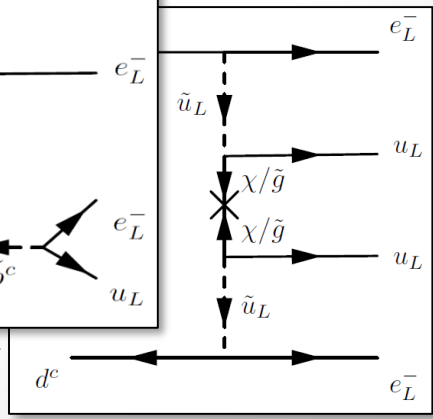
Left-Right Symmetry



$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$



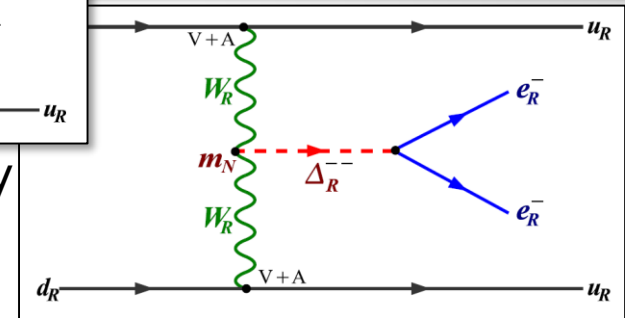
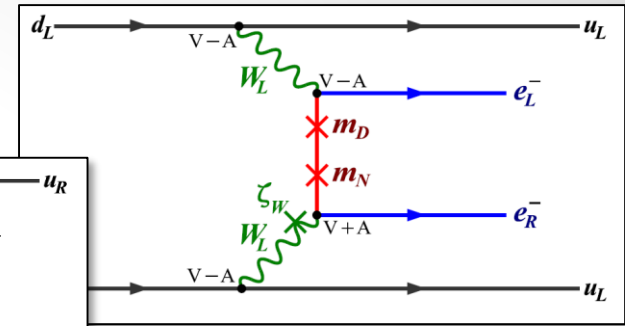
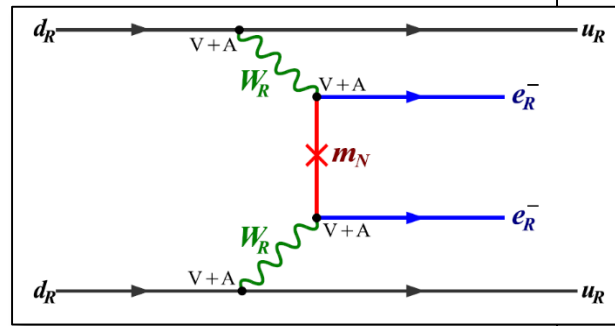
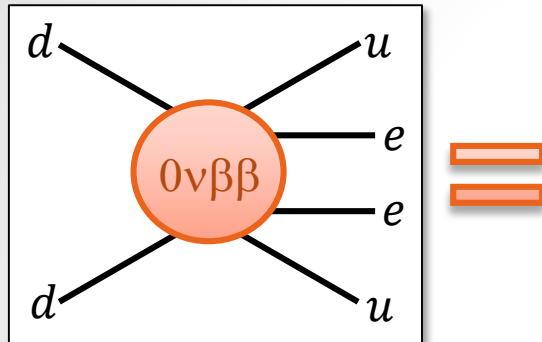
R-Parity Violating SUSY



- Extra Dimensions
- Majorons
- Leptoquarks
- ...

New Physics and $0\nu\beta\beta$

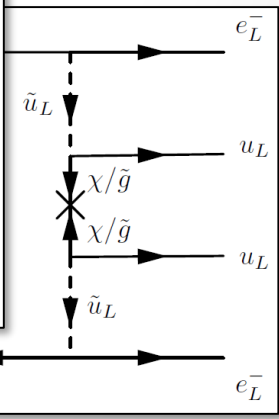
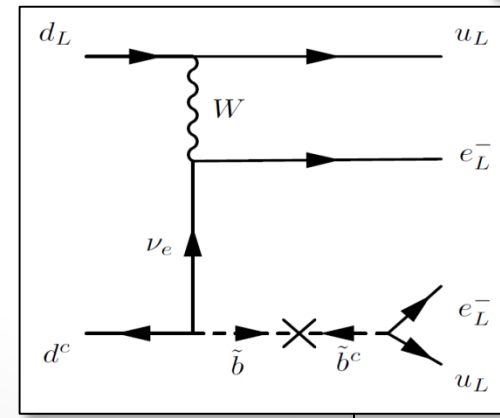
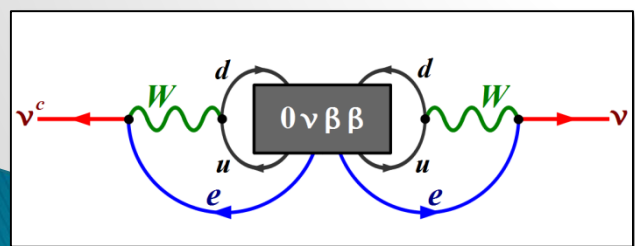
- ▶ Plethora of New Physics scenarios



Left-Right Symmetry

$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$

- ▶ Neutrinos still Majorana

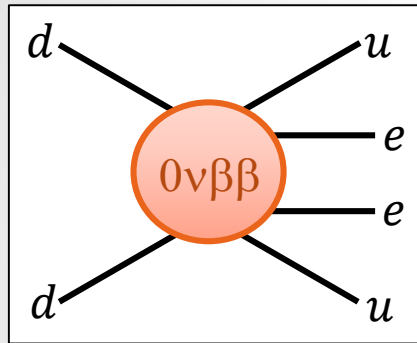


R-Parity Violating SUSY

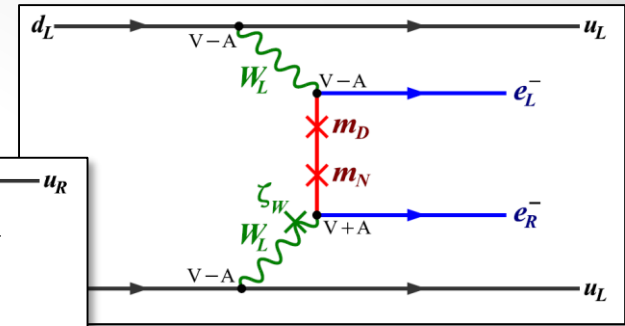
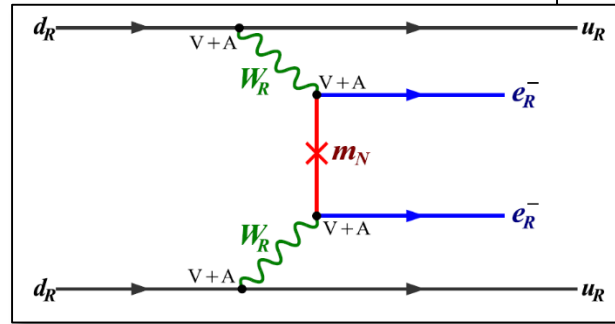
- Extra Dimensions
- Majorons
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- ...

New Physics and $0\nu\beta\beta$

Examples in Left-Right Symmetry



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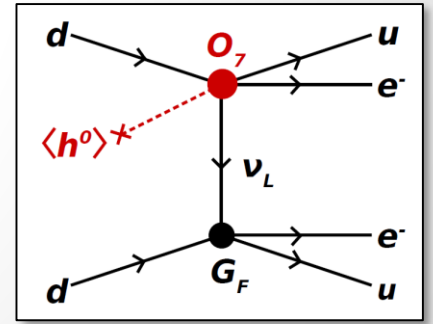
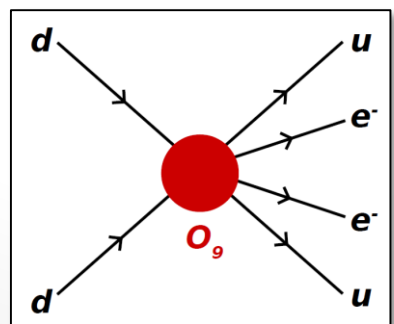


$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$

$$\epsilon_3^{RRZ} = \sum_{i=1}^3 V_{ei}^2 \frac{m_p}{m_N} \frac{m_W^4}{m_{WR}^4} \approx \frac{10^{-8}}{(\Lambda/1 \text{ TeV})^5}$$

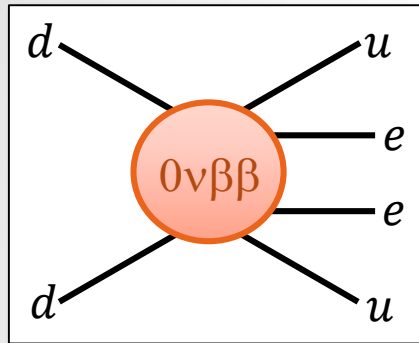
$$\epsilon_{V-A}^{V+A} = \sum_{i=1}^3 U_{ei} W_{ei} \tan \zeta_W \approx \frac{10^{-9}}{(\Lambda/10 \text{ TeV})^3}$$

- ▶ $0\nu\beta\beta$ probes the TeV scale
- ▶ Limits on 6D and 9D eff. operators

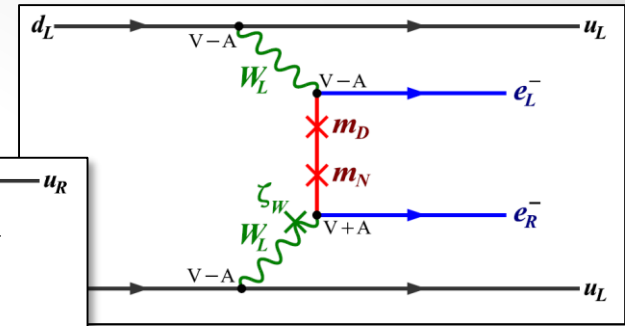
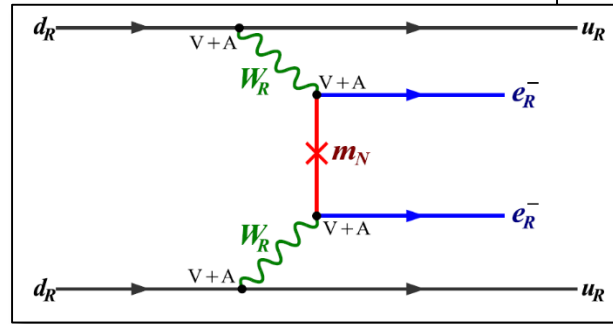


New Physics and $0\nu\beta\beta$

Examples in Left-Right Symmetry



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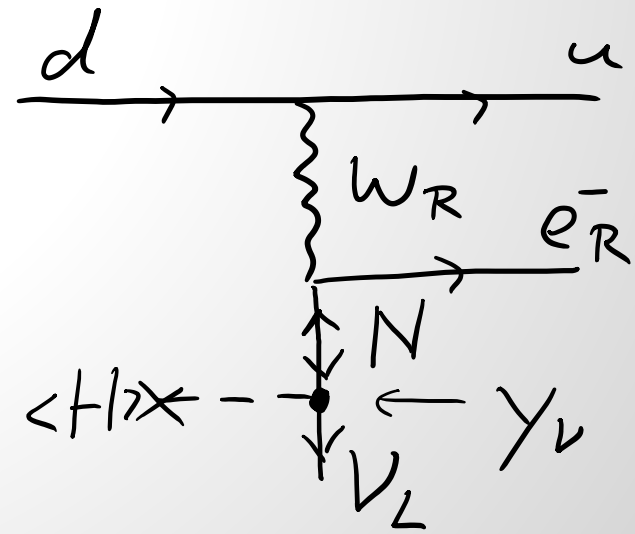
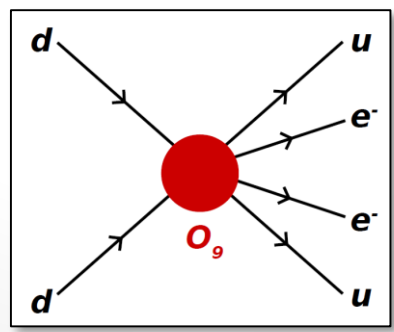


$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$

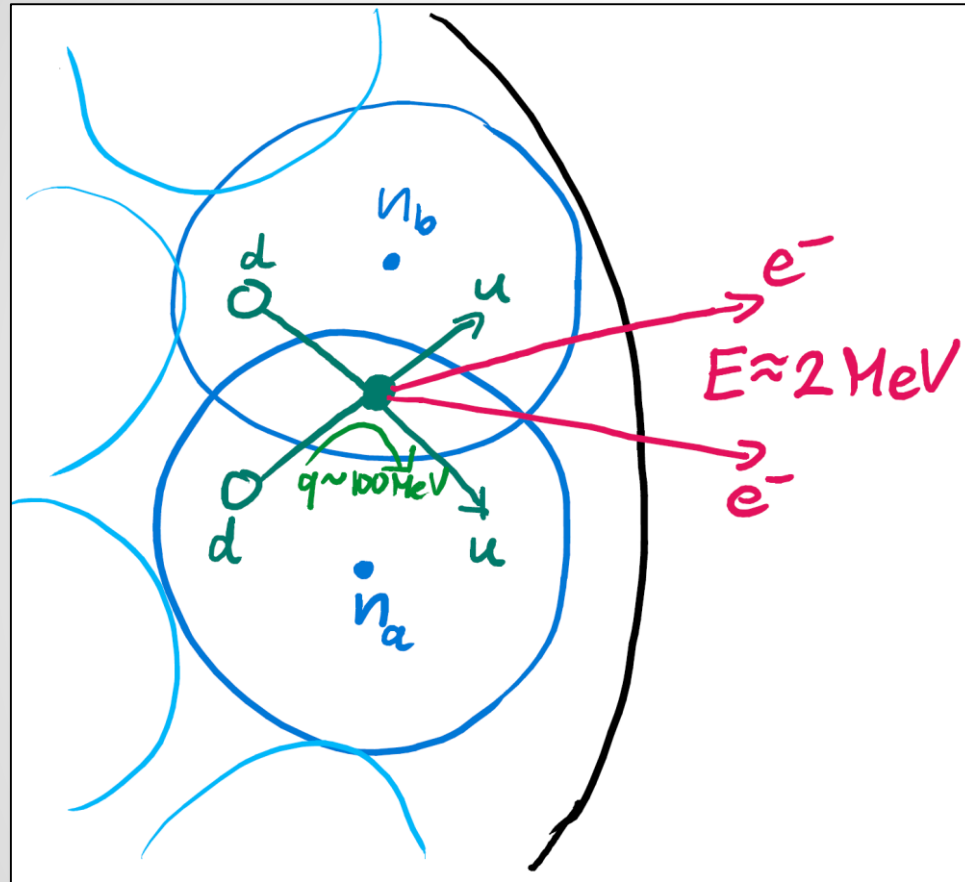
- ▶ $0\nu\beta\beta$ probes the TeV scale
- ▶ Limits on 6D and 9D eff. operators

$$\epsilon_3^{RRZ} = \sum_{i=1}^3 V_{ei}^2 \frac{m_p}{m_N} \frac{m_W^4}{m_{WR}^4} \approx \frac{10^{-8}}{(\Lambda/1 \text{ TeV})^5}$$

$$\epsilon_{V-A}^{V+A} = \sum_{i=1}^3 U_{ei} W_{ei} \tan \zeta_W \approx \frac{10^{-9}}{(\Lambda/10 \text{ TeV})^3}$$



Short-Range Mechanisms



► Evaluation of limits on short-range operators

(Graf, FFD, Iachello, Kotila, PRD 98, 095023)

- General parton level operators (Paes et al. '01)
- Nucleon currents

$$\langle p | \bar{u}(1 \pm \gamma_5)d | n \rangle = \bar{N}\tau^+ [F_S(q^2) \pm F_{PS}(q^2)\gamma_5] N',$$

$$\langle p | \bar{u}\gamma^\mu(1 \pm \gamma_5)d | n \rangle = \bar{N}\tau^+ \left[F_V(q^2)\gamma^\mu - i\frac{F_W(q^2)}{2m_p}\sigma^{\mu\nu}q_\nu \right] N'$$

$$\pm \bar{N}\tau^+ \left[F_A(q^2)\gamma^\mu\gamma_5 - \frac{F_P(q^2)}{2m_p}\gamma_5q^\mu \right] N',$$

- Form factors with enhancement for

$$F_{PS}(q^2) = \frac{g_{PS}}{(1 + q^2/m_{PS}^2)^2} \frac{1}{1 + q^2/m_\pi^2}, \quad g_{PS} = 349$$

$$F_P(q^2) = \frac{g_A}{(1 + q^2/m_A^2)^2} \frac{1}{1 + q^2/m_\pi^2} \frac{4m_p^2}{m_\pi^2} \left(1 - \frac{m_\pi^2}{m_A^2} \right)$$

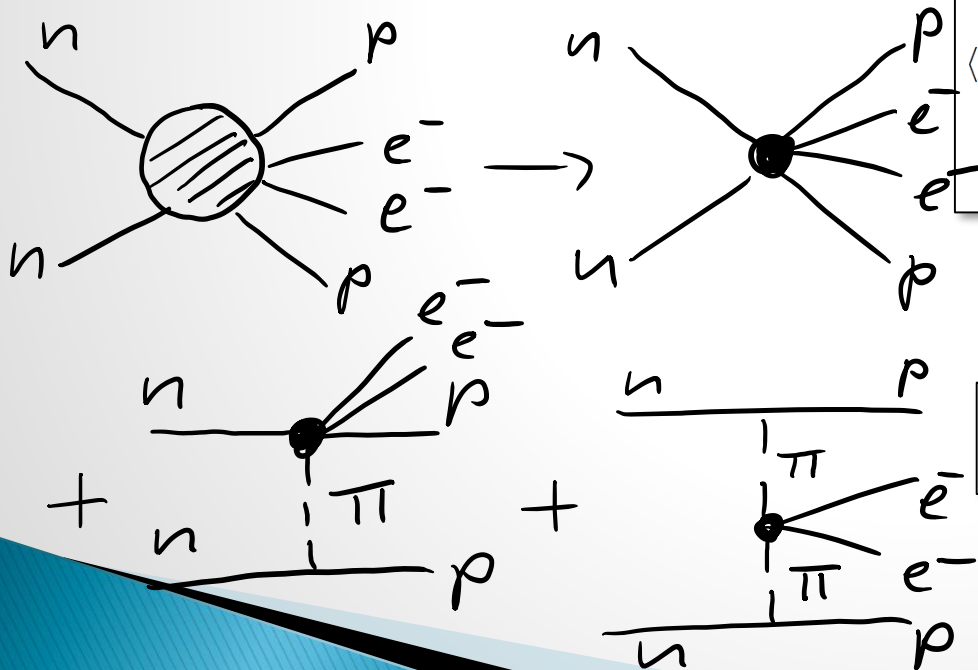
Short-Range Mechanisms

▶ Pion-mediated contributions

- R-parity violating SUSY
(Faessler, Kovalenko, Simkovic, Schwieger, Phys.Rev.Lett. 78 (1997) 183)
- Chiral EFT with Pion-operators from Lattice QCD
(Cirigliano, Dekens, de Vries, Graesser, Mereghetti, JHEP 1812 (2018) 097)

▶ Evaluation of limits on short-range operators

- (Graf, FFD, Iachello, Kotila, PRD 98, 095023)
- General parton level operators (Paes et al. '01)
- Nucleon currents



$$\langle p | \bar{u}(1 \pm \gamma_5)d | n \rangle = \bar{N}\tau^+ [F_S(q^2) \pm F_{PS}(q^2)\gamma_5] N',$$

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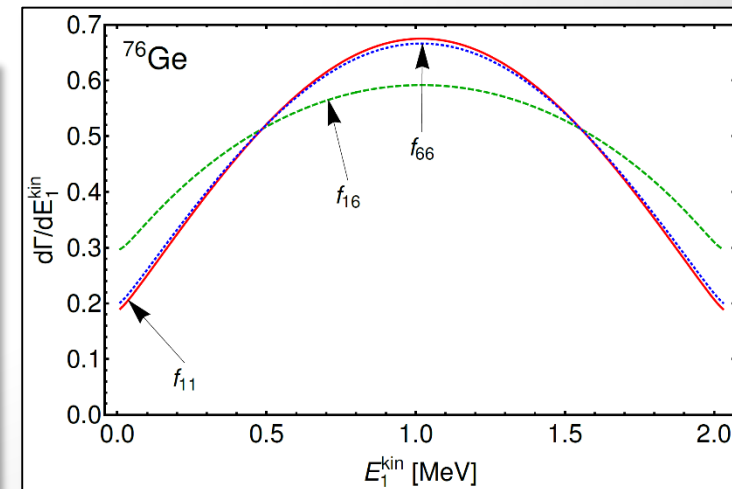
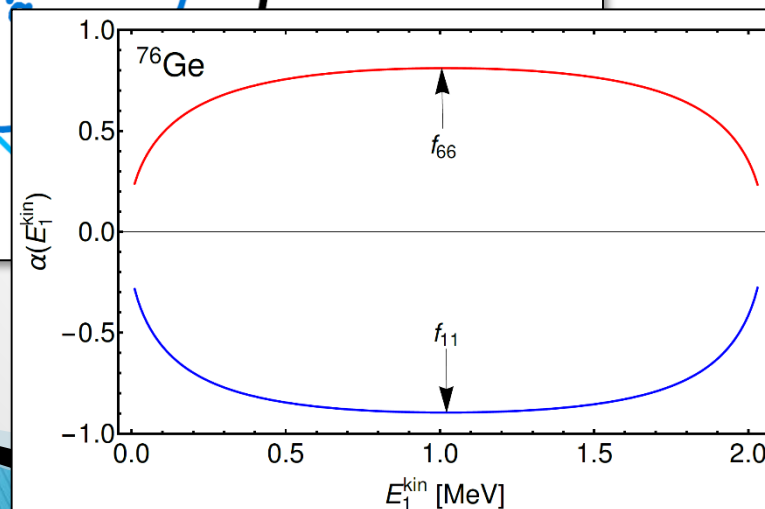
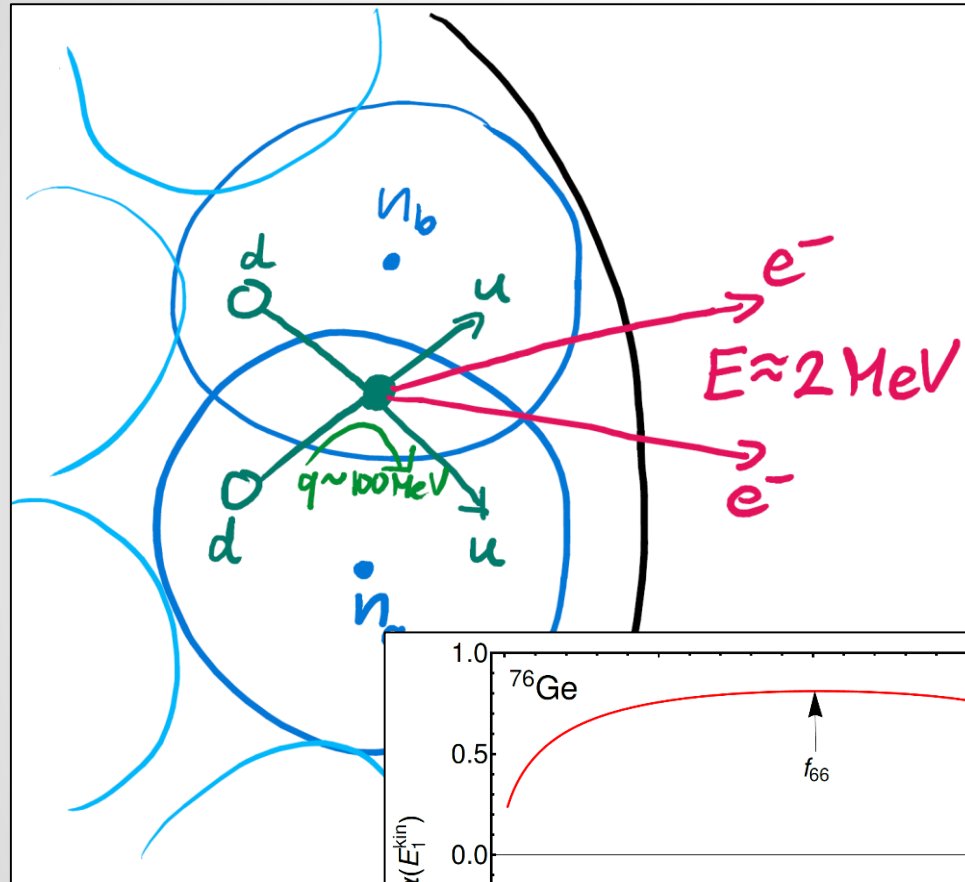
- Form factors with enhancement for

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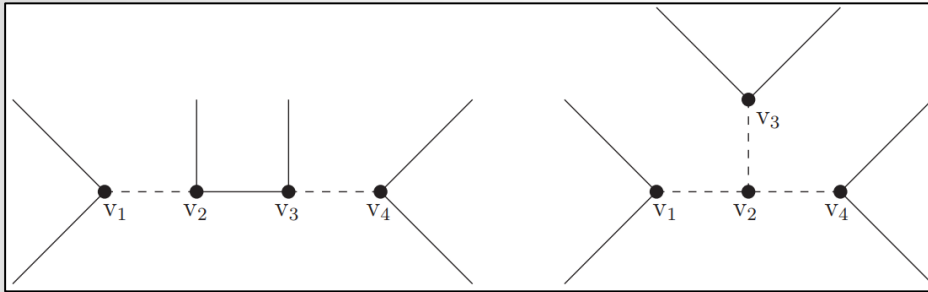
$$F_P(q^2) = \frac{g_A}{(1 + q^2/m_A^2)^2} \frac{1}{1 + q^2/m_\pi^2} \frac{4m_p^2}{m_\pi^2} \left(1 - \frac{m_\pi^2}{m_A^2} \right)$$

Short-Range Mechanisms

- ▶ Evaluation of limits on short-range operators
 - (Graf, FFD, Iachello, Kotila, PRD 98, 095023)
 - Evaluation of additional NMEs
 - Numerical determination of e^- wavefunctions (nuclear Coulomb potential and e^- cloud screening \rightarrow e^- energy and angular distribution)



Short-Range Mechanisms



#	Decomposition	Long Range?	Mediator ($U(1)_{em}, SU(3)_c$) S or V_ρ	ψ	S' or V'_ρ
1-i	$(\bar{u}d)(\bar{e})(\bar{e})(\bar{u}d)$	(a)	$(+1, \mathbf{1})$	$(0, \mathbf{1})$	$(-1, \mathbf{1})$
1-ii-a	$(\bar{u}d)(\bar{u})(d)(\bar{e}\bar{e})$		$(+1, \mathbf{8})$ $(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(0, \mathbf{8})$ $(+5/3, \mathbf{3})$ $(+5/3, \mathbf{3})$	$(-1, \mathbf{8})$ $(+2, \mathbf{1})$ $(+2, \mathbf{1})$
1-ii-b	$(\bar{u}d)(d)(\bar{u})(\bar{e}\bar{e})$		$(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(+4/3, \mathbf{3})$ $(+4/3, \mathbf{3})$	$(+2, \mathbf{1})$ $(+2, \mathbf{1})$

- Evaluation of limits on short-range operators (Graf, FFD, Iachello, Kotila, PRD 98, 095023)
 - Improved limits on effective interactions and NP scales

$$\frac{1}{\Lambda_{NP}^5} = \frac{G_F^2}{2m_p} \epsilon_i$$

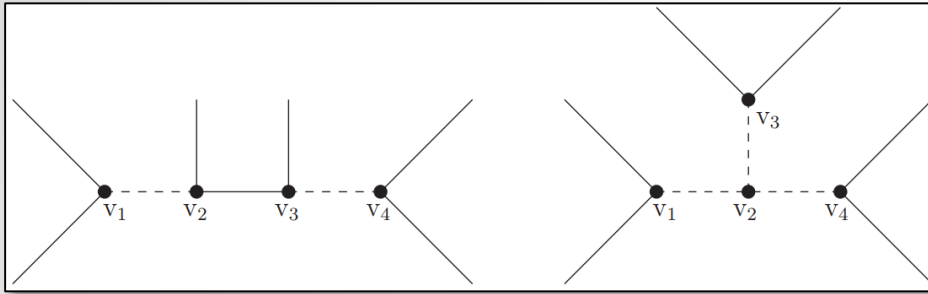
including QCD RGE effects (Mahajan, PRL 112, 031804; Gonzalez, Hirsch, Kovalenko, PRD 93, 013017)

Bonnet, Hirsch, Ota, Winter, JHEP 1303 (2013) 055

$T_{1/2}^{\text{exp}} [y]$	JJj		$J_{\mu\nu}J^{\mu\nu}j$		$J_\mu J^\mu j$		$J^\mu J_{\mu\nu}j^\nu$		$J_\mu Jj^\mu$	
	$ c_1^{XX} $	$ c_1^{LR} $	$ c_2^{XX} $	$ c_3^{XX} $	$ c_3^{LR} $	$ c_4^{XX} $	$ c_4^{LR} $	$ c_5^{XX} $	$ c_5^{RL,LR} $	
^{76}Ge 5.3×10^{25} [71]	0.62	0.36	88	160	260	580	400	25	12	
^{130}Te 2.8×10^{24} [72]	1.4	0.83	200	350	580	1300	880	59	28	
^{136}Xe 1.1×10^{26} [73]	0.24	0.14	32	72	130	250	190	9.6	4.7	

$\times 10^{-10}$

Short-Range Mechanisms



- ▶ Evaluation of limits on short-range operators (Graf, FFD, Iachello, Kotila, PRD 98, 095023)
 - Improved limits on effective interactions and NP scales

#	Decomposition	Long Range?	Mediator ($U(1)_{em}, SU(3)_c$) S or V_ρ	ψ	S' or V'_ρ	Models/Refs./Comments
1-i	$(\bar{u}d)(\bar{e})(\bar{e})(\bar{u}d)$	(a)	$(+1, \mathbf{1})$	$(0, \mathbf{1})$	$(-1, \mathbf{1})$	Mass mechan., RPV [58–60], LR-symmetric models [39], Mass mechanism with ν_S [61], TeV scale seesaw, e.g., [62, 63] [64]
1-ii-a	$(\bar{u}d)(\bar{u})(d)(\bar{e}\bar{e})$		$(+1, \mathbf{8})$ $(+1, \mathbf{1})$	$(0, \mathbf{8})$ $(+5/3, \mathbf{3})$	$(-1, \mathbf{8})$ $(+2, \mathbf{1})$	Bonnet, Hirsch, Ota, Winter, JHEP 1303 (2013) 055
1-ii-b	$(\bar{u}d)(d)(\bar{u})(\bar{e}\bar{e})$		$(+1, \mathbf{8})$ $(+1, \mathbf{1})$	$(+4/3, \mathbf{3})$ $(+4/3, \mathbf{3})$	$(+2, \mathbf{1})$ $(+2, \mathbf{1})$	
2-i-a	$(\bar{u}d)(d)(\bar{e})(\bar{u}\bar{e})$		$(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(+4/3, \mathbf{3})$ $(+4/3, \mathbf{3})$	$(+1/3, \mathbf{3})$ $(+1/3, \mathbf{3})$	RPV [58–60], LQ [65, 66]
2-i-b	$(\bar{u}d)(\bar{e})(d)(\bar{u}\bar{e})$	(b)	$(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(0, \mathbf{1})$ $(0, \mathbf{8})$	$(+1/3, \mathbf{3})$ $(+1/3, \mathbf{3})$	
2-ii-a	$(\bar{u}d)(\bar{u})(\bar{e})(d\bar{e})$		$(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(+5/3, \mathbf{3})$ $(+5/3, \mathbf{3})$	$(+2/3, \mathbf{3})$ $(+2/3, \mathbf{3})$	RPV [58–60], LQ [65, 66]
2-ii-b	$(\bar{u}d)(\bar{e})(\bar{u})(d\bar{e})$	(b)	$(+1, \mathbf{1})$ $(+1, \mathbf{8})$	$(0, \mathbf{1})$ $(0, \mathbf{8})$	$(+2/3, \mathbf{3})$ $(+2/3, \mathbf{3})$	
2-iii-a	$(d\bar{e})(\bar{u})(d)(\bar{u}\bar{e})$	(c)	$(-2/3, \mathbf{3})$ $(-2/3, \mathbf{3})$	$(0, \mathbf{1})$ $(0, \mathbf{8})$	$(+1/3, \mathbf{3})$ $(+1/3, \mathbf{3})$	RPV [58–60]
2-iii-b	$(d\bar{e})(d)(\bar{u})(\bar{u}\bar{e})$		$(-2/3, \mathbf{3})$ $(-2/3, \mathbf{3})$	$(-1/3, \mathbf{3})$ $(-1/3, \mathbf{6})$	$(+1/3, \mathbf{3})$ $(+1/3, \mathbf{3})$	RPV [58–60]

$$= \frac{G_F^2}{2m_p} \epsilon_i$$

QCD RGE effects
[112, 031804; Gonzalez, Panko, PRD 93, 013017]

$$J_\mu J^\mu$$

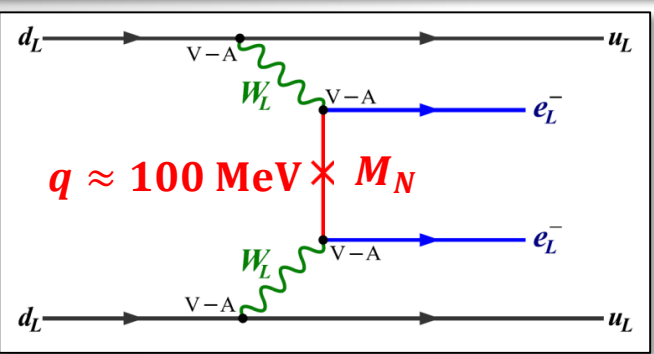
c_5^{XX}	$ c_5^{RL,LR} $
25	12
59	28
9.6	4.7

$\times 10^{-10}$

Heavy Sterile Neutrinos

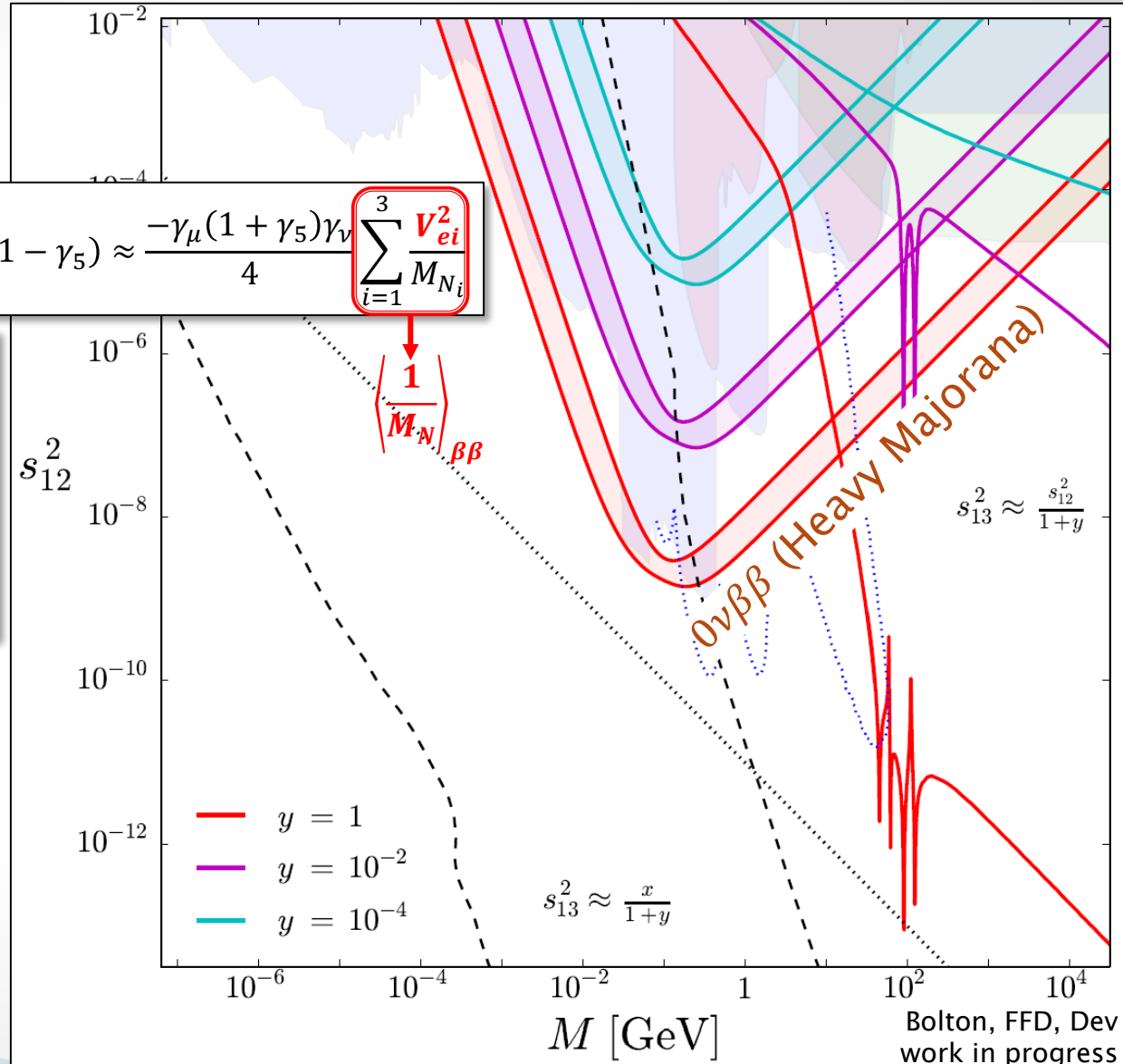
- ▶ with masses larger than ≈ 100 MeV

$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 V_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + M_{N_i}}{q^2 - M_{N_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{-\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4} \sum_{i=1}^3 \frac{V_{ei}^2}{M_{N_i}}$$



Different nuclear matrix elements

- Short-range operator



Interference with $m_{\beta\beta}$

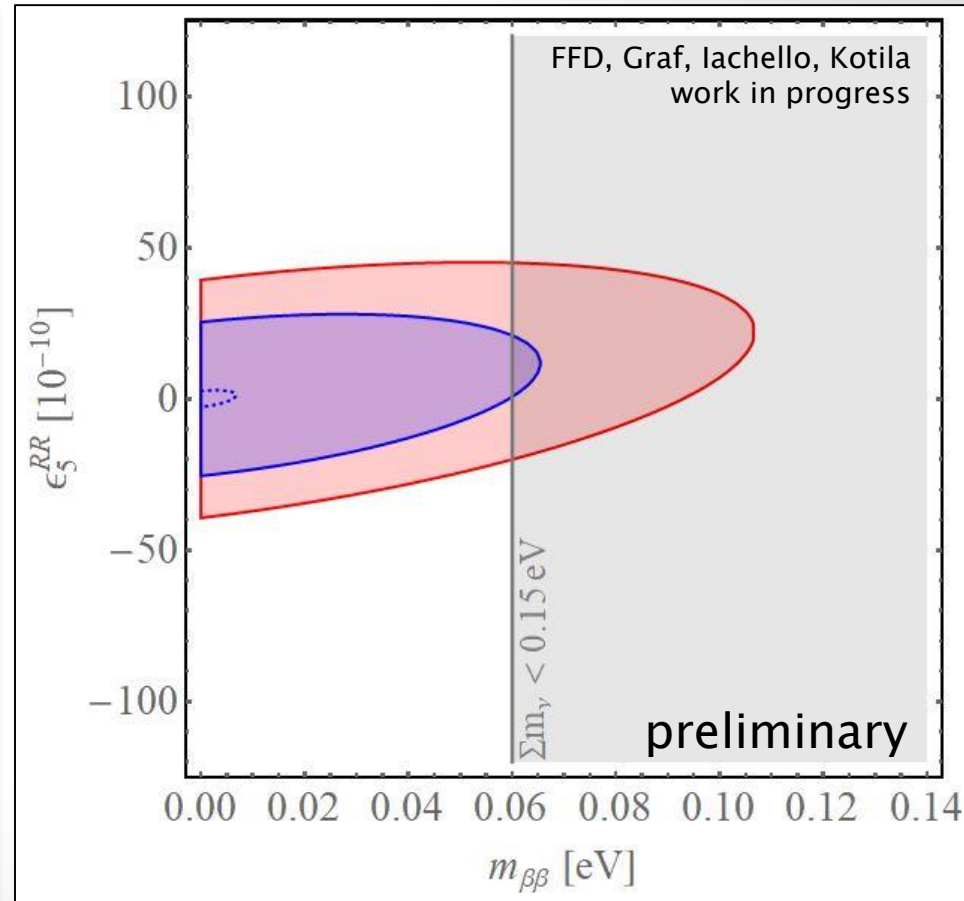
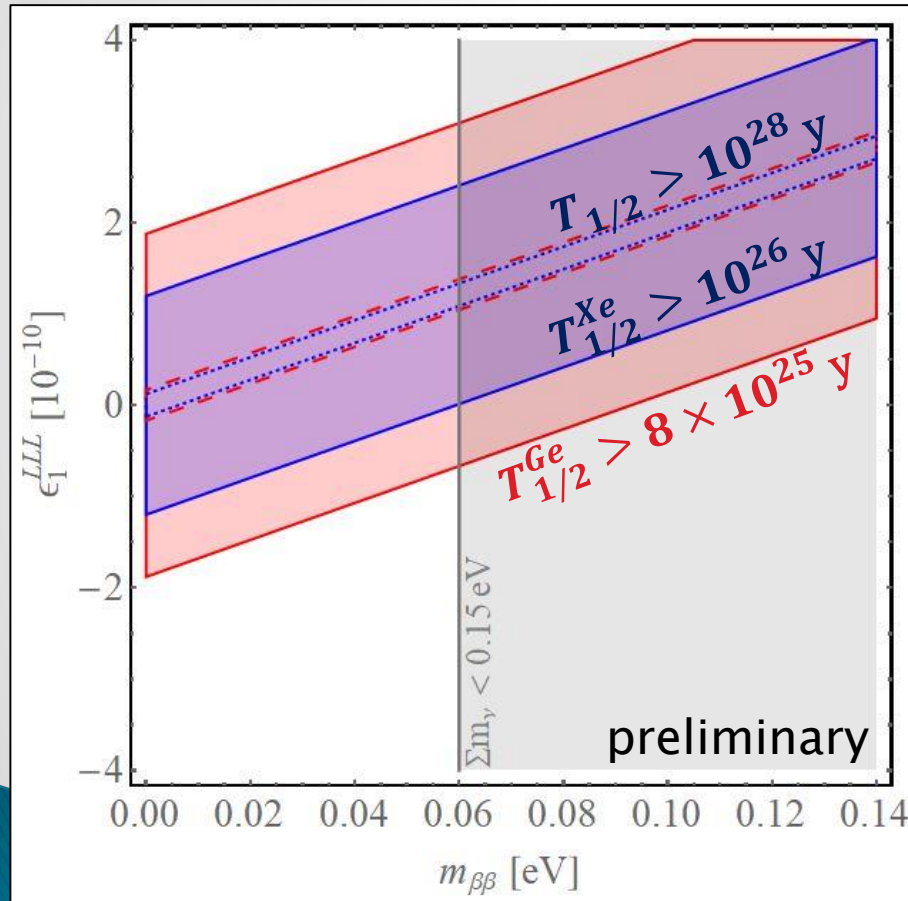
Same lepton current

$$T_{1/2}^{-1} = G_\nu \left| \frac{m_{\beta\beta}}{m_e} \mathcal{M}_\nu + \epsilon \mathcal{M}_\epsilon \right|^2$$

e.g. heavy neutrino exchange
Mitra, Pascoli, Wong, Phys. Rev. D 90 (2014) 093005

Different currents

$$T_{1/2}^{-1} = \frac{|m_{\beta\beta}|^2}{m_e^2} |\mathcal{M}_\nu|^2 G_\nu + |\epsilon|^2 |\mathcal{M}_\epsilon|^2 G_\epsilon + 2\text{Re} \left[\frac{m_{\beta\beta}}{m_e} \epsilon^* \mathcal{M}_\nu \mathcal{M}_\epsilon^* \right] G_{\nu\epsilon}$$



Interference with $m_{\beta\beta}$

Same lepton current

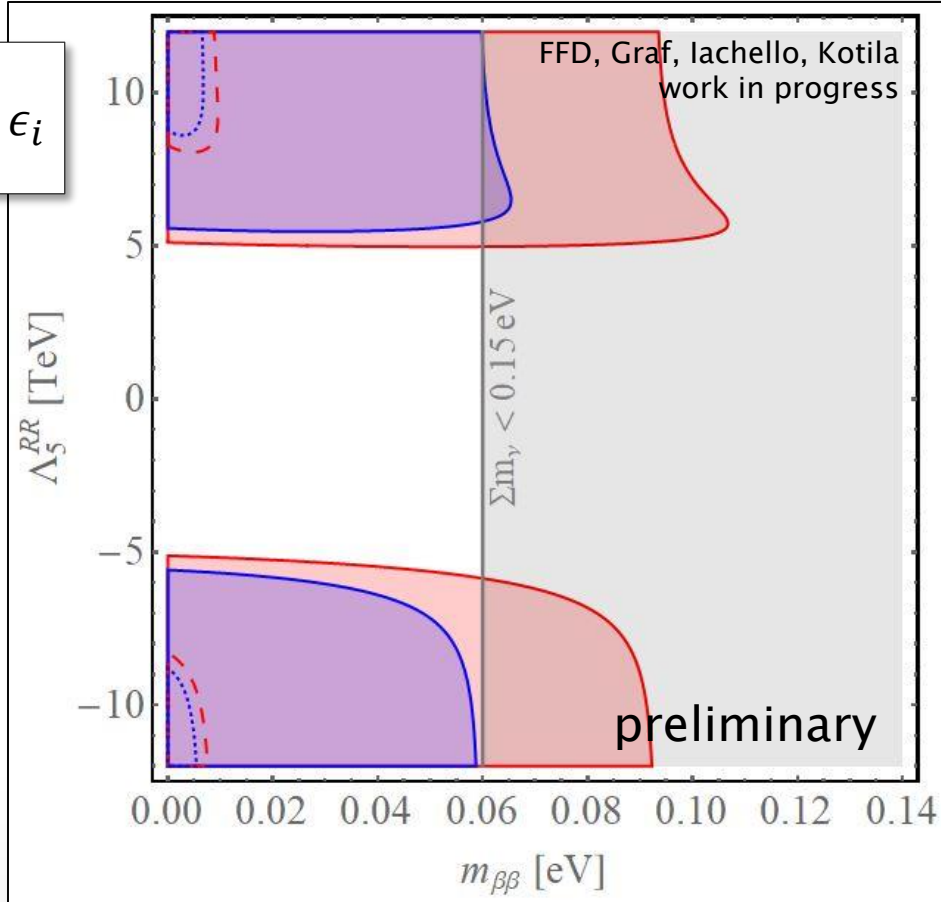
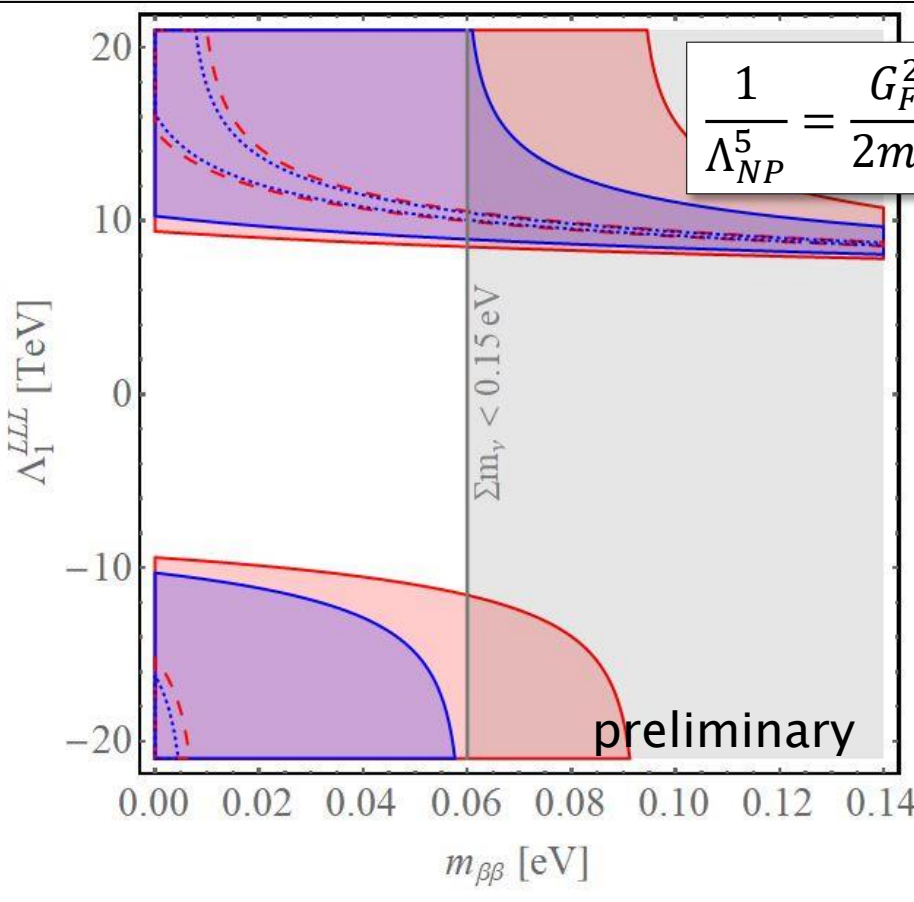
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Different currents

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$$\frac{1}{\Lambda_{NP}^5} = \frac{G_F^2}{2m_p} \epsilon_i$$



Conclusion

▶ Neutrinos much lighter than other fermions

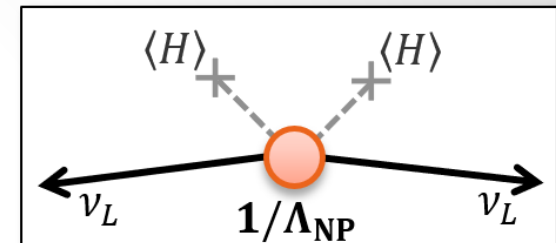
- Dirac or Majorana? Lepton Number Violation?
- Natural suppression of charged LFV?
- Determination of absolute mass scale

▶ $0\nu\beta\beta$ is crucial probe for BSM physics

- New LNV physics at the LHC scale?
- Standard Mass Mechanism?
 - 5-dim operator from LNV at GUT scale
- Experimentally and theoretically challenging

▶ Importance of probing LNV around the TeV scale

- E.g. searches for heavy neutral leptons at the LHC
- Can we rule out mechanisms of neutrino mass generation?
- Impact on baryon asymmetry of the Universe
(FFD, Harz, Hirsch, Phys. Rev. Lett. 112 (2014) 221601)



$$\frac{T_{1/2}^{0\nu\beta\beta}}{10^{28} \text{ y}} \approx \left(\frac{\Lambda_{NP}}{10^{15} \text{ GeV}} \right)^2$$