

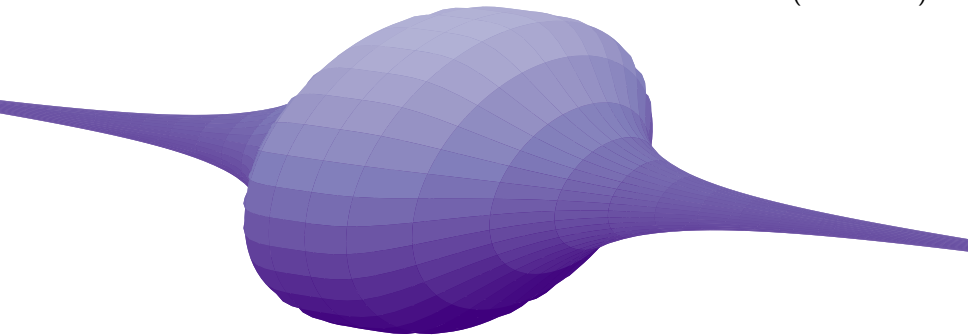
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The University of Manchester

The Geometry of Inflation

arXiv:1806.02431 (PRD 2018)

arXiv:1812.07095 (PRD 2019)



Kieran Finn

with Apostolos Pilaftsis and Sotirios Karamitsos

Monday 1st July 2019

"God always geometrizes" - Plato

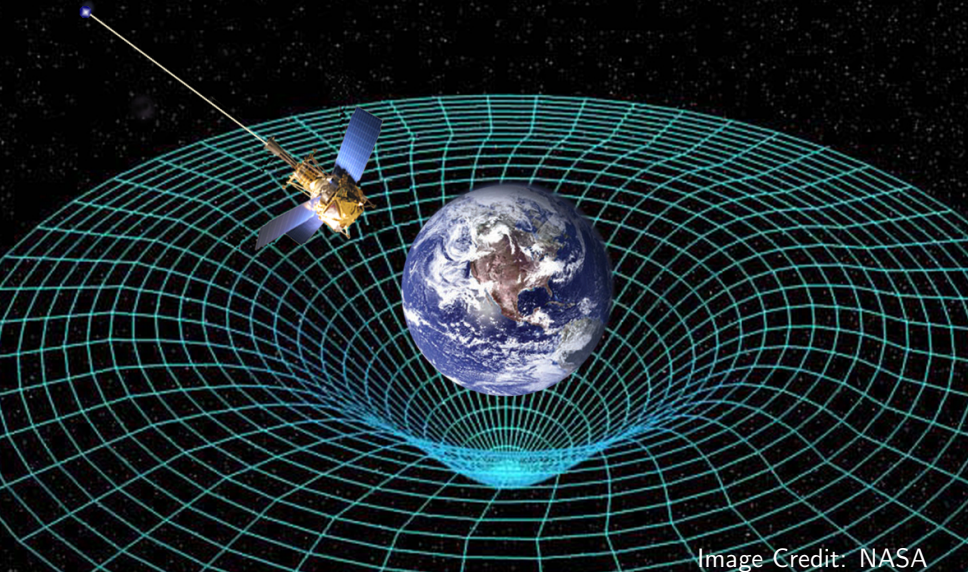


Image Credit: NASA

Geometric Interpretation of Homogeneous Field Theories

Lagrangian

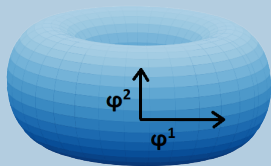
- $L = \frac{1}{2} k_{ab}(\varphi) \dot{\varphi}^a \dot{\varphi}^b$

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Field Space

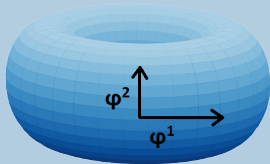


Geometric Interpretation of Homogeneous Field Theories

Lagrangian

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Field Space



Equations of Motion

- $\ddot{\varphi}^a + \Gamma_{bc}^a \dot{\varphi}^b \dot{\varphi}^c = 0$
- $\Gamma_{bc}^a \equiv \frac{1}{2} k^{ad} (k_{bd,c} + k_{cd,b} - k_{bc,d})$

Noether Symmetries

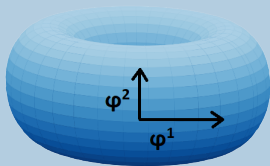
- $\varphi^a \rightarrow \varphi^a + \xi^a(\varphi)$
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Geometric Interpretation of Homogeneous Field Theories

Lagrangian

- $L = \frac{1}{2} k_{ab}(\varphi) \dot{\varphi}^a \dot{\varphi}^b - V(\varphi)$

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Equations of Motion

- $\ddot{\varphi}^a + \Gamma_{bc}^a \dot{\varphi}^b \dot{\varphi}^c = -k^{ab} V_{,b}$
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Noether Symmetries

- $\varphi^a \rightarrow \varphi^a + \xi^a(\varphi)$
- $\nabla_b \xi_c + \nabla_c \xi_b = 0$
- $V_{,a} \xi^a = 0$

The Eisenhart Lift for Field Theories

L.P. Eisenhart: 1928

KF, S. Karamitsos and A. Pilaftsis: 2018 (arXiv:1806.02431)

Original Theory

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Lifted Theory

- $L = \frac{1}{2}k_{ab}(\varphi)\dot{\varphi}^a\dot{\varphi}^b + \frac{1}{2}\frac{1}{V(\varphi)}\dot{\chi}^2$

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Extended Field Space

- $\phi^A = \{\varphi^a, \chi\}$
- $L = \frac{1}{2}G_{AB}\dot{\phi}^A\dot{\phi}^B$

$$G_{AB} = \begin{pmatrix} k_{ab} & 0 \\ 0 & \frac{1}{V(\varphi)} \end{pmatrix}$$

Extension to Non-Homogeneous Field theories

Original Theory

- $\mathcal{L} = \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} k_{ab}(\varphi) \partial_\mu \varphi^a \partial_\nu \varphi^b - V(\varphi) \right)$
- $\nabla^\mu \nabla_\mu \varphi^a + \Gamma_{bc}^a \partial_\mu \varphi^b \partial^\mu \varphi^c = -k^{ab} V_{,b}$

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Extended Theory

- $\mathcal{L} = \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} k_{ab}(\varphi) \partial_\mu \varphi^a \partial_\nu \varphi^b + \frac{1}{2} \frac{1}{V(\varphi)} \nabla_\mu B^\mu \nabla_\nu B^\nu \right)$
- $\nabla^\mu \nabla_\mu \varphi^a + \Gamma_{bc}^a \partial_\mu \varphi^b \partial^\mu \varphi^c = -\frac{1}{2} \left(\frac{\nabla_\nu B^\nu}{V(\varphi)} \right)^2 k^{ab} V_{,b}$
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The Eisenhart Lift for Inflation

- $S = \int d^4x \sqrt{-g} \left(-\frac{1}{2}R + \frac{1}{2}(\partial^\mu \varphi)(\partial_\mu \varphi) - V(\varphi) \right)$
- $g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2), \quad \varphi = \varphi(t)$

Homogeneous Inflation

- $L = -3a\dot{a}^2 + \frac{1}{2}a^3\dot{\varphi}^2 - a^3V(\varphi)$

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- $L = -3a\dot{a}^2 + \frac{1}{2}a^3\dot{\varphi}^2 + \frac{1}{2} \frac{1}{a^3V(\varphi)} \dot{\chi}^2$

Field Space Metric for Inflation

- $G_{AB} = \begin{pmatrix} -6a & 0 & 0 \\ 0 & a^3 & 0 \\ 0 & 0 & \frac{1}{a^3V(\varphi)} \end{pmatrix}$

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- $\ddot{\varphi} + 3H\dot{\varphi} + \frac{1}{2} \left(\frac{\dot{\chi}}{a^3V} \right)^2 V'(\varphi) = 0$
- $H^2 + 2\frac{\ddot{a}}{a} = \frac{1}{2}\dot{\varphi}^2 + \frac{1}{2} \left(\frac{\dot{\chi}}{a^3V} \right)^2 V(\varphi)$
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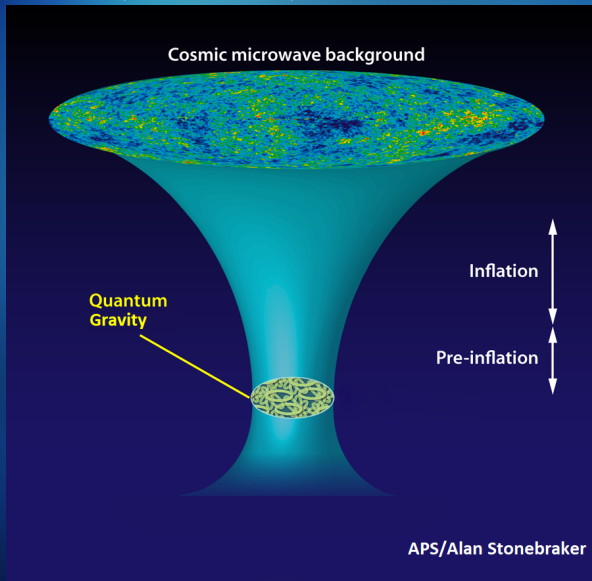
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Eisenhart Lifted Inflation

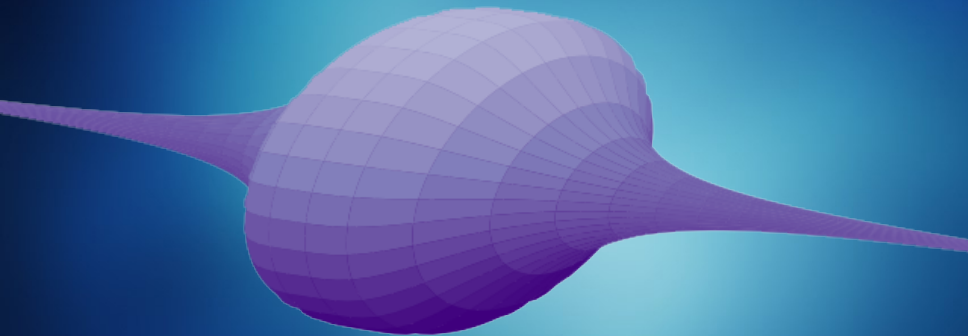
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Initial Conditions for Inflation

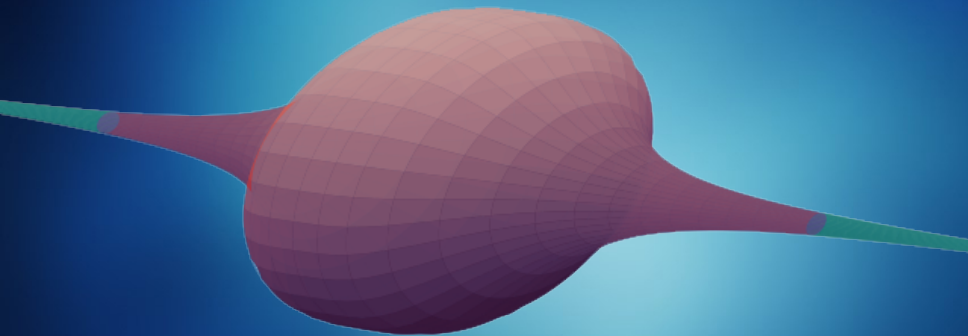
KF and S. Karamitsos: 2019 (arXiv:1812.07095)



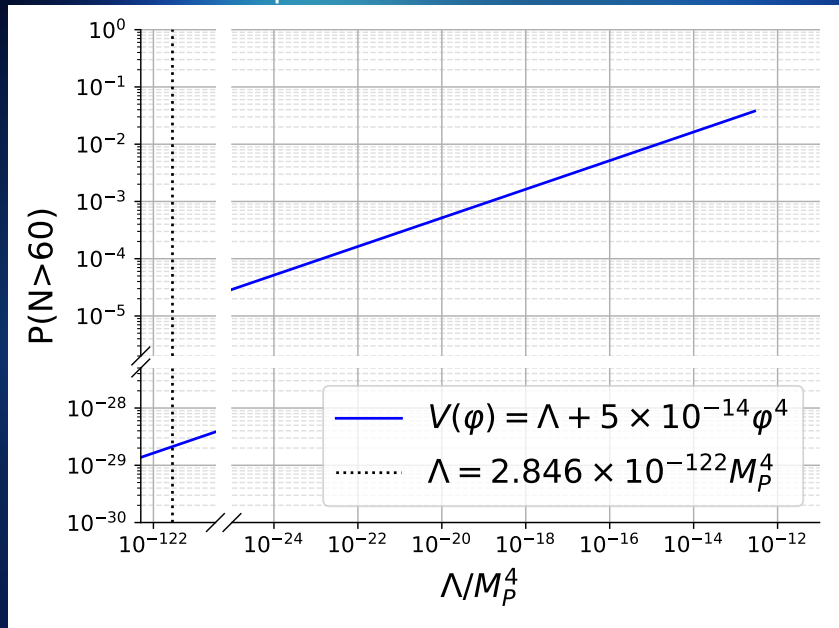
Phase Space Manifold



Phase Space Manifold



Fraction of Phase Space that allows Inflation

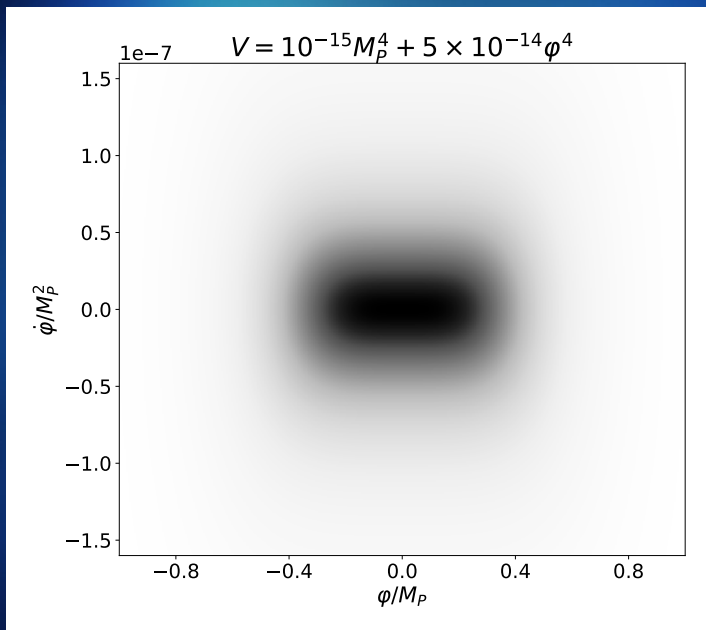


Summary

- The Eisenhart lift allows us to describe any scalar field theory geometrically
- For homogeneous theories we must add an additional scalar field
- For non homogeneous theories we must add a vector field
- We can use this to construct a geometric description of inflation
- We can use this to study the initial conditions problem

Backup Slides

Distribution of the Inflaton Field



Distribution of the Inflaton Field

