

Minimal Vector (Isotriplet) Dark Matter

Alexander Belyaev



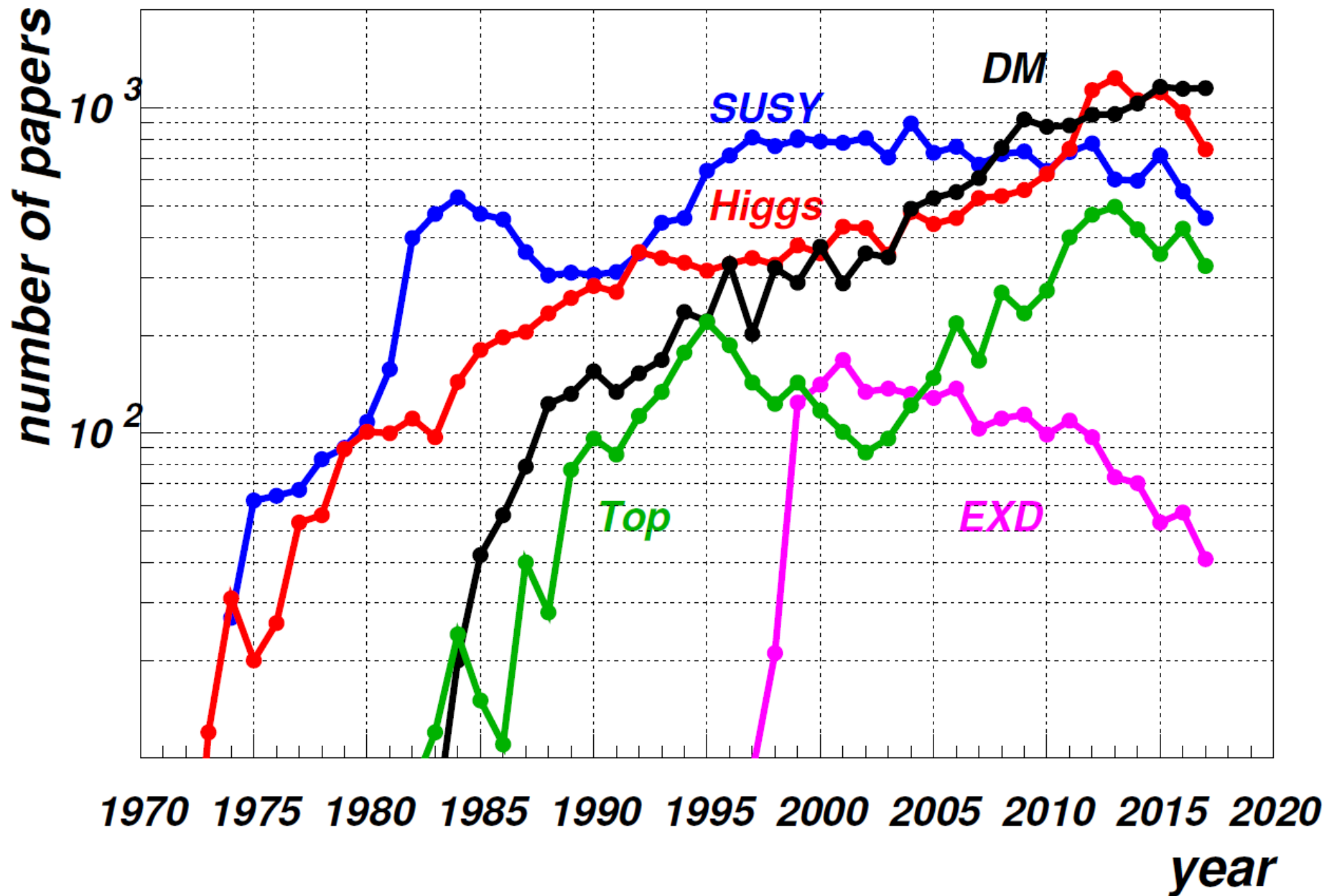
Southampton University & Rutherford Appleton Laboratory

G.Cacciapaglia, J.McKay, D. Marin, A.Zerwekh, AB – arXiv:**1808.10464**, PRD

An abstract graphic on the left side of the slide, featuring a warm color palette of orange, red, and yellow. It contains several clusters of thin, multi-colored lines radiating from central points, resembling particle tracks or fiber-like structures.

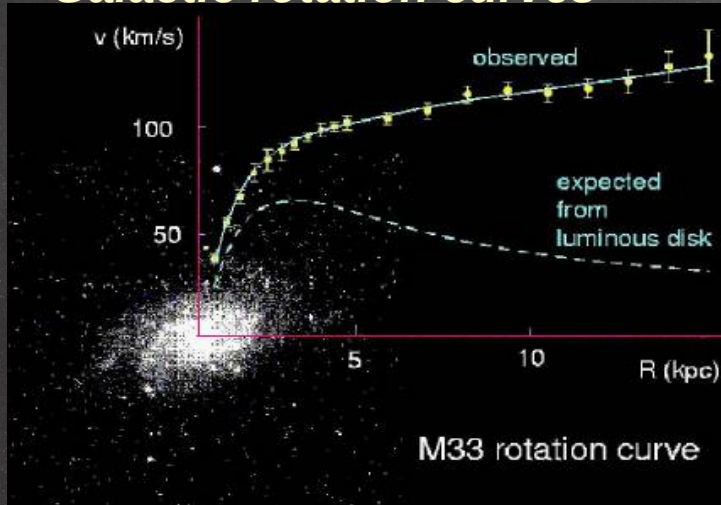
PASCOS 2019
XXV International Symposium

Why we are so keen to study DM?

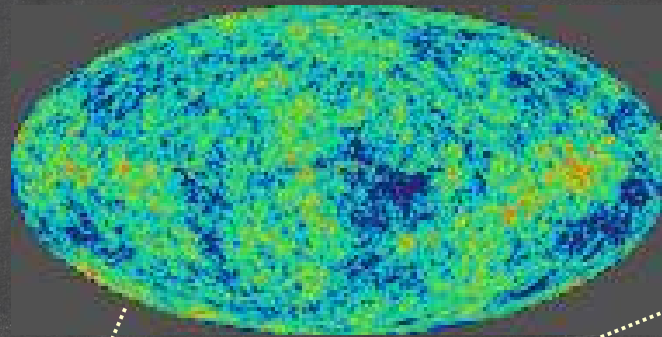


Because the existence of DM is the strongest evidence for BSM!

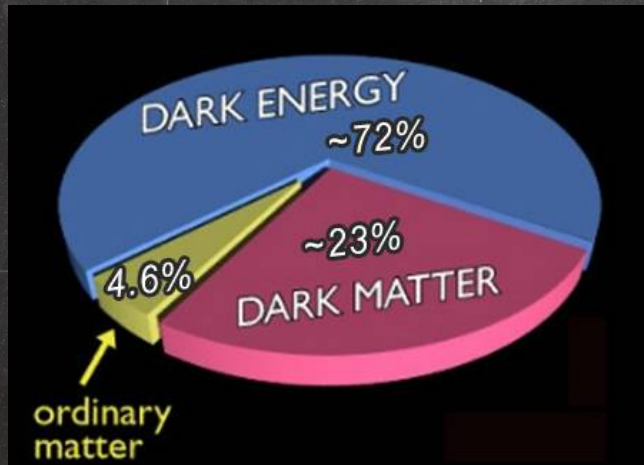
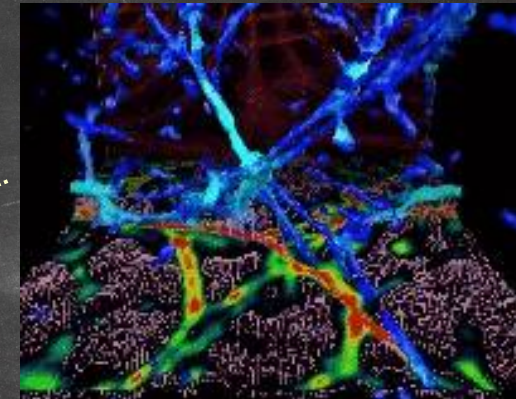
Galactic rotation curves



CMB: WMAP and PLANCK



Large Scale Structures



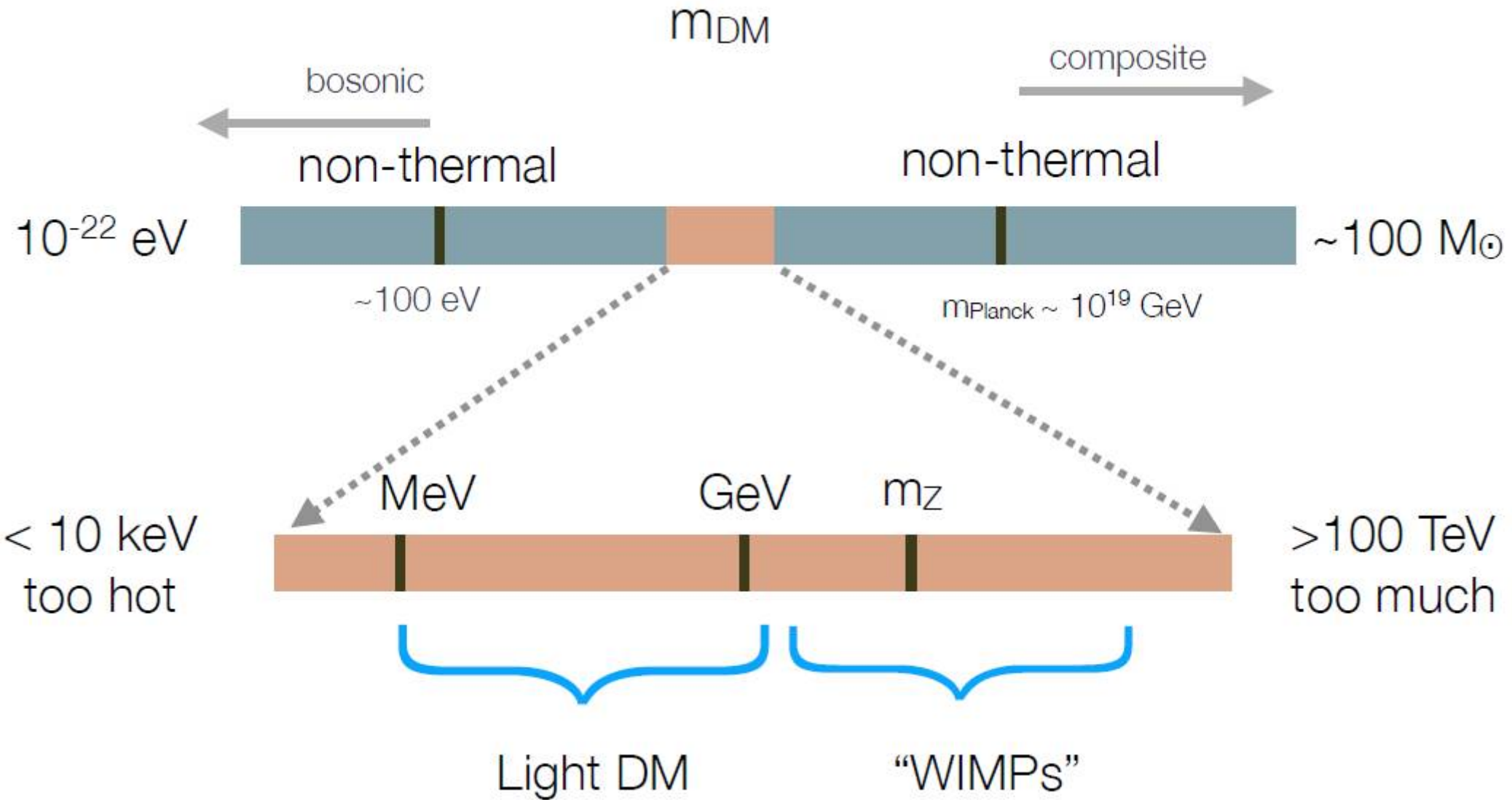
Gravitational lensing



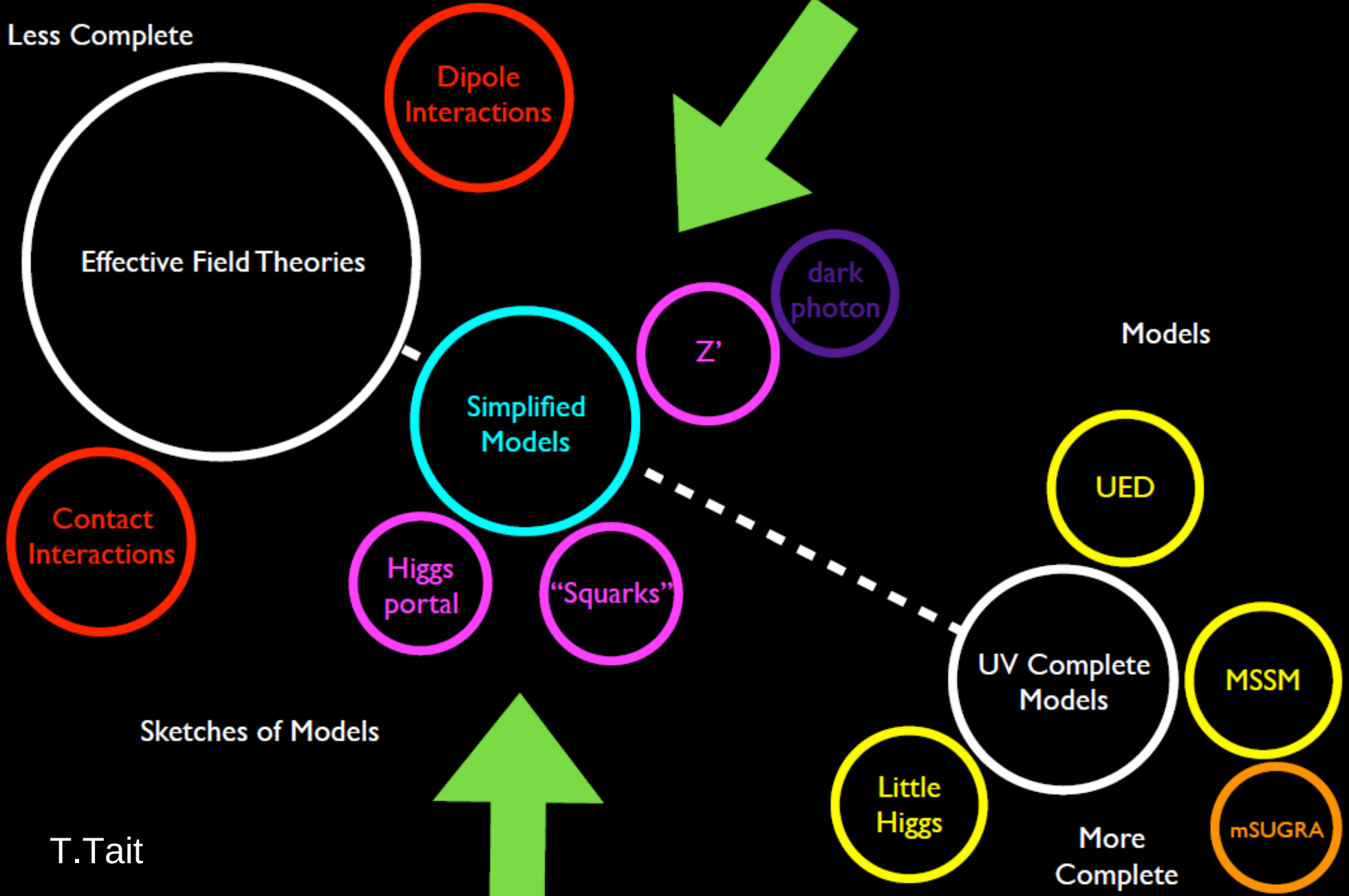
Bullet cluster



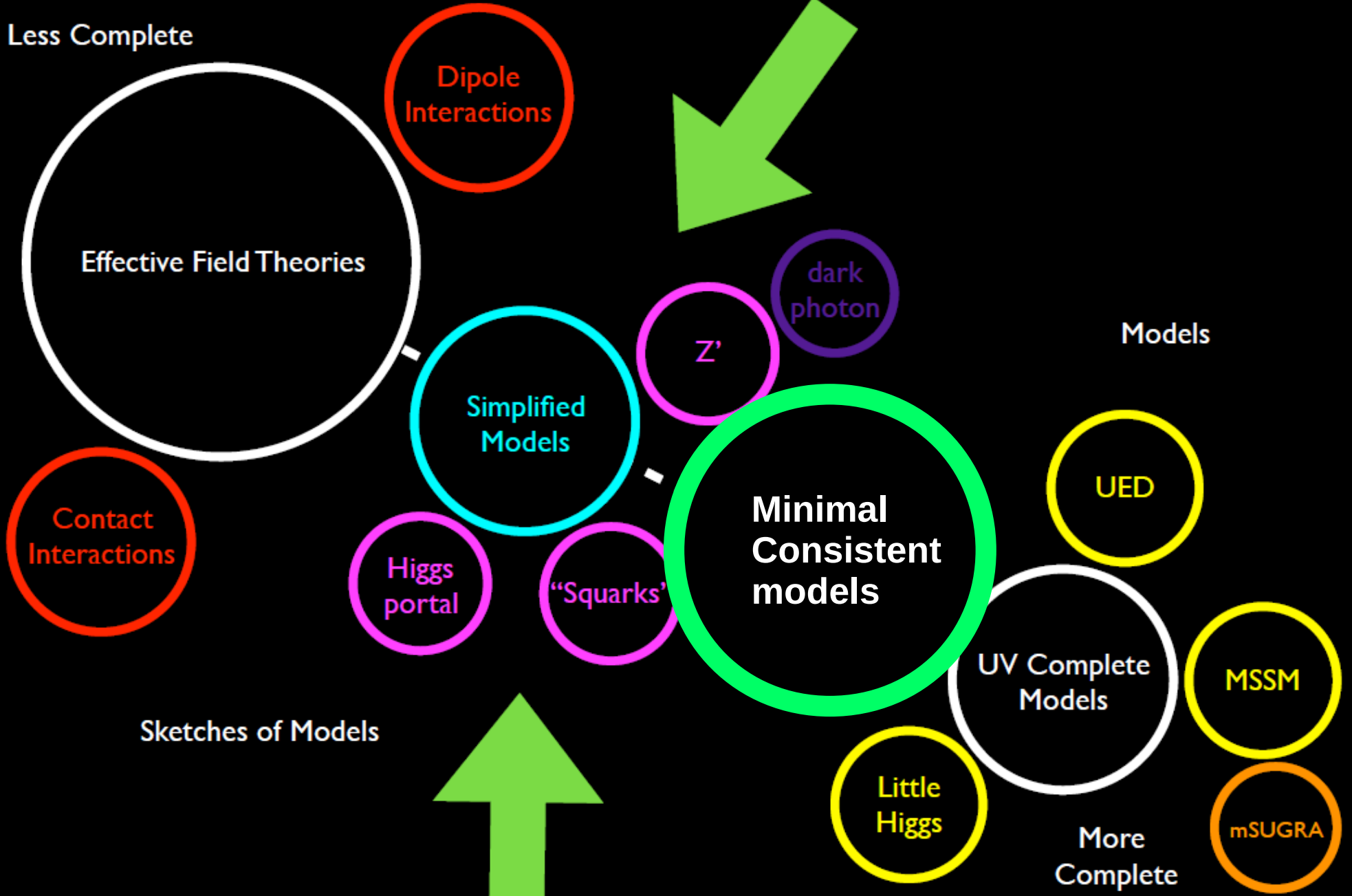
Mass range for thermal DM



Spectrum of Theory Space



Spectrum of Theory Space

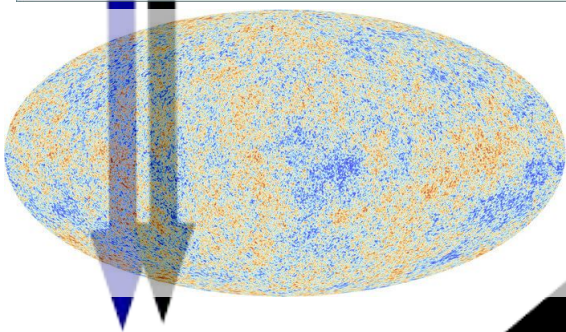


DM

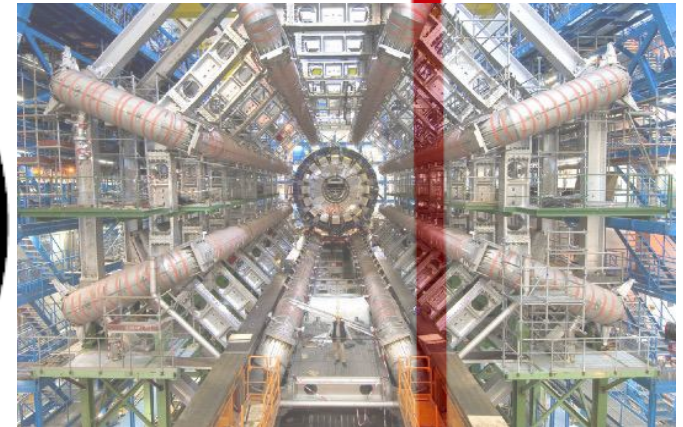
DM

Correct Relic density: efficient
(co) annihilation at the time
of early Universe

*Efficient
annihilation now:
Indirect Detection*



Dark Matter (DM) Signatures



*Efficient production
at colliders*

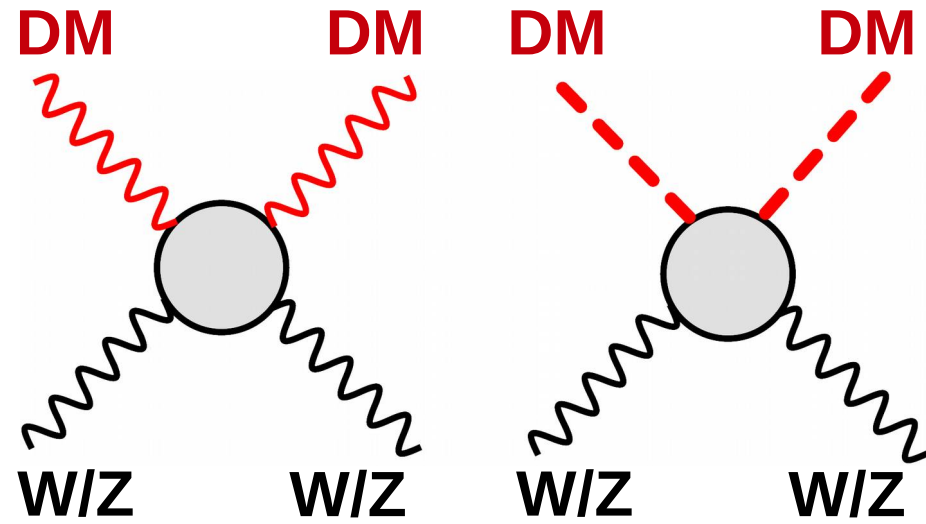
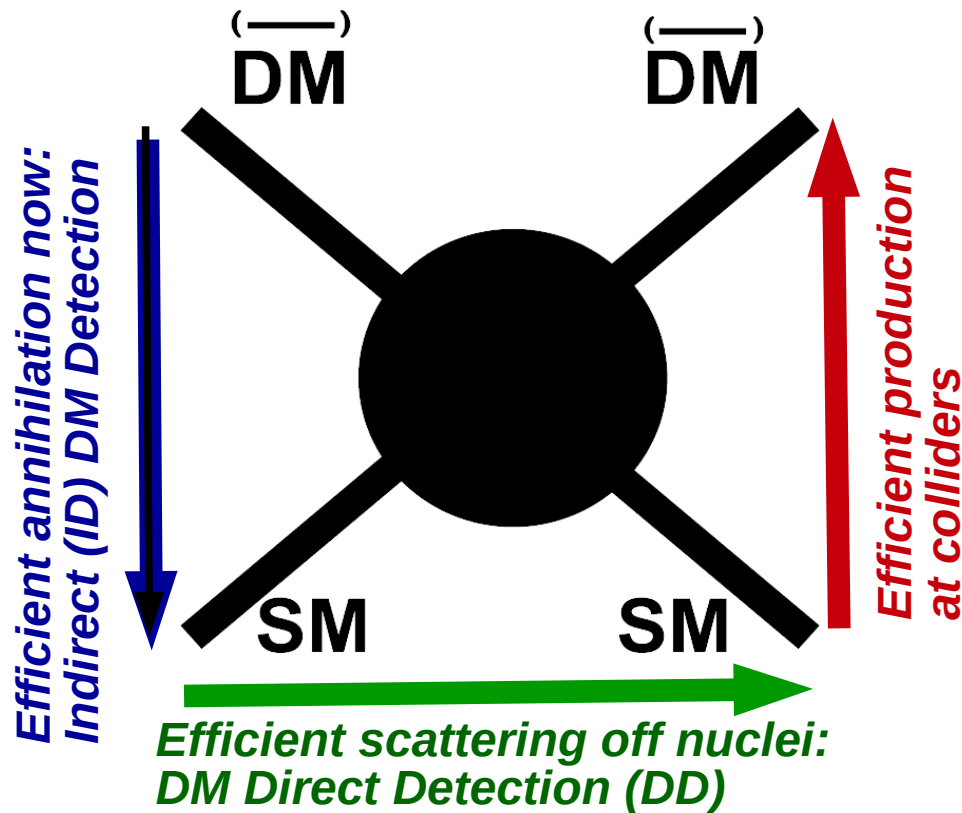
SM



SM

Efficient scattering off nuclei: Direct Detection

Complementarity of DM searches



Example of DM interactions with negligible/suppressed DD rates

Important: there is no 100% correlation between signatures above. E.g. the high rate of annihilation does not always guarantee high rate for DD!

Actually there is a great complementarity in this:

- In case of NO DM Signal – we can efficiently exclude DM models
- In case of DM signal – we have a way to determine the nature of DM

General Lagrangian with massless gauge field and massive vector (matter) field

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{2} \text{Tr} \{ G_{\mu\nu} G^{\mu\nu} \} \\
 & -\text{Tr} \{ D_\mu V_\nu D^\mu V^\nu \} + (1 + a) \text{Tr} \{ D_\mu V_\nu D^\nu V^\mu \} \\
 & + a_1 \text{Tr} \{ (D_\mu V_\nu - D_\nu V_\mu) [V^\mu, V^\nu] \} \\
 & + \frac{a_2}{2} \text{Tr} \{ [V_\mu, V_\nu] [V^\mu, V^\nu] \} \\
 & + i a_3 \text{Tr} \{ G_{\mu\nu} [V^\mu, V^\nu] \} + M^2 \text{Tr} \{ V_\nu V^\nu \}
 \end{aligned}$$

$$\begin{aligned}
 G_\mu & \rightarrow U G_\mu U^{-1} - \frac{1}{g} (\partial_\mu U) U^{-1} \\
 V_\mu & \rightarrow U V_\mu U^{-1} \\
 D_\mu V_\nu & = \partial_\mu V_\nu - i [W_\mu, V_\nu]
 \end{aligned}$$

A. Zerwekh
Int.J.Mod.Phys. A28 (2013)1350054
arXiv:1207.5233

G^μ is the gauge field, while

V_μ transforms homogeneously under the gauge transformations, heaving the properties of the matter field

It turns out that this theory can be unitary!

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{2} \text{Tr} \{G_{\mu\nu} G^{\mu\nu}\} \\
 & -\text{Tr} \{D_\mu V_\nu D^\mu V^\nu\} + (1 + a) \text{Tr} \{D_\mu V_\nu D^\nu V^\mu\} & a & = 0 \\
 & ~~+ a_1 \text{Tr} \{(D_\mu V_\nu - D_\nu V_\mu) [V^\mu, V^\nu]\}~~ & a_1 & = 0 \\
 & -\frac{g^2}{2} \text{Tr} \{[V_\mu, V_\nu] [V^\mu, V^\nu]\} & a_2 & = -g^2 \\
 & -ig \text{Tr} \{G_{\mu\nu} [V^\mu, V^\nu]\} + M^2 \text{Tr} \{V_\nu V^\nu\} & a_3 & = -g
 \end{aligned}$$

$$\begin{aligned}
 G_\mu & \rightarrow U G_\mu U^{-1} - \frac{1}{g} (\partial_\mu U) U^{-1} \\
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G^μ is the gauge field, while

V_μ transforms homogeneously under the gauge transformations, having the properties of the matter field

Theory with massive vector field in the adjoint representation interacting with non-abelian gauge field

$$\begin{aligned}\mathcal{L} = & -\frac{1}{2}\text{Tr} \{G_{\mu\nu}G^{\mu\nu}\} - \text{Tr} \{D_\mu V_\nu D^\mu V^\nu\} + \text{Tr} \{D_\mu V_\nu D^\nu V^\mu\} \\ & -\frac{g^2}{2}\text{Tr} \{[V_\mu, V_\nu][V^\mu, V^\nu]\} - ig\text{Tr} \{G_{\mu\nu} [V^\mu, V^\nu]\} + M^2\text{Tr}\{V_\nu V^\nu\}\end{aligned}$$

- there is a consistent gauge theory for a massive spin-1 field without scalar degrees of freedom
- Unitarity requires no trilinear V interactions: the massive spin-1 field is odd under a new **Z_2 symmetry**
- The massive vector is not a force carrier but a true matter field
- The theory has bi-gauge origin (see backup slides for details)

The Minimal Isotriplet Vector Dark Matter Model (MI VDM)

$$\begin{aligned}\mathcal{L} = & \mathcal{L}_{SM} - \text{Tr} \{D_\mu V_\nu D^\mu V^\nu\} + \text{Tr} \{D_\mu V_\nu D^\nu V^\mu\} \\ & - \frac{g^2}{2} \text{Tr} \{[V_\mu, V_\nu] [V^\mu, V^\nu]\} - ig \text{Tr} \{W_{\mu\nu} [V^\mu, V^\nu]\} + \tilde{M}^2 \text{Tr} \{V_\nu V^\nu\}\end{aligned}$$

$$D_\mu V_\nu = \partial_\mu V_\nu - i[W_\mu, V_\nu]$$

The Minimal Isotriplet Vector Dark Matter Model (MI VDM)

$$\mathcal{L} = \mathcal{L}_{SM} - \text{Tr} \{D_\mu V_\nu D^\mu V^\nu\} + \text{Tr} \{D_\mu V_\nu D^\nu V^\mu\} \\ - \frac{g^2}{2} \text{Tr} \{[V_\mu, V_\nu] [V^\mu, V^\nu]\} - ig \text{Tr} \{W_{\mu\nu} [V^\mu, V^\nu]\} + \tilde{M}^2 \text{Tr} \{V_\nu V^\nu\}$$

$$+ a (\Phi^\dagger \Phi) \text{Tr} \{V_\nu V^\nu\}$$

we must add interactions with Φ
Higgs doublet

$$D_\mu V_\nu = \partial_\mu V_\nu - i[W_\mu, V_\nu]$$

$$M_V^2 = \tilde{M}^2 + \frac{1}{2} a v^2$$

only two parameters: **a** and **M**

The Minimal Isotriplet Vector Dark Matter Model (MI VDM)

$$\mathcal{L} = \mathcal{L}_{SM} - Tr \{ D_\mu V_\nu D^\mu V^\nu \} + Tr \{ D_\mu V_\nu D^\nu V^\mu \} - \frac{g^2}{2} Tr \{ [V_\mu, V_\nu] [V^\mu, V^\nu] \} - ig Tr \{ W_{\mu\nu} [V^\mu, V^\nu] \} + \tilde{M}^2 Tr \{ V_\nu V^\nu \}$$

$$+ a (\Phi^\dagger \Phi) Tr \{ V_\nu V^\nu \}$$

we must add interactions with Φ Higgs doublet

$$D_\mu V_\nu = \partial_\mu V_\nu - i[W_\mu, V_\nu]$$

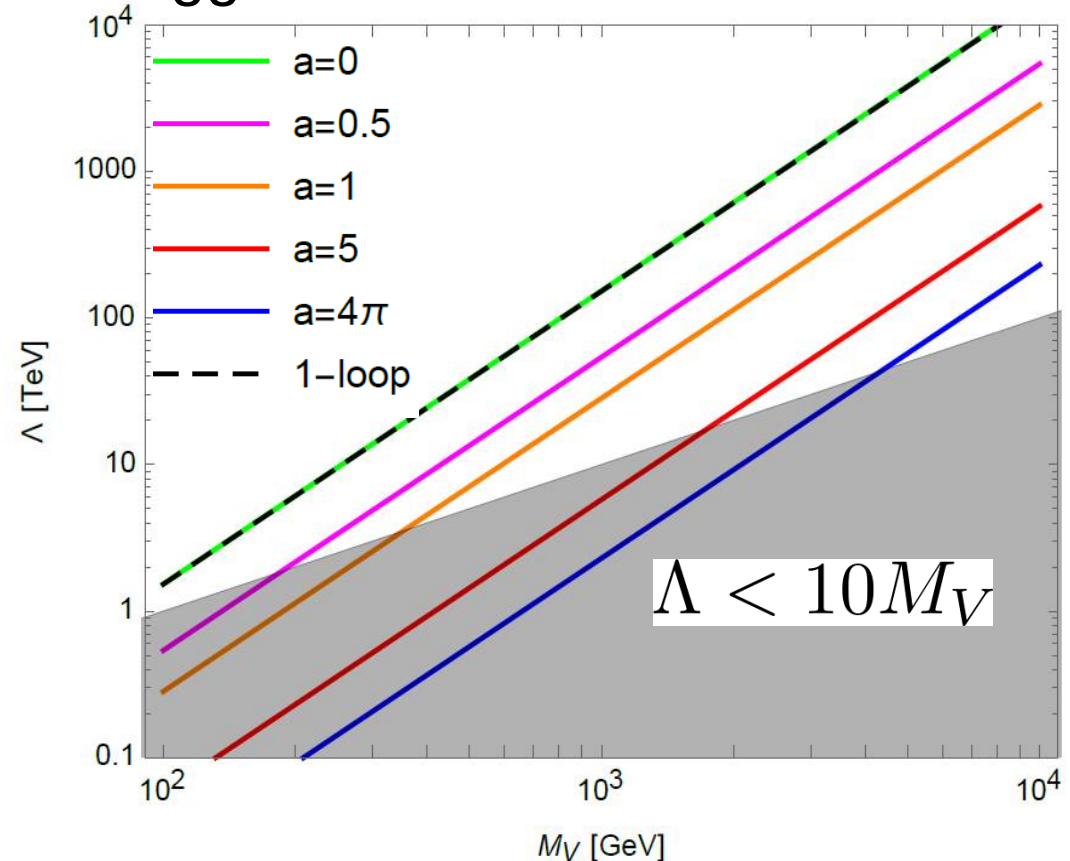
$$M_V^2 = \tilde{M}^2 + \frac{1}{2} a v^2$$

only two parameters: **a** and **M**

- due to the V-H interaction, Unitarity became an issue:

$$V_L^+ V_L^- \rightarrow V_L^+ V_L^- \quad |a_0| < 1/2$$

$$\Lambda \approx \frac{8\sqrt{\pi} M_V^2}{\sqrt{4a^2 v^2 + 3g^2 M_W^2}}$$

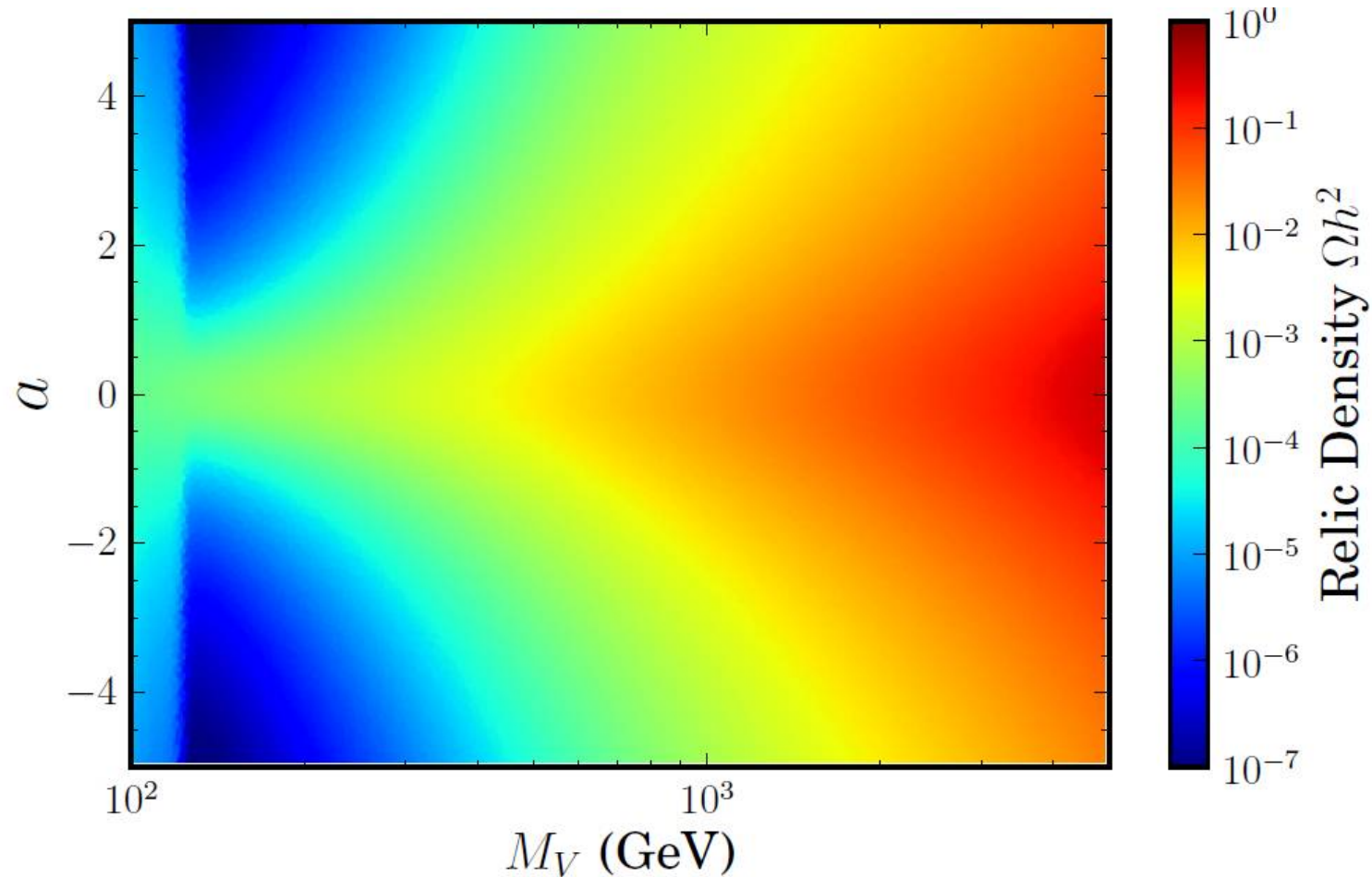


For $M_V > 2\text{TeV}$, unitarity violation occurs at sufficiently high scales

Phenomenology of MI VDM

- The model is implemented into CalcHEP package using LanHEP
 - ➔ DM observables with micrOMEGAs
 - ➔ DM Collider phenomenology with CalcHEP
 - ➔ Available at HEPMDB (hepmdb.soton.ac.uk) as `hepmdb:0118.0283` model
- Main relevant observables/constraints:
 - ➔ DM relic density
 - ➔ DM direct detection (XENON1T)
 - ➔ Collider observables
 - disappearing charged tracks from V^+ due to the small V^0 - V^+ mass split
 - from $H \rightarrow \gamma\gamma$ decay branching ratio

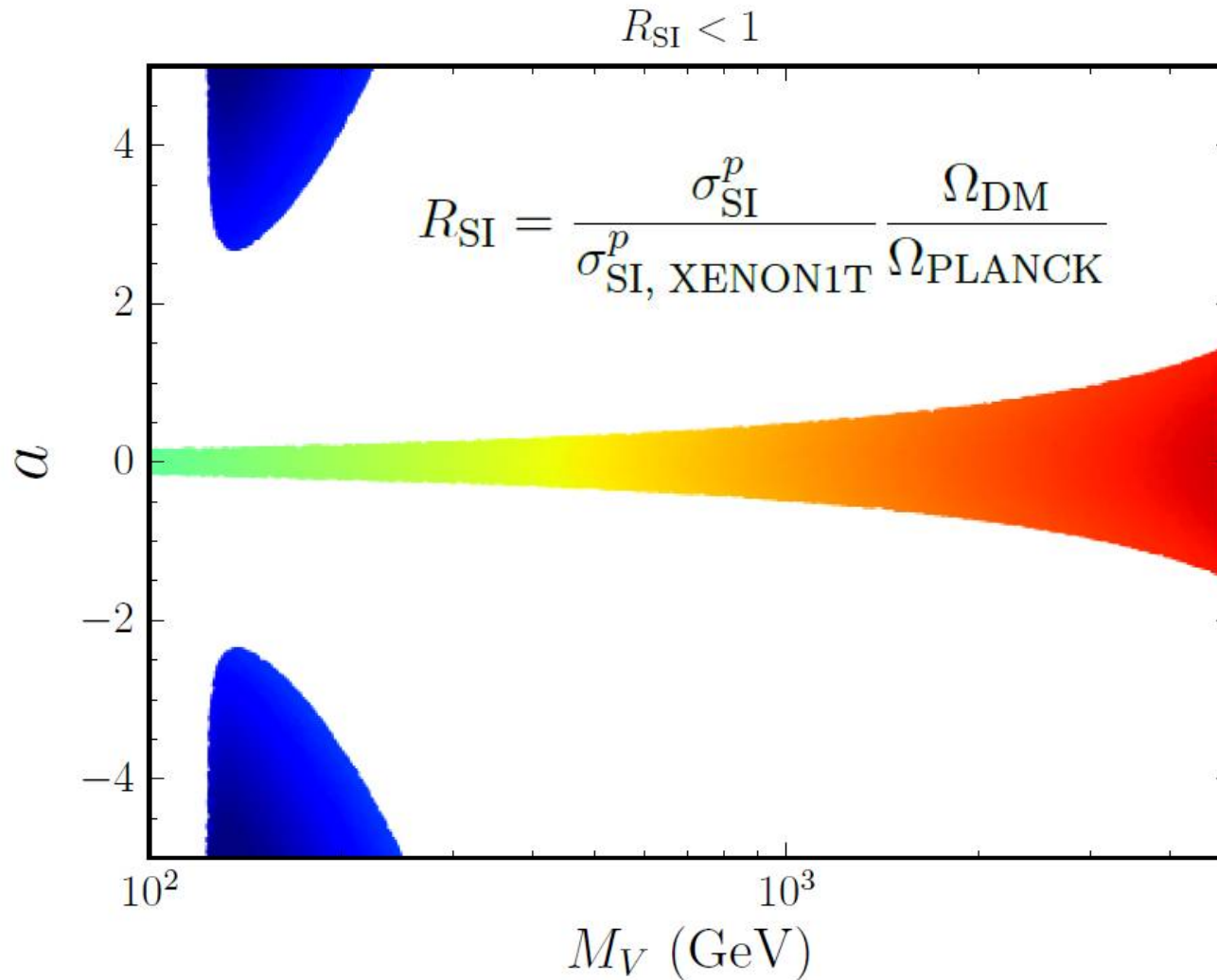
The probe of MI VDM parameter space



The relic density map in M_V - a parameter space

The probe of MI VDM parameter space

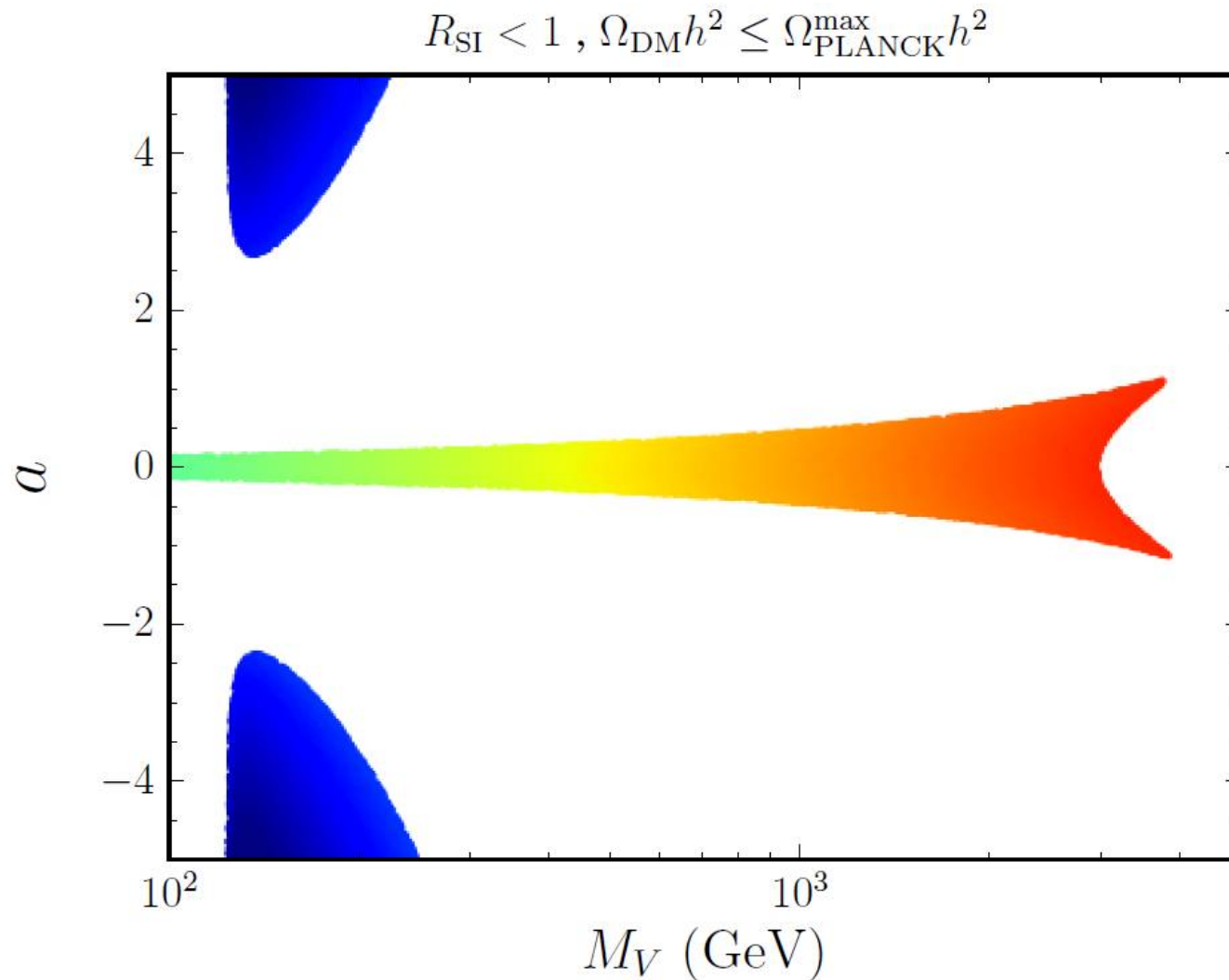
The relic density map in M_V - a parameter space



DM DD constraints from XENON1T: limiting HVV interactions

The probe of MI VDM parameter space

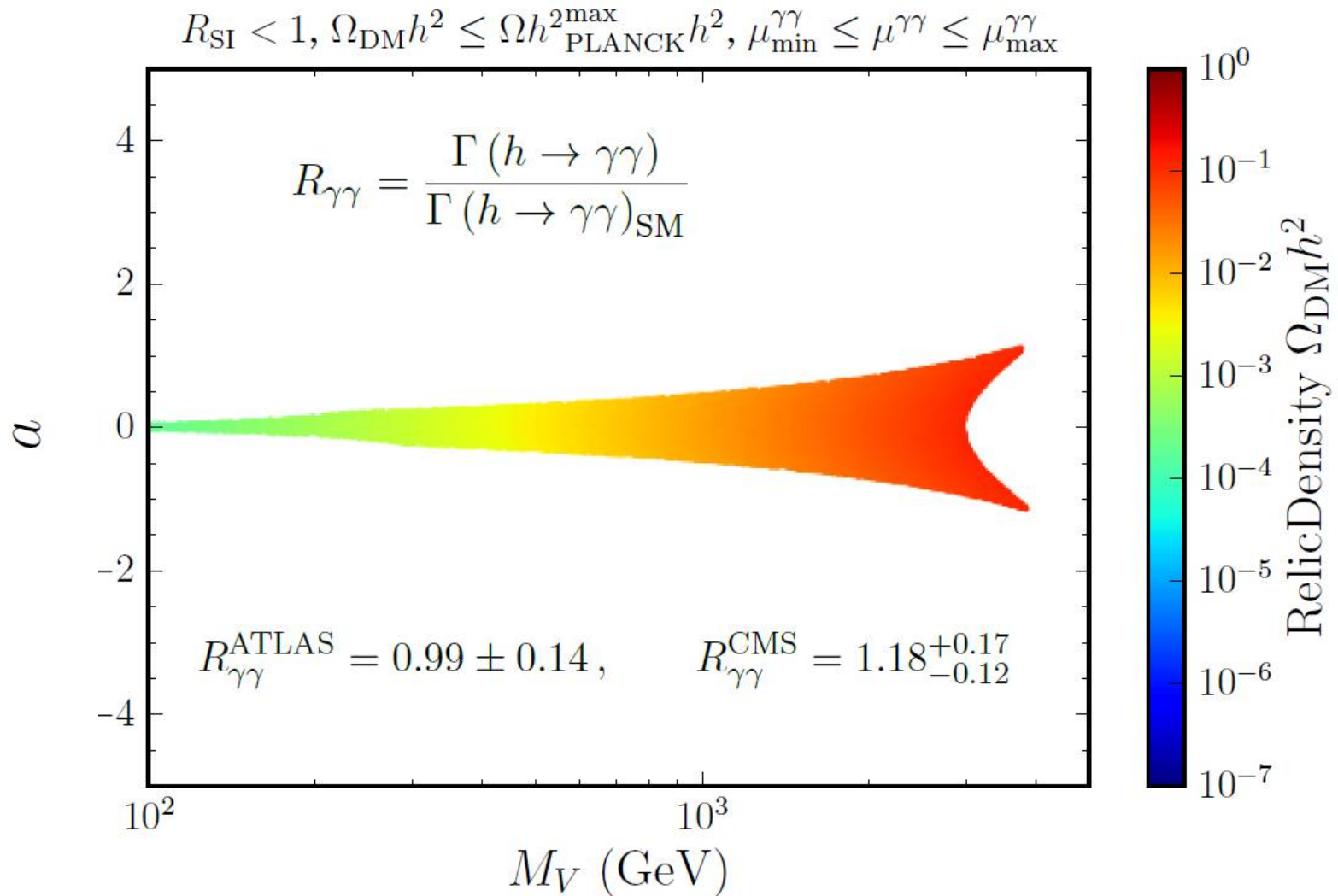
The relic density map in M_V - a parameter space



Relic density constraints from PLANCK: an upper limit on DM mass

The probe of MI VDM parameter space

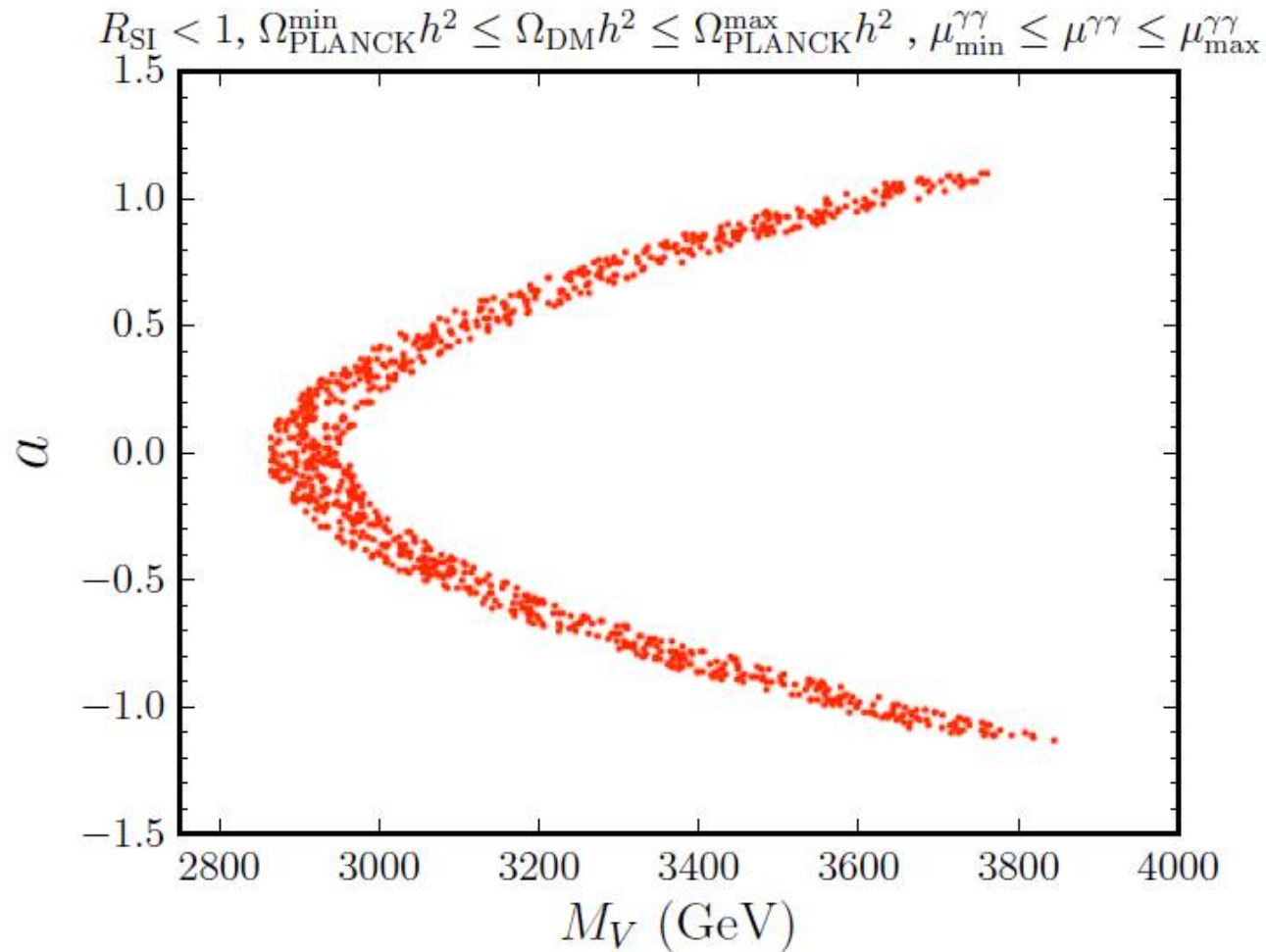
The relic density map in M_V - a parameter space



$H \rightarrow \gamma\gamma$ constraints from LHC: eliminates large a & low M_V region

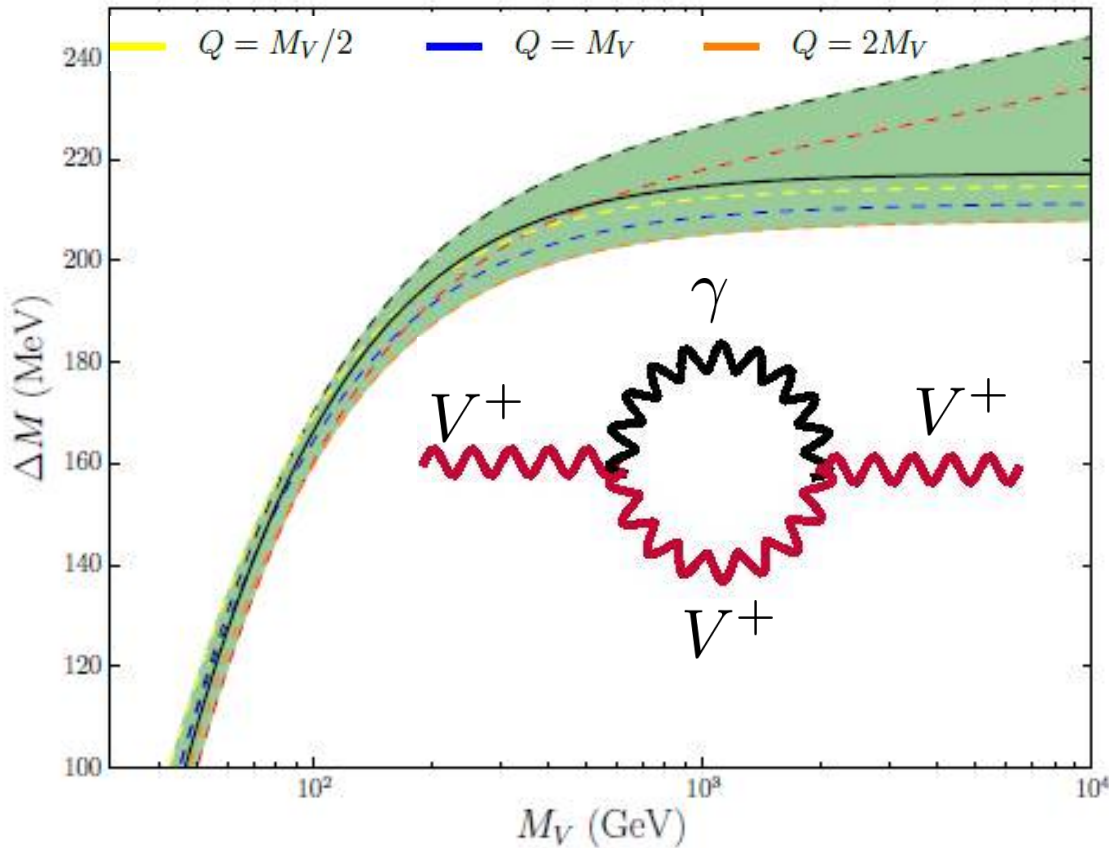
Upper and lower limit on DM mass!

application of upper and lower PLANCK limits on relic density

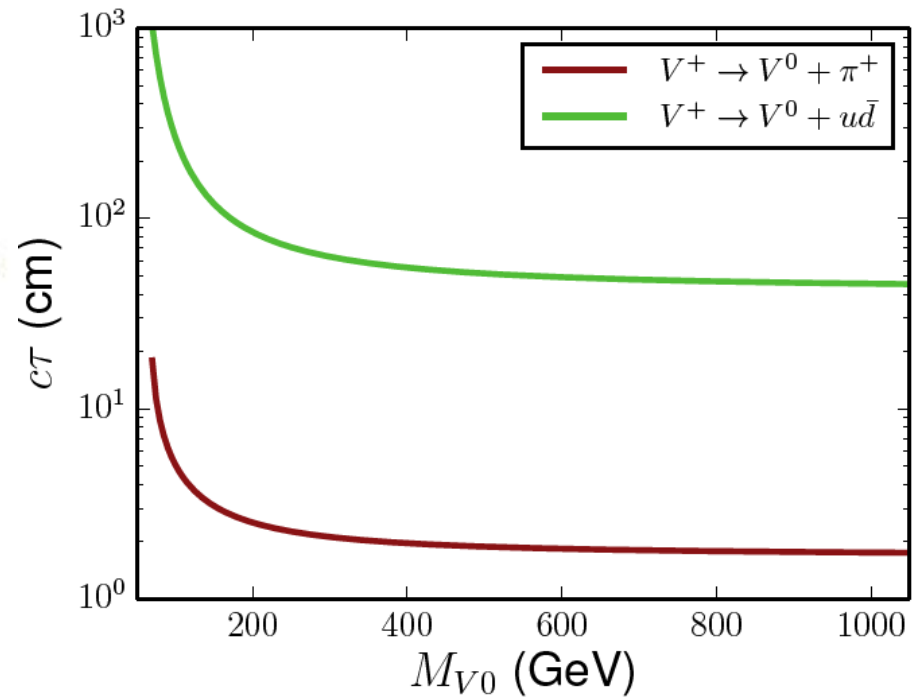


we have found narrow region for viable DM mass: 2.8-3.8 TeV
Can we probe it?

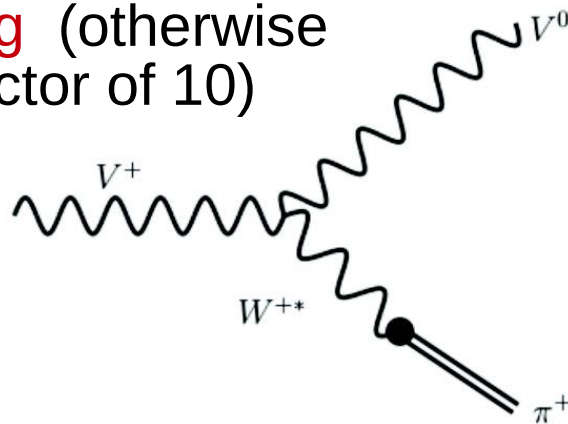
Disappearing Charged tracks from MI VDM



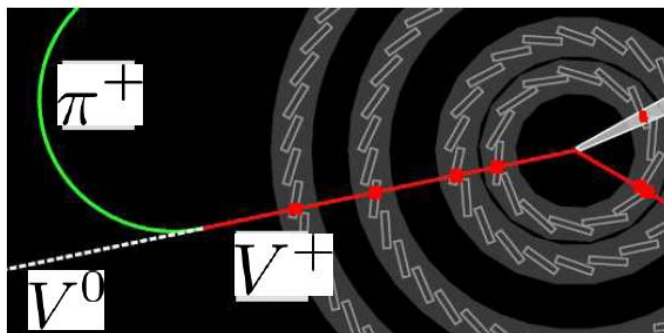
The small mass gap due to radiative mass corrections ($\sim 200\text{MeV} \sim \text{pion mass}$) between DM and its charged partner will lead to the **disappearing charge tracks** signatures



The life-time should be properly evaluated using **W-pion mixing** (otherwise overestimated by factor of 10)

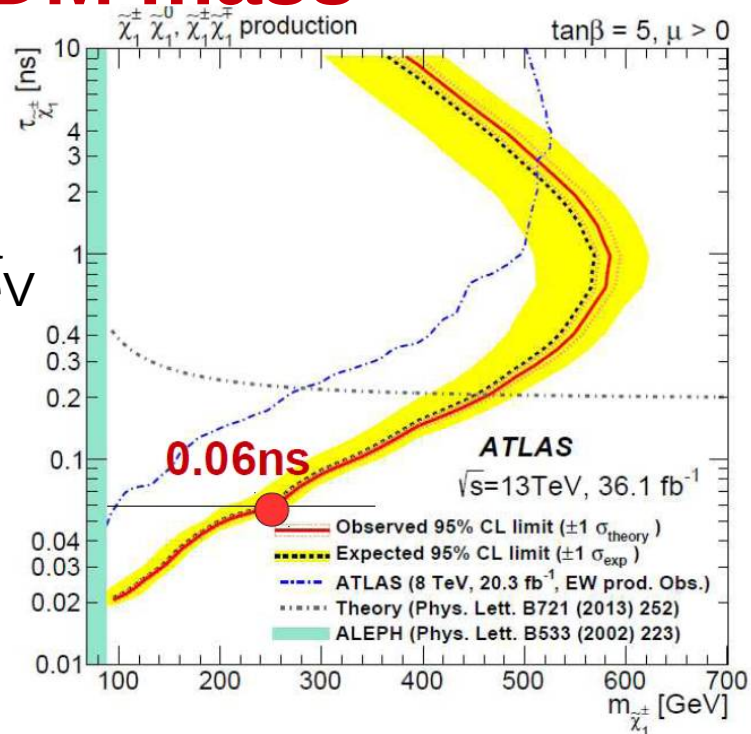


Collider sensitivity to VDM mass



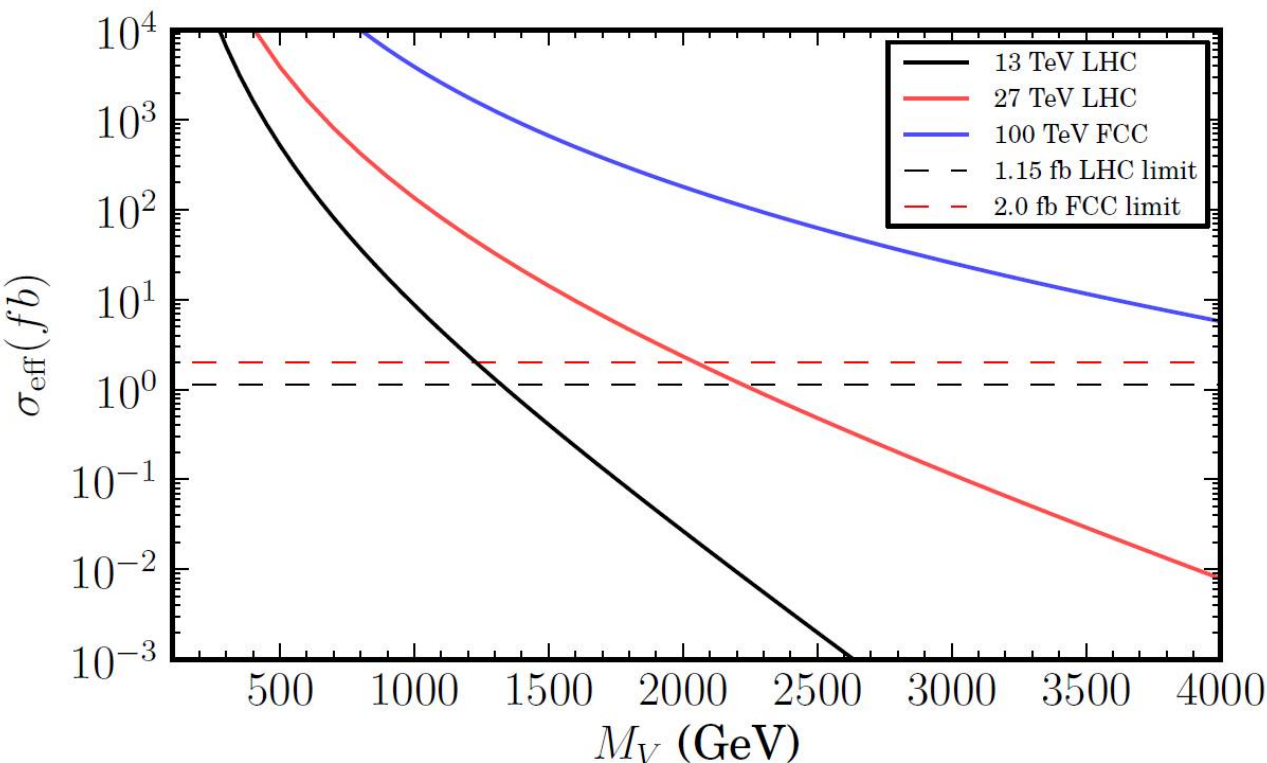
using ATLAS
 arXiv:1712.02118 for LHC
 interpretation and
 Mahbubani, Schwaller, Zurita
 arXiv:1703.05327 for 100 TeV
 FCC projections

AB, Cacciapaglia, McKay, Martin, Zerwekh arXiv:1808.10464
 LHC@13, @27TeV and FCC@100 TeV constraints from LLP searches



Current bound from LHC on DM mass from the minimal vector triplet model: **1.3 TeV !**

100 TeV FCC will cover DM mass **beyond 4 TeV:**
 will discover or close the model



Summary

- ⇒ we have build an extension of the SM with the new massive spin-one isotriplet field with only two new parameters: M_V and V coupling to H
- ⇒ perturbative unitarity, requires new Z_2 symmetry (so it is motivated by theory and not imposed by hand), making V^0 a DM candidate
- ⇒ If DM has a mass in the range 2.8-3.8TeV, the model explains the measured DM relic density and simultaneously satisfies complementary current experimental constraints
- ⇒ Spectacular feature of V^+ – long lifetime due to a small (radiative) mass split with V^0 , providing disappearing tracks signature
- ⇒ future hadron collider with a centre-of-mass energy of 100TeV will probe masses beyond 4TeV – completely covering the allowed parameter space of the model with disappearing tracks signature.

So, we have an opportunity to discover or exclude the model

Thank you!

Backup Slides

A bi-gauge origin of the theory with massive vector field

A simple rotation

$$G = \frac{1}{\sqrt{2}} (A_1 + A_2)$$

$$V = \frac{1}{\sqrt{2}} (A_1 - A_2)$$

$$\mathcal{L} = -\frac{1}{2} \text{Tr} [F_{1\mu\nu} F_1^{\mu\nu}] - \frac{1}{2} \text{Tr} [F_{2\mu\nu} F_2^{\mu\nu}] + \frac{M^2}{2} \text{Tr} [(A_{1\mu} - A_{2\mu})^2]$$

$$F_{1\mu\nu} = \partial_\mu A_{1\nu} - \partial_\nu A_{1\mu} - i\sqrt{2}g [A_{1\mu}, A_{1\nu}]$$

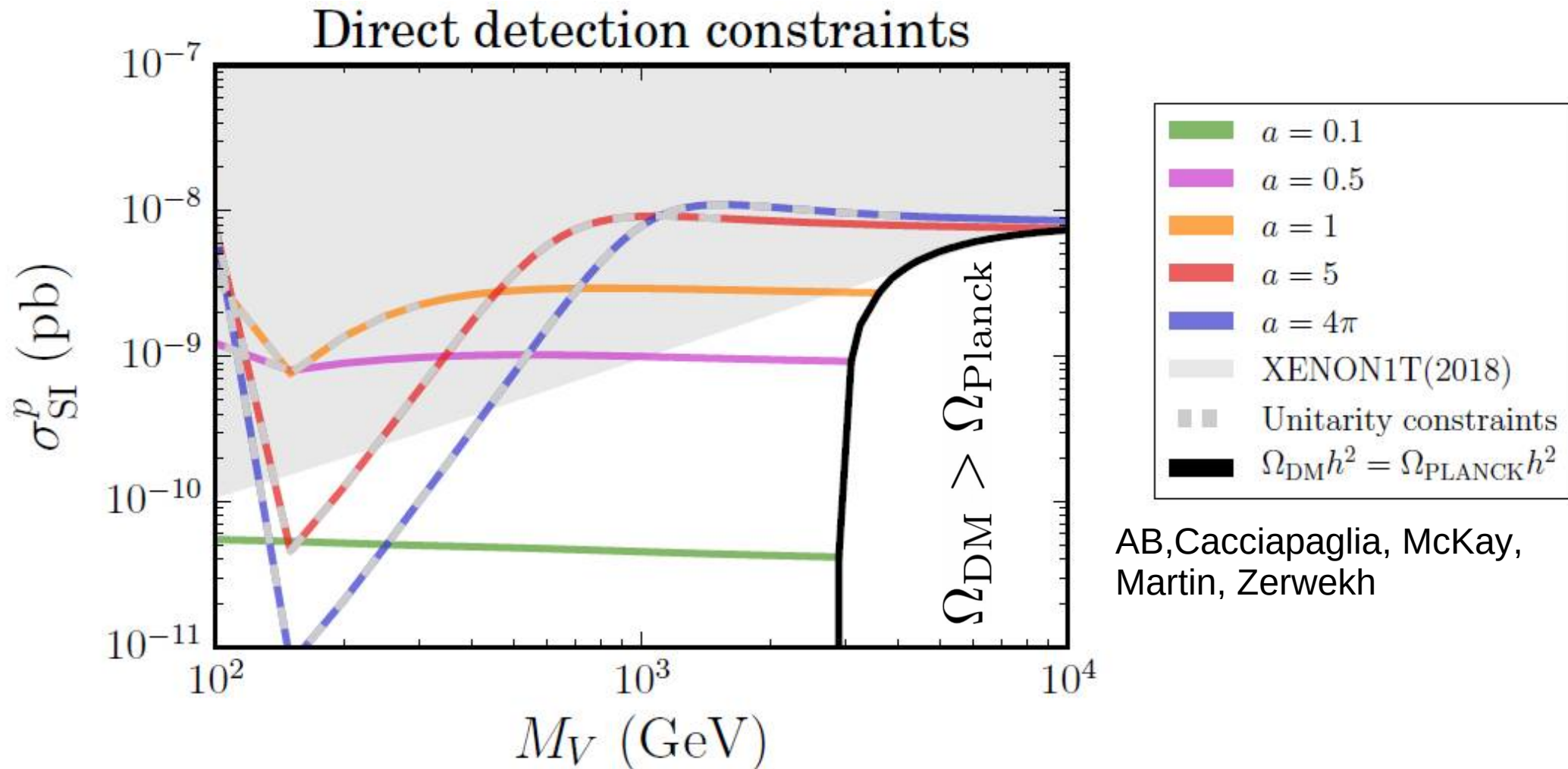
Same coupling constant

$$F_{2\mu\nu} = \partial_\mu A_{2\nu} - \partial_\nu A_{2\mu} - i\sqrt{2}g [A_{2\mu}, A_{2\nu}]$$

$$A_{i\mu} \rightarrow U A_{i\mu} U^{-1} - \frac{1}{\sqrt{2}g} (\partial_\mu U) U^{-1} \quad (i = 1, 2)$$

Power of DM DD to rule out theory space

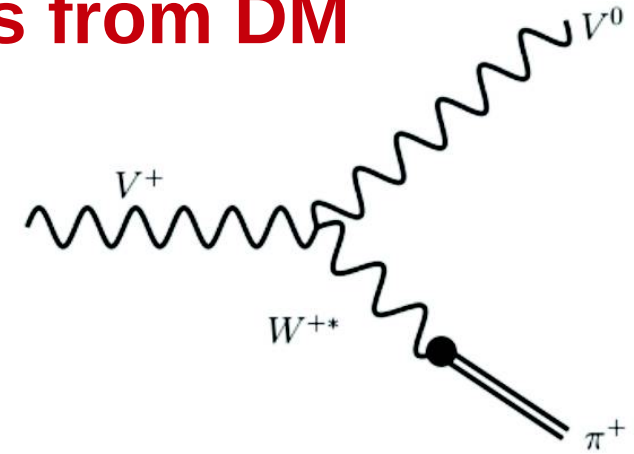
Vector DM Model



- ZENON 1T excludes **both** large $HV_{DM} V_{DM}$ couplings and large M_{DM}
- The **lower masses** (rest of space) can be covered at future colliders

Disappearing Charged Tracks from DM

The small mass gap between (\sim pion mass) DM and its charged partner will lead to the **disappearing charge tracks**



The life-time should be properly evaluated using **W-pion mixing**

$$\mathcal{L}_{\pi^- V^+ V^0} = \frac{g^2 f_\pi}{2\sqrt{2} M_W^2} [(p_{V^+} - p_{V^0})_\alpha g_{\beta\gamma} + p_{V^0}{}_\beta g_{\alpha\gamma} - p_{V^+}{}_\gamma g_{\alpha\beta}] p_{\pi^-}^\alpha \pi^- V^{+\beta} V^{0\gamma}$$

$pp \rightarrow V^+ V^- / V^\pm V^0$ @ 13 TeV

