

Asymptotic Scale Invariance and its Consequences

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Phys. Rev. D 99, 103528 with Mikhail Shaposhnikov

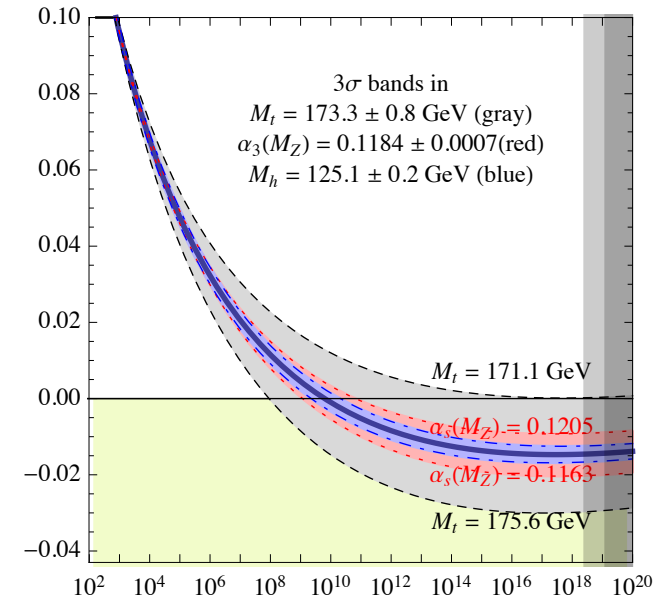
PASCOS2019, Manchester, July 1

Introduction

- ✓ SM works very well (to explain collider expt.).
- ✓ Higgs boson with $m_h \simeq 125 \text{ GeV}$
- ✓ SM itself can be extended up to...

Running Higgs quartic coupling

D. Buttazzo, G. Degrassi, P.P. Giardino,
G.F. Giudice, F. Sala, A. Salvio, A. Strumia (2014)

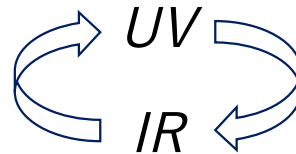


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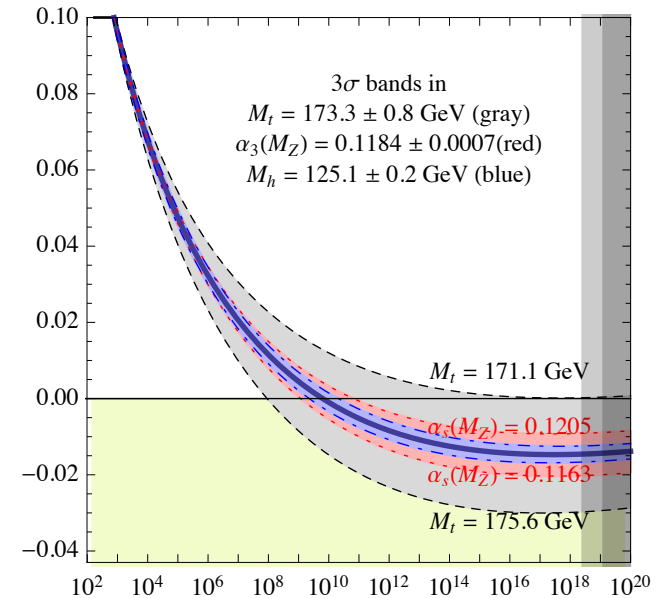
Running Higgs quartic coupling

Any information on UV completion?
(extracted from boundary condition)



EW vacuum meta-stability?

D. Buttazzo, G. Degrassi, P.P. Giardino, G.F. Giudice, F. Sala, A. Salvio, A. Strumia (2014)



- New physics can easily alter the running.

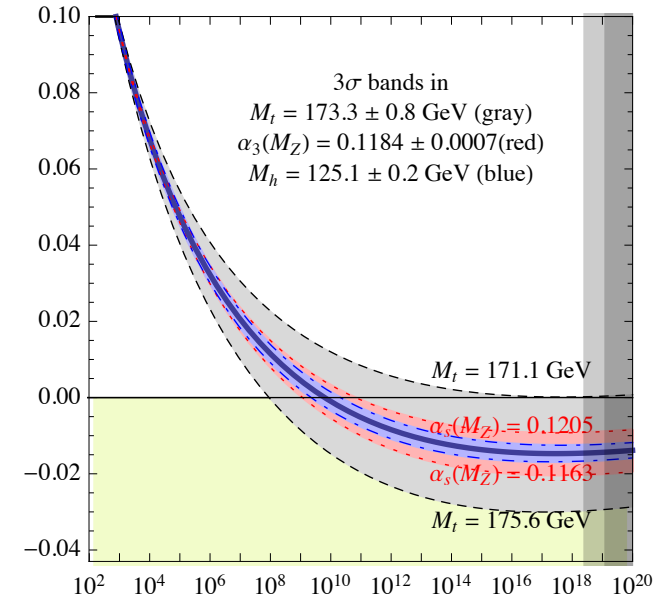
motivated by unsolved issues (DM, neutrino mass, BAU, etc.)

Introduction

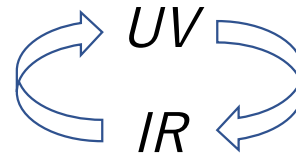
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EW vacuum meta-stability?

□ Impact of **symmetry in low-energy effective theory**

Quantum Scale invariance

(Assumption)
 following from unknown UV completion

Outline

Scale invariance (classical/quantum)

Asymptotic SI and vacuum meta-stability issue
without any technical aspects

Self-consistency and Validity of EFT approach
with some technical aspects

Summary
(regularization/renormalization)

Scale invariance (classical)

Invariance under $\left\{ \begin{array}{l} x^\mu \rightarrow \sigma^{-1} x^\mu \\ \Phi(x) \rightarrow \sigma^{d_\Phi} \Phi(x) \end{array} \right.$ d_Φ : Mass dimension of dynamical fields Φ

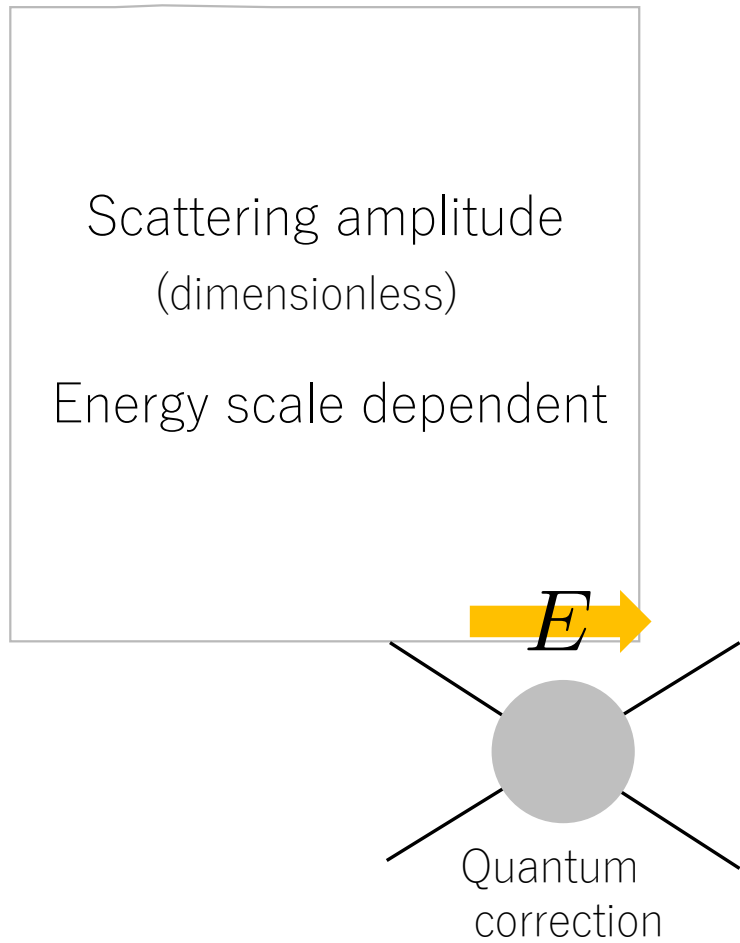
Explicit mass scale breaks SI :

$$V = -\frac{\mu_{\text{EW}}^2}{2} h^2 + \frac{\lambda}{4} h^4$$

In the SM of particle physics sector,
SI is broken only by the Higgs mass term.

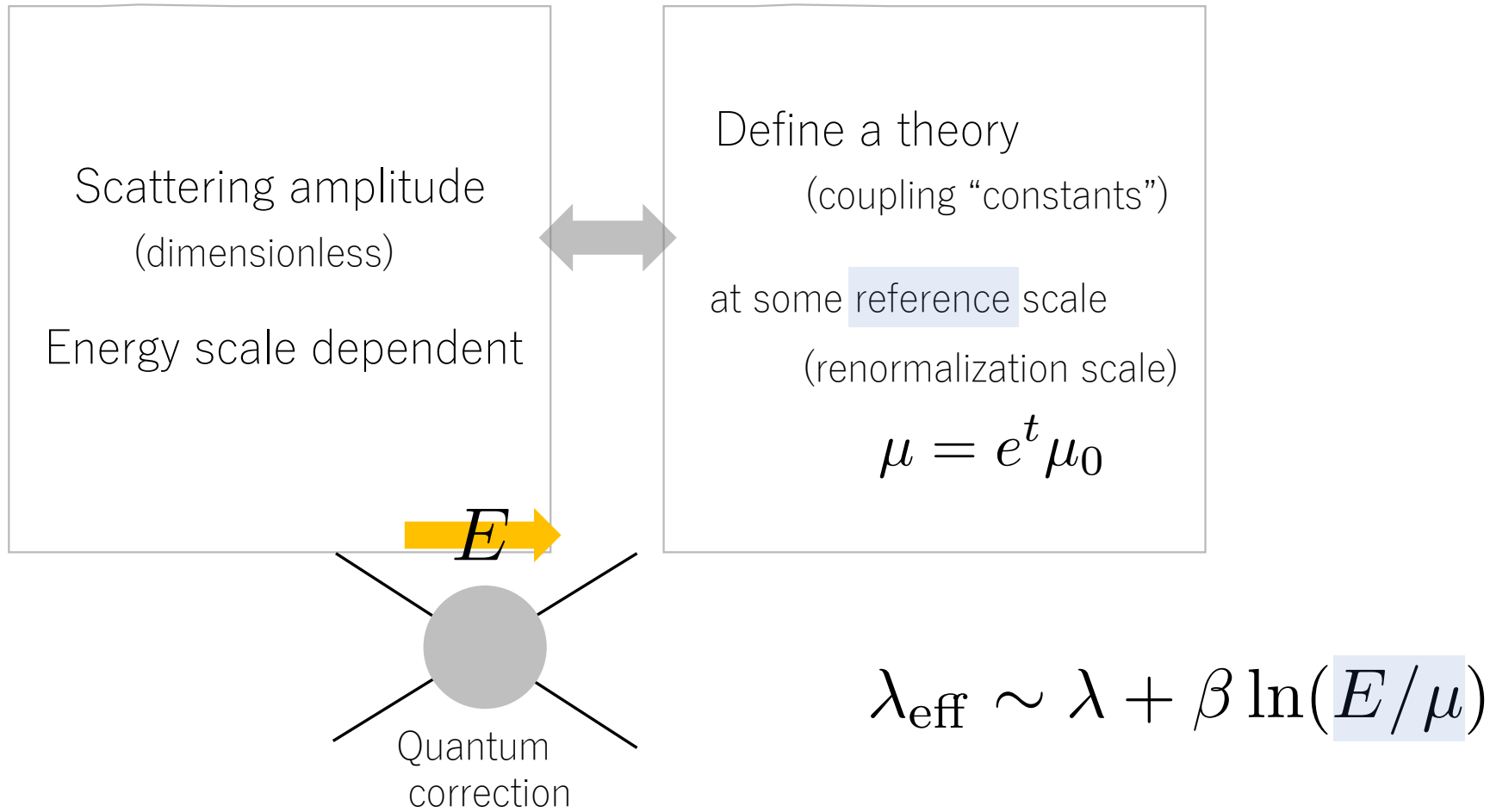
Scale-invariant for $h \gg \mu_{\text{EW}}$.
(approximately)

Is SI necessarily anomalous?

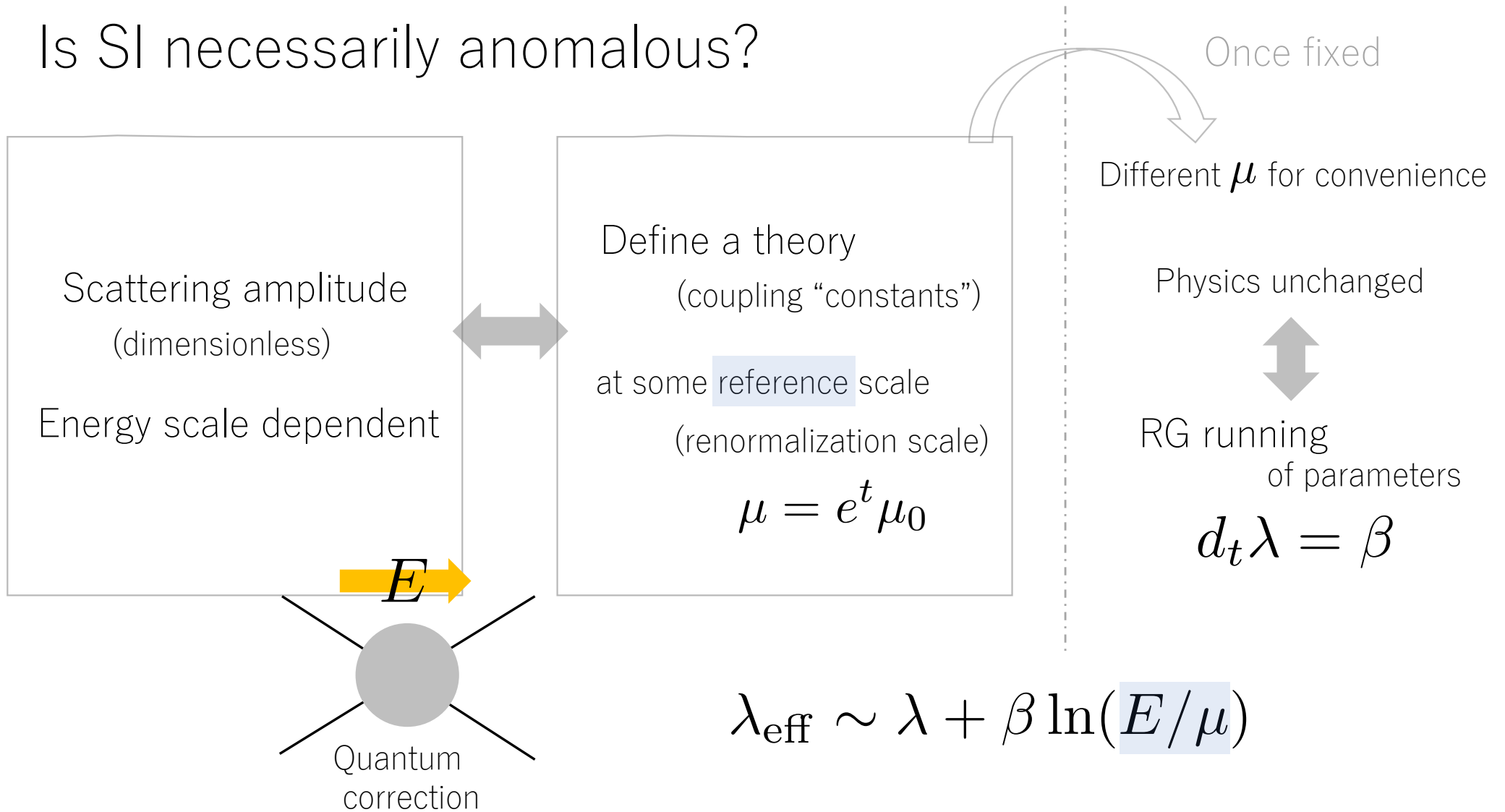


$$\lambda_{\text{eff}} \sim \lambda + \beta \ln(E/\mu)$$

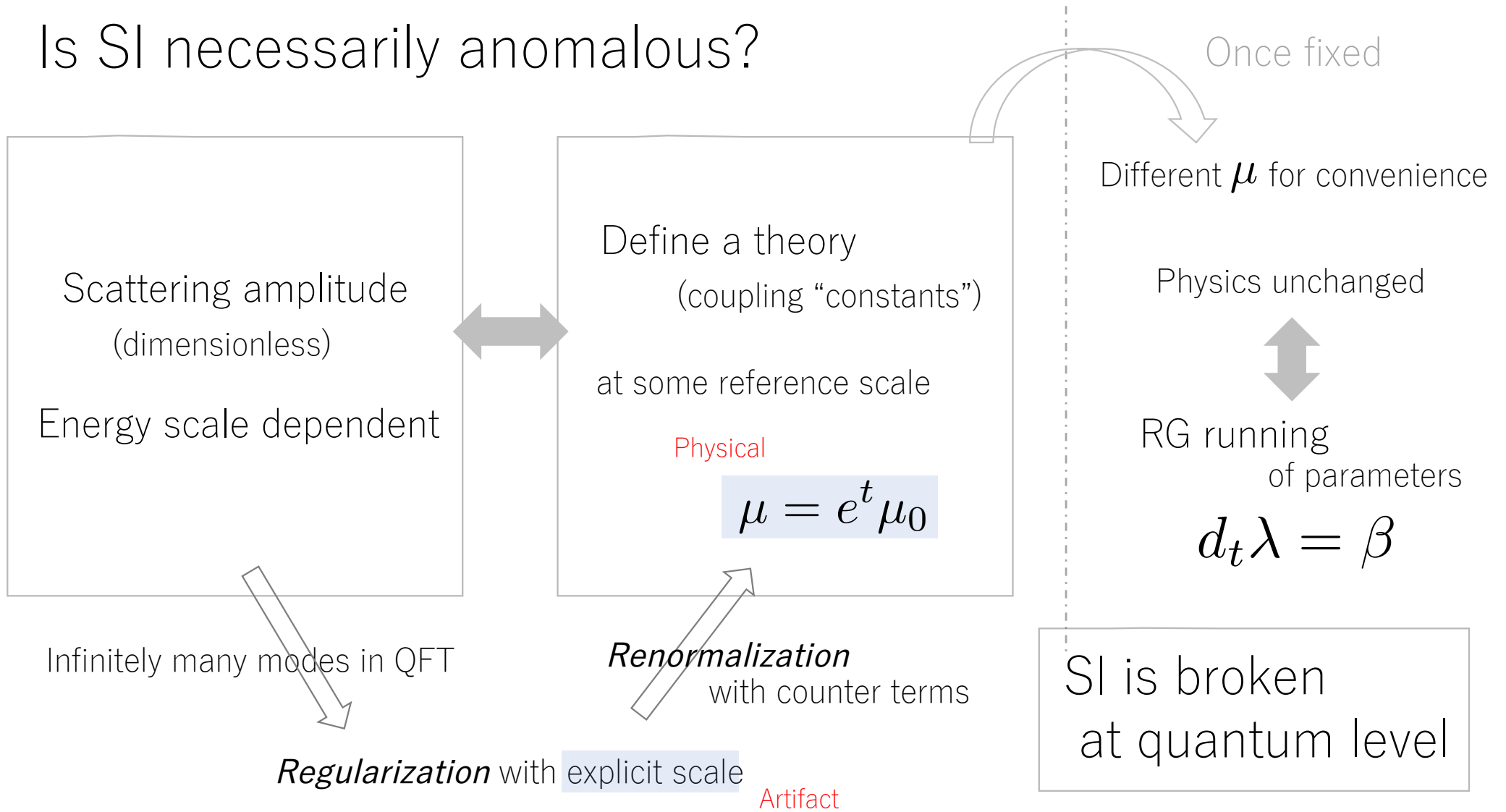
Is SI necessarily anomalous?



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Is SI necessarily anomalous?

Not anomalous if there is no explicit scale

➔ **Dynamical** “reference” scale to define a theory (↔ Different UV completion)

$$\omega = \phi \times f(\hbar/\phi, \dots)$$



Quantum Scale Inv.
with $\beta \neq 0$

F.Englert, C.Truffin, R.Gastmans (1976)

M.Shaposhnikov, D.Zenhausern (2009)

C.Tamarit (2013)

D.M.Ghilencea, Z.Lalak, P.Olszewski (2016)

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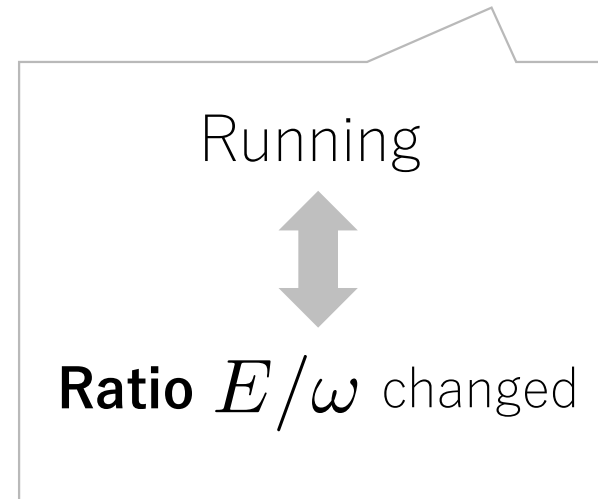
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Asymptotic scale invariance

Assume the simplest form

$$\omega^2 \propto \xi h^2 + \zeta \phi^2 + \dots \quad \text{Dynamical fields} \quad (\text{Exactly SI})$$



Dimensionless constants

Asymptotic scale invariance

Assume the simplest form

$$\omega^2 \propto \xi h^2 + \zeta \phi^2 + \dots \quad \text{Dynamical fields}$$

➔ $h^2 + \mu_\star^2$

SM Higgs

- Maybe regarded as a toy model of the exactly SI one.
- No new d.o.f involved.
- μ_\star is “hidden” at tree-level

Scale-invariant for $h \gg \mu_\star$

Asymptotic scale invariance

Couplings still run : $\beta \neq 0$

What's the consequence of scale invariance?

Coleman-Weinberg one-loop correction

Effective
Higgs potential

$$V_{\text{eff}}(h) = \frac{\lambda h^4}{4} + \sum_i (-1)^{F_i} \frac{m_i^4}{(4\pi)^2} \ln \frac{m_i^2}{\omega^2} + \dots$$

SM particle masses $m_i^2 \propto h^2$

: Relevant energy scale
in vacuum diagrams

Reference scale

Compared with

Asymptotic scale invariance

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Reference scale

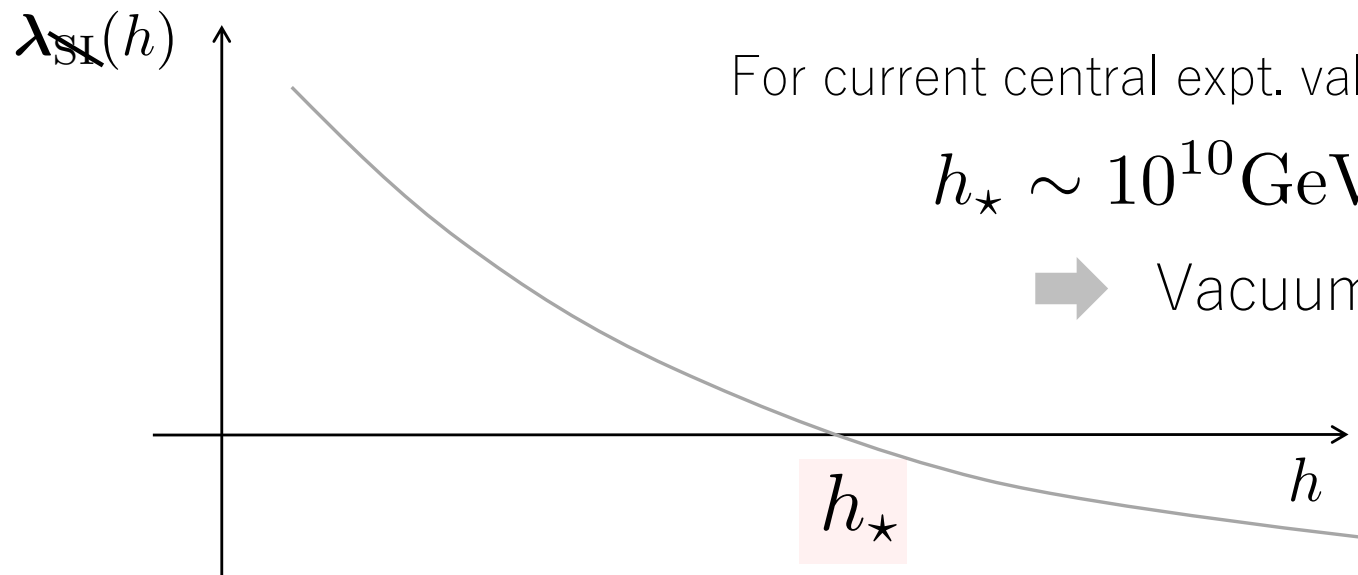
→ $\frac{V_{\text{eff}}}{h^4/4} \equiv \boldsymbol{\lambda}(h) = \lambda + B \ln \frac{h^2}{\omega^2} + \dots$ “Running”

Asymptotic scale invariance

When SI is *anomalous*,

$$\lambda_{\text{SI}}(h) = \lambda + B \ln \frac{h^2}{\mu^2} + \dots$$

Explicit breaking



For current central expt. value of m_t ,

$$h_* \sim 10^{10} \text{ GeV}$$

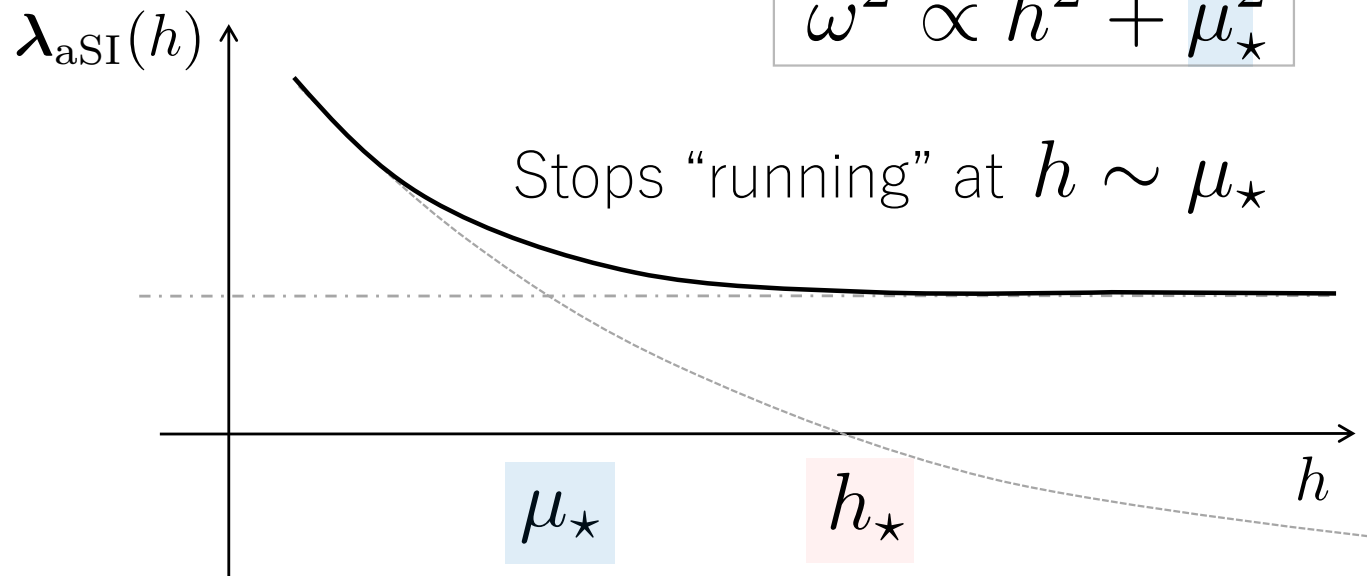
➔ Vacuum meta-stability

Asymptotic scale invariance

$$E/\omega \Rightarrow \text{Constant}$$

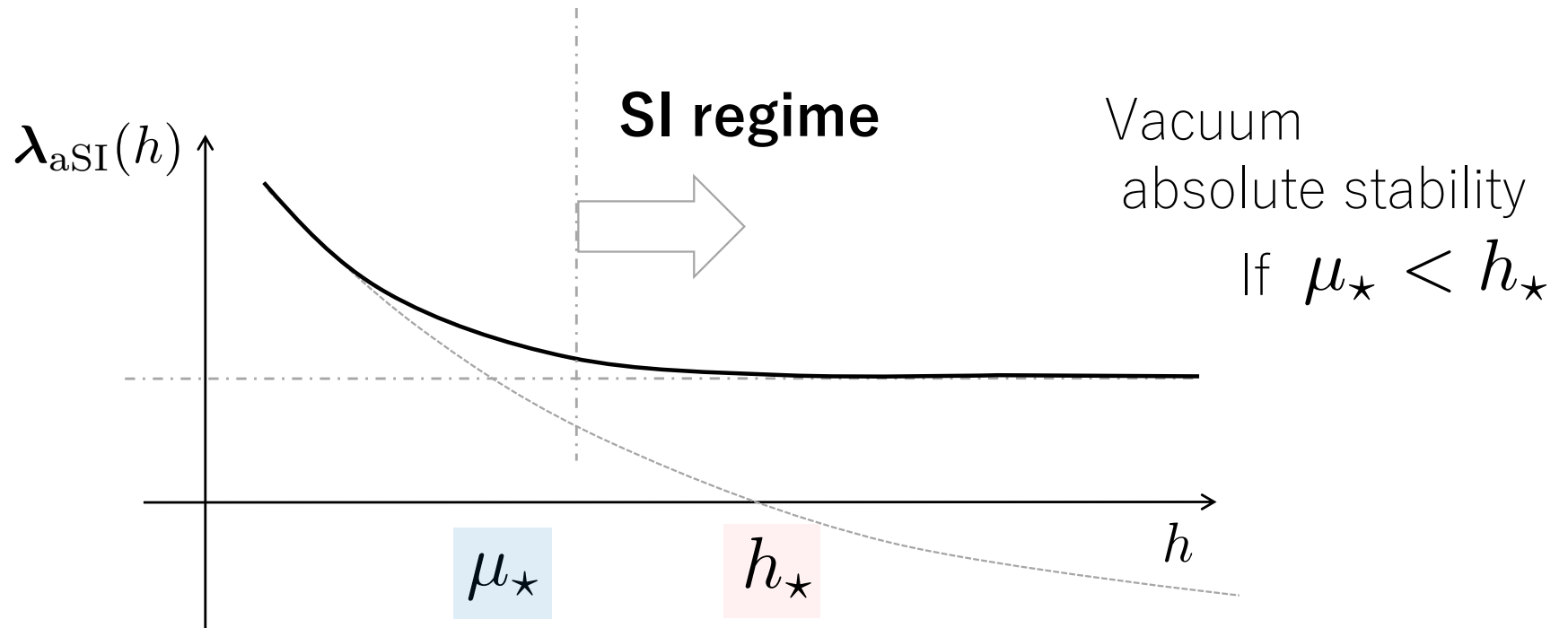
$$\lambda_{\text{aSI}}(h) = \lambda + B \ln \frac{h^2}{\omega^2} + \dots \quad \Rightarrow \quad \lambda_{\text{aSI}}(\infty) = \lambda_{\text{SI}}(\mu_\star)$$

$$\omega^2 \propto h^2 + \mu_\star^2$$



Asymptotic scale invariance

$V_{\text{eff}}(h) \propto h^4$ for $h \gtrsim \mu_\star$ as if there were no quantum correction



Self-consistency

Regularization
Renormalization } both respect SI for $\hbar \gg \mu_\star$.

Asymptotic SI is manifest
at each order of perturbative computation.

Self-consistency

Regularization
Renormalization } both respect SI for $h \gg \mu_*$.

Dimensional regularization ($n = 4 - 2\varepsilon$)

$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4} \quad : \text{mass dimension } n$$

$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} = \omega \Big|_{\text{BG}}^{\frac{2\varepsilon}{1-\varepsilon}} \times \left(1 + \frac{2\varepsilon h}{\mu_*^2 + h^2} \delta h + \dots \right)$$

$$\omega^2 \propto \mu_*^2 + h^2$$

Fluctuation

$$h \rightarrow h + \delta h$$

Background

Self-consistency

Regularization
Renormalization } both respect SI for $h \gg \mu_*$.

Dimensional regularization ($n = 4 - 2\varepsilon$)

$$\omega^2 \propto \mu_*^2 + h^2$$

$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4} : \text{Non-renormalizable}$$

➔ Non-polynomial operators for renormalization

$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{h^{4+2k}}{(\mu_*^2 + h^2)^k}$$

Asymptotically
scale-invariant

$$(h^{1/d_h})^n \text{ for } h \gg \mu_*$$

Self-consistency

Regularization
Renormalization } both respect SI for $h \gg \mu_*$.

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$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4} : \text{Non-renormalizable}$$

$$\omega^2 \propto \mu_*^2 + h^2$$

➔ Non-polynomial operators for renormalization

Up to what energy scale is this **effective theory** valid?

Validity of EFT approach

Non-polynomial operators $\frac{h^{4+2k}}{(\mu_\star^2 + h^2)^k}$ required

Unitarity bound

N -particle amplitude

$$\mathcal{M}_N \sim E^{4-N}$$

at most

J.M.Cornwall, D.N.Levin, G.Tiktopoulos (1974)

Tree unitarity violation

$$\text{at } \Lambda \sim \sqrt{\mu_\star^2 + h^2}$$

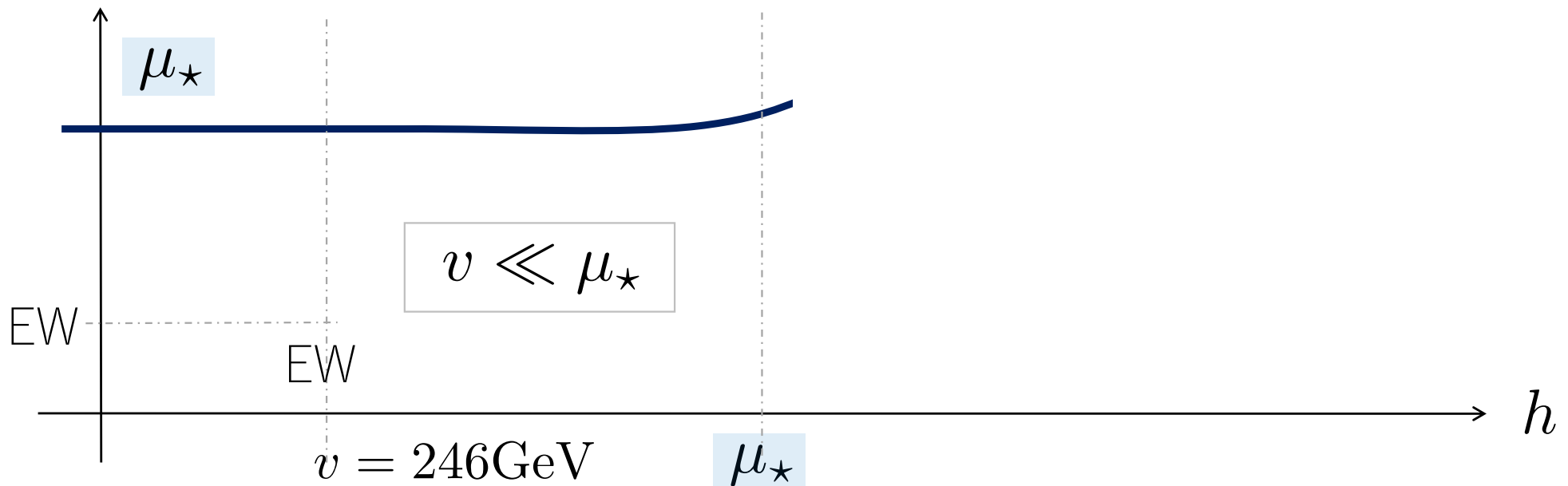
{ Strong coupling
or
New physics

Validity of EFT approach

Non-polynomial operators $\frac{h^{4+2k}}{(\mu_\star^2 + h^2)^k}$ required

$\Lambda \sim \sqrt{\mu_\star^2 + h^2}$: Field dependent

F.Bezrukov, A.Magnin, M.Shaposhnikov, S.Sibiryakov (2010) Higgs inflation

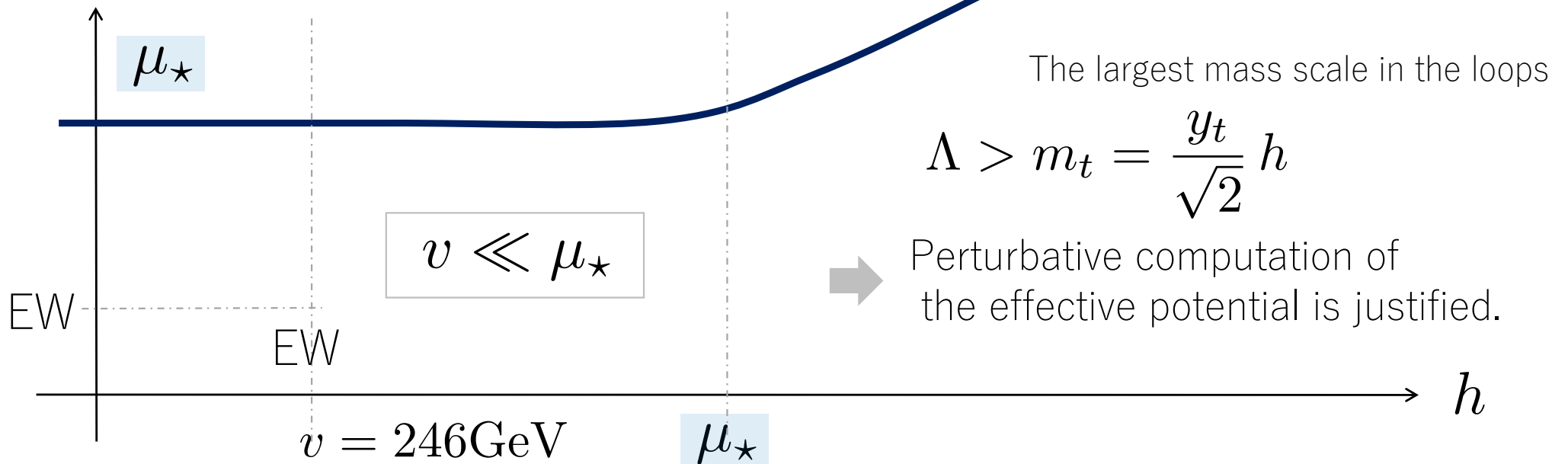


Validity of EFT approach

Non-polynomial operators $\frac{h^{4+2k}}{(\mu_\star^2 + h^2)^k}$ required

$$\Lambda \sim \sqrt{\mu_\star^2 + h^2} \quad : \text{Field dependent}$$

$\sim h$ **SI regime**



The largest mass scale in the loops

$$\Lambda > m_t = \frac{y_t}{\sqrt{2}} h$$

Perturbative computation of the effective potential is justified.

Some comments

Tree unitarity violation scale

$$\Lambda \sim \sqrt{\mu_\star^2 + h^2} \neq$$

Heavy particle

$$M \sim \Lambda$$

Cf. Planck scale

as tree unitarity violation scale

➔ does not necessarily mean large radiative correction to Higgs mass.

F.Bezrukov, M.Shaposhnikov (2007)

J.Garcia-Bellido, J.Rubio, M.Shaposhnikov, D.Zenhausern (2011) Higgs-Dilaton model

Asymptotic SI in *Higgs inflation* (prescription 1)

➔ Asymptotically flat potential
(Einstein frame)

$$\omega^2 \propto M_{\text{P,eff}}^2 = M_{\text{P}}^2 + \xi h^2$$

Effective Planck mass in Jordan frame

In general, $\omega^2 \not\propto M_{\text{P,eff}}^2$

➔ More variety of
Higgs potential shapes
(before asymptotic flatness)

Summary

Asymptotically scale-invariant model

- ✓ Quantum Scale invariance for $h \gg \mu_*$.
- ✓ Perturbative realization within EFT with dimensional regularization.
- ✓ Field dependent tree unitarity violation scale $\Lambda \sim \sqrt{\mu_*^2 + h^2}$.
- ✓ Effective potential is computable without knowing UV completion.

Cosmology based on effective potential

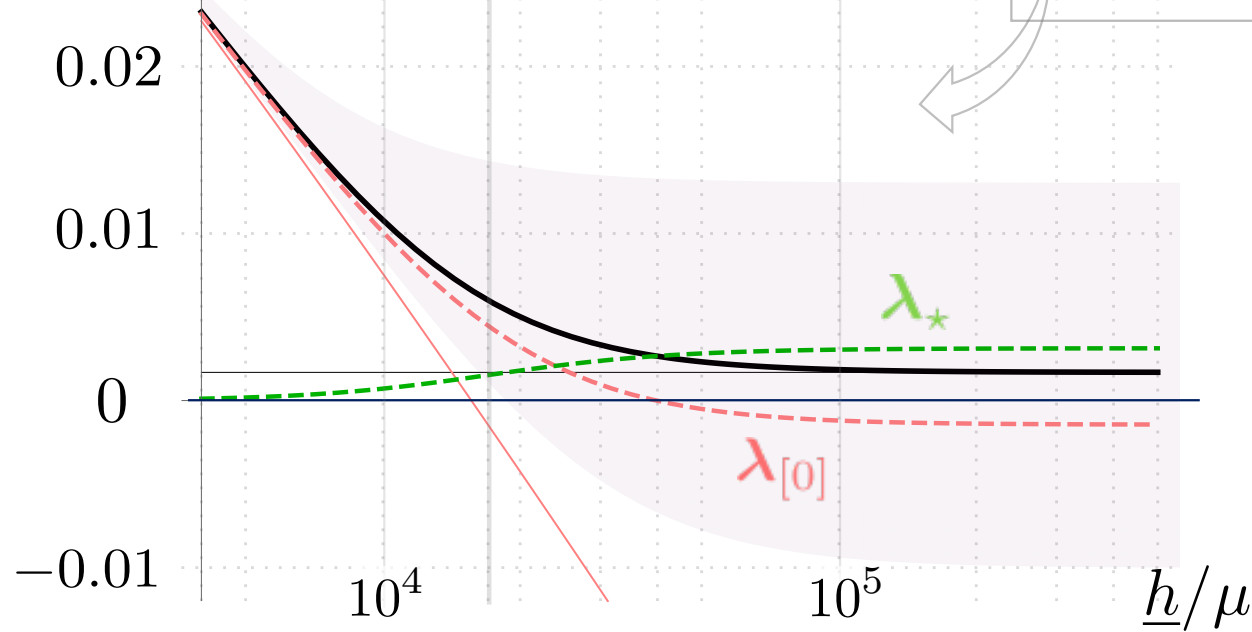
- ✓ Asymptotic scale invariance can be responsible for
absolute stability of our EW vacuum.

Thank you

Asymptotic scale invariance

$$V_{\text{eff}}(h) = \sum_{k=0}^{\infty} \frac{\lambda_{[k]}}{4} \frac{h^{4+2k}}{(\mu_{\star}^2 + h^2)^k}$$

$$\lambda_{\text{aSI}}(\underline{h}; \mu_{\star}; \mu)$$



Two-loop computation with assumption
 No non-polynomial operator at tree-level
 ↓
 Correction with higher power k
 → More loop-suppressed

Asymptotic value

$$\sum_{k=0}^{\infty} \lambda_{[k]}$$

Finiteness (scale invariance) constraints UV completion.

Asymptotic SI and Higgs inflation

$$\frac{\mathcal{L}}{\sqrt{-g}} = \frac{M_{\text{P,eff}}^2}{2} R - \frac{g^{\mu\nu}}{2} \partial_\mu h \partial_\nu h - V(h) + \dots$$

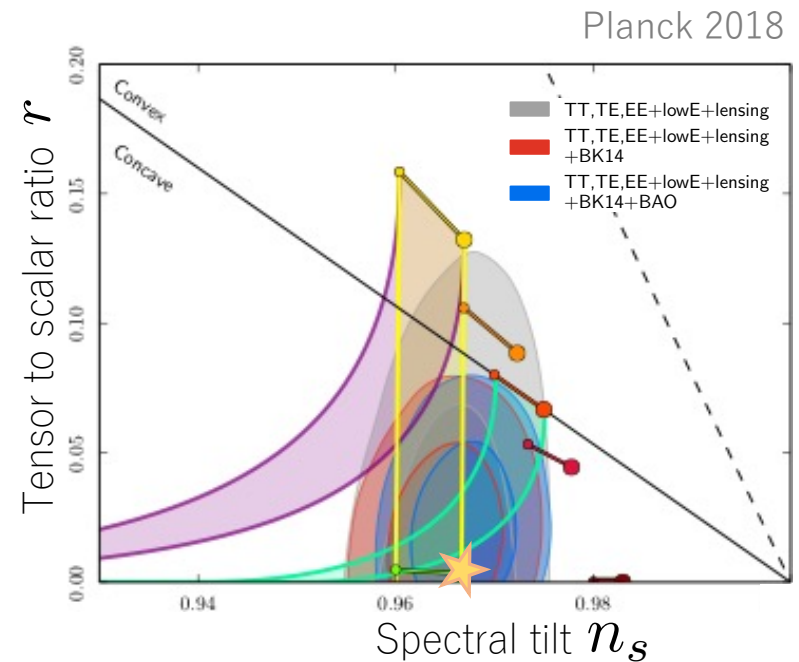
Effective
Planck mass

$$M_{\text{P,eff}}^2 = M_{\text{P}}^2 + \xi h^2$$

F.Bezrukov, M.Shaposhnikov (2007)

$$\xi \sim 10^4 \sqrt{\lambda} \quad \text{Large non-minimal coupling}$$

$$\rightarrow A_s \simeq 2.2 \times 10^{-9}$$



Asymptotic SI and Higgs inflation

Renormalization prescription	$\omega^2 \propto$	
I	$M_{\text{P}}^2 + \xi h^2 = M_{\text{P},\text{eff}}^2$	F.Bezrukov, M.Shaposhnikov (2007)
II	M_{P}^2 (constant)	A.O.Barvinsky, A.Y.Kamenshchik, A.A.Starobinsky (2008)

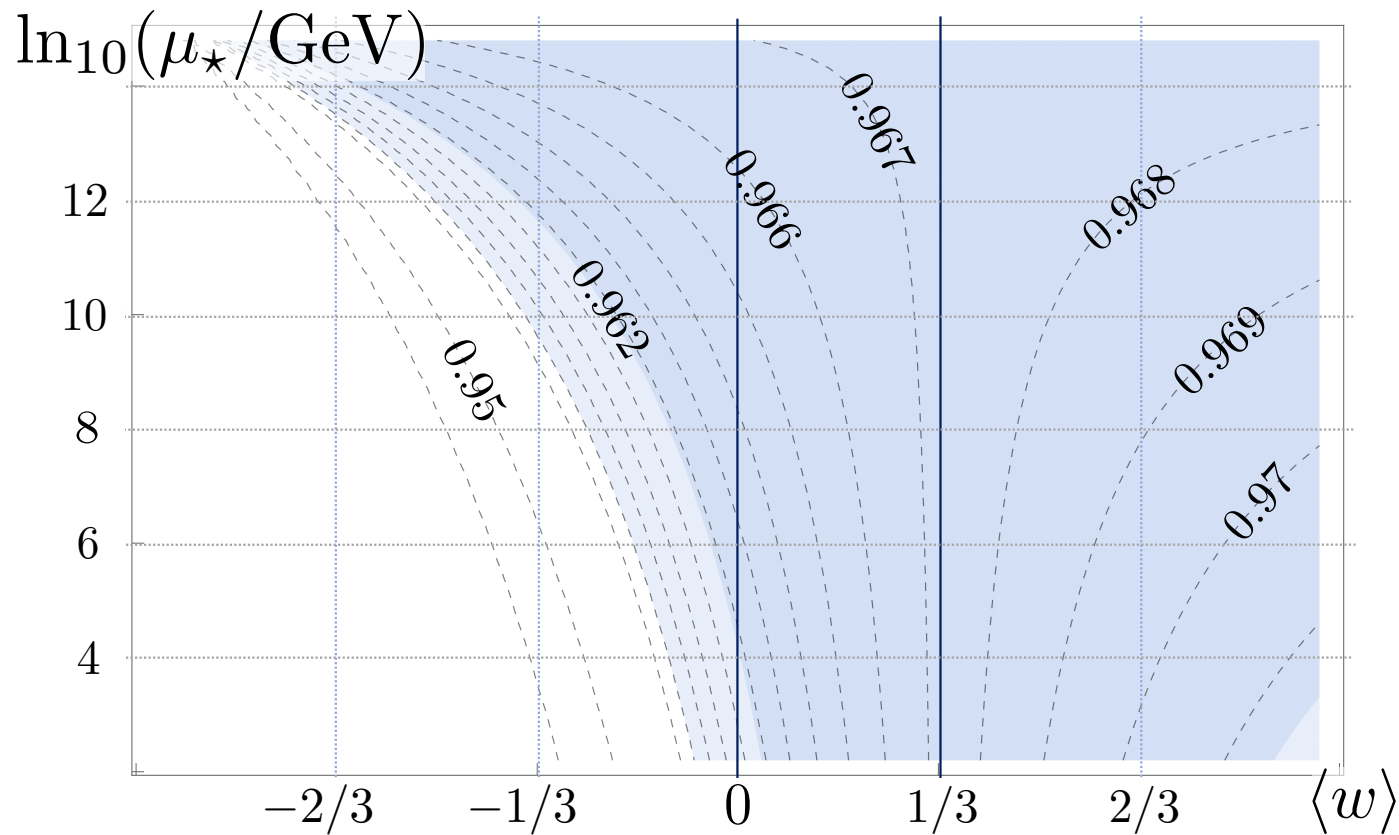
$$\frac{\lambda h^4}{4} \Rightarrow \omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4}$$

→ $\mu_{\star}^2 + h^2 \propto M_{\text{P}}^2 + \xi_{\star} h^2$

$$\xi_{\star} = M_{\text{P}}^2 / \mu_{\star}^2 \neq \xi$$

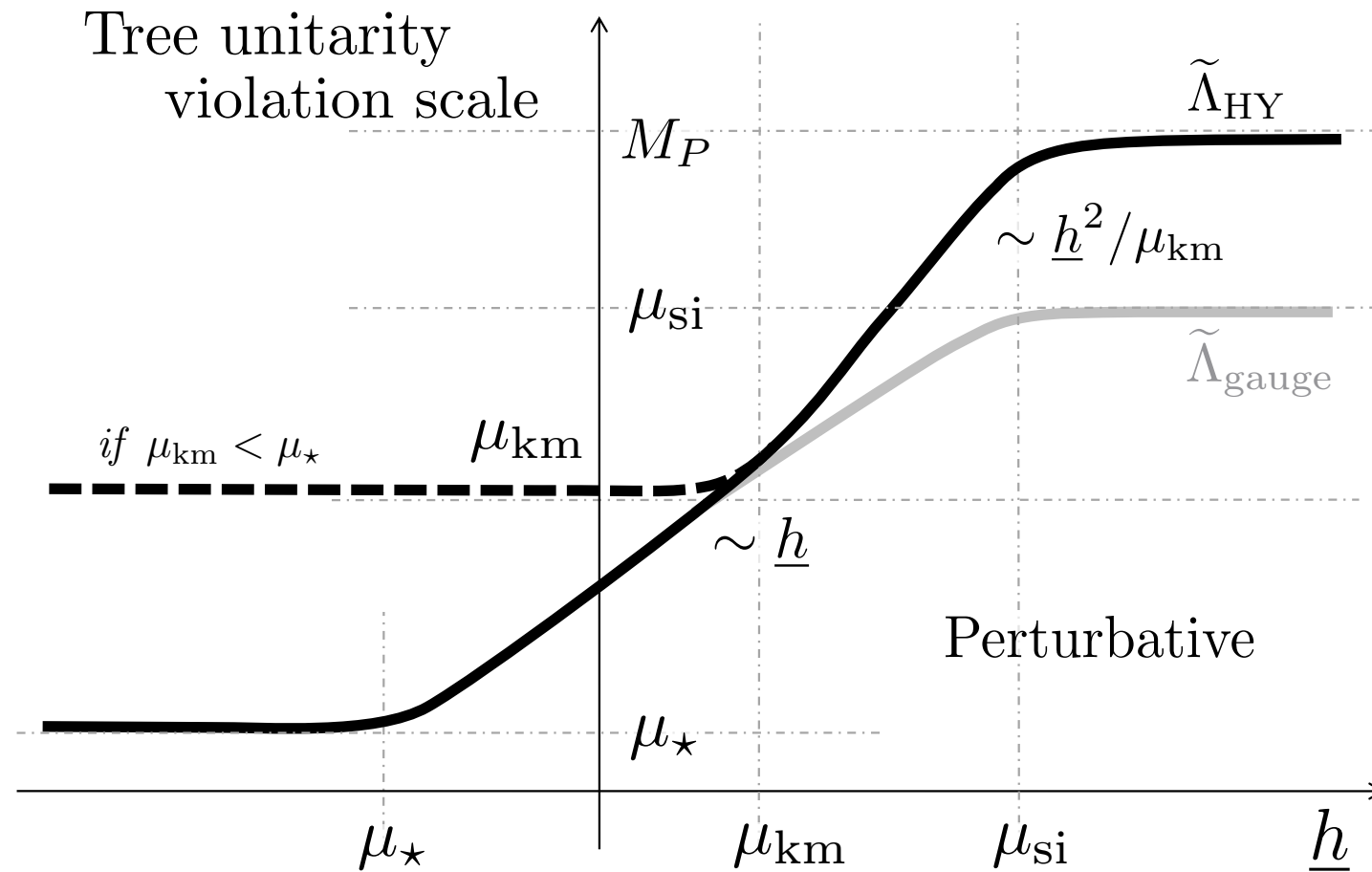
Asymptotic SI and Higgs inflation

Contour plot of spectral tilt n_s



e-folding-averaged equation of state during high-energy phase

Validity of EFT approach with nonminimal coupling



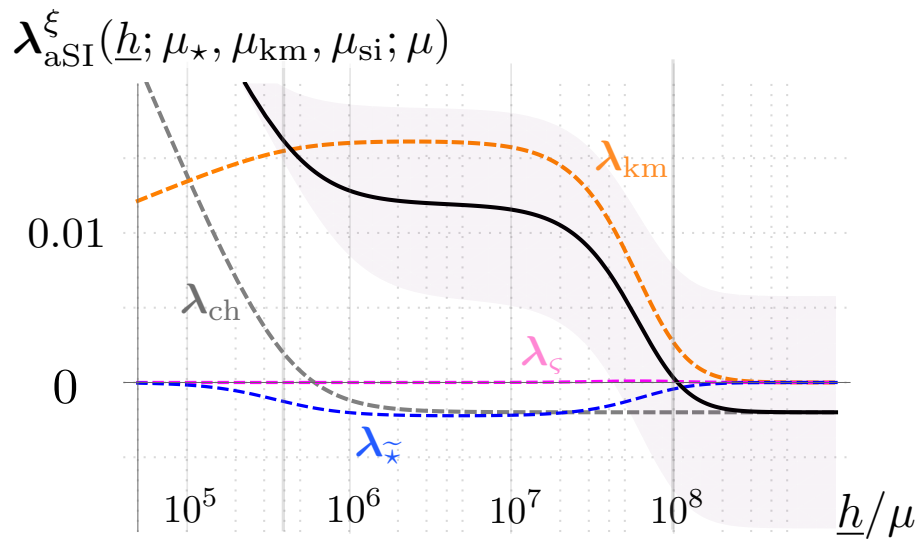
$$\mu_{\text{si}} \sim M_P / \sqrt{\xi}$$

$$\mu_{\text{km}} \sim M_P / \xi$$

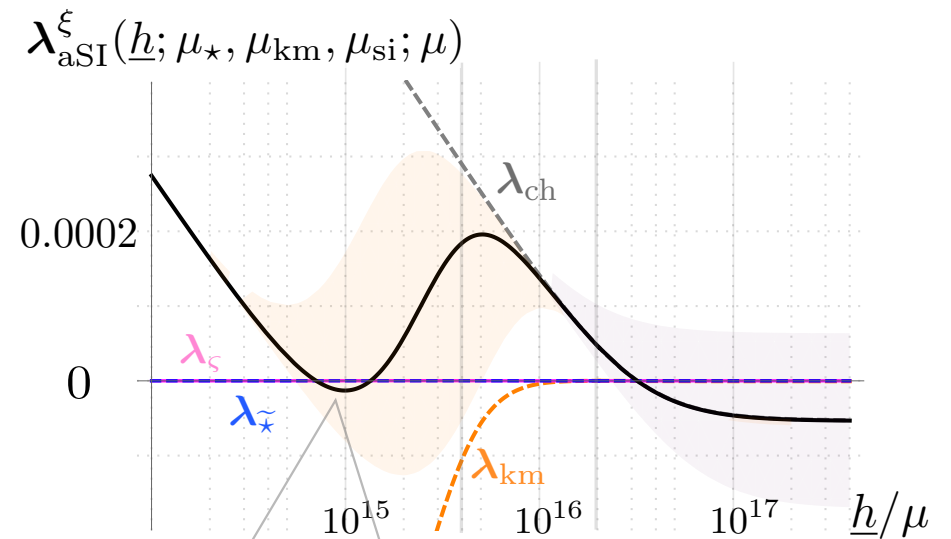
$$\text{for } \xi \gg 1$$

λ stops “running” before/after it jumps

$$\mu_\star \ll \mu_{\text{km}}$$



$$\mu_{\text{km}} \ll \mu_\star$$



$$h \sim \mu_{\text{km}}$$

Higgs-graviton mixing becomes significant. (Jordan frame) \rightarrow Jump

Jump makes Higgs inflation possible even with vacuum meta-stability.

F.Bezrukov, J.Rubio, M.Shaposhnikov (2015)