

Forbidden Freeze-In

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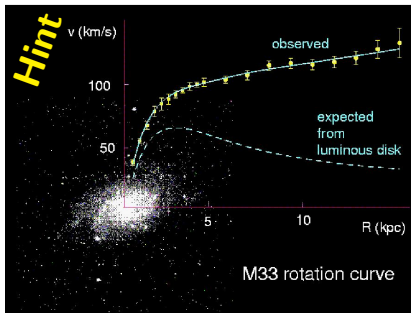
In collaboration with L. Darmé, A. Hryczuk, L. Roszkowski

[arXiv:1907.XXXXX](https://arxiv.org/abs/1907.XXXXX)

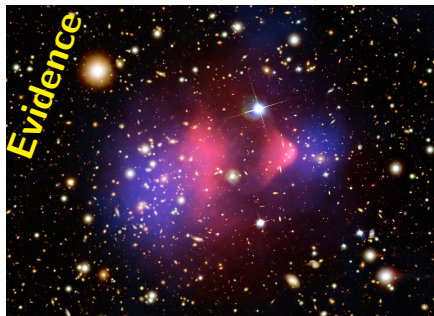
- 1 Introduction
 - Why Particle Dark Matter
 - The Dark Matter Particle
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 - Standard Freeze-in
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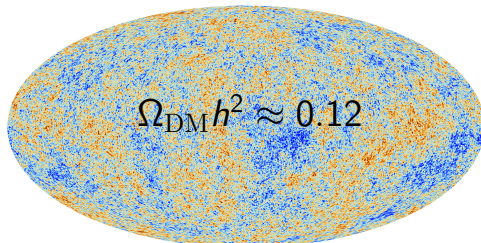
Why Particle Dark Matter



E. Corbelli and P. Salucci, *Mon. Not. Roy. Astron. Soc.* **311** 441 (2000),
arXiv:astro-ph/9909252.



M. Markevitch, *ESA Spec. Publ.* **604** (2006) 723, astro-ph/0511345.
Clowe, Bradac, et. al. *Astrophys. J.* **648**, L109 (2006), astro-ph/0608407



N. Aghanim et al. [Planck Collaboration], arXiv:1807.06209 [astro-ph.CO].

The Dark Matter Particle

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- Cold/Warm and non-relativistic today.

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Dark Matter production via Freeze-in: ¹

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L. Covi, H. B. Kim, J. E. Kim and L. Roszkowski, JHEP **0105**, 033 (2001), hep-ph/0101009

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arXiv:0911.1120 [hep-ph]

F. Elahi, C. Kolda and J. Unwin, JHEP **1503**, 048 (2015), arXiv:1410.6157 [hep-ph]

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- Number of particles stabilizes at $T_{FI}^{IR} \sim m_S$ or $T_{FI}^{UV} \sim T_{RH}$.

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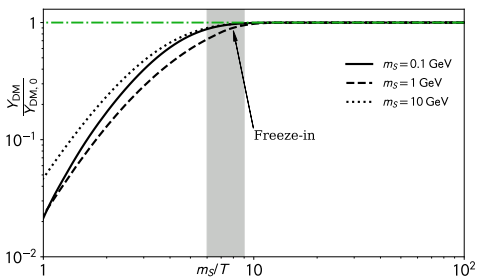
F. Elahi, C. Kolda and J. Unwin, JHEP **1503**, 048 (2015), arXiv:1410.6157 [hep-ph]

IR and UV Freeze-in

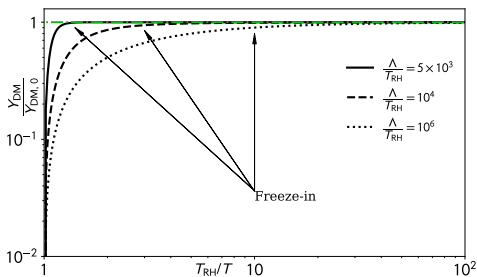
Consider χ the Dark Matter particle, and S in equilibrium with the plasma.

$$\mathcal{L}_{\text{int}} = -yS\bar{\chi}\chi$$

$$\mathcal{L}_{\text{int}} = -\frac{1}{\Lambda}SS\bar{\chi}\chi$$



IR



UV

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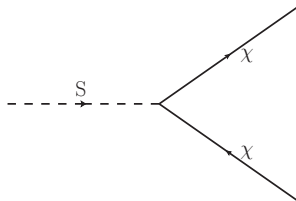
Behavior was noted in:

A. Strumia, JHEP **1006**, 036 (2010),
arXiv:1003.5847 [hep-ph].

M. J. Baker, M. Breitbach, J. Kopp and
L. Mittnacht, JHEP **1803**, 114 (2018),
arXiv:1712.03962 [hep-ph].

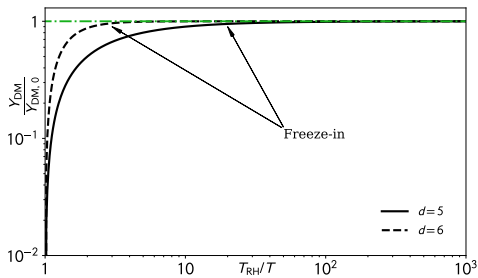
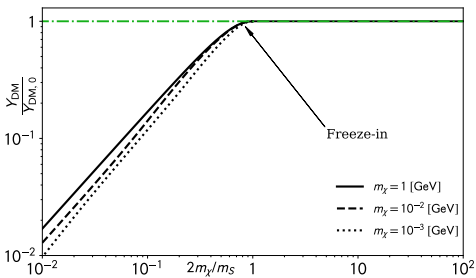
L. Bian and Y. L. Tang, JHEP **1812**, 006
(2018), arXiv:1810.03172 [hep-ph].

IR and UV Forbidden Freeze-in



IR

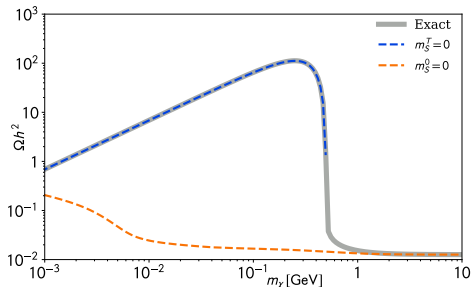
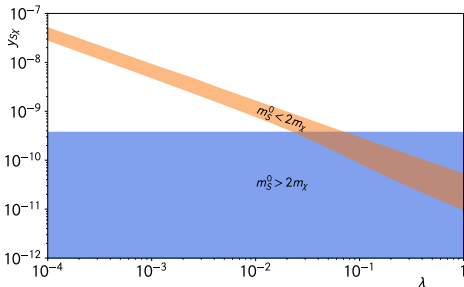
UV



Standard vs Forbidden Freeze-in

Comparing the two types of Freeze-in.

$$\mathcal{L} \supset -y_{S\chi} S \bar{\chi}\chi - \frac{\lambda}{4!} S^4 - \frac{1}{2} m_S^{02} S^2 - m_\chi \bar{\chi}\chi + (\text{SH - terms}),$$
$$m_S^2(T) \approx m_S^{02} + \frac{\lambda}{24} T^2.$$



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Assume Dark Matter (χ) a Dirac particle, and portal particle (S) that couples to the Higgs:

$$\mathcal{L}^{\text{DM}} = \bar{\chi} (i\gamma_\mu D^\mu - \mu_\chi) \chi + \frac{1}{2} (D^\mu S)(D_\mu S) - y_{XS} S \bar{\chi} \chi - V_{HS} ,$$

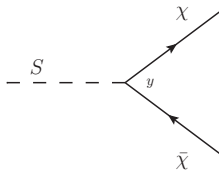
$$V_{HS} = \frac{\mu_S^2}{2} S^2 + \frac{\lambda_S}{4!} S^4 + A S H^\dagger H + \lambda_{HS} S^2 H^\dagger H .$$

After Electroweak phase transition $S \rightarrow S - v_S$ ($v_S \ll \mu_S, v$), and S and h mix by $\theta \ll 1$. This induces interaction of the form

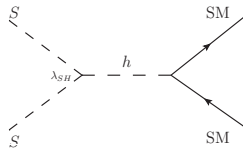
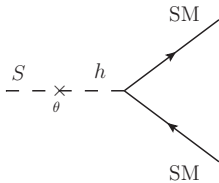
$$\mathcal{L}^{\text{YSM}} = \theta \frac{m_f}{v} S \bar{f} f .$$

Relevant Processes

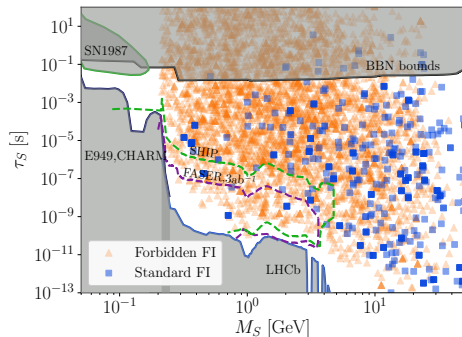
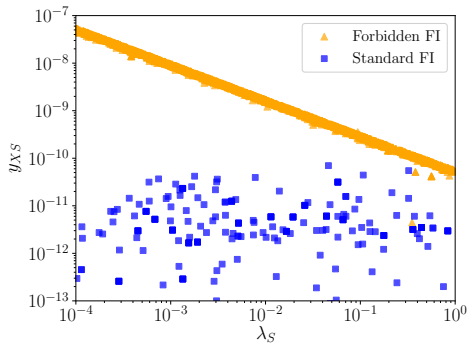
Dark Matter is produced by decays of S .



S is kept in equilibrium via its interactions with the SM.



Searching for observed relic using MultiNest ²:



²F. Feroz, M. P. Hobson and M. Bridges Mon. Not. Roy. Astron. Soc. **398**, 1601 (2009) [arXiv:0809.3437 [astro-ph]].

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Future direction:

- Study *a lot* of models under Forbidden Freeze-in.

Thank you

Backup

$$s = \frac{2\pi^2}{45} h(T) T^3$$

$$\delta_h = 1 + \frac{1}{3} \frac{d \log(h)}{d \log(T)}$$

$$H = \sqrt{\frac{4\pi^3}{45 m_P^2} g(T) T^2},$$

$$x = \frac{m_S}{T},$$

$$\frac{dY_{\text{DM}}}{dx} = \left(\frac{\Gamma_\chi(m_S, m_\chi)}{5.93 \times 10^{-19} \text{ GeV}} \right) \left(\frac{1 \text{ GeV}}{m_S} \right)^2 \frac{K_1(x) x^3}{\sqrt{g} h} \delta_h,$$

$$Y_{\text{DM},0}^{\text{UV}} \approx \frac{4^n n! (n+1)!}{2.3 \times 10^{-15} g_\star^{3/2}(T_{\text{RH}})} \left(\frac{\text{GeV}}{m_S} \right) \left(\frac{m_S}{\Lambda} \right)^{2n} \frac{x_{\text{RH}}^{1-2n}}{2n-1},$$

$$Y_{\text{DM},0}^{\text{IR}} \approx \left(\frac{\Gamma_\chi(m_S, m_\chi)}{1.6 \times 10^{-36} \text{ GeV}} \right) \left(\frac{1 \text{ GeV}}{m_S} \right)^2 m_\chi \left(\frac{1}{\sqrt{g} h} \right) \Big|_{x=\langle x \rangle},$$

$$\langle x \rangle \equiv \frac{\int_0^\infty dx x^3 K_1(x) \times x}{\int_0^\infty dx x^3 K_1(x)} \approx 3.4.$$

$$z \equiv \frac{2m_\chi}{\alpha T} ,$$

$$\Gamma_\chi \sim \frac{\gamma_{S\chi}}{16\pi} m_S \left(\frac{m_S}{\Lambda}\right)^{2n} = \frac{\gamma_{S\chi}}{16\pi} \alpha^{2n+1} \left(\frac{T}{\Lambda}\right)^{2n} T ,$$

$$\frac{dY_{\text{DM}}}{dz} = \left(\frac{\Gamma_\chi}{5.93 \times 10^{-19} \text{ GeV}}\right) \left(\frac{1 \text{ GeV}}{2m_\chi}\right)^2 \frac{\alpha^4 K_1(\alpha)}{\sqrt{g} h} \delta_h z^3 .$$

$$Y_{\text{DM},0}^{\text{UV}} = \frac{z_{\text{RH}}^{1-2n} - 1}{2n - 1} \left(\frac{\alpha^4 K_1(\alpha) \gamma_{S\chi}}{2.96 \times 10^{-17}}\right) \left(\frac{2m_\chi}{\Lambda}\right)^{2n}$$

$$\left(\frac{1 \text{ GeV}}{2m_\chi g_\star^{3/2}(T_{\text{RH}})}\right) ,$$

$$Y_{\text{DM},0} = \left(\frac{\alpha^2 y_{S\chi}}{7.1 \times 10^{-9}}\right)^2 \left(\frac{\text{GeV}}{m_\chi}\right) K_1(\alpha) \left(\frac{1}{\sqrt{g} h}\right)_{z=\langle z \rangle} ,$$

$$\langle z \rangle \equiv \frac{\int_0^1 dz (1 - z^2)^{3/2} \times z}{\int_0^1 dz (1 - z^2)^{3/2}} \approx 0.34 .$$