Sbottoms as probes to MSSM with nonholomorphic soft interactions

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[Ref.: UC, **Abhishek Dey**, JHEP 1610 (2016) 027, arXiv:1604.06367],
[UC, **AseshKrishna Datta, Samadrita Mukherjee, Abhaya Kumar Swain**, JHEP 1810 (2018) 202, arXiv: 1809.05438]

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MSSM

- MSSM Superpotential and soft SUSY breaking terms::

$$\mathcal{W} = \mu H_D.H_U - Y^e_{ij} H_D.L_i \bar{E}_j - Y^d_{ij} H_D.Q_i \bar{D}_j - Y^u_{ij} Q_i.H_U \bar{U}_j$$

$$A.B = \epsilon_{\alpha\beta} A^\alpha B^\beta$$

$$\begin{aligned} -\mathcal{L}_{soft} &= [\tilde{q}_{iL}.hu(A_u)_{ij}\tilde{u}^*_{jR} + h_d.\tilde{q}_{iL}(A_d)_{ij}\tilde{d}^*_{jR} + h_d.\tilde{l}_{iL}(A_e)_{ij}\tilde{e}^*_{jR} + h.c.] \\ &+ (B\mu h_d.h_u + h.c.) + m_d^2|h_d|^2 + m_u^2|h_u|^2 \\ &+ \tilde{q}^*_{iL}(M^2_{\tilde{q}})_{ij} + \tilde{u}^*_{iR}(M^2_{\tilde{u}})_{ij}\tilde{u}_{jR} + \tilde{d}^*_{iR}(M^2_{\tilde{d}})_{ij}\tilde{d}_{jR} + \tilde{l}^*_{iL}(M^2_{\tilde{l}})_{ij}\tilde{l}_{jL} \\ &+ \text{gaugino mass terms} \end{aligned}$$

- Possible origin of soft terms: SUSY breaking parametrized by vev of F-term of a chiral superfield X, so that $< X >= \theta\theta < F >\equiv \theta\theta F$. X couples to Φ and a gauge strength superfield W^a_α.

Type	Term	Naive Suppression	Origin
soft	$\phi\phi^*$	$\frac{\|F\|^2}{M^2} \sim m_W^2$	$\frac{1}{M^2}[XX^*\Phi\Phi^*]_D$
	ϕ^2	$\frac{\mu F}{M} \sim \mu m_W$	$\frac{\mu}{M}[X\Phi^2]_F$
	ϕ^3	$\frac{F}{M} \sim m_W$	$\frac{1}{M}[X\Phi^3]_F$
	$\lambda\lambda$	$\frac{F}{M} \sim m_W$	$\frac{1}{M}[XW^\alpha W_\alpha]_F$

- Are there any more possible soft terms ?

Nonholomorphic soft SUSY breaking terms

- **S. Martin, Phys. Rev D., 2000; Possible non-holomorphic soft SUSY breaking terms:**

Type	Term	Naive Suppression	Origin
"maybe soft"	$\phi^2\phi^*$	$\frac{\lvert F\rvert^2}{M^3} \sim \frac{m_W^2}{M}$	$\frac{1}{M^3}[XX^*\Phi^2\Phi^*]_D$
	$\psi\psi$	$\frac{\lvert F\rvert^2}{M^3} \sim \frac{m_W^2}{M}$	$\frac{1}{M^3}[XX^*D^\alpha\Phi D_\alpha\Phi]_D$
	$\psi\lambda$	$\frac{\lvert F\rvert^2}{M^3} \sim \frac{m_W^2}{M}$	$\frac{1}{M^3}[XX^*D^\alpha\Phi W_\alpha]_D$

- "maybe soft": *In the absence of a gauge singlet field* the above non-holomorphic terms are of soft SUSY breaking in nature. But, these have mass scale suppression by M.
- A gauge singlet scalar field would have tadpole contributions causing hard SUSY breaking [Bagger and Poppitz PRL 1993].
- NHSSM: Also known for non-standard soft terms: MSSM + NH terms like $\phi^2\phi^*$ along with a higgsino mass soft term of $\psi\psi$ type.

$$-\mathcal{L}'_{soft} = h_d^c.\tilde{q}_{iL}(A'_u)_{ij}\tilde{u}^*_{jR} + \tilde{q}_{iL}.h_u^c(A'_d)_{ij}\tilde{d}^*_{jR} + \tilde{l}_{iL}.h_u^c(A'_e)_{ij}\tilde{e}^*_{jR} + \mu'\tilde{h}_u.\tilde{h}_d + h.c.$$

Higgs fields are replaced with their conjugates: h_d going with up-type of squarks etc. We consider the new parameters to be of unknown origin while having strengths similar to other soft terms.

- $V_{\rm Higgs}$ is unaffected. But, the potential involving charged and colored scalar fields needs a separate study for CCB [**J. Beuria and A. Dey**, JHEP 2017].

Nonholomorphic terms: A partial list of related analyses and our present work

- **Early mentions: Girardello, Grisaru 1982, Hall and Randall 1990"** labelled as hard SUSY breaking terms while considering gauge singlets in the picture. But, MSSM does not have a gauge singlet. **Jack and Jones**, PRD 2000: Quasi IF fixed points and RG invariant trajectories; **Jack and Jones** PLB 2004: General analyses with NH terms involving RG evolutions.
- Works performed under Constrained MSSM (CMSSM)/minimal supergravity(mSUGRA) setup for studying the Higgs mass and observables like $\mathrm{Br}(B \to X_s + \gamma)$ etc.: **Hetherington** JHEP 2001, **Solmaz et. al.** PRD 2005, PLB 2008, PRD 2015. The analyses involve mixed type of inputs given at the unification and electroweak scales.
- **Ross, Schmidt-Hoberg, Staub** PLB 2016, JHEP 2017. Focused on fine-tuning and higgsino DM, stressed the importance of the bilinear higgsino term and performed RGE.
- **UC, A. Dey** JHEP 2016: No specific mechanism for SUSY breaking: all the parameters are given at the low scale. Impact on muon $g-2$ apart from EW fine-tuning, Higgs mass etc.
  UC, D. Das, S. Mukherjee, JHEP 2018: On GMSB type of realization of NHSSM.
  J. Beuria and A. Dey, JHEP 2017, CCB effects in NHSSM
  UC, A. Datta, S. Mukherjee, A. K. Swain: JHEP 2018, Sbottom phenomenology.

NHSSM: scalars and electroweakinos

- $$\textit{Squarks}: \quad M_{\tilde{u}}^2 = \begin{bmatrix} m_{\tilde{Q}}^2 + (\frac{1}{2} - \frac{2}{3}\sin^2\theta_W)M_Z^2\cos 2\beta + m_u^2 & -m_u(A_u - (\mu + A_u')\cot\beta) \\ -m_u(A_u - (\mu + A_u')\cot\beta) & m_{\tilde{u}}^2 + \frac{2}{3}\sin^2\theta_W M_Z^2 \cos 2\beta + m_u^2 \end{bmatrix},$$

  Sleptons (off-diagonal): $-m_\mu[A_\mu - (\mu + A_\mu')\tan\beta] \Rightarrow A_\mu' \tan\beta$ **potentially enhances** $(g-2)_\mu^{\rm SUSY}$**, particularly affecting the** $\tilde{\chi}_1^0 - \tilde{\mu}$ **loop contributions.**

- $$\text{Higgs mass corrections}: \Delta m_{h,top}^2 = \frac{3g_2^2\bar{m}_t^4}{8\pi^2 M_W^2}\left[\ln\left(\frac{m_{\tilde{t}_1}m_{\tilde{t}_2}}{\bar{m}_t^2}\right) + \frac{X_t^2}{m_{\tilde{t}_1}m_{\tilde{t}_2}}\left(1 - \frac{X_t^2}{12 m_{\tilde{t}_1}m_{\tilde{t}_2}}\right)\right],$$

  Here, $X_t = A_t - (\mu + A_t')\cot\beta \Rightarrow$ influence on m_h.

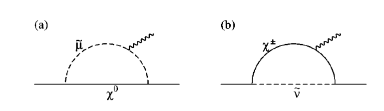
- $$\textit{Charginos}: M_{\tilde{\chi}^\pm} = \begin{pmatrix} M_2 & \sqrt{2}M_W\sin\beta \\ \sqrt{2}M_W\cos\beta & -(\mu + \mu') \end{pmatrix},$$

  $m_{\tilde{\chi}_1^\pm} \gtrsim 100$ GeV $\Rightarrow |\mu + \mu'| \gtrsim 100$ GeV. Neutralino matrix: similar changes from μ to $\mu + \mu'$.

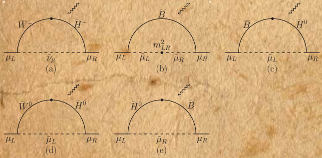
- The Higgs potential at the tree level is independent of μ'. If $|(\mu + \mu')| << M_1, M_2 \Rightarrow \tilde{\chi}_1^0$ is higgsino-like. It is possible to have an acceptable higgsino-like LSP (with mass $\tilde{1}$ TeV) with very small μ ($\sim$i.e. small electroweak fine-tuning.)

Muon anomalous magnetic moment: $(g-2)_\mu$ in MSSM

- Large discrepancy from the SM (more than 3σ): $a_\mu^{exp} - a_\mu^{SM} = (29.3 \pm 8) \times 10^{-10}$
- MSSM contributions to muon (g-2): Diagrams involving charginos and neutralinos



Gauge Eigenstate basis:



$$\Delta a_\mu(\tilde{W}, \tilde{H}, \tilde{\nu}_\mu) \simeq 15 \times 10^{-9} \left(\frac{\tan\beta}{10}\right)\left(\frac{(100\,\mathrm{GeV})^2}{M_2\mu}\right)\left(\frac{f_C}{1/2}\right),$$

$$\Delta a_\mu(\tilde{W}, \tilde{H}, \tilde{\mu}_L) \simeq -2.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right)\left(\frac{(100\,\mathrm{GeV})^2}{M_2\mu}\right)\left(\frac{f_N}{1/6}\right),$$

$$\Delta a_\mu(\tilde{B}, \tilde{H}, \tilde{\mu}_L) \simeq 0.76 \times 10^{-9} \left(\frac{\tan\beta}{10}\right)\left(\frac{(100\,\mathrm{GeV})^2}{M_1\mu}\right)\left(\frac{f_N}{1/6}\right),$$

$$\Delta a_\mu(\tilde{B}, \tilde{H}, \tilde{\mu}_R) \simeq -1.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right)\left(\frac{(100\,\mathrm{GeV})^2}{M_1\mu}\right)\left(\frac{f_N}{1/6}\right),$$

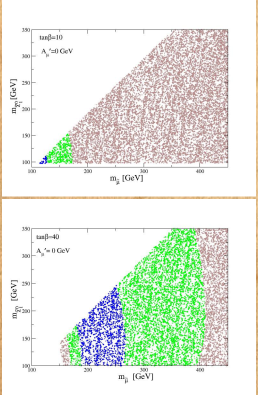
$$\Delta a_\mu(\tilde{\mu}_L, \tilde{\mu}_R, \tilde{B}) \simeq 1.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right)\left(\frac{(100\,\mathrm{GeV})^2}{m^2_{\tilde{\mu}_L} m^2_{\tilde{\mu}_R}/M_1\mu}\right)\left(\frac{f_N}{1/6}\right).$$

- Slepton L-R mixing in MSSM: $m_\mu(A_\mu - \mu\tan\beta)$
- The mixing influences the last item of Δa_μ shown in blue. Typically, A_μ is quite smaller than $\mu\tan\beta$, especially for large $\tan\beta$.
- In NHSSM: $m_\mu[(A_\mu - A'_\mu \tan\beta) - \mu\tan\beta]$ A'_μ **effect is enhanced by** $\tan\beta$ causing a significant change in Δa_μ.

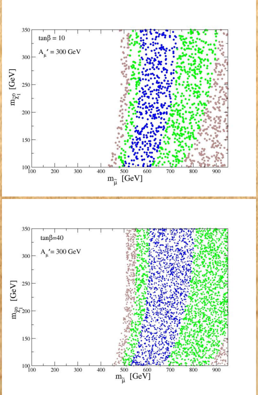
[Ref. arXiv 1303.4256 by Endo, Hamaguchi, Iwamoto, Yoshinaga]

Results of muon g-2 in MSSM

Ref: UC, ***A Dey****, JHEP 1610 (2016) 027, arXiv:1604.06367*



MSSM: for $\tan\beta = 10$ and 40. Blue, green and brown are 1σ, 2σ and 3σ levels of $g-2$.

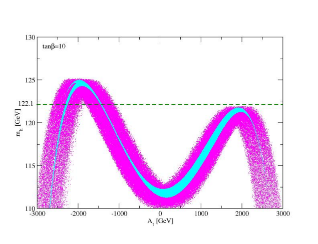


NHSSM: for $\tan\beta = 10$ and 40.

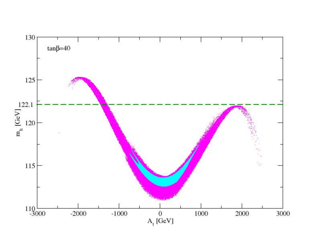
Non-vanishing A'_μ in NHSSM pushes up the smuon mass scale significantly.

Impact of non-holomorphic soft parameters on m_h

A 2 to 3 GeV change in m_h can be possible via A'_t in NHSSM. The effect is larger for a smaller $\tan\beta$. Results with $\mathrm{Br}(B \to X_s + \gamma)$.
A 3 GeV uncertainty in SUSY Higgs boson mass m_h is assumed. Cyan:MSSM, Magenta:NHSSM



- **Correct m_h possible for significantly smaller $|A_t|$.**
- Scanning: $0 \leqslant \mu \leqslant 1$ TeV, $-2 \leqslant \mu' \leqslant 2$ TeV, $-3 \leqslant A'_t \leqslant 3$ TeV.



- Since the effect of A'_t is suppressed by $\tan\beta$ [$X_t = A_t - (\mu + A'_t)\cot\beta$], m_h is affected rather marginally.
- Unlike MSSM, large $|A_t|$ regions in NHSSM are valid via $\mathrm{Br}(B \to X_s + \gamma)$ causing m_h to be acceptable for large A_t. $\mathrm{Br}(B_s \to \mu^+\mu^-)$ limits are not important once $\mathrm{Br}(B \to X_s + \gamma)$ constraint is imposed.

Probing NHSSM via sbottom decay at the LHC: Outline

Ref: UC, AseshKrishna Datta, Samadrita Mukherjee, Abhaya Kumar Swain, JHEP 1810 (2018) 202, arXiv: 1809.05438

- $\tilde{b}_1$ pair production and decay leading to $2b + \not{E}_T$ via $\tilde{b}_1 \to b + \widetilde{\chi}_1^0$. Relevant other decay modes: $\tilde{b}_1 \to t + \widetilde{\chi}_1^-$ and $\tilde{b}_1 \to \tilde{t}_1 + W^-$. Kinematic elimination used for $\tilde{b}_1 \to \tilde{t}_1 + W^-$.
- Higgsino dominated $\widetilde{\chi}_1^0$ ($\mu \leq 350$ GeV) is considered for naturalness preference. $\mu' = 0$ is chosen in the main body of the analysis for simplicity.
- We keep the left and the right mass parameters $m_{\tilde{b}_L}$ and $m_{\tilde{b}_R}$ to be the same in the main body of the work. $\Rightarrow$ For no mixing, $\tilde{b}_1$ and $\tilde{b}_2$ are very close to L and R like respectively with essentially equal masses. With $A_b = 0$, mixing occurs via $(\mu + A'_b)$ that itself is associated with a $\tan\beta$ enhancement.
- Bottom mass radiative corrections may change y_b significantly in NHSSM. This, not only may affect the L-R mixing but may also change couplings concerned with the above electroweakinos and quarks in the $\tilde{b}_i$ decay modes. This is used to distinguish the models.
- Parton-level yields for $(\sigma_{\tilde{b}_1\tilde{b}_1} \times \mathrm{BR}[\tilde{b}_1 \to b\widetilde{\chi}_1^0]^2)$ in the final state $2b + \not{E}_T$ arising from pair-produced $\tilde{b}_1$ at the 13 TeV run are compared for NHSSM and MSSM for varying A'_b. Parameter space explored for large yield zones.

Outline contd.

- Analysis is extended to involve $\tilde{b}_2$. Comparison made with MSSM with proper ratio of yields involving $\tilde{b}_1$ and $\tilde{b}_2$ pair productions and decay into $2b + \not{E}_T$.
- Analysis is extended to varying $m_{\tilde{b}_L}$ and $m_{\tilde{b}_R}$.
- Implications on stop searches are probed in relation to appearance of large y_b via radiative effects that may affect $\tilde{t}_1 \to b\tilde{\chi}_1^+$.

We try to understand how the relevant couplings behave while A_b' changes. Consequently, our study investigates how the branching ratios $\mathrm{Br}(\tilde{b}_1 \to b + \widetilde{\chi}_1^0)$ and $\mathrm{Br}(\tilde{b}_1 \to t + \widetilde{\chi}_1^{\pm})$ are affected and finally how the cross sections leading to yields vary for changing A_b'. We try to isolate the interesting zones of parameter space in relation to LHC for any distinct feature with respect to MSSM .

MSSM and NHSSM corrections to y_b

$$\Delta m_b^{(\tilde{g})}{}_{\mathrm{MSSM}} = \frac{2\alpha_3}{3\pi} m_{\tilde{g}} \mu y_b \frac{v_u}{\sqrt{2}} I(m^2_{\tilde{b}_1}, m^2_{\tilde{b}_2}, m^2_{\tilde{g}}),$$

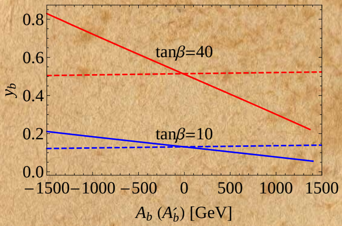
$$\Delta m_b^{\tilde{h}^+}{}_{\mathrm{MSSM}} = \frac{y_t y_b}{16\pi^2} \mu A_t y_t \frac{v_u}{\sqrt{2}} I(m^2_{\tilde{t}_1}, m^2_{\tilde{t}_2}, \mu^2),$$

$$\Delta m_b^{(\tilde{g})}{}_{\mathrm{NHSSM}} = \frac{2\alpha_3}{3\pi} m_{\tilde{g}} A'_b y_b \frac{v_u}{\sqrt{2}} I(m^2_{\tilde{b}_1}, m^2_{\tilde{b}_2}, m^2_{\tilde{g}}),$$

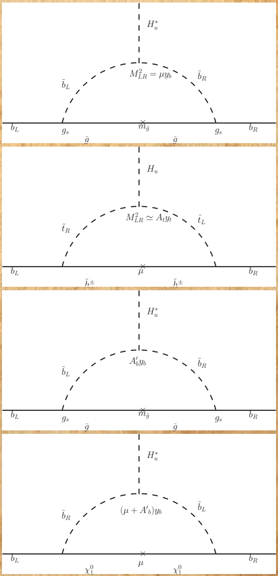
$$\Delta m_b^{\tilde{h}^0}{}_{\mathrm{NHSSM}} = \frac{y_b^2}{16\pi^2} \mu(\mu + A'_b) y_b \frac{v_u}{\sqrt{2}} I(m^2_{\tilde{b}_1}, m^2_{\tilde{b}_2}, \mu^2).$$

where, $I(a,b,c) = -\frac{ab\ \ln(a/b) + bc\ \ln(b/c) + ca\ \ln(c/a)}{(a-b)(b-c)(c-a)}$.

- Δm_b(or Δy_b) is proportional to $\tan\beta$. **Large y_b for negative A'_b and large $\tan\beta$.**



MSSM (dashed line), NHSSM (Solid line);
$\mu, M_1, M_2 = 200, 500, 1100$ GeV, $\mu' = 0$.



Sbottom-electroweakino couplings

$\tilde{b}_i \to b\widetilde{\chi}_1^0$ and $\tilde{b}_i \to t\widetilde{\chi}_1^-$ couplings:
$C_L P_L + C_R P_R$
The decay rates $\propto C_L^2 + C_R^2$
For $\tilde{b}_i$-b-$\tilde{\chi}_1^0$:

$$
\begin{aligned}
C_L &= -\frac{i}{6}(-3\sqrt{2}g_2 N_{12}^* Z_{i3}^d + 6N_{13} y_b Z_{i6}^d \\
&+ \sqrt{2}g_1 N_{11} Z_{i3}^d), \\
C_R &= -\frac{i}{3}(3y_b Z_{i3}^d N_{13} + \sqrt{2}g_1 Z_{i6}^d N_{11}).
\end{aligned}
$$

For $\tilde{b}_i$-t-$\tilde{\chi}_1^-$:

$$
\begin{aligned}
C_L &= i(y_t Z_{i3}^d V_{12}), \\
C_R &= i(-g_2 U_{11}^* Z_{i3}^d + U_{12}^* y_b Z_{i6}^d).
\end{aligned}
$$

N_{ij} are neutralino diagonalizing matrix elements. N_{13}, N_{14} will be large for higgsino dominated LSP. Z_{ij}'s are for squark diagonalizing matrix elements where large Z_{13} and Z_{16} would mean large L and R-components in $\tilde{b}_1$.

- We consider higgsino like $\widetilde{\chi}_1^0$ and $\widetilde{\chi}_1^{\pm}$.
- With above higgsino domination, for $\tilde{b}_i \to b\widetilde{\chi}_1^0$ both C_L and C_R are approximately proportional to y_b.
- For $\tilde{b}_1 \to t\widetilde{\chi}_1^-$, couplings for L-type $\tilde{b}_1$ is $\propto y_t$ and the same for R-like $\tilde{b}_1$ is $\propto y_b$.
- A left like $\tilde{b}_1$ will largely decay via $t\widetilde{\chi}_1^-$. Thus it will have a smaller BR for $b\tilde{\chi}_{1,2}^0$. All the above that are generic in MSSM are also true for NHSSM.
- Significantly different behavior in NHSSM results from non-vanishing A_b' that may cause a large y_b for large $\tan\beta$ via radiative effects.
- We ignored $\tilde{b}_1 \to \tilde{t}_1 W^-$ kinematically by the choice of $m_{\tilde{b}_1} < m_{\tilde{t}_1} + m_W$.

Sbottom-electroweakino couplings

For $\tilde{b}_i$-b-$\tilde{\chi}_1^0$ coupling:

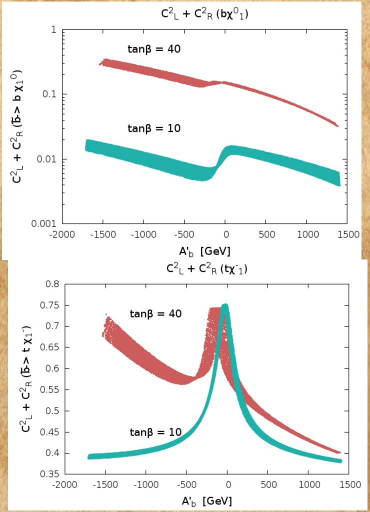
$$
\begin{aligned}
C_L &= -\frac{i}{6}(-3\sqrt{2}g_2 N_{12}^* Z_{i3}^d + 6N_{13}y_b Z_{i6}^d \\
&+ \sqrt{2}g_1 N_{11} Z_{i3}^d), \\
C_R &= -\frac{i}{3}(3y_b Z_{i3}^d N_{13} + \sqrt{2}g_1 Z_{i6}^d N_{11}).
\end{aligned}
$$

For $\tilde{b}_i$-t-$\tilde{\chi}_1^-$ coupling:

$$
\begin{aligned}
C_L &= i(y_t Z_{i3}^d V_{12}), \\
C_R &= i(-g_2 U_{11}^* Z_{i3}^d + U_{12}^* y_b Z_{i6}^d).
\end{aligned}
$$

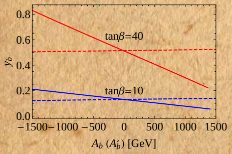
The decay rates $\propto C_L^2 + C_R^2$.

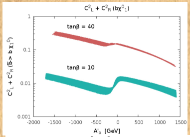
- Spread appears due to $100 < \mu < 350$ GeV. Region around $\mu + A_b' \simeq 0$ refers to scenarios of $\tilde{b}_1$, $\tilde{b}_2$ to be Left and Right like with negligible mixing. Large non-vanishing A_b' zones refer to larger mixing.
- For $\tilde{b}_1 \to b\tilde{\chi}_1^0$, a change of sign of Z_{13}^d occurs near the no mixing zone. y_b enhancement/suppression occurs for negative/positive A_b' especially for large $\tan\beta$.
- In the small A_b' zone $\tilde{b}_1$ is Left like. Thus $\tilde{b}_1 \to t\tilde{\chi}_1^-$, is peaked due to y_t irrespective of $\tan\beta$. For large negative A_b' and large $\tan\beta$ y_b enhancement effect is seen in the left zone. For the small $\tan\beta$ case, y_t dominates over y_b.

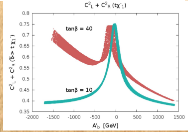


A_b' dependence for $\tilde{b}_1$ decays

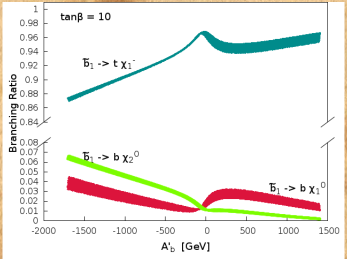
Branching ratios

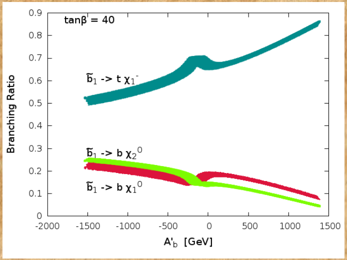




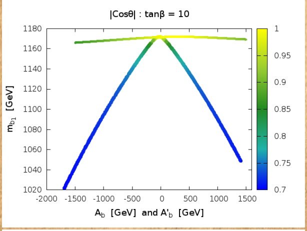


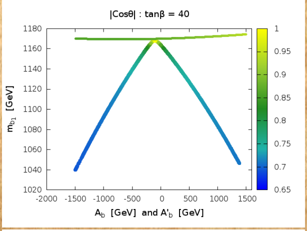
- Br($\tilde{b}_1 \to b\widetilde{\chi}^0_1$) essentially follows the profile of the couplings.
- When $\tilde{b}_1$ is left dominated (central region) the $\tilde{b}_1 \to t\widetilde{\chi}^-_1$ decay rate is driven by y_t, hence becomes large.
- Difference of rates gets smaller for increase in y_b i.e. large negative A'_b and large $\tan\beta$ cases where competition sets in between the modes.

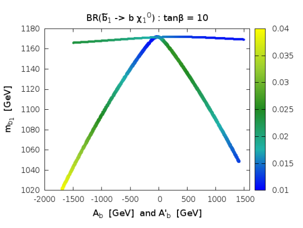


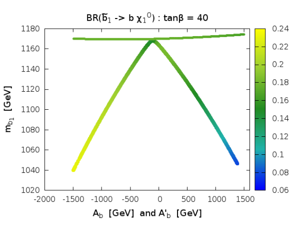


$\mu, M_1, M_2 = 200, 500, 1100$ GeV
$m_{\tilde{Q}_3} = m_{\tilde{t}_L} = m_{\tilde{b}_L} = m_{\tilde{b}_R}(m_{\tilde{D}_3})$=1.2 TeV and $m_{\tilde{t}_R}(m_{\tilde{U}_3}) = 1.5$ TeV
Large mixing in NHSSM cases compared to MSSM (top flatter lines).



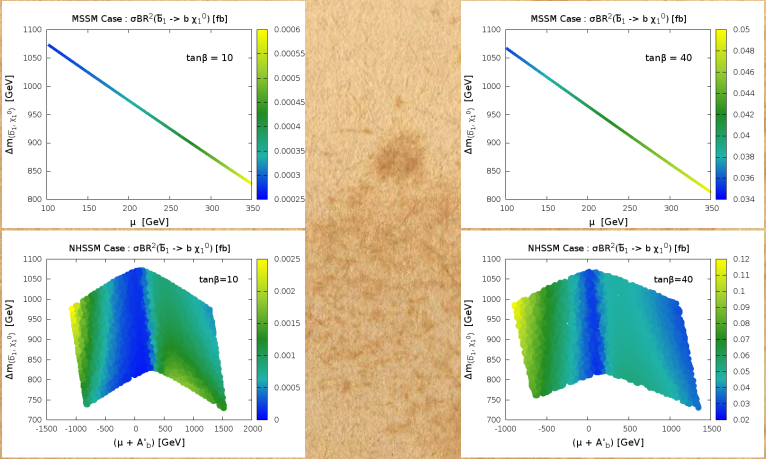






Signal Strength: Parton-level yields

$pp \to \tilde{b}_1 \tilde{b}_1^* $ at LHC 13 TeV, $\tilde{b}_1 \to b\widetilde{\chi}_1^0$ leading to $2b + \not{E}_T$.



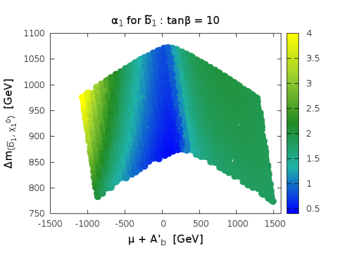
With $A_b = 0$, $(\mu + A_b')\tan\beta$ controls the L-R mixing in NHSSM. Blue central regions refer to larger $\mathrm{Br}(\tilde{b}_1 \to t\widetilde{\chi}_1^-)$ since $\tilde{b}_1 \sim \tilde{b}_L \Rightarrow$ suppressed $\mathrm{Br}(\tilde{b}_1 \to b\widetilde{\chi}_1^0)$. NHSSM can give a much higher yield for a large $\tan\beta$.

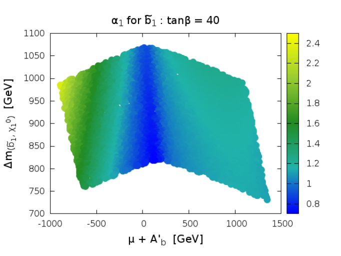
Ratio of yields for $\tilde{b}_1$

$100 < \mu < 350$ GeV and $|A'_b| < 1.2$ TeV.

$$\alpha_1(A'_b) = \frac{[(\sigma_{\tilde{b}_1\tilde{b}_1} \times \mathrm{BR}[\tilde{b}_1 \to b\tilde{\chi}^0_1]^2)]^{\mathrm{NHSSM}}}{[(\sigma_{\tilde{b}_1\tilde{b}_1} \times \mathrm{BR}[\tilde{b}_1 \to b\tilde{\chi}^0_1]^2)]^{\mathrm{MSSM}}}$$

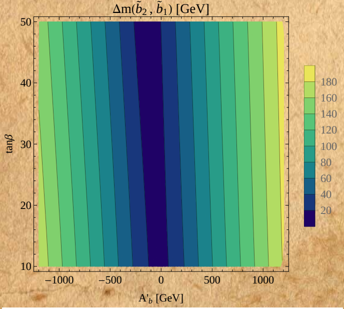
Ratio α_1 refers to a μ value for MSSM that would be equivalent to $\mu + A'_b$ used in NHSSM. There is about an 8-fold increase from the lowest to the highest value for $\tan\beta = 10$ and around a 6-fold increase for $\tan\beta = 40$. Largest regions of α_1 fall in the negative large A'_b zone due to y_b-enhancement.

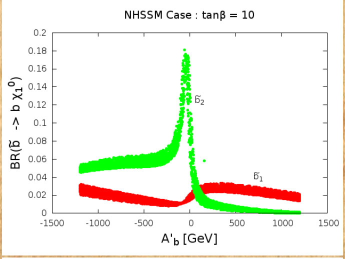


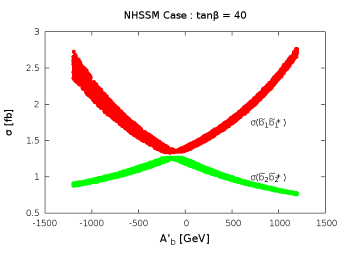


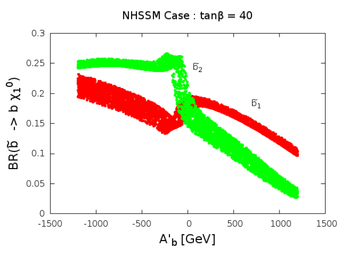
Including $\tilde{b}_2$ in the picture

Because of the parameter choice, the mass difference of $\tilde{b}_1$ and $\tilde{b}_2$ is hardly very large.





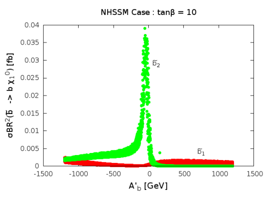


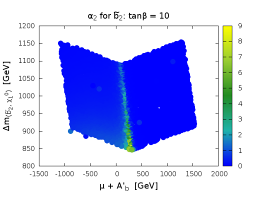


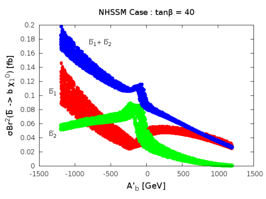
$\tilde{b}_2 \simeq \tilde{b}_R$ near $A'_b = 0$ that suppresses $\mathrm{Br}(\tilde{b}_2 \to t\widetilde{\chi}_1^-)$ in favour of $\mathrm{Br}(\tilde{b}_2 \to b\widetilde{\chi}_1^0)$.

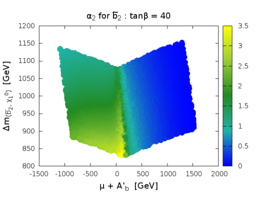
Ratio of yields for $\tilde{b}_2$

$$\alpha_2(A'_b) = \frac{\left[(\sigma_{\tilde{b}_2\tilde{b}_2} \times \mathrm{BR}[\tilde{b}_2 \to b\tilde{\chi}_1^0]^2)\right]^{\mathrm{NHSSM}}}{\left[(\sigma_{\tilde{b}_2\tilde{b}_2} \times \mathrm{BR}[\tilde{b}_2 \to b\tilde{\chi}_1^0]^2)\right]^{\mathrm{MSSM}}}$$





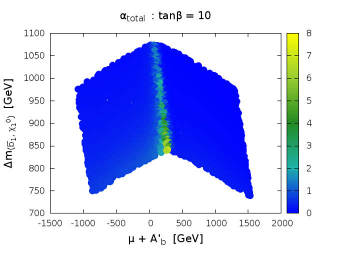


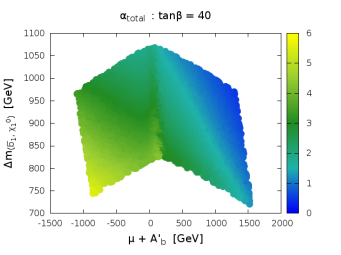


For large and negative A'_b and large $\tan\beta$, the $\tilde{b}_2 \to b\tilde{\chi}_1^0$ decay rate is larger because of R-domination in $\tilde{b}_2$. This is especially true for the $A'_b \simeq 0$ zone. This in turn suppresses $\tilde{b}_2 \to t\tilde{\chi}_1^-$ in favour of $\tilde{b}_2 \to b\tilde{\chi}_1^0$. This indeed compensates the relatively smaller production cross section for $\tilde{b}_2$ leading to similar yields as in the case of $\tilde{b}_1$. The combined yields for $\tilde{b}_1$ plus $\tilde{b}_2$ could reach to few tens of events at 300 fb^{-1} for the favourable parameter zone.

Ratio of yields for $\tilde{b}_1$ plus $\tilde{b}_2$

$$\alpha_{\mathrm{total}}(A'_b) = \frac{\sum_{i=1,2}\left[(\sigma_{\tilde{b}_i\tilde{b}_i} \times \mathrm{BR}[\tilde{b}_i \to b\tilde{\chi}_1^0]^2)\right]^{\mathrm{NHSSM}}}{\sum_{i=1,2}\left[(\sigma_{\tilde{b}_i\tilde{b}_i} \times \mathrm{BR}[\tilde{b}_i \to b\tilde{\chi}_1^0]^2)\right]^{\mathrm{MSSM}}}$$

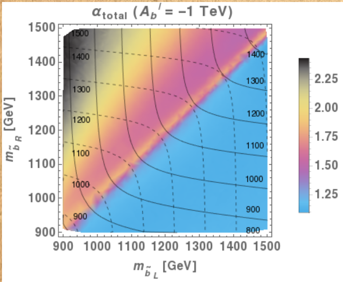




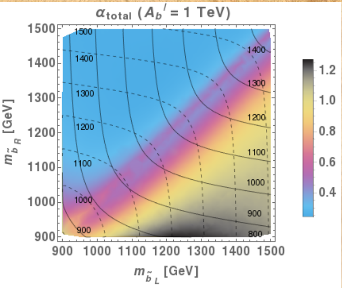
Up to eight-fold (six-fold) increased rates could be possible for $\tan\beta = 10$ (40) over the expected MSSM rates in the final state under consideration.

α_{total} for varying L and R sbottom masses

Variations of α_{total} in the $m_{\tilde{b}_L} - m_{\tilde{b}_R}$ plane for $A_b' = -1$ TeV (left) and $A_b' = 1$ TeV (right) and for fixed values of $\tan\beta$ (=40) and μ (=200 GeV). Contours of constant $m_{\tilde{b}_1}$ ($m_{\tilde{b}_2}$) are overlaid with solid (dashed) lines along the right (left) edges of the plots.

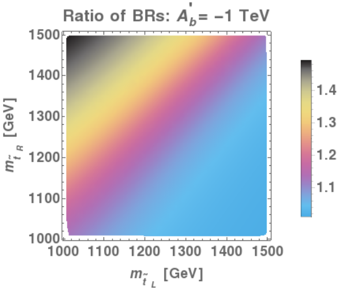


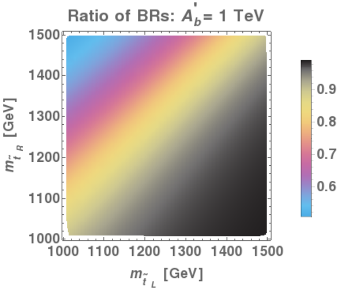
With A_b' large and negative the relative yield is maximum in the upper left corner. Here, the denominator for the MSSM value goes to a minimum. This happens via $\tilde{b}_1$ becoming further $\tilde{b}_L$-like in MSSM than that of NHSSM (where the mixing effect is larger due to A_b'). The effect of larger y_b for NHSSM also comes into play.



α_{total} can rather be marginally above unity for positive A_b'. It maximizes (black region) when $\tilde{b}_1$ is heavily $\tilde{b}_R$-like.

Generally, NHSSM involves a $\tan\beta$ suppression for stop mixing. However, $\tilde{t}_i - b - \widetilde{\chi}_1^+$ vertex would indicate that a Left like stop would couple to a higgsino like chargino and a b-quark with strength y_b. Hence its decay rate would be different from that of MSSM depending on $\tan\beta$ and A'_b.
$\tan\beta = 40$ and $\mu = 200$ GeV.





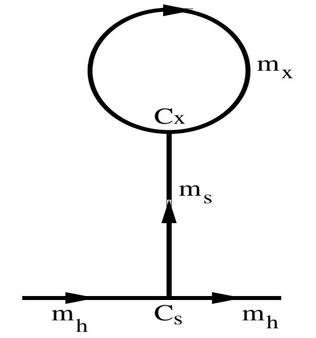
Conclusion

- Non-holomorphic MSSM is a simple extension of MSSM with a few virtues like it is able to isolate the electroweakino sector to some degree from the scalar sector. Hence, it is able to reduce the EW fine-tuning while allowing a higgsino type of $\widetilde{\chi}_1^0$ to be a single component DM candidate.
- It can accommodate muon $g-2$ result rather easily for some region of parameter space.
- It has unique signatures for the scalar sector especially for the down type of quarks and sleptons and it has some degree of influence on the Higgs sector too. It may have interesting signature on flavor physics.
- Distinguishing the signatures of NHSSM from MSSM can be challenging. However, the bottom Yukawa coupling may receive large radiative corrections and thus it may have some interesting consequences.
- A suitably designed multi-channel study may illuminate useful ways to distinguish the scenario from MSSM more effectively.
- Implications may be studied for suitable models by going beyond MSSM.

THANK YOU FOR YOUR ATTENTION

Backup pages

Tadpole correction



S: a singlet field. m_X: a very heavy scalar mass

Tadpole contribution: $\sim C_S C_X \frac{m_X^2}{m_S^2} ln(\frac{m_X^2}{m_h^2})$

If $m_S << m_X$ the tadpole contribution becomes very large.

For discussions: Ref. Hetherington, JHEP 2001

Hard SUSY breaking terms

S. Martin, Phys. Rev D., 2000; Possible non-holomorphic hard SUSY breaking terms:

Type	Term	Naive Suppression	Origin
	ϕ^4	$\frac{F}{M^2} \sim \frac{m_W}{M}$	$\frac{1}{M^2}[X\Phi^4]_F$
	$\phi^3\phi^*$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\Phi^3\Phi^*]_D$
	$\phi^2\phi^{*2}$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\Phi^2\Phi^{*2}]_D$
	$\phi\psi\psi$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\,\Phi D^\alpha\Phi D_\alpha\Phi]_D$
hard	$\phi^*\psi\psi$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\,\Phi^* D^\alpha\Phi D_\alpha\Phi]_D$
	$\phi\psi\lambda$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\,\Phi D^\alpha\Phi W_\alpha]_D$
	$\phi^*\psi\lambda$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\,\Phi^* D^\alpha\Phi W_\alpha]_D$
	$\phi\lambda\lambda$	$\frac{F}{M^2} \sim \frac{m_W}{M}$	$\frac{1}{M^2}[X\Phi W^\alpha W_\alpha]_F$
	$\phi^*\lambda\lambda$	$\frac{\lvert F\rvert^2}{M^4} \sim \frac{m_W^2}{M^2}$	$\frac{1}{M^4}[XX^*\,\Phi^* W^\alpha W_\alpha]_D$

Electroweak Fine-tuning Components

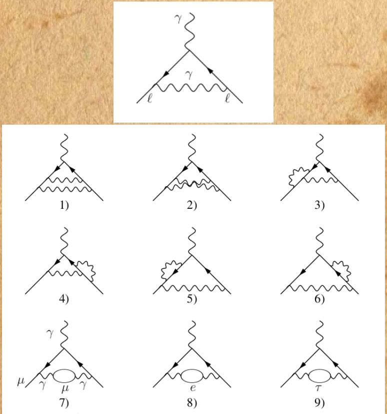
$$
\begin{aligned}
\Delta(\mu) &= \frac{4\mu^2}{m_Z^2}\left(1+\frac{m_A^2+m_Z^2}{m_A^2}\tan^2 2\beta\right), \\
\Delta(b) &= \left(1+\frac{m_A^2}{m_Z^2}\right)\tan^2 2\beta, \\
\Delta(m_{H_u}^2) &= \left|\frac{1}{2}\cos 2\beta+\frac{m_A^2}{m_Z^2}\cos^2\beta-\frac{\mu^2}{m_Z^2}\right| \times \left(1-\frac{1}{\cos 2\beta}+\frac{m_A^2+m_Z^2}{m_A^2}\tan^2 2\beta\right), \\
\Delta(m_{H_u}^2) &= \left|-\frac{1}{2}\cos 2\beta+\frac{m_A^2}{m_Z^2}\sin^2\beta-\frac{\mu^2}{m_Z^2}\right| \times \left|1+\frac{1}{\cos 2\beta}+\frac{m_A^2+m_Z^2}{m_A^2}\tan^2 2\beta\right|,
\end{aligned}
$$

$$
\Delta_{Total} = \sqrt{\textstyle\sum_i \Delta_{p_i}^2}, \tag{1}
$$

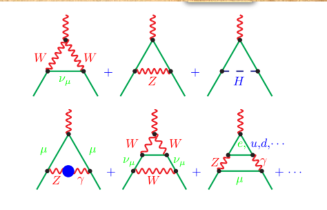
Ref. Perelstein, Spethmann: JHEP 2007, hep-ph/0702038

SM contributions: a_μ^{SM}

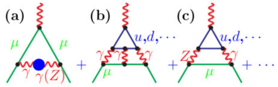
1 and 2-loop QED:



Weak contributions:



hadronic contributions:



(a) Hadronic vacuum polarization $O(\alpha^2), O(\alpha^3)$
(b) Hadronic light-by-light scattering $O(\alpha^3)$
(c) Hadronic effects in 2-loop EWRC $O(\alpha G_F m_\mu^2)$

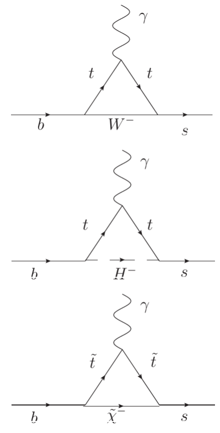
Light quark loops
↓
Hadronic "blobs"

$\mathrm{Br}(B \to X_s + \gamma)$ in MSSM

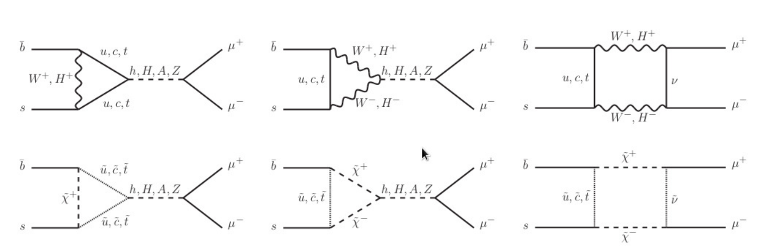
- SM contribution (almost saturates the experimental value) $\to t - W^{\pm}$ loop.
- MSSM contribution:
  1. $\widetilde{\chi}^{\pm} - \tilde{t}$ loop:
  $BR(b \to s\gamma)|_{\widetilde{\chi}^{\pm}} = \mu A_t tan\beta f(m_{\tilde{t}_1}, m_{\tilde{t}_2}, m_{\widetilde{\chi}^{\pm}}) \frac{m_b}{v(1+\Delta m_b)}$
  2. $H^{\pm} - t$ loop:
  $BR(b \to s\gamma)|_{H^{\pm}} = \frac{m_b(y_t cos\beta - \delta y_t sin\beta)}{vcos\beta(1+\Delta m_b)} g(m_{H^{\pm}}, m_t)$
  where,

$$
\begin{aligned}
\delta y_t &= y_t \frac{2\alpha_s}{3\pi} \mu M_{\tilde{g}} tan\beta[\cos^2\theta_t I(m_{\tilde{s}_L}, m_{\tilde{t}_2}, M_{\tilde{g}}) \\
&+ sin^2\theta_t I(m_{\tilde{s}_L}, m_{\tilde{t}_1}, M_{\tilde{g}})]
\end{aligned}
$$

- Destructive interference for $A_t\mu < 0 \to$ preferred.
- NLO contributions (from squark-gluino loops: due to the corrections of top and bottom Yukawa couplings) become important at large μ or large $\tan\beta$.



$B_s \to \mu^+\mu^-$ in MSSM



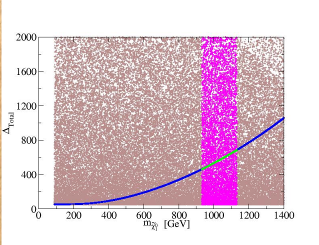
- Dominant SM contribution from : Z penguin top loop & W box diagram.
- SM value : $BR(B_s \to \mu^+\mu^-)$=3.23 $\pm$ 0.27$\times 10^{-9}$.
- LHCb result : $3.2^{+1.4}_{-1.2}(stat.)^{+0.5}_{-0.3}(syst.) \to$ no room for large deviation.
- $BR(B_s \to \mu^+\mu^-)_{SUSY} \propto \frac{tan^6\beta}{m_A{}^4}$

NH terms affecting or not affecting muon g-2 in two benchmark points where $\tilde{\chi}_1^0$ is bino-like

Table 1. *Benchmark points for NHSSM. Masses are shown in GeV. Only the two NHSSM benchmark points shown satisfy the phenomenological constraint of Higgs mass, dark matter relic density along with direct detection cross section, muon anomaly, $\mathrm{Br}(B \to X_s + \gamma)$ and $\mathrm{Br}(B_s \to \mu^+\mu^-)$. The associated MSSM points are only given for comparison and do not necessarily satisfy all the above constraints.*

Parameters	MSSM	NHSSM	MSSM	NHSSM
$m_{1,2,3}$	472, 1500, 1450	472, 1500, 1450	243, 250, 1450	243, 250, 1450
$m_{\tilde{Q}_3}/m_{\tilde{U}_3}/m_{\tilde{D}_3}$	1000	1000	1000	1000
$m_{\tilde{Q}_2}/m_{\tilde{U}_2}/m_{\tilde{D}_2}$	1000	1000	1000	1000
$m_{\tilde{Q}_1}/m_{\tilde{U}_1}/m_{\tilde{D}_1}$	1000	1000	1000	1000
$m_{\tilde{L}_3}/m_{\tilde{E}_3}$	2236	2236	1000	1000
$m_{\tilde{L}_2}/m_{\tilde{E}_2}$	592	592	500	500
$m_{\tilde{L}_1}/m_{\tilde{E}_1}$	592	592	500	500
A_t, A_b, A_τ	-1500, 0, 0	-1500, 0, 0	-1368.1, 0, 0	-1368.1, 0, 0
A'_t, A'_μ, A'_τ	0, 0, 0	2234, 169, 0	0, 0, 0	3000, 200, 0
$\tan\beta$	**10**	**10**	**40**	**40**
μ	500	500	390.8	390.8
μ'	0	-175	0	1655.5
m_A	1000	1000	1000	1000
$m_{\tilde{g}}$	1438.9	1439.1	1438.9	1438.9
$m_{\tilde{t}_1}, m_{\tilde{t}_2}$	894.4, 1151.2	865.5, 1154.9	907.8, 1137.5	903.4, 1141.4
$m_{\tilde{b}_1}, m_{\tilde{b}_2}$	1032.4, 1046.2	1026.3, 1045.1	1013.8, 1051.2	1017.7, 1056.5
$m_{\tilde{\mu}_L}, m_{\tilde{\nu}_\mu}$	596.4, 596.3	573.5, 595.9	502.0, 497.1	465.8, 496.3
$m_{\tilde{\tau}_1}, m_{\tilde{\nu}_\tau}$	2237.1, 2238.5	2237.1, 2238.5	985.4, 997.2	988.5, 998.8
$m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_2^\pm}$	504.2, 1483.6	677.6, 1484.7	244.6, 421.0	262.3, 1255.2
$m_{\tilde{\chi}_1^0}, m_{\tilde{\chi}_2^0}$	448.6, 509.0	464.0, 680.6	231.3, 249.9	240.9, 262.1
$m_{\tilde{\chi}_3^0}, m_{\tilde{\chi}_4^0}$	522.6, 1483.5	683.2, 1484.7	400.7, 421.0	1253.3, 1253.7
$m_{H^\pm}$	1011.9	1005.8	955.7	1011.6
m_H, m_h	1008.1, 121.4	984.8, 122.8	948.0, 122.4	990.2, 122.8
$\mathrm{Br}(B \to X_s + \gamma)$	3.00×10^{-4}	3.01×10^{-4}	2.01×10^{-4}	4.05×10^{-4}
$\mathrm{Br}(B_s \to \mu^+\mu^-)$	3.40×10^{-9}	3.45×10^{-9}	5.06×10^{-9}	1.65×10^{-9}
a_μ	1.94×10^{-10}	22.3×10^{-10}	34.8×10^{-10}	35.8×10^{-10}
$\Omega_{\tilde{\chi}_1^0} h^2$	0.035	0.095	0.0114	0.122
$\sigma^{\mathrm{SI}}_{\tilde{\chi}_1^0 p}$ in pb	4.01×10^{-9}	3.47×10^{-10}	6.79×10^{-9}	3.15×10^{-12}

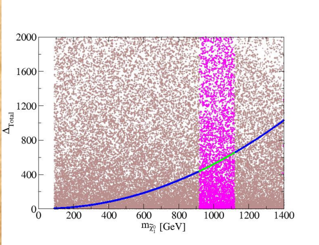
Electroweak fine-tuning and higgsino dark matter



Δ_{Total} vs $m_{\tilde{\chi}_1^0}$ for $\tan\beta = 10$

MSSM (i.e. with $\mu' = A'_t = 0$): Thin blue line and partly green line in the middle. Δ_{Total} is little above 400.

NHSSM: brown and magenta. Consistent region satisfying a 3σ level of WMAP/PLANCK constraints are shown. EWFT in NHSSM ranges from too high to too low (~ 50).



Δ_{Total} vs $m_{\tilde{\chi}_1^0}$ for $\tan\beta = 40$

EWFT in NHSSM can be vanishingly small.

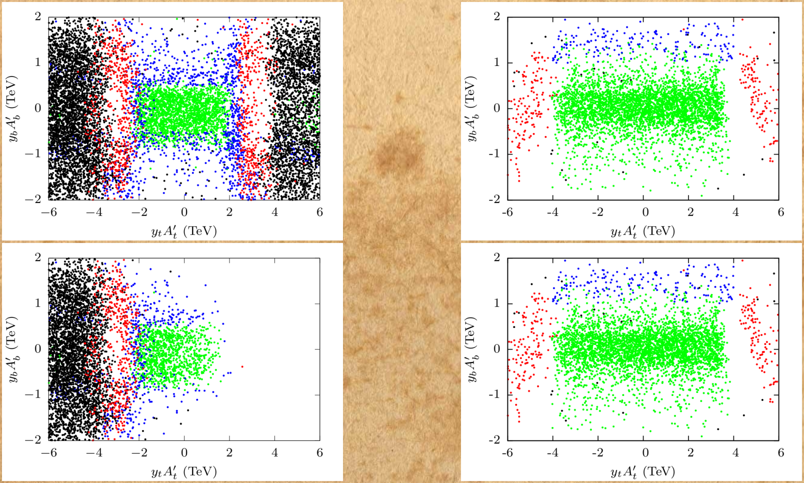
$-3 \text{ TeV} < \mu, \mu' < 3 \text{ TeV}$

$-3 \text{ TeV} < \mathrm{A_t}, \mathrm{A'_t} < 3 \text{ TeV}$

EW fine-tuning differs from FT estimate in UV complete scenario like CMSSM with NH terms. There, an FT expression would depend on NH parameters. The FT related low scale parameters p_i are no longer independent. NH+CMSSM still has FT estimate dominantly controlled by μ^2 (Ross *et. al.* 2016, 2017).

NHSSM: Limiting trilinears with Charge and Color Breaking Constraints

Jyotiranjan Beuria and Abhishek Dey, JHEP 2017; "Exploring Charge and Color Breaking vacuum in Non-Holomorphic MSSM"



tan $\beta = 10$; Without and with Higgs mass constraint.

tan $\beta = 40$; Without and with Higgs mass constraint.

Color codes: Green: Stable vacuum, Blue: Long lived, Red: Thermally excluded, Black: Unstable

NHSSM: Limiting trilinears with Charge and Color Breaking Constraints

Jyotiranjan Beuria and Abhishek Dey, JHEP 2017; "Exploring Charge and Color Breaking vacuum in Non-Holomorphic MSSM"

$$\begin{aligned} \mathrm{V}|_{\mathrm{tree}} \quad = \quad & m_2^2 H_u^2 + m_1^2 H_d^2 + m_{\tilde{t}_L}^2 \tilde{t}_L^2 + m_{\tilde{t}_R}^2 \tilde{t}_R^2 - 2B_\mu H_d H_u + 2y_t A_t H_u \tilde{t}_R \tilde{t}_L \\ & -2y_t(\mu + A_t')\tilde{t}_L \tilde{t}_R H_d + y_t^2(H_u^2 \tilde{t}_L^2 + H_u^2 \tilde{t}_R^2 + \tilde{t}_R^2 \tilde{t}_L^2) + \frac{g_1^2}{8}(H_u^2 - H_d^2 + \frac{\tilde{t}_L^2}{3} - \frac{4\tilde{t}_R^2}{3})^2 \\ & + \frac{g_2^2}{8}(H_u^2 - H_d^2 - \tilde{t}_L^2)^2 + \frac{g_3^2}{6}(\tilde{t}_L^2 - \tilde{t}_R^2)^2, \qquad (2) \end{aligned}$$

Only stops receiving vevs apart from up and down Higgses:

$$\{|A_t| + |\mu| + |A_t'|\}^2 < 3\left(m_1^2 + m_2^2 + m_{\tilde{t}_L}^2 + m_{\tilde{t}_R}^2 - 2B_\mu\right). \qquad (3)$$

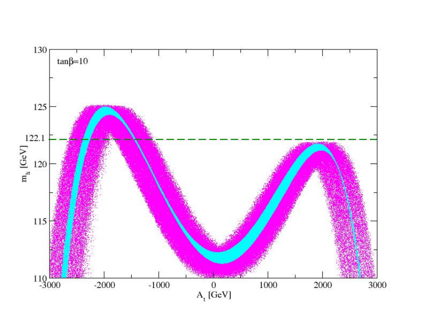
Only sbottoms receiving vevs apart from up and down Higgses:

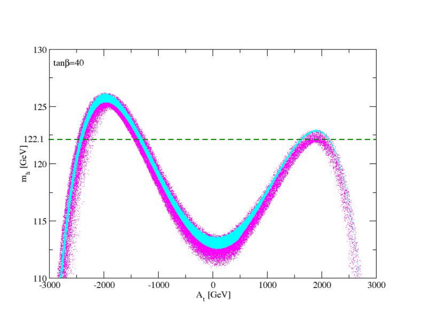
$$\{|A_b| + |\mu| + |A_b'|\}^2 < 3\left\{1 - \frac{g_1^2 + g_2^2}{24 y_b^2}\right\}\left(m_1^2 + m_2^2 + m_{\tilde{b}_L}^2 + m_{\tilde{b}_R}^2 - 2B_\mu\right). \qquad (4)$$

Analytically derived constraints have limited scopes. Apart from the scenario of many scalars receiving vevs, one needs to consider long lived vacuum, thermal stability of vacuum etc. $\Rightarrow$ Code: Vevacious.

Impact of non-holomorphic soft parameters on m_h

A 2 to 3 GeV change in m_h can be possible via A_t'. The effect is larger for a smaller $\tan\beta$.
Cyan:MSSM, Magenta:NHSSM





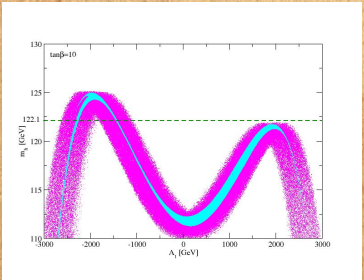
m_h is enhanced/decreased by 2-3 GeV due to non-holomorphic terms.
- **Correct m_h possible for significantly smaller $|A_t|$.**

- Since A_t' is associated with a suppression by $\tan\beta$ [off-diag term in stop sector: $X_t = A_t - (\mu + A_t')\cot\beta$], m_h is affected only marginally.

- $0 \leqslant \mu \leqslant 1\ \mathrm{TeV}$, $-2 \leqslant \mu' \leqslant 2\ \mathrm{TeV}$, $-3 \leqslant A_t' \leqslant 3\ \mathrm{TeV}$.
- A 3 GeV uncertainty in computation of m_h in SUSY is assumed.

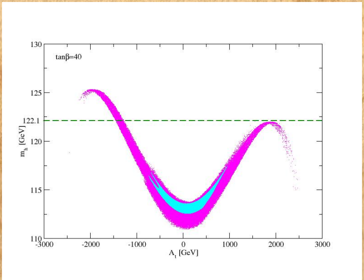
Imposing $\mathrm{Br}(B \to X_s + \gamma)$ and $\mathrm{Br}(B_s \to \mu^+\mu^-)$ constraints

$2.77 \times 10^{-4} \leqslant \mathrm{Br}(B \to X_s + \gamma) \leqslant 4.09 \times 10^{-4}$, $0.8 \times 10^{-9} \leqslant \mathrm{Br}(\mathbf{B_s} \to \mu^+\mu^-) \leqslant 5 \times 10^{-9}$ [both at 3σ]



m_h vs A_t for $\tan\beta = 10$ with the above constraints.

$\Rightarrow$ Essentially unaltered results for a low $\tan\beta$ like 10.



m_h vs A_t for $\tan\beta = 40$.

$\Rightarrow$ $\mathrm{Br}(B \to X_s + \gamma)$ that increases with $\tan\beta$ takes away large $|A_t|$ zones of MSSM (cyan). Large $|A_t|$ with $\mu A_t < 0$ is discarded via the lower bound and vice versa. Thus m_h does not reach the desired limit beyond $|A_t| \sim 1$ TeV in MSSM.

NHSSM: The effect of A_t' is via L-R mixing: $[A_t \to A_t - (\mu + A_t')\cot\beta]$. Thus large $|A_t|$ regions are valid via $\mathrm{Br}(B \to X_s + \gamma)$ and m_h may stay above the desired limit. $\mathrm{Br}(B_s \to \mu^+\mu^-)$ limits are not important once $\mathrm{Br}(B \to X_s + \gamma)$ constraint is imposed.

Electroweak fine-tuning in MSSM

EWSB conditions out of minimization of $V_{\rm Higgs}$:

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 - m_{H_u}^2 \tan^2\beta}{\tan^2\beta - 1} - |\mu|^2, \qquad \sin 2\beta = \frac{2b}{m_{H_d}^2 + m_{H_u}^2 + 2|\mu|^2} \tag{5}$$

Electroweak Fine-tuning:

$$\Delta_{p_i} = \left| \frac{\partial \ln m_Z^2(p_i)}{\partial \ln p_i} \right|, \qquad \Delta_{Total} = \sqrt{\sum_i \Delta_{p_i}^2}, \text{ where } p_i \equiv \{\mu^2, b, m_{H_u}, m_{H_d}\}$$

- For $\tan\beta$ and μ both not too small the most important term is $\Delta(\mu) \simeq \frac{4\mu^2}{m_Z^2}$. For a moderately large $\tan\beta$, a small μ means a small Δ_{Total}.
- NH soft terms do not contribute to $V_{\rm Higgs}$ at the tree level. Possibility of small μ with a larger higgsino LSP mass $\sim |\mu + \mu'|$ satisfying the DM data (as a single component one). This is unlike MSSM.

Bilinear Higgsino soft term

- One may try to absorb μ' in the superpotential sector that may give rise to its F-terms of the potential involving Higgs scalars. It appears that the following reparametrization of μ, μ' and Higgs scalar mass parameters may evade the need of a bilinear higgsino soft term. $\mu \to \mu + \delta$, $\mu' \to \mu' + \delta$, and $m^2_{H_{U,D}} \to m^2_{H_{U,D}} - 2\mu\delta - \delta^2$
- A reparametrization would however involve *ad-hoc* correlations between unrelated parameters [Jack and Jones 1999, Hetherington 2001 etc.].
- Such correlations are arbitrary, at least in view of fine-tuning. In particular, there may be a scenario where definite SUSY breaking mechanisms generate bilinear higgsino soft terms whereas it may keep the scalar sector unaffected. [Ross *et. al.* 2016, 2017, Antoniadis *et. al.* 2008, Perez *et. al.* 2008 etc].
- The μ' term that is traditionally retained, isolates a fine-tuning measure (typically $\sim \mathrm{factor} \times \mu^2/M_Z^2$) from the higgsino mass $(\mu + \mu')$: $\Rightarrow$ Possibility of a large higgsino mass (like a TeV satisfying DM relic limits) while having a small fine-tuning.

In a general standpoint we acknowledge the importance of trilinear and bilinear NH soft terms, irrespective of a suppression predicted by a *given* model. Unlike other analyses, we will use a pMSSM type of work on Non-holomorphic supersymmetric SM (NHSSM).