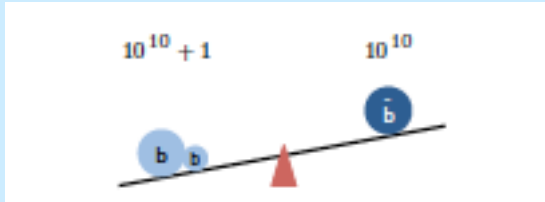


Spontaneous CPT Violation and Quantum Anomalies in a model for Matter-Antimatter Asymmetry in the Cosmos

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Baryogenesis (BAU)



Small difference between number of baryons and anti baryons

$$\beta \equiv \frac{n_B}{s} = (8.59 \pm 0.11) \times 10^{-11}$$

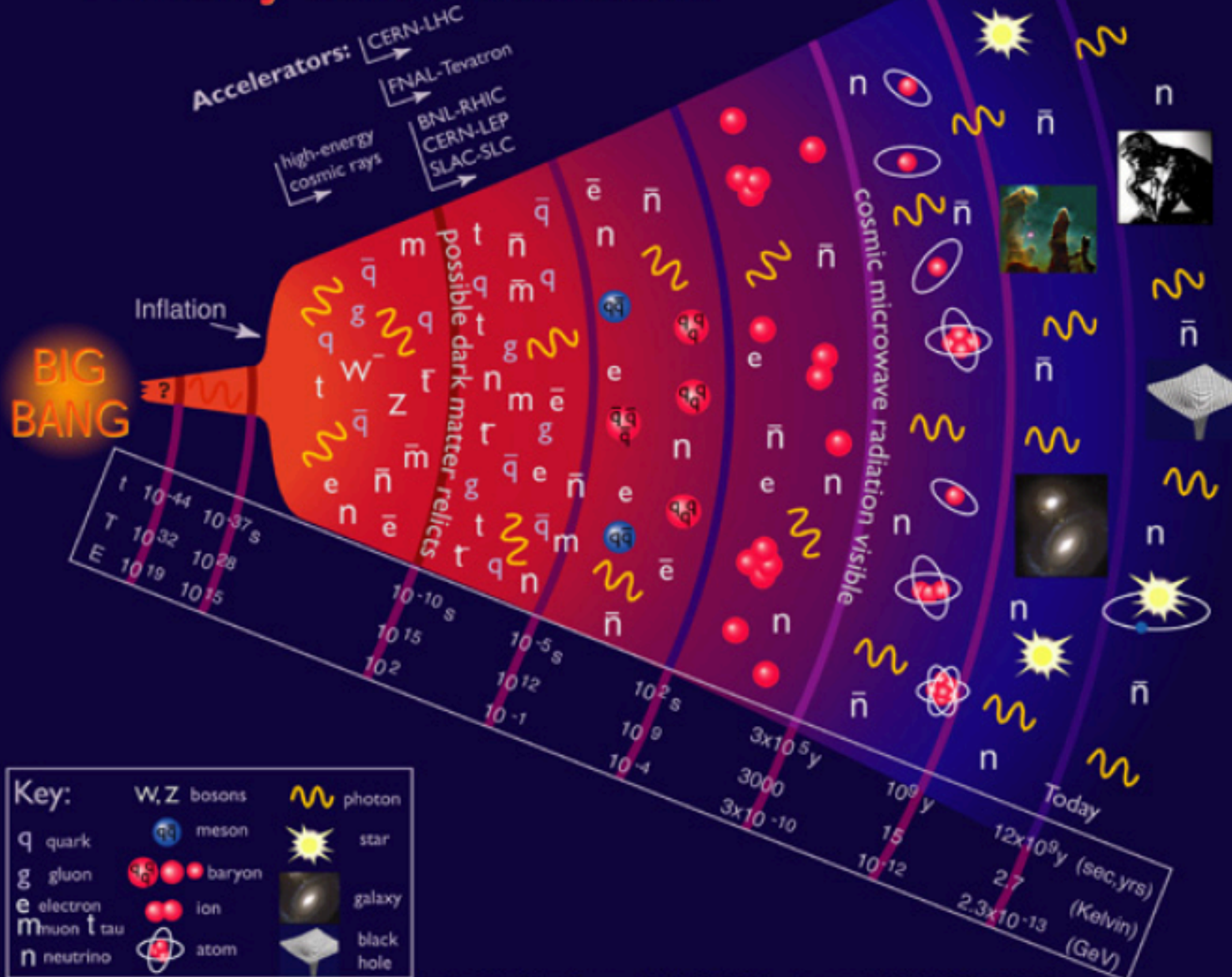
$$n_B = n_b - n_{\bar{b}}$$

n_b is the baryon number density
 $n_{\bar{b}}$ is the anti-baryon number density
 s is the entropy density $\approx 7.04n_\gamma$

For $T \geq GeV$ $n_b \sim n_{\bar{b}} \sim n_\gamma$

- **BAU cannot be due to an initial condition**
- **Galaxy and anti-galaxy: large gamma ray emissions not observed**
- **BAU needs BSM physics: many proposals; many satisfy Sakharov criteria**

History of the Universe



Particle Data Group, LBNL, © 2000. Supported by DOE and NSF

Sakharov criteria (Sakharov JETP Lett 5 (1967) 24)

- **Baryon number violating process (BVP)**
- **C and CP violations**
- **Out of equilibrium (assumes CPT invariance)**

- **BVP not present at perturbative level in SM**

- **Second requirement**

Since $\Gamma(X \rightarrow Y+b) \neq \Gamma(\bar{X} \rightarrow \bar{Y}+\bar{b})$ **and**

$$X \rightarrow Y+b \quad \Delta B = 1$$

$$\bar{X} \rightarrow \bar{Y}+\bar{b} \quad \Delta B = -1$$

- **The third condition avoids**

$$\Gamma(X \rightarrow Y+b) = \Gamma(Y+b \rightarrow X)$$

which otherwise would lead to $\Delta B = 0$

- **Expanding universe leads to non-equilibrium**

Standard model: electroweak BAU

- **Baryon number violating process: sphaleron process**
- **C and P violations: chiral gauge theory and CP phase**
- **Out of equilibrium: *first order* electroweak phase transition**

(Kuzmin, Rubakov, Shaposhnikov, Cohen, Kaplan, Nelson ...)

- **Attractive scenario but **not viable** since phase transition is second order for Higgs mass $M_H > 70\text{GeV}$ in the SM**
- **An alternative mechanism** is baryogenesis through leptogenesis using a background induced ***CPTV*** and ***anomalies***
- (see de Cesare, Mavromatos, Sarkar: *Eur Phys J C* 75 (2015) 514,
Bossingham, Mavromatos, Sarkar: *Eur. Phys. J C* 78 (2018) 113
Bossingham, Mavromatos, Sarkar: *Eur. Phys. J C* 79 (2019) 50)

CPT symmetry

- **C** is charge conjugation
- **P** is parity (space reflection)
- **T** is time-reversal
- **CPT theorem:** For any *Lorentz invariant Lagrangian* $L(x)$ in flat space times

$$\Theta L(x) \Theta^{-1} = L^\dagger(-x)$$

where $\Theta \equiv CPT$.

based on unitarity and locality of interactions
(Schwinger, Pauli, Luders , Jost and Bell)

Assumptions of CPT theorem may not hold in the early universe, e.g. Lorentz invariance

- **An approach to CPTV :Standard Model Extension *effective* field theory with ψ_f a generic fermion**

$$L \supset \dots + \bar{\psi}^f \left(i\gamma^\mu \nabla_\mu - m_f \right) \psi^f + a_\mu \bar{\psi}^f \gamma^\mu \psi^f + b_\mu \bar{\psi}^f \gamma^\mu \gamma^5 \psi^f + \dots$$



Lorentz and CPT from spontaneous violation

(See: Collady, Kosteletsky Phys Rev D55 (1997) 676,
Bluhm et al, Phys Rev Lett 84 (2000) 1098,
Kosteletsky, Russell Rev. Mod. Phys. 83 (2011) 11)

- $\Delta m = m - \bar{m}$ for quarks and antiquarks because of CPTV but in equilibrium has to be too large for BAU (Dolgov 2009)

CPT violation in a String Inspired Model

- CPT violation \longrightarrow violation of Lorentz invariance (Greenberg 2002)
- A CPTV possibility:
 1. Bosonic gravitational multiplet of strings consists of graviton $g_{\mu\nu}$, spin 0 scalar field the dilaton Φ , spin 1 antisymmetric gauge field $B_{\mu\nu}$ which has gauge symmetry

$$B_{\mu\nu} \rightarrow B_{\mu\nu} + \partial_{\mu} \theta_{\nu} - \partial_{\nu} \theta_{\mu}$$

in closed string sector

2. Gauge invariant field strength

$$H_{\mu\nu\rho} = \partial_{[\mu} B_{\nu\rho]}$$

3. Bosonic part of (3 + 1) dimensional effective action

$$S_B = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left(R - e^{-4\Phi} H_{\lambda\mu\nu} H^{\lambda\mu\nu} - \Omega \right) + \dots$$

Role of H-field as Torsion

- **Effective gravitational action in string low energy limit**

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{2\kappa^2} R - \frac{1}{6} H_{\mu\nu\rho} H^{\mu\nu\rho} \right) = \int d^4x \sqrt{-g} \left(\frac{1}{2\kappa^2} \bar{R} \right)$$

where

$$\bar{\Gamma}_{\nu\rho}^{\mu} = \Gamma_{\nu\rho}^{\mu} + \frac{\kappa}{\sqrt{3}} H_{\nu\rho}^{\mu} \neq \bar{\Gamma}_{\rho\nu}^{\mu}$$



contorsion

- **Generalised curvature** $\bar{R}(\bar{\Gamma})$

Dual pseudoscalar in 3 + 1 dimensions:

- $b(x)$ Kalb-Ramond axion

$$B^\mu = \partial^\mu b = -\frac{1}{4} e^{-2\Phi} \epsilon_{abc}{}^\mu H^{abc}$$

Split quantum field into a *background* and fluctuation:

$$b(x) = \bar{b}(x) + \tilde{b}(x)$$

Possible backgrounds:

$$\bar{b}(x) \propto t$$

such backgrounds exist in bosonic non-critical strings

(Antoniadis, Bachas, Ellis and Nanopoulos, Nucl. Phys B 328 (1989) 117)

Fermion coupling to H-torsion using vielbeins

$$g_{\mu\nu} = e^a{}_\mu \eta_{ab} e^b{}_\nu$$

$$\omega_{bca} = e_{b\lambda} \left(\partial_a e_c^\lambda + \Gamma_{\gamma\mu}^\lambda e_c^\gamma e_a^\mu \right)$$

Spin connection

$$\bar{D}_a = \left(\partial_a - \frac{i}{4} \bar{\omega}_{bca} \sigma^{bc} \right) \quad \text{where} \quad \bar{\omega}_{ab\mu} = \omega_{ab\mu} + K_{ab\mu}$$

and

$$K_{abc} = \frac{1}{2} \left(H_{cab} - H_{abc} - H_{bca} \right)$$

$$S_\psi = \frac{i}{2} \int d^4x \sqrt{-g} \left(\bar{\psi} \gamma^\mu \bar{D}_\mu \psi - \bar{D}_\mu \bar{\psi} \gamma^\mu \psi \right)$$

$$S_\psi \supset \int d^4x B_a \bar{\psi} \gamma^a \gamma^5 \psi$$

where

$$B^d \sim \epsilon^{abcd} H_{bca}$$

Effective action for axions

- **Let** $S_{eff} = S + S_{\psi}$
- **The axion terms in** S_{eff} **are**

$$\sim \int d^4x \sqrt{-g} \left[\frac{4}{3\kappa^2} \partial_{\mu} b \partial^{\mu} b - \partial_{\nu} b J^{5\nu} \right]$$

The classical homogeneous isotropic background can be

$$\partial_t \left[\sqrt{-g} \left(\frac{8}{3\kappa^2} \partial_0 \bar{b} - J_0^5 \right) \right] = 0 \quad J_{\mu}^5 = \sum_i \bar{\psi}_i \gamma_{\mu} \gamma^5 \psi_i$$

An approximate solution is plausibly obtained on using from thermal average

$$\langle J_0^5 \rangle = 0$$

$$\frac{d}{dT} \left(a^3(T) B_0 \right) \simeq 0 \Rightarrow B_0(T) = AT^3$$

Dispersion relations for (Dirac) fermions and anti-fermions in presence of background

$$E = \sqrt{\vec{p}^2 + m^2 + B_0^2 + 2|\vec{p}|B_0}$$

$$\bar{E} = \sqrt{\vec{p}^2 + m^2 + B_0^2 - 2|\vec{p}|B_0}$$

$$n - \bar{n} = \frac{g}{(2\pi)^3} \int d^3p \left(\frac{1}{1 + \exp\left(\frac{E}{T}\right)} - \frac{1}{1 + \exp\left(\frac{\bar{E}}{T}\right)} \right) \neq 0$$

Is direct baryogenesis possible?

BUT

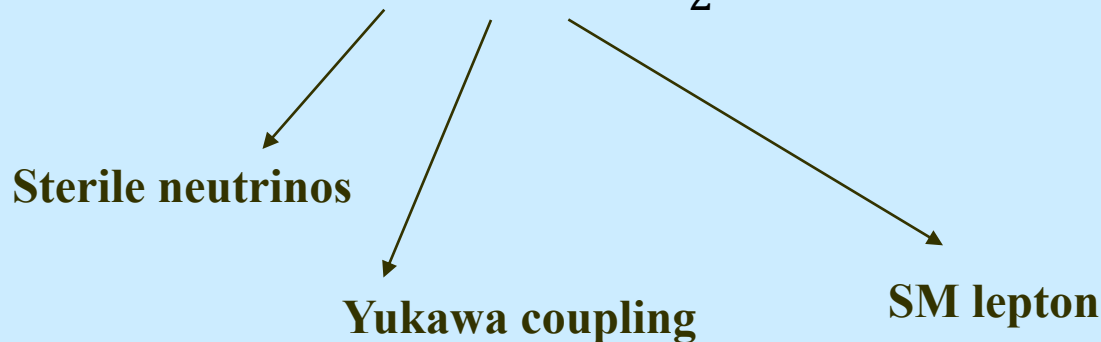
Detailed mechanisms for freeze out?

Phenomenologically relevant baryogenesis need non-minimal higher derivative LV and CPTV fermionic interactions

Leptogenesis model \supset Standard Model has baryogenesis

Existing such model (seesaw mechanism):
Produces light neutrino masses relevant for neutrino oscillations
SM extension with heavy N_I extra right-handed neutrinos ($N_f = 2, 3$)

$$L = L_{SM} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - f_{\alpha I} \bar{L}_\alpha N_I \tilde{\phi} - \frac{M_I}{2} \bar{N}_I^c N_I + h.c.$$



(Minkowski, Yanagida, Mohapatra, Senjanovic, Sechter, Valle, ...)

$$\tilde{\varphi}_i = \varepsilon_{ij} \varphi_j$$

$$B_0 = \text{const} \neq 0$$

$$L = L_{SM} + i\bar{N}\not{\partial}N - \frac{m_N}{2}(\bar{N}^c N + \bar{N}N^c) - \bar{N}B\gamma^5 N - \sum_k y_k \bar{L}_k \tilde{\varphi} N + h.c.$$

$$B_i = 0, \quad i = 1, 2, 3$$

CPTV leptogenesis model motivated by seesaw model

- Kalb-Ramond background model for leptogenesis
- **N is a Majorana spinor; is the adjoint of the Higgs field**
- L_k is the lepton doublet of the k th generation of the standard model
- Restrict to generation $k = 1$
- homogeneous and isotropic
- The model breaks Lorentz invariance and CPT
- Gives microscopic justification of standard model extension
- Decay of heavy right-handed Majorana neutrino

-411 De Cesare, NEM & Sarkar, Eur.Phys.J. C75 (2015) no.10, 514
Bossingham, NEM & Sarkar, Eur.Phys.J. C78 (2018) no.2, 113 & arXiv:1810.13384 [hep-ph].

Leptogenesis implies baryogenesis in model \supset SM

- SM Lagrangian has global $U(1)$ chiral symmetry
- Classically: B and L conservation for *individual* generations
- Quantum mechanically: $B+L$ is anomalous, $B - L$ is exact symmetry
- So non-conservation of $L \rightarrow$ non-conservation of B
- C is explicitly broken in SM
- CP is broken by background Kalb-Ramond
- The expansion of the Universe is an out of equilibrium situation: freeze out
- Prediction of BAU from $B - L$ conservation

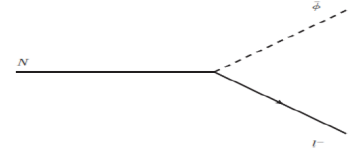
$$N \rightarrow l^- \phi^+, \nu \phi^0$$

$$N \rightarrow l^+ \phi^-, \bar{\nu} \phi^0$$

Channels

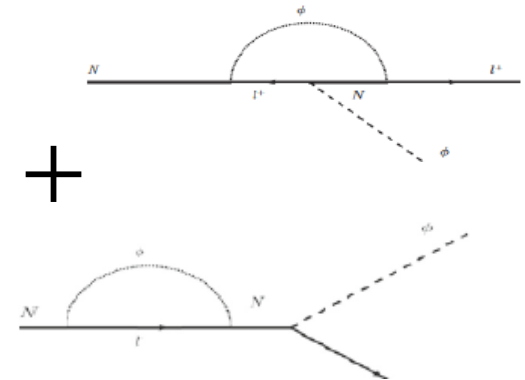
Tree level

Lepton number and CP violation



Contrast with Fukugita-Yanagida leptogenesis

Early universe $T > 10^5 \text{ GeV}$



Boltzmann equations

- Boltzmann equations to find the out equilibrium abundances

$$xHs \frac{dY_1}{dx} - \lambda I = - \sum \left[\gamma^{eq}(1, 2... \rightarrow 3, 4...) \frac{Y_1}{Y_1^{eq}} \frac{Y_2}{Y_2^{eq}} \dots - \gamma^{eq}(3, 4... \rightarrow 1, 2...) \frac{Y_3}{Y_3^{eq}} \frac{Y_4}{Y_4^{eq}} \dots \right]$$

RHS are all the different interactions that take place

Only need to consider the tree level decays of the heavy neutrinos into leptons/anti-leptons and the reverse processes



$$x \equiv \frac{m_N}{T}, \qquad Y_\chi \equiv \frac{n_\chi}{s}$$

Vacuum energy

$$\rho_{B_0}^{DE}(x) = \frac{4}{3\kappa^2} \partial_\mu \bar{b} \partial^\mu \bar{b}$$

$$= \frac{M_P^2}{6\pi} B_0^2 = \frac{M_P^2 \Phi^2}{6\pi} x^{-6}$$

$$\rho_{B_0}^{DE}(x_0) \simeq 10^{-158} M_P^2$$

$$\frac{\Delta L^{total}}{s} = q \frac{\Phi}{m_N} \simeq 8 \times 10^{-11}$$

$$B_i < 10^{-31} \text{ GeV}$$

$$B_0 < .01 \text{ eV}$$

$$T_D = O(100) \text{ TeV}$$

$$B_0(x) \simeq \Phi x^{-3}$$

$$\Phi \simeq (0.36 - 0.74) \text{ keV}$$

Conclusions

B_0 DOES NOT contribute to the CME in the presence of a magnetic field even though it looks like a chiral chemical potential (Hull, Mavromatos, Dvornikov) gravitational anomaly

$$\begin{aligned} \text{Tr}(\bar{R}(\bar{\omega}) \wedge \bar{R}(\bar{\omega})) &= \text{Tr}(R(\omega) \wedge R(\omega)) \\ &+ d \left[\text{Tr} \left(H \wedge R + H \wedge DH + \frac{2}{3} H \wedge H \wedge H \right) \right] \end{aligned}$$

Dark energy contribution less than cosmological constant