

# Bayesian and frequentist approaches to resonance searches

arXiv:1902.03243

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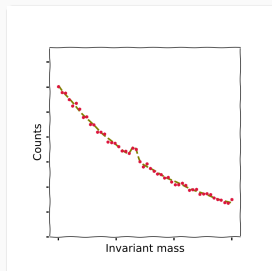
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2. Interpretations
3. Results from toy Higgs search
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## **Background**

# What is that?

A new particle? or just a fluctuation?

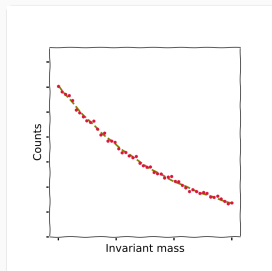
How can we characterise our uncertainty?



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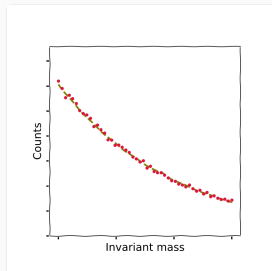
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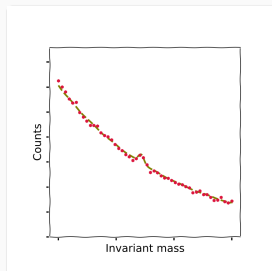
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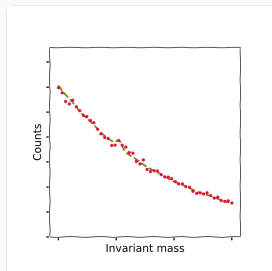
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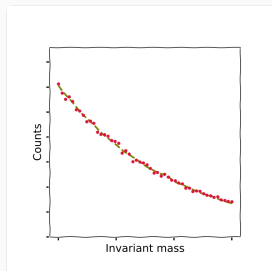




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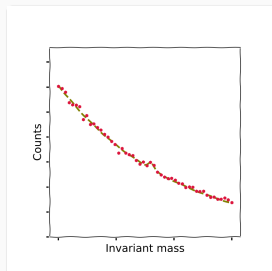
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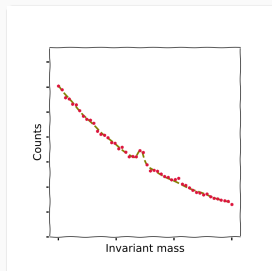
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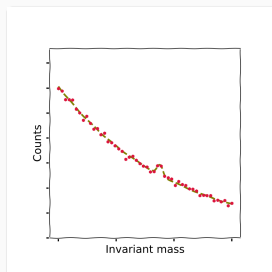
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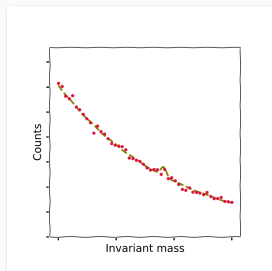
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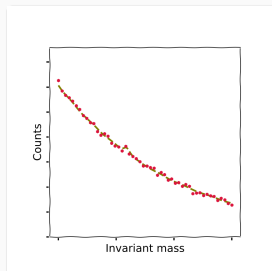
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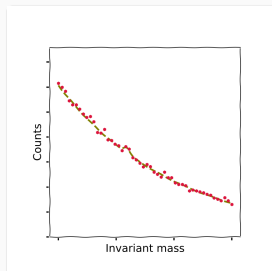
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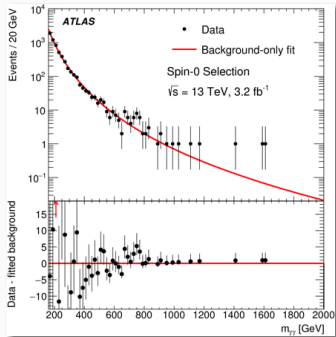
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## Another 750 GeV?



or something real? Should you write a paper about it? Announce a press conference? Start writing your Nobel prize speech?



# Interpretations

## Frequentist: what is probability?

Probabilities **are not** degrees of certainty or belief.

Probabilities **are** frequencies at which events occur in identical repeat experiments.

$$P(A) = \lim_{N \rightarrow \infty} \frac{n_A}{N}$$

## Frequentist: what can we do?

We **cannot** quantify our uncertainty about the resonance.

We **can** attempt to control the frequency at which we would make a type-1 error.

Type-1 error: Reject null hypothesis when it is true.

We must specify a null hypothesis,  $H_0$ , and a desired type-1 error rate,  $\alpha$ .

We reject  $H_0$  at a pre-chosen significance  $\alpha$  or we do not.

The rate  $\alpha$  (implicitly) chosen to be about  $10^{-7}$  ( $5\sigma$ ) in particle physics.

## Frequentist: how do we do it? i

We construct a **test-statistic** that measures discrepancies between data and the null hypothesis, e.g. the log-likelihood ratio,

$$q \equiv -2 \ln \frac{\max_{\theta_1} P(D | M_1, \theta_1)}{\max_{\theta_2} P(D | M_0, \theta_2)}$$

This involves numerical optimisation of the likelihood function.

We calculate the  $p$ -value.

$p$ -value: probability of obtaining a test-statistic at least as extreme as the one we saw, if the null hypothesis was true.

The observed  $p$ -value is not a continuous measure of our confidence in  $H_0$ . The  $p$ -value was a means to controlling the type-1 error rate.

It is common nevertheless to interpret  $p$  as a measure of our confidence in  $H_0$ .

## Frequentist: global or local?

If the data had been different, we would have constructed a resonance model with a different mass to match the different data.

We would have **looked elsewhere**.

**Global  $p$ -values** account for this **look-elsewhere effect**.

We calculated **global  $p$ -values** with Gross-Vitells [1] and Monte-Carlo simulations.

## Bayesian: what is probability?

Probabilities **are** degrees of belief about any proposition.

There is a unique rule for updating them in light of information — **Bayes' theorem**.

$$P(A|B) = \frac{P(B|A)P(A)}{P(B)}$$

Bayesian statistics  $\Leftrightarrow$  probability theory

## Bayesian: what can we do?

We can simply update our belief in the signal + background model relative to the background only model.

The factor that updates our belief is a **Bayes factor**.

$$\begin{aligned}\text{Bayes factor} &= \frac{\text{Relative belief after data}}{\text{Relative belief before data}} \\ &= \frac{P(D | M_1)}{P(D | M_2)}\end{aligned}$$



## Bayesian: how do we do it?

The numerator and denominator are so-called Bayesian evidences. For a model with parameters  $\theta$ ,

$$P(D|M) = \int P(D|M, \theta) p(\theta|M) d\theta$$

To compare with the  $p$ -value, we calculate the posterior of the background model, assuming equal prior odds,

$$P(M_0|D) = \frac{1}{1+B}$$

This is the **plausibility of the background model in light of data.**

## Likelihood function

A component of Bayesian and frequentist analysis. The probability of obtaining data given a particular model and parameters.

Our data is binned. The likelihood is a product of Poissons, one for each bin.

$$P(D|M, \boldsymbol{\theta}) = \prod_i \frac{e^{-\lambda_i} \lambda_i^{o_i}}{o_i!},$$

where the expected number of events depends on the model parameters,  $\lambda = \lambda(\boldsymbol{\theta})$ .

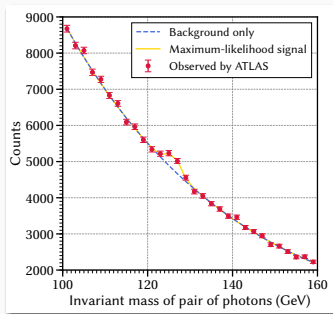
## **Results from toy Higgs search**

From quantum mechanics, we learned an antidote to disputes about interpretations.

Shut up and calculate.

## Toy problem

To make calculations, let's pick a toy problem to study. The search for the Higgs in the diphoton channel by ATLAS with 25/fb [2].

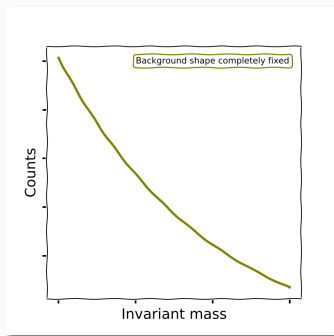


An important search for the discovery of the Higgs.

## Background model

There is a monotonically falling background.

We could describe it by a basis of polynomials (e.g. Bernstein) but so that we can perform many calculations, we just use a **fixed background** and neglect parametric uncertainties in it.



We model the signal predicted by a Higgs as a Gaussian centred at  $m_h$ .

The **width** was the experimental resolution of about 1.5 GeV.

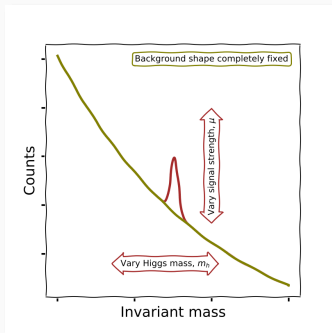
We specified the **strength** relative to the Standard Model prediction (at 125 GeV),

$$\mu = \frac{\text{efficiency} \times \text{cross section}}{(\text{efficiency} \times \text{cross section})_{\text{SM @ 125 GeV}}}$$

This is an approximation as we did not model dependence of efficiency or cross section as functions of Higgs mass.

## Signal model ii

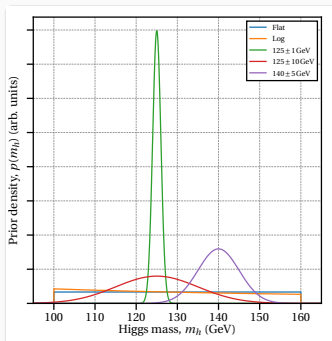
There were thus two unknown parameters describing the location and strength of the resonance,  $m_h$  and  $\mu$ .





# Priors

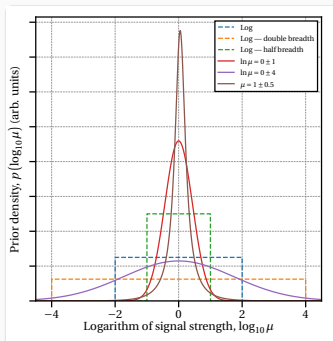
For our Bayesian calculations, we must place priors on  $m_h$  and  $\mu$ . We experiment with several choices.



Broad priors (log and flat) and narrow ones representing specific prior knowledge.

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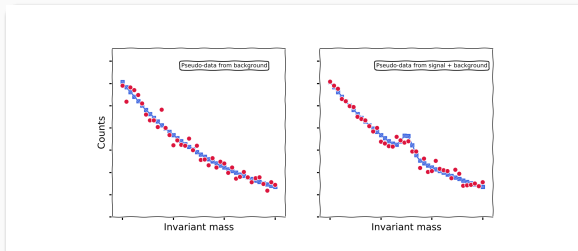
We vary the breadth of the log prior for the signal strength, and the shape of the prior.

We use the real 25/fb collected by ATLAS [2].

We sample our own **pseudo-data** from the background model and the signal + background model with  $\mu = 1$ ,  $m_h = 125$  GeV.

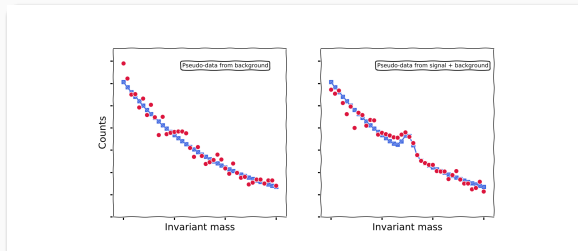
The models tell us the expected number of counts in each bin for a particular integrated luminosity.

We sample pseudo-data at many integrated luminosities by drawing counts from Poisson distributions in each bin.



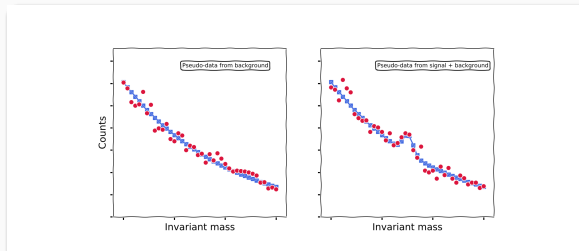
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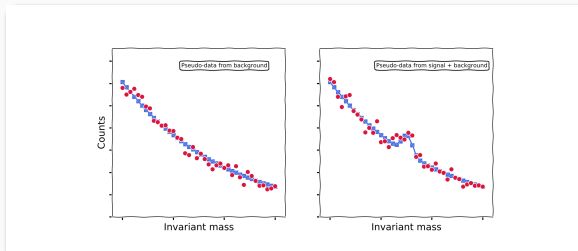
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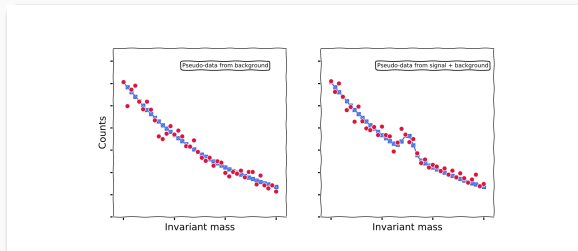
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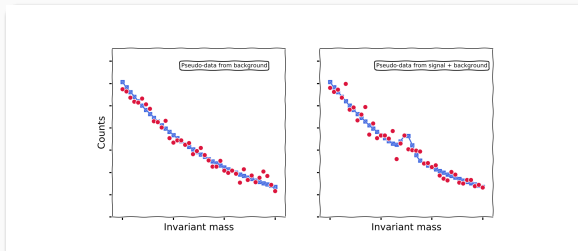
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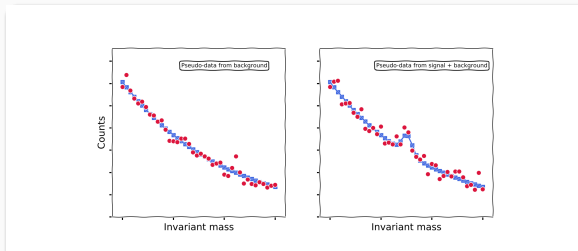
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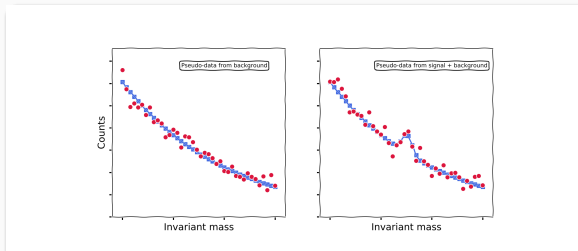
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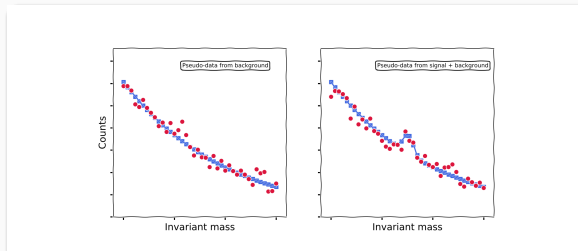
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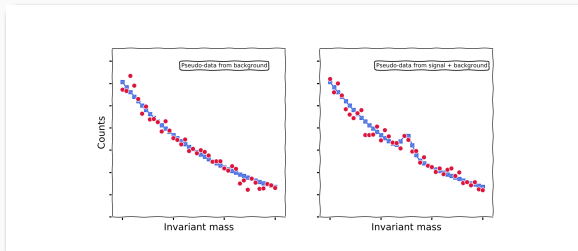
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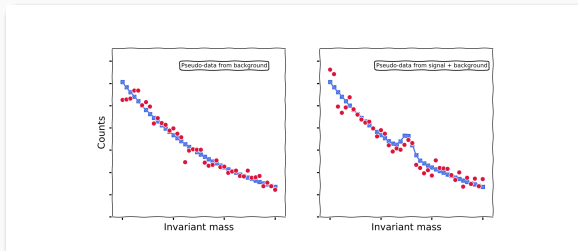
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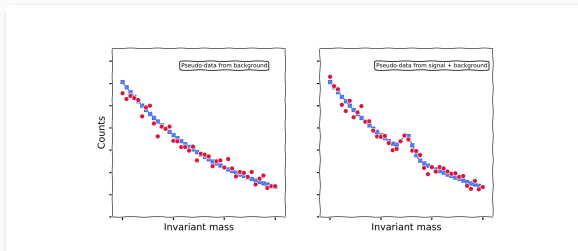
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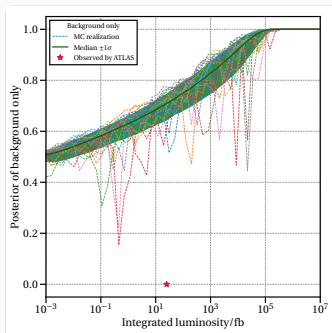


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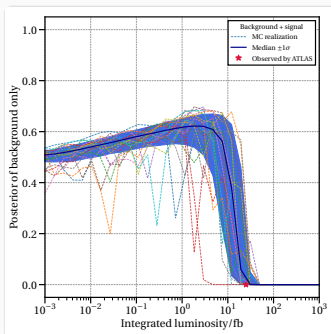
## Evolution of $p$ -value and posterior as we collect data



The posterior slowly approaches 1 when the background model is correct

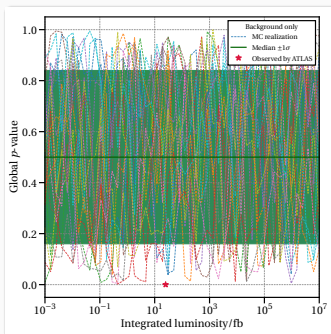


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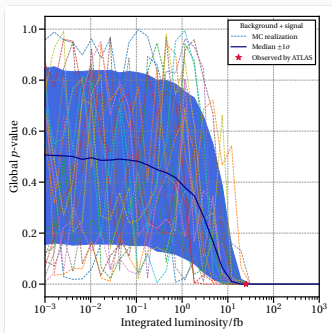
and zero when the signal model is correct, though in this case there is an extremely mild preference for the background model until about 10/fb.

## Evolution of $p$ -value and posterior as we collect data



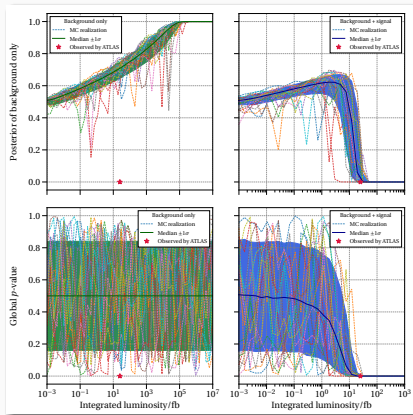
The  $p$ -value makes a random walk between 0 and 1 when the background model is correct

## Evolution of $p$ -value and posterior as we collect data



and when the signal model is correct, it makes a (noisy) walk towards zero.

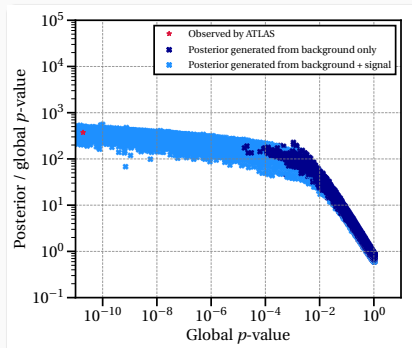
# Evolution of $p$ -value and posterior as we collect data



Bayesian (top)/frequentist (bottom). Background model true (left)/signal model true (right).

## Comparison between $p$ -value and posterior

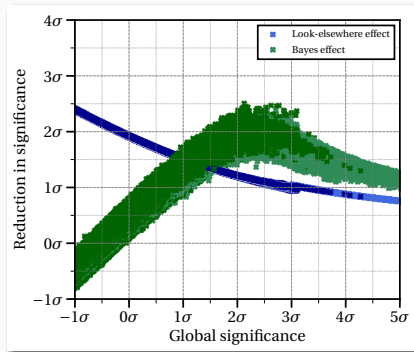
We performed about a million pseudo-experiments.



The posterior of the background model about  $10^2 - 10^3$  times greater than global  $p$ -value!

# The Bayes effect

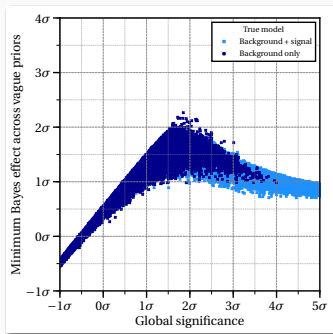
The magnitude of the effect greater than the well-known look-elsewhere effect.



Global significances reduced by  $1 - 2\sigma$ .

## Prior dependence

We checked many priors. The effect could be reduced but remained important.



See paper [3] for full discussion about prior dependence of this effect.

## **Conclusions**



1. First detailed comparison of Bayesian and frequentist methods in resonance searches
2. Posterior ultimately converged to 0 or 1;  $p$ -value makes random walk if  $H_0$  correct
3.  $p$ -values overstate evidence against the null!  $p$ -value  $\lll$  posterior of background model
4. Checked that the effect was robust with respect to several choices of prior
5. When looking at an anomaly, we must remember the look-elsewhere effect and the Bayes effect

- <sup>1</sup> E. Gross and O. Vitells, “Trial factors for the look elsewhere effect in high energy physics,” *Eur. Phys. J.* **C70**, 525–530 (2010), [arXiv:1005.1891](#).
- <sup>2</sup> G. Aad et al., “Measurements of Higgs boson production and couplings in diboson final states with the ATLAS detector at the LHC,” *Phys. Lett.* **B726**, [Erratum: *Phys. Lett.*B734,406(2014)], 88–119 (2013), [arXiv:1307.1427](#).
- <sup>3</sup> A. Fowlie, “Bayesian and frequentist approaches to resonance searches,” (2019), [arXiv:1902.03243](#).