

asymptotic in the spotlight

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2019

International Symposium

U Manches

1 Jul 2019

sium

ic safety spotlight

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PASCOS 2019
XXV Internation

**asymptotic safety
in the spotlight**

asymptotic safety in the spotlight

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asymptotic safety in the spotlight

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standard model

local QFT for fundamental interactions

strong nuclear force

weak force

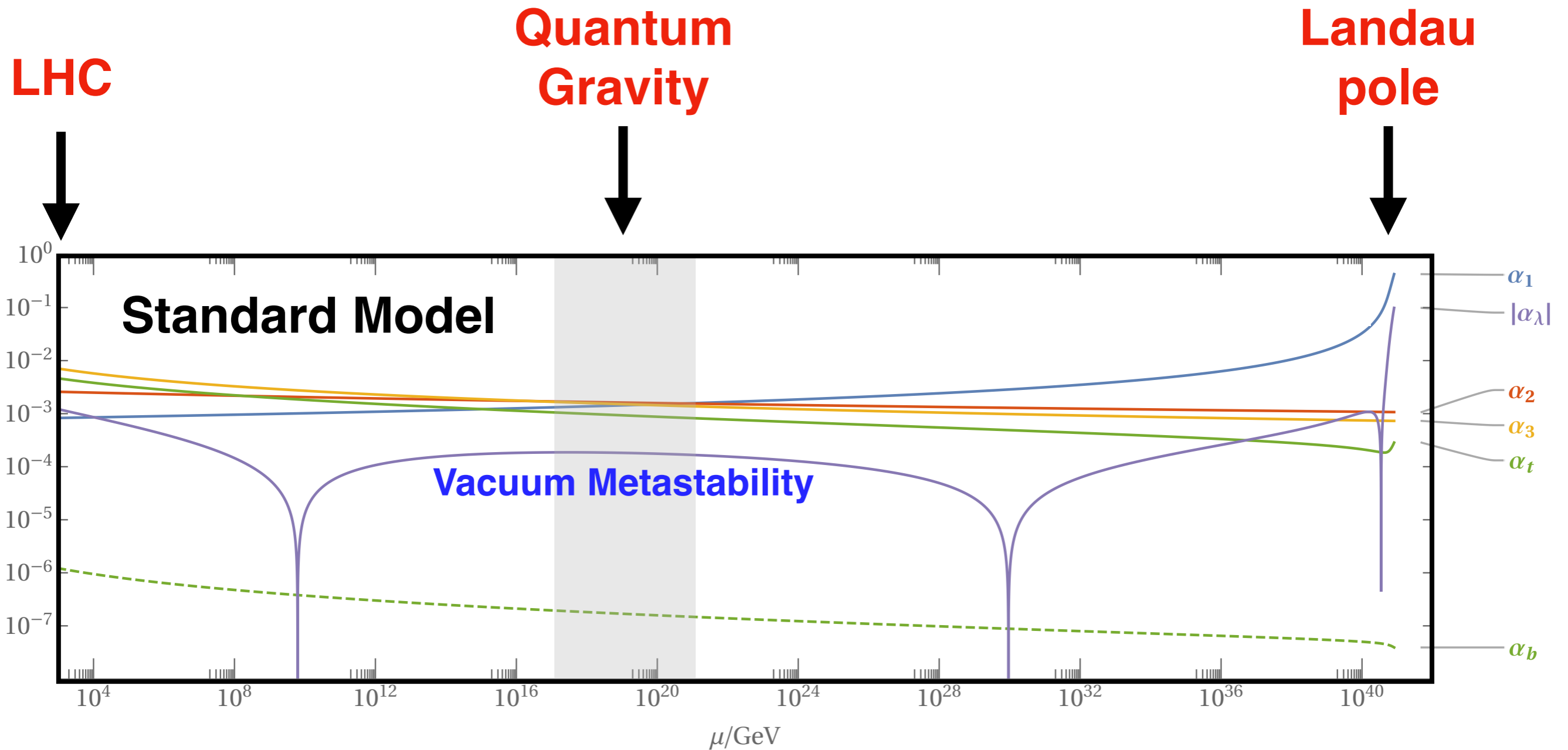
electromagnetic force

open challenges

what comes **beyond the SM**?

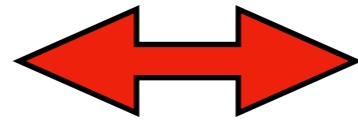
and how does **gravity** fit in?

where are we ?



what is asymptotic safety?

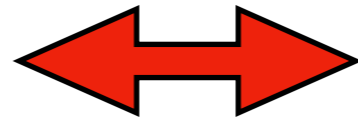
fundamental QFT



UV fixed point

Wilson '71

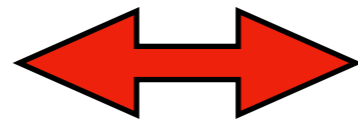
asymptotic freedom



non-interacting
UV fixed point

Gross, Wilzcek '73 , Politzer '73

asymptotic safety



interacting
UV fixed point

Weinberg '79

exact asymptotic safety

infinite-N non-linear sigma

Brezin, Zinn-Justin '76
Bardeen, Lee, Shrock '76

2+eps infinite-NF Gross-Neveu
quantum gravity

Gawedzki, Kupiainen '85
Christensen, Duff '78
Gastmans, Kallosh, Truffin '78
Weinberg '79

3d infinite-N scalars
infinite-NF Gross-Neveu

Pisarski '82
Bardeen, Moshe, Bander '84
Rosenstein, War, Park' 89
de Calan, Faria da Veiga, Magnen, de Seneor '91

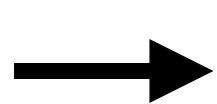
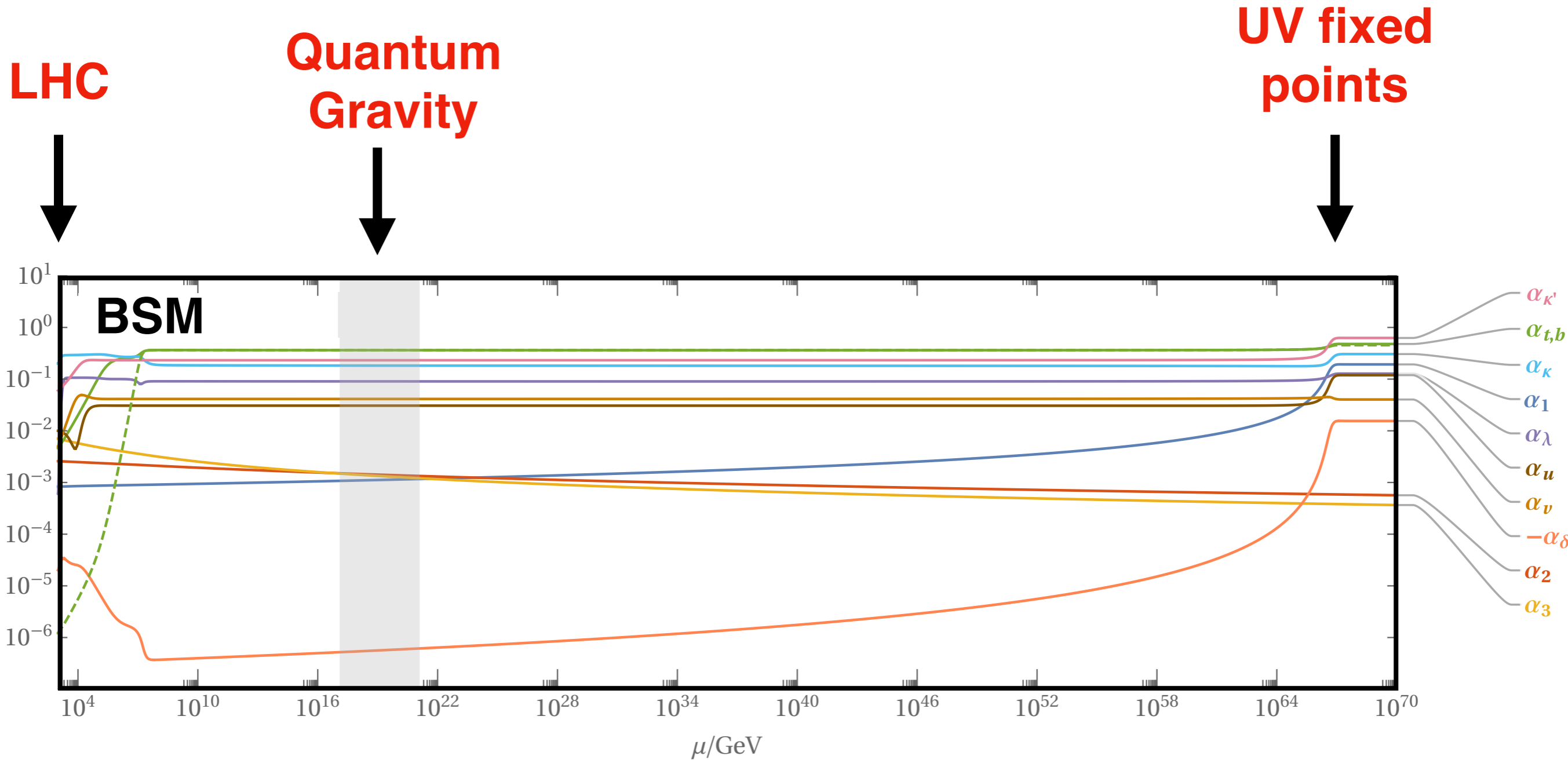
4d gauge + matter

Litim, Sannino '14
Bond, Litim '16, '17, '18

4d quantum gravity

Reuter '96
Litim '03

why asymptotic safety?



AS models beyond the SM and their phenomenology:
Talk by Tom Steudtner (today's parallel session 1530)

today:

understand **asymptotic safety**
in general weakly-coupled 4d QFTs

asymptotic safety
and models beyond the SM

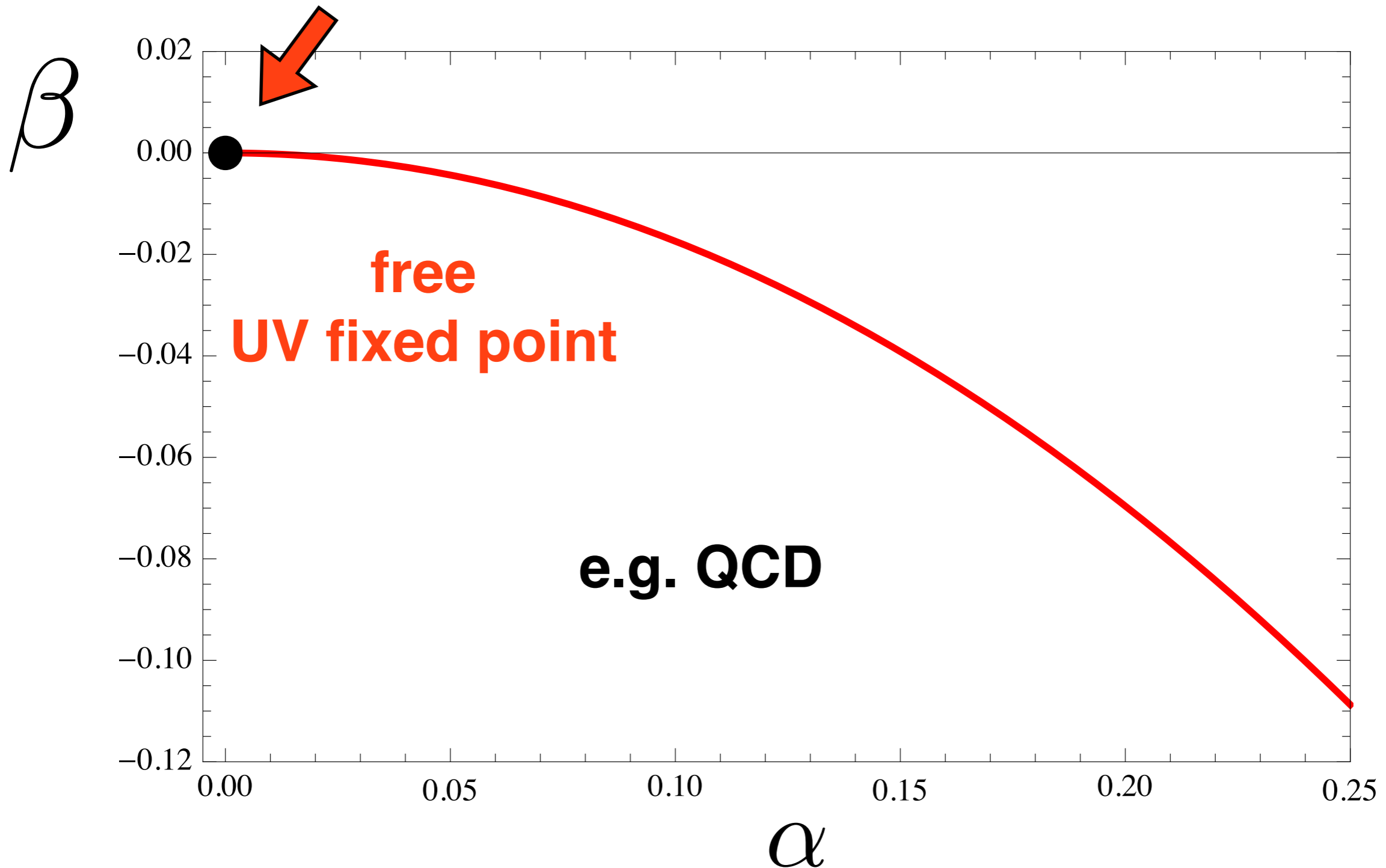
signatures of **asymptotic safety**
in **quantum gravity**

asymptotic safety in gauge-matter theories

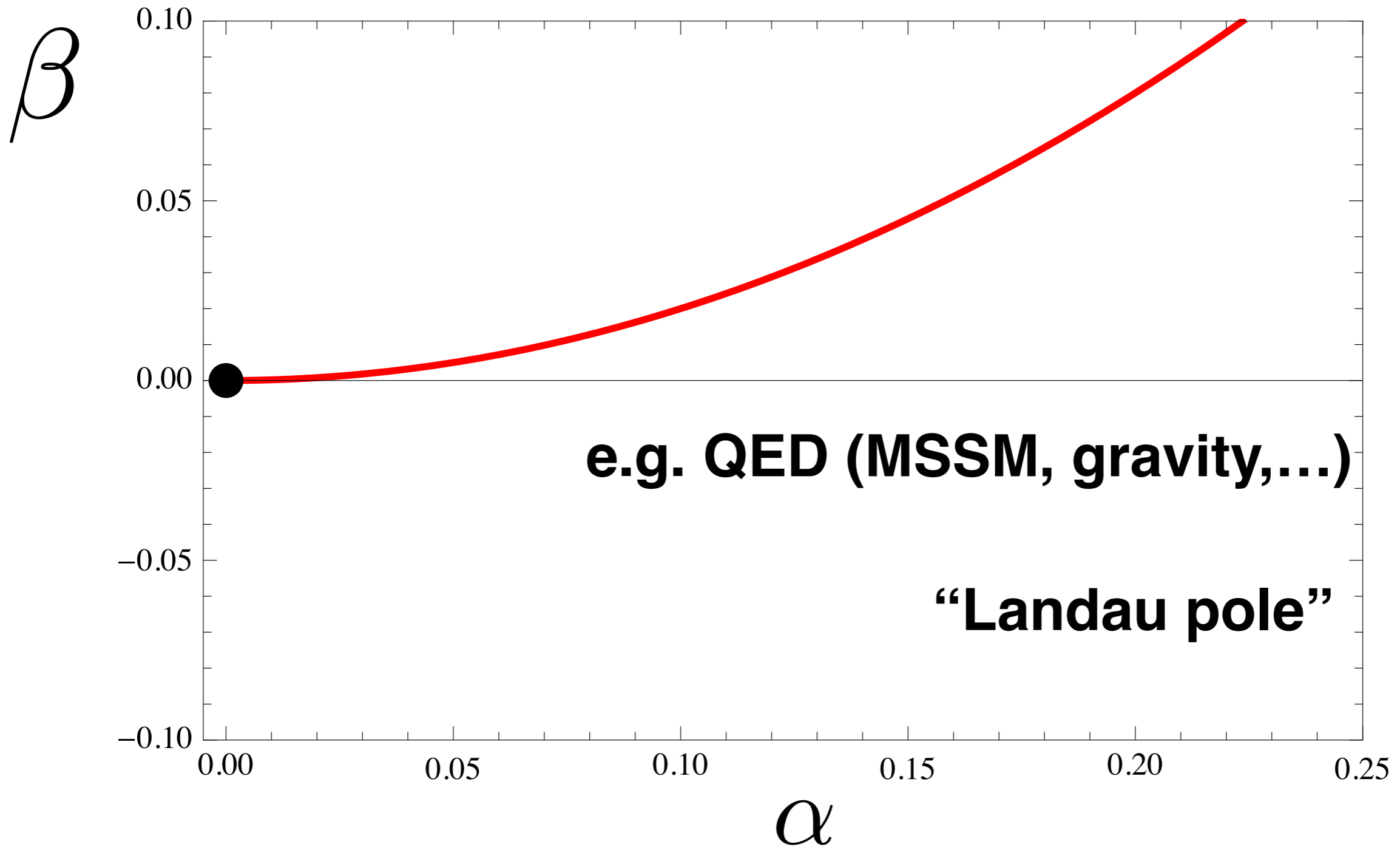
**AD Bond, DF Litim, 1608.00519/EPJC
1707.04217/PRD
1709.06953/PRL
1710.07612/PRD
1801.08527/PRL**

$$\beta = \frac{d\alpha}{d \ln \mu}$$

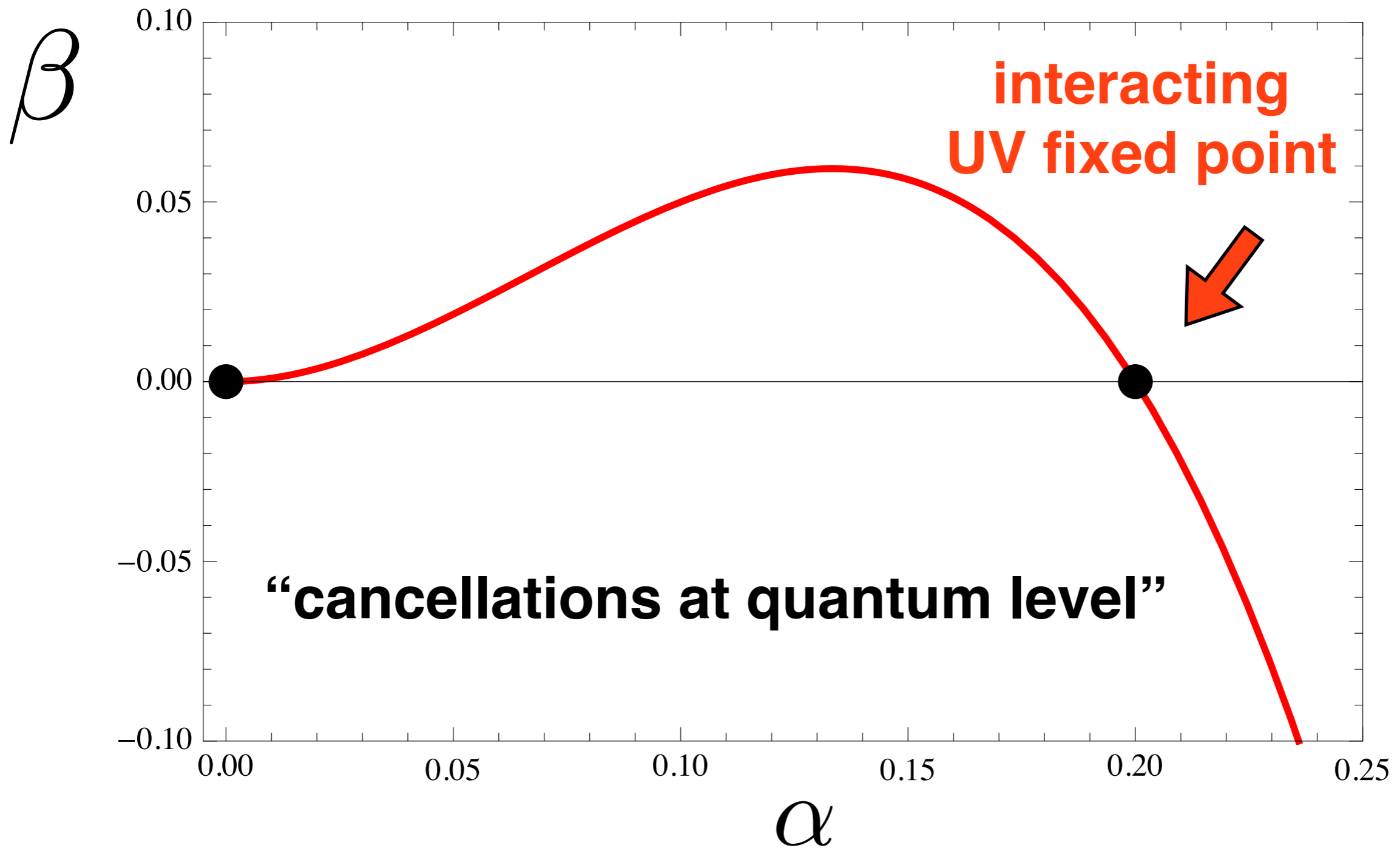
asymptotic freedom



infrared freedom



asymptotic safety



fields

vectors A_μ^a , **fermions** ψ_I , **scalars** ϕ^A

path integral

$$Z[J] = \exp -i \int d^4x (L + L_{\text{gf}} + L_{\text{gh}} + J^i \Phi_i)$$

action

$$L = \frac{1}{4g_a^2} \text{Tr} F_{\mu\nu}^a F_a^{\mu\nu} + i\psi_I \not{D}\psi_I + \frac{1}{2} (D_\mu \phi^A)^2$$

$$+ \frac{1}{2} Y^A_{IJ} \phi^A \psi_I \xi \psi_J + \frac{1}{4!} \lambda_{ABCD} \phi^A \phi^B \phi^C \phi^D$$

fields

vectors A_μ^a , **fermions** ψ_I , **scalars** ϕ^A

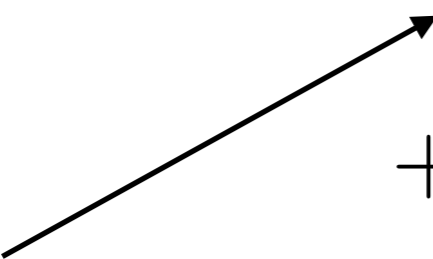
path integral

$$Z[J] = \exp -i \int d^4x (L + L_{\text{gf}} + L_{\text{gh}} + J^i \Phi_i)$$

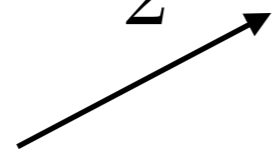
action

$$L = \frac{1}{4g_a^2} \text{Tr} F_{\mu\nu}^a F_a^{\mu\nu} + i\psi_I \not{D}\psi_I + \frac{1}{2} (D_\mu \phi^A)^2$$

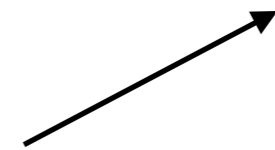
gauge



Yukawa

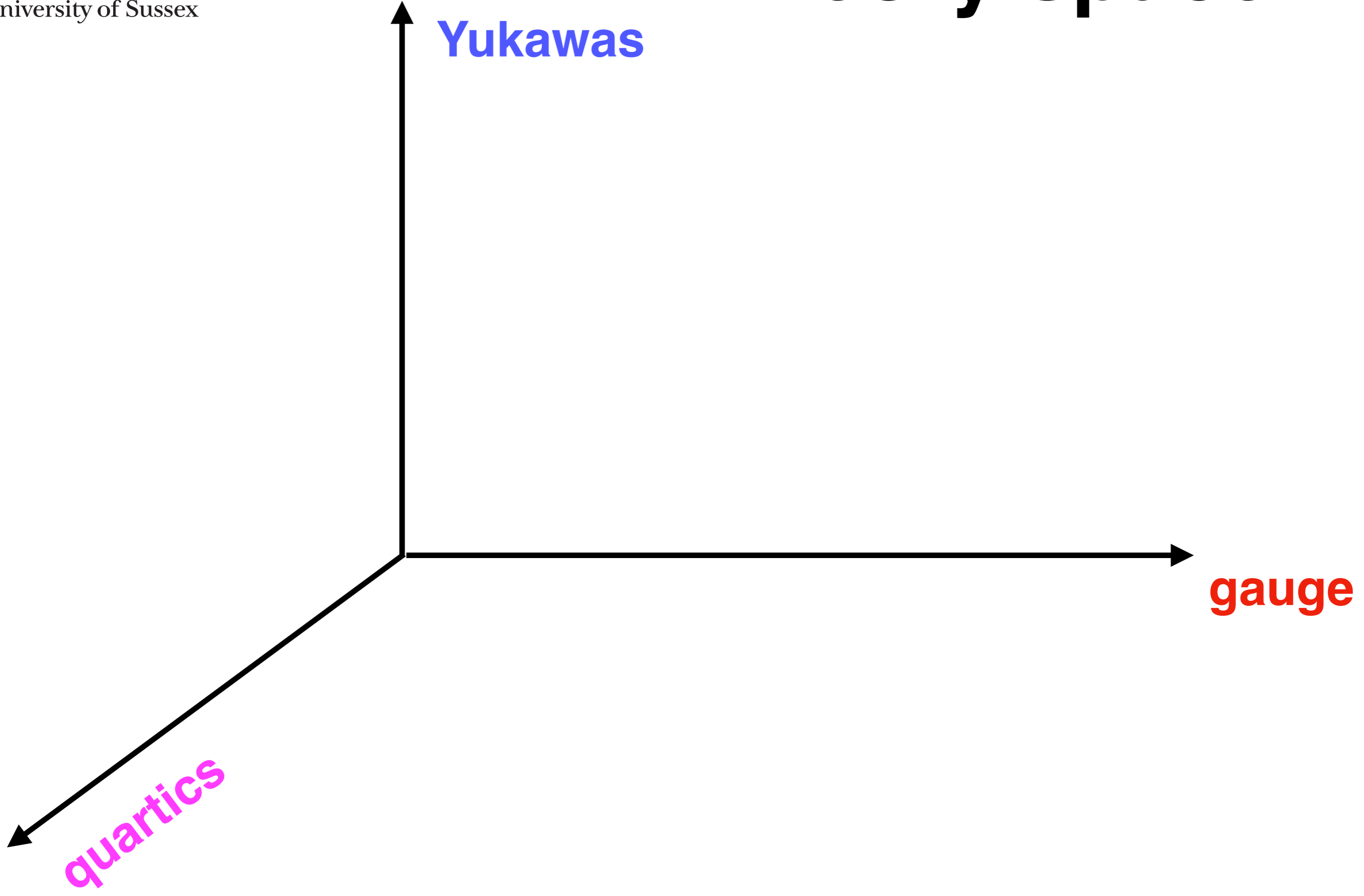


quartics

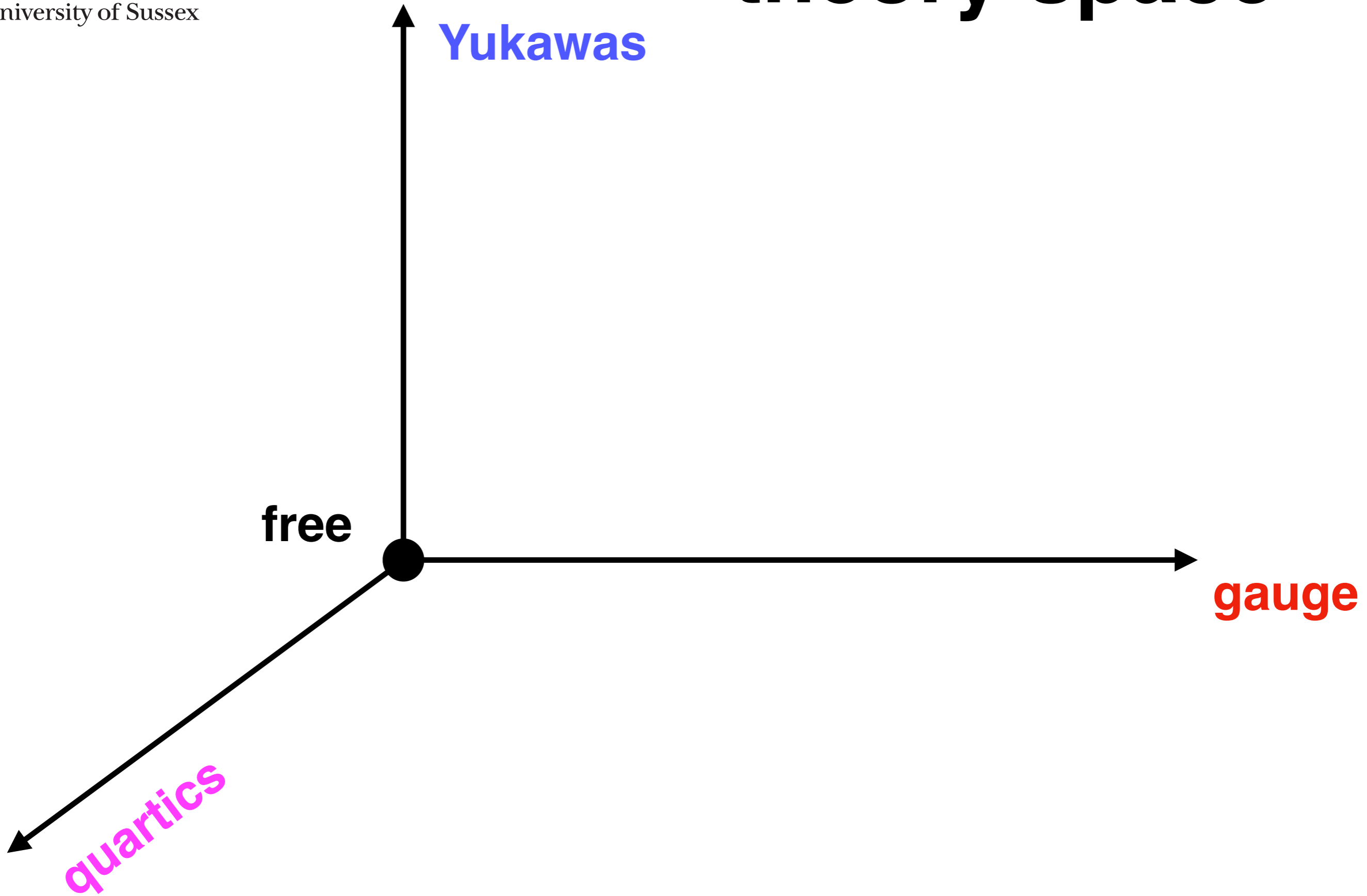


$$+ \frac{1}{2} Y^A_{IJ} \phi^A \psi_I \xi \psi_J + \frac{1}{4!} \lambda_{ABCD} \phi^A \phi^B \phi^C \phi^D$$

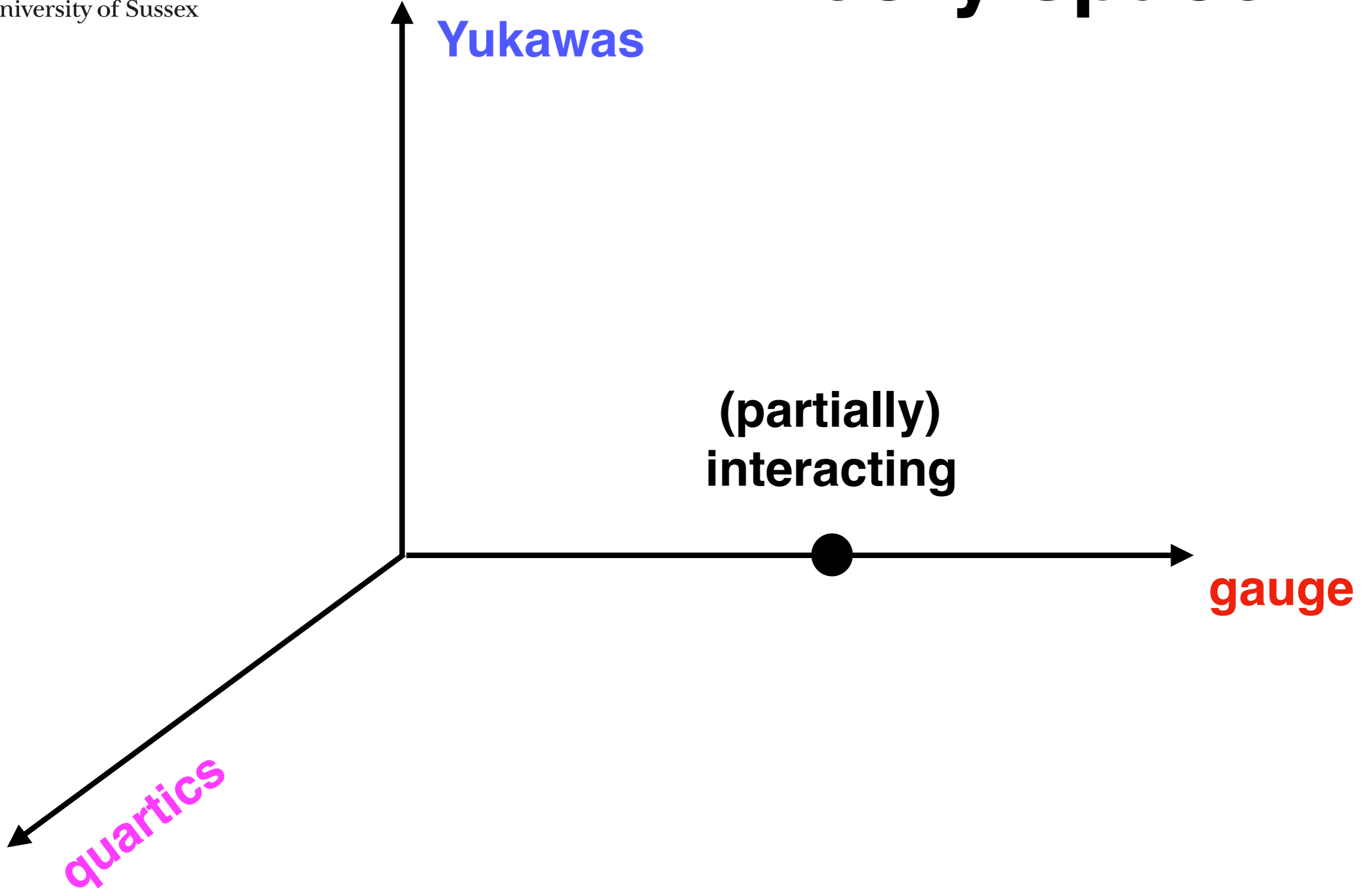
“theory space”



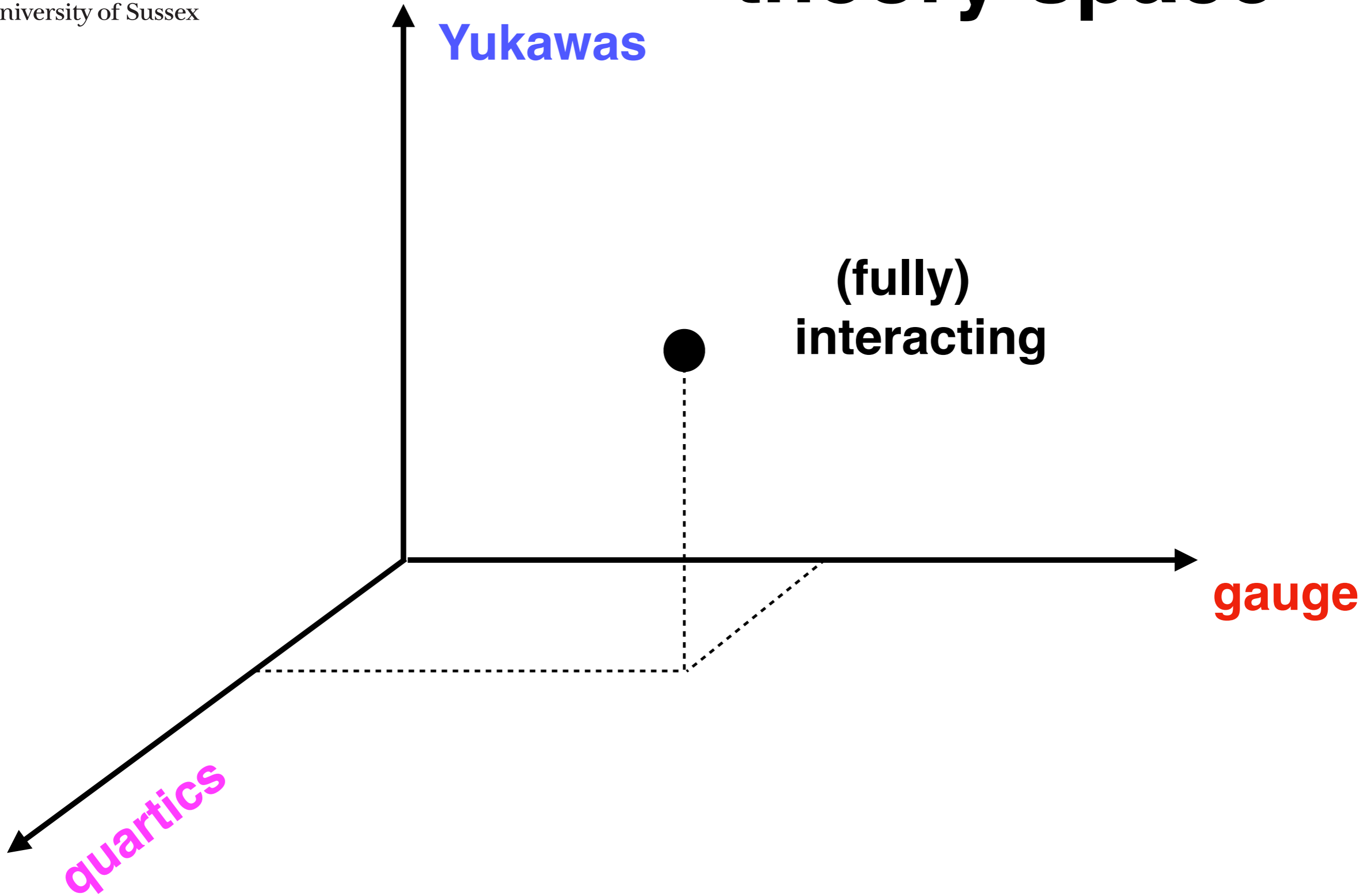
“theory space”



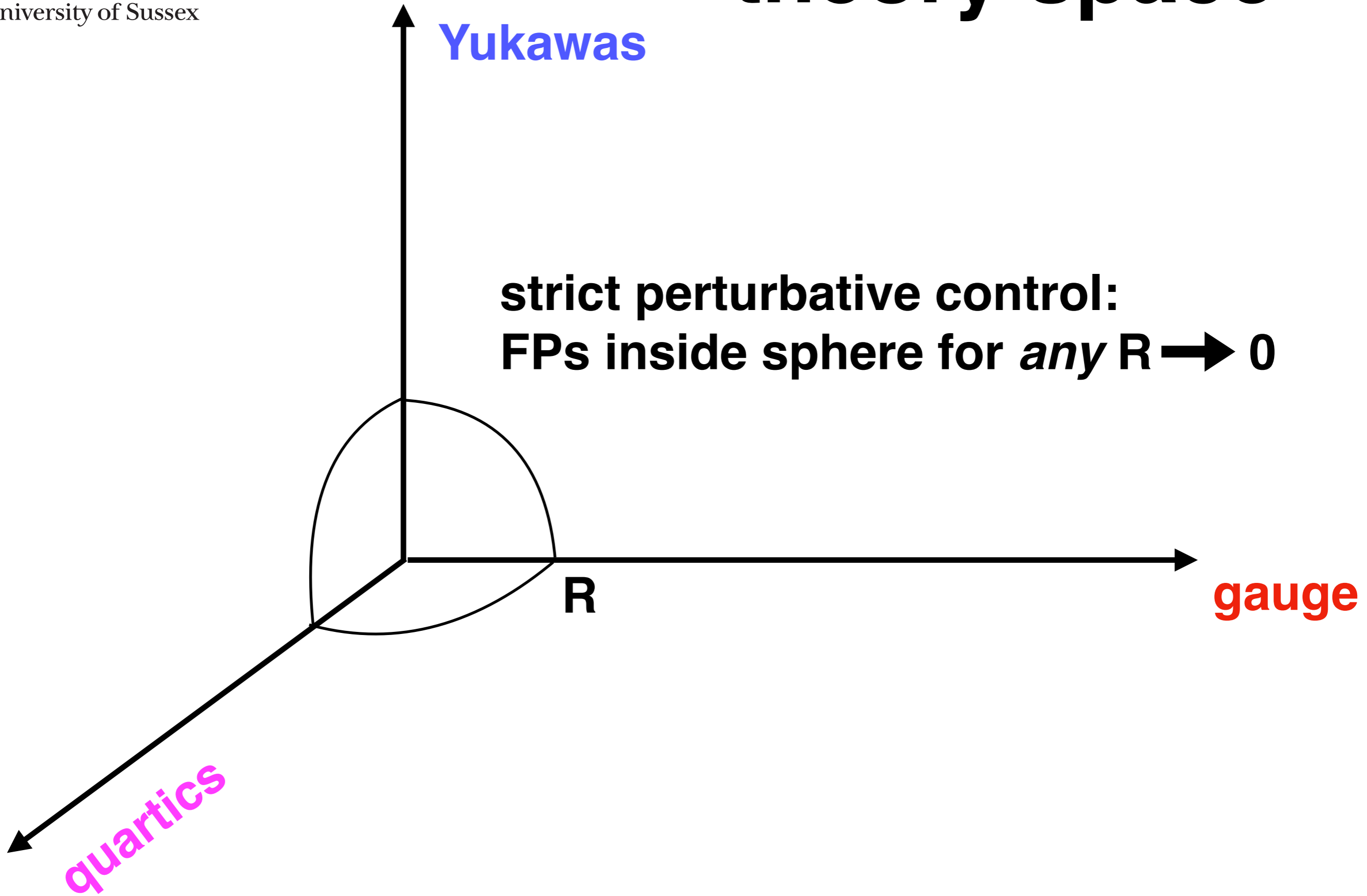
“theory space”



“theory space”



“theory space”



interacting fixed point

gauge

Y Y Y N N

Yukawas

N N Y N Y

quartics

N Y Y Y Y

“Banks
Zaks”

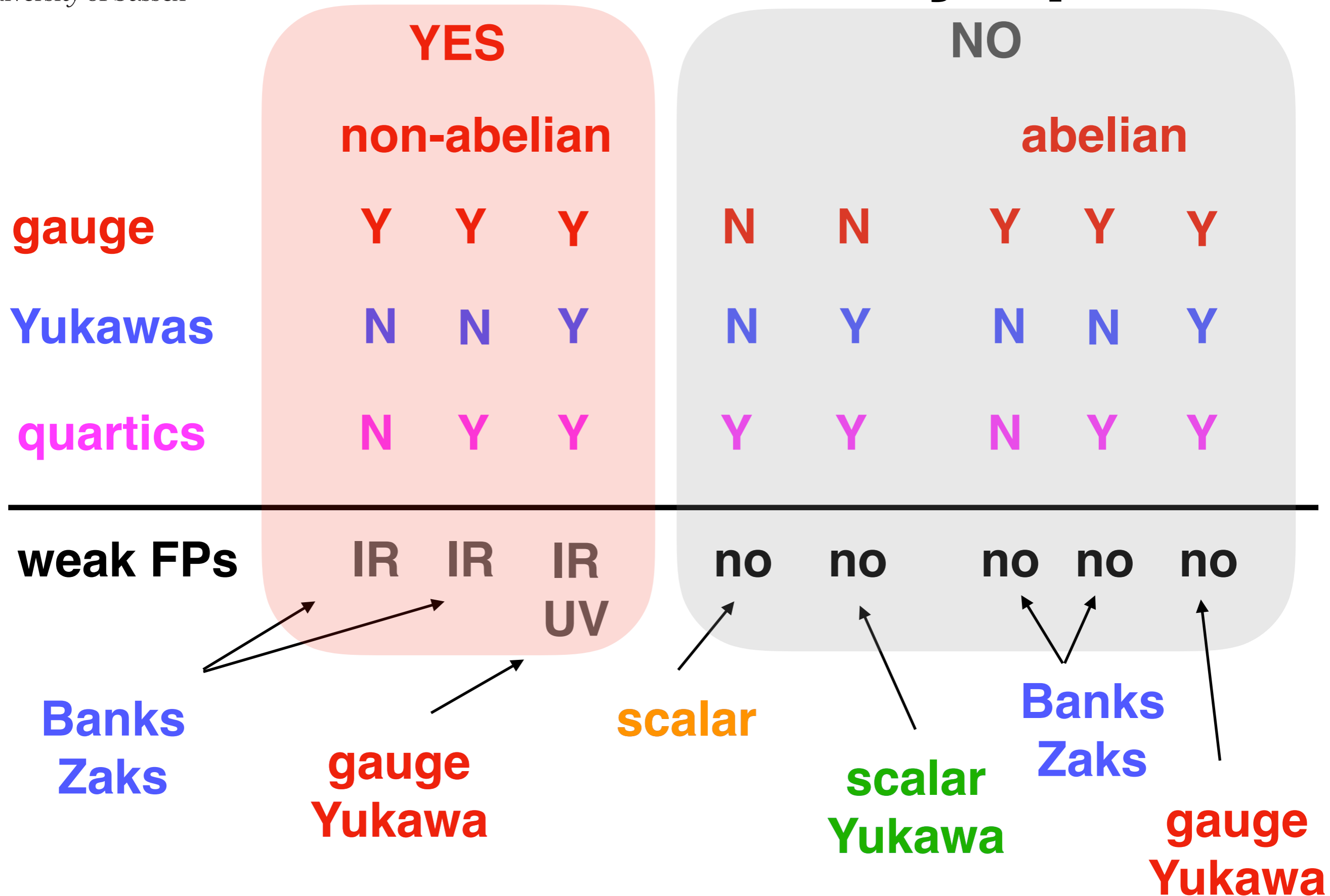
“gauge
Yukawa”

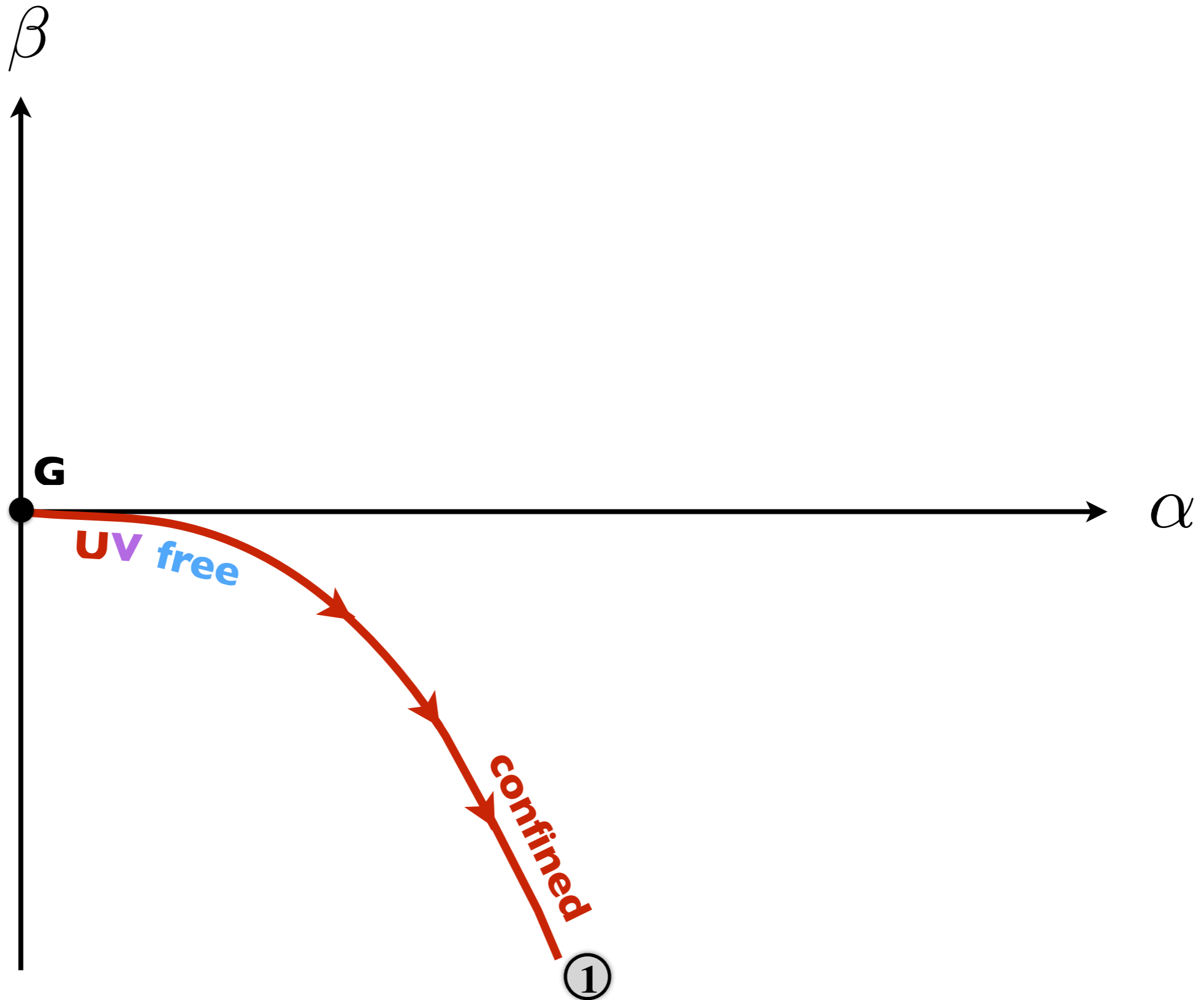
“scalar”

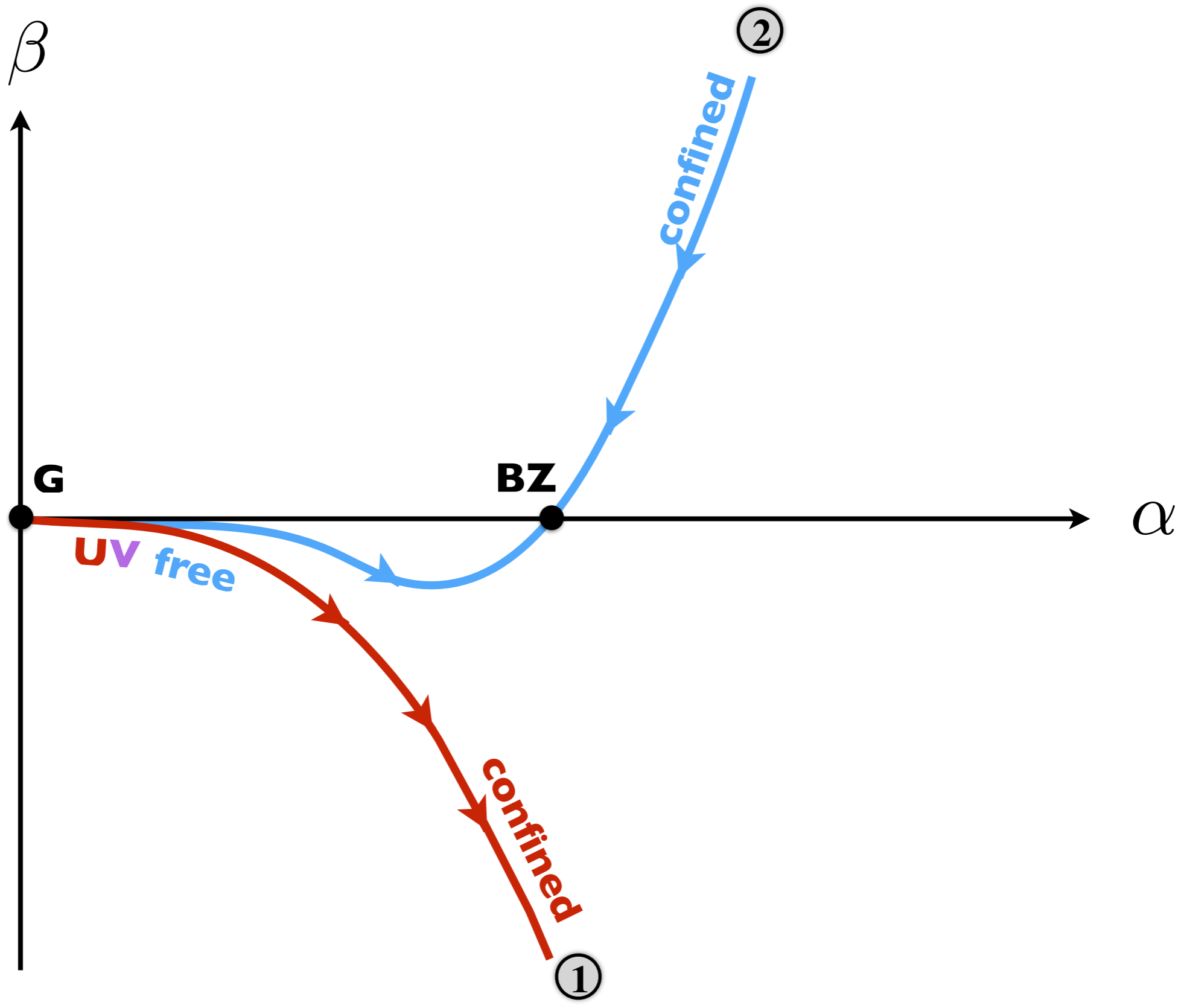
“scalar
Yukawa”

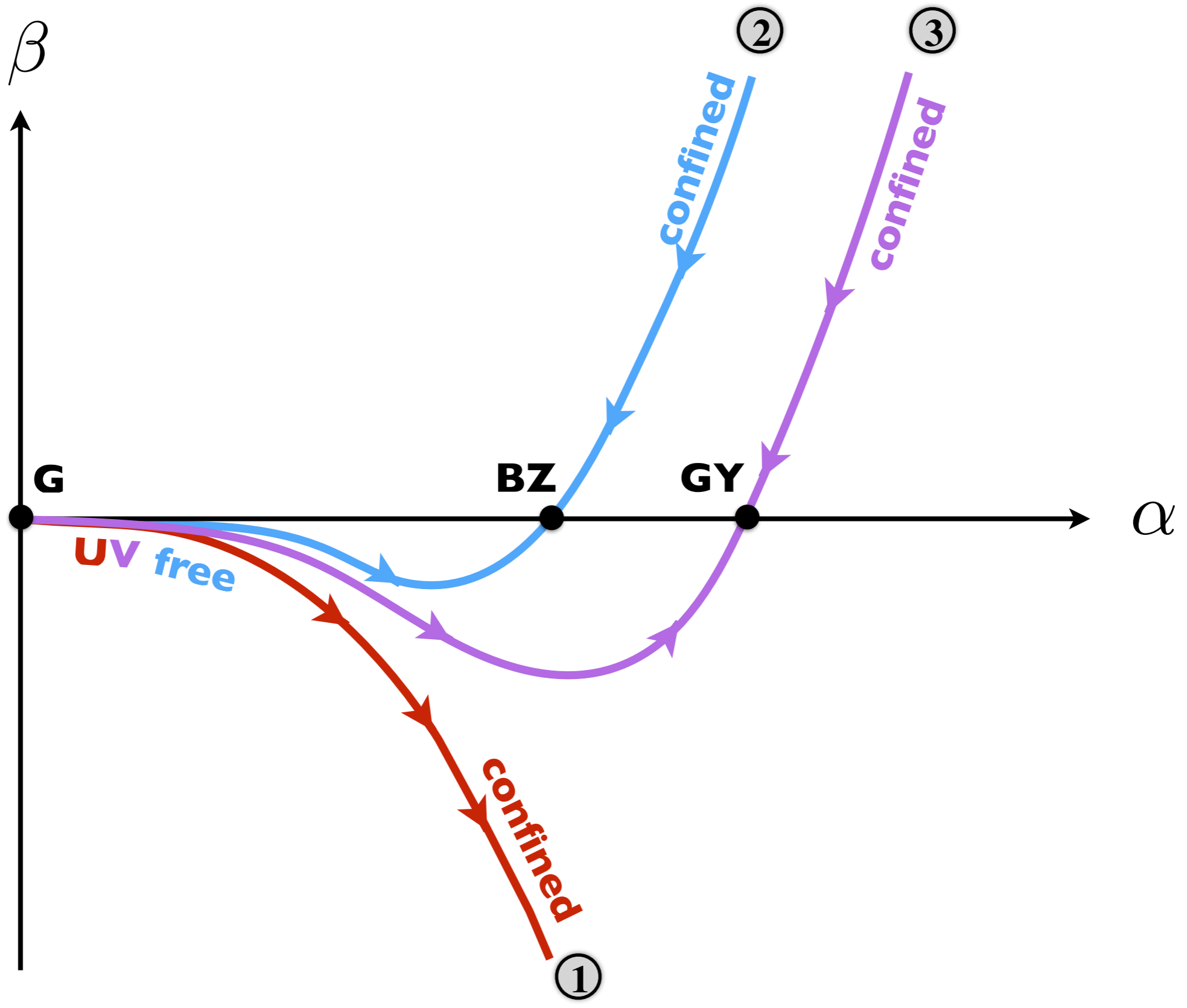


“theory space”

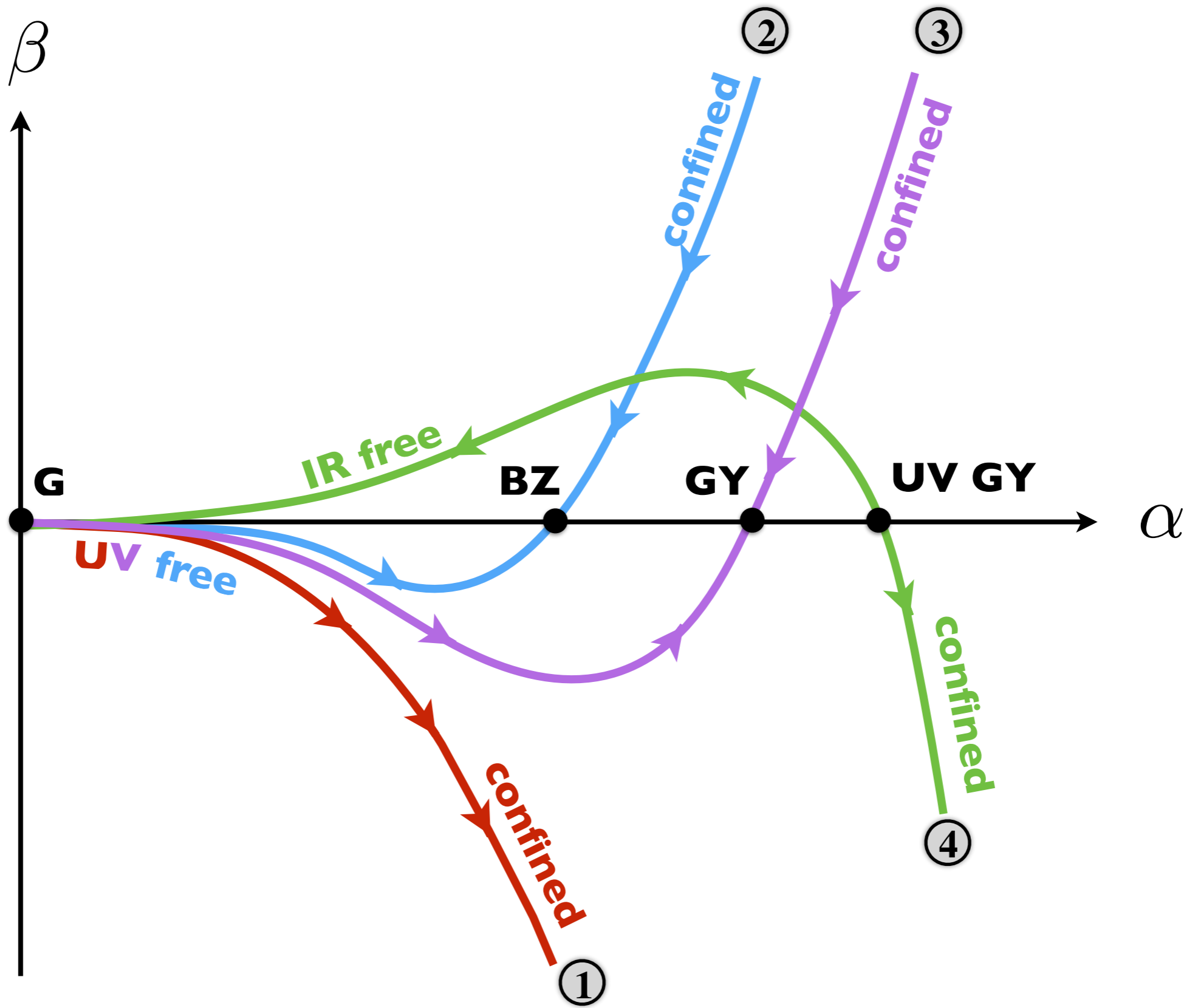


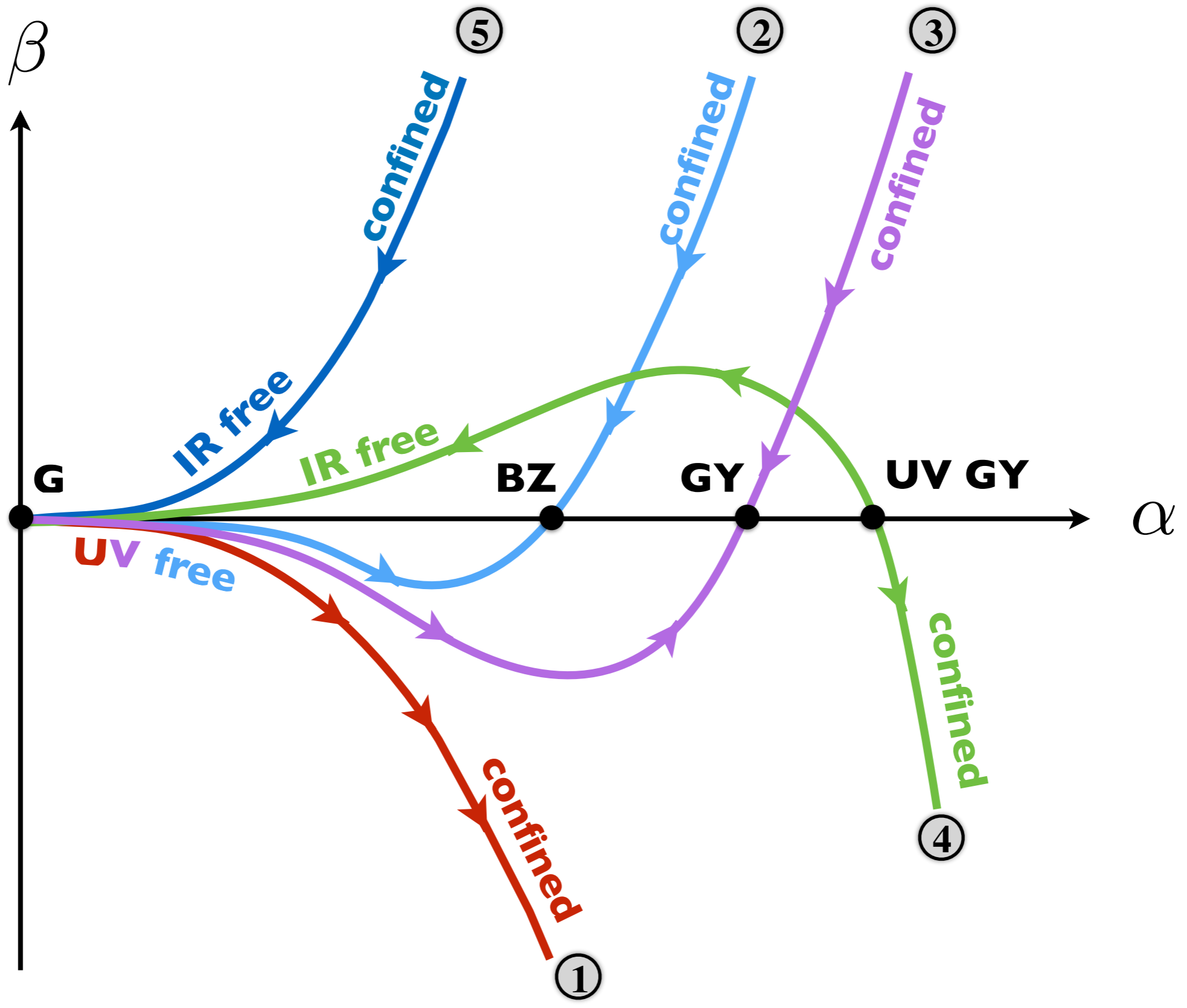


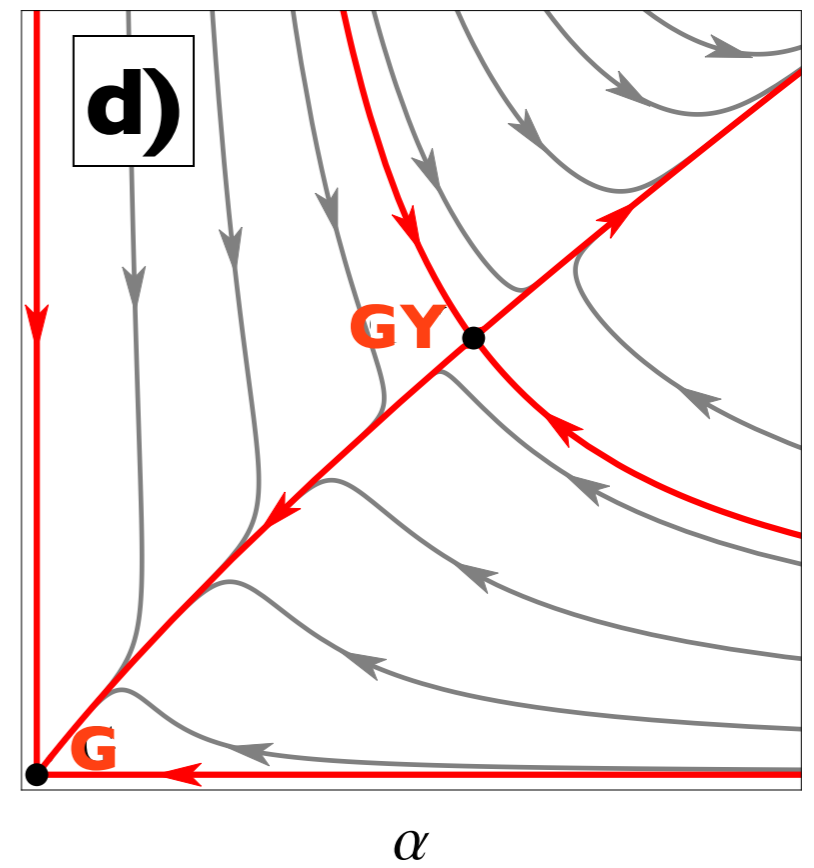
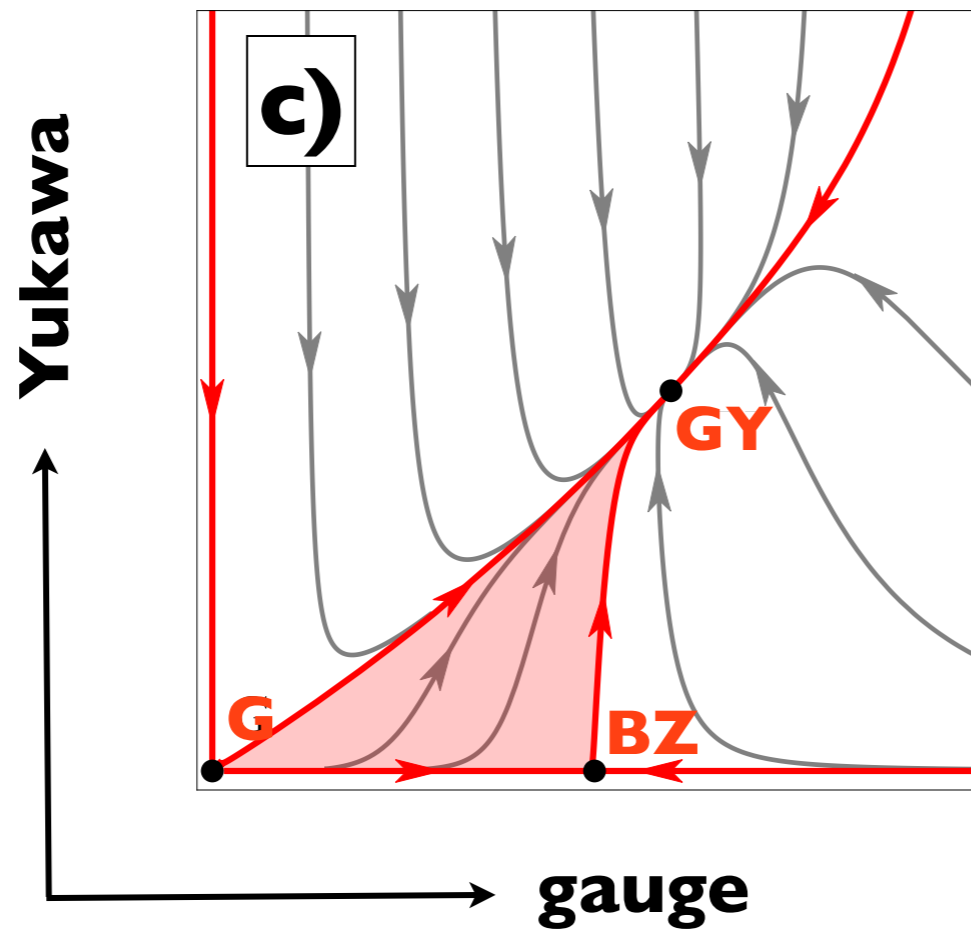
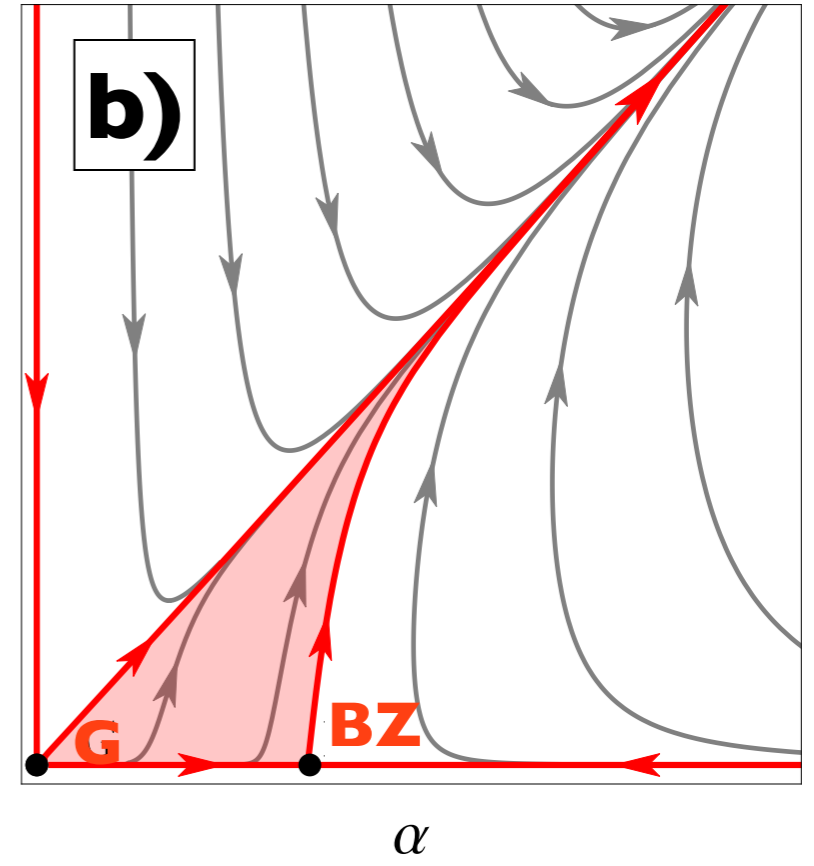
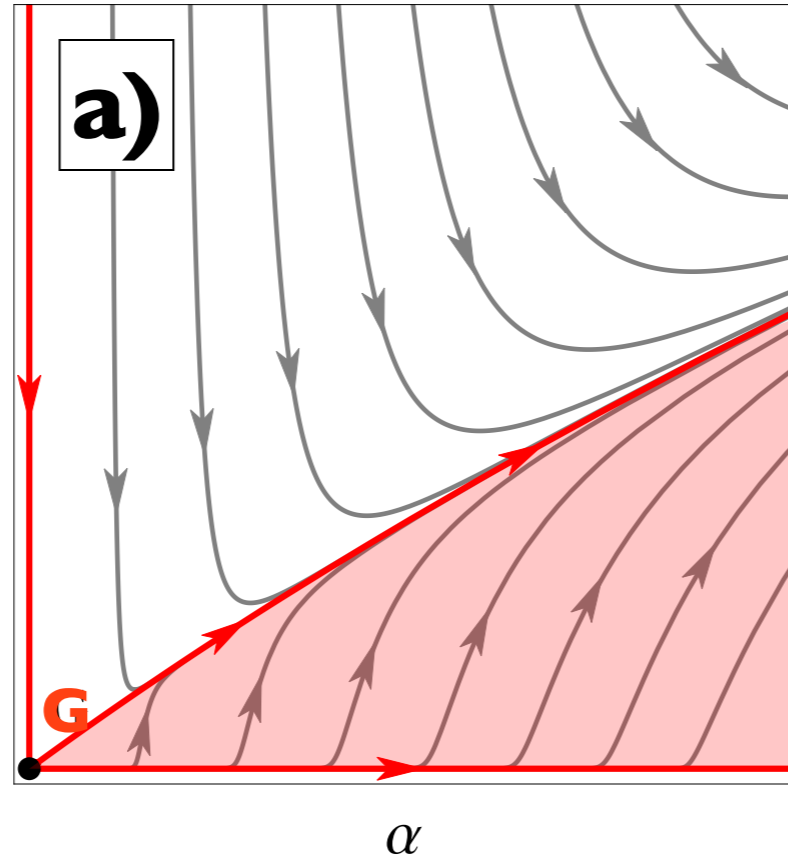




renormalisation group

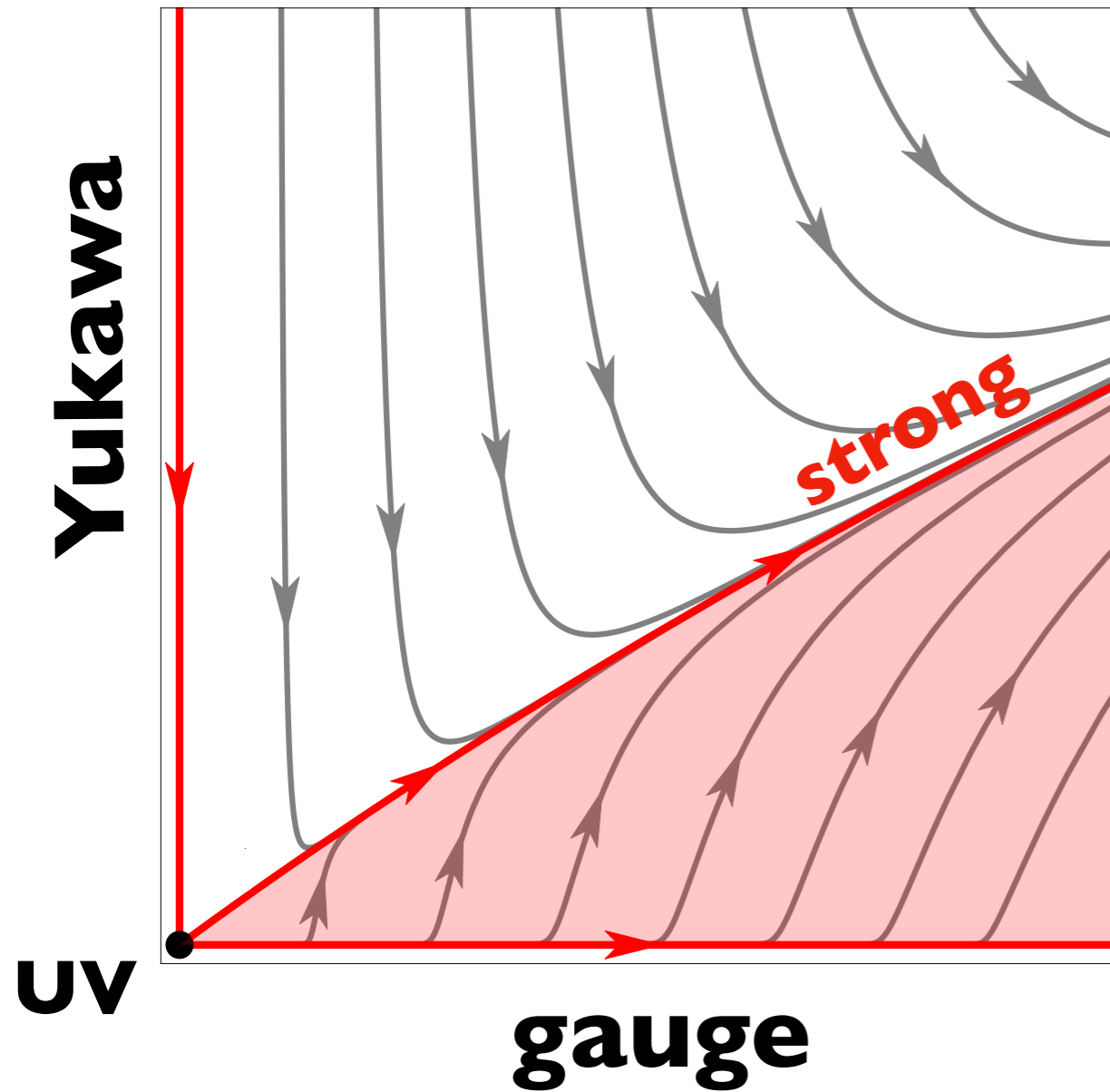






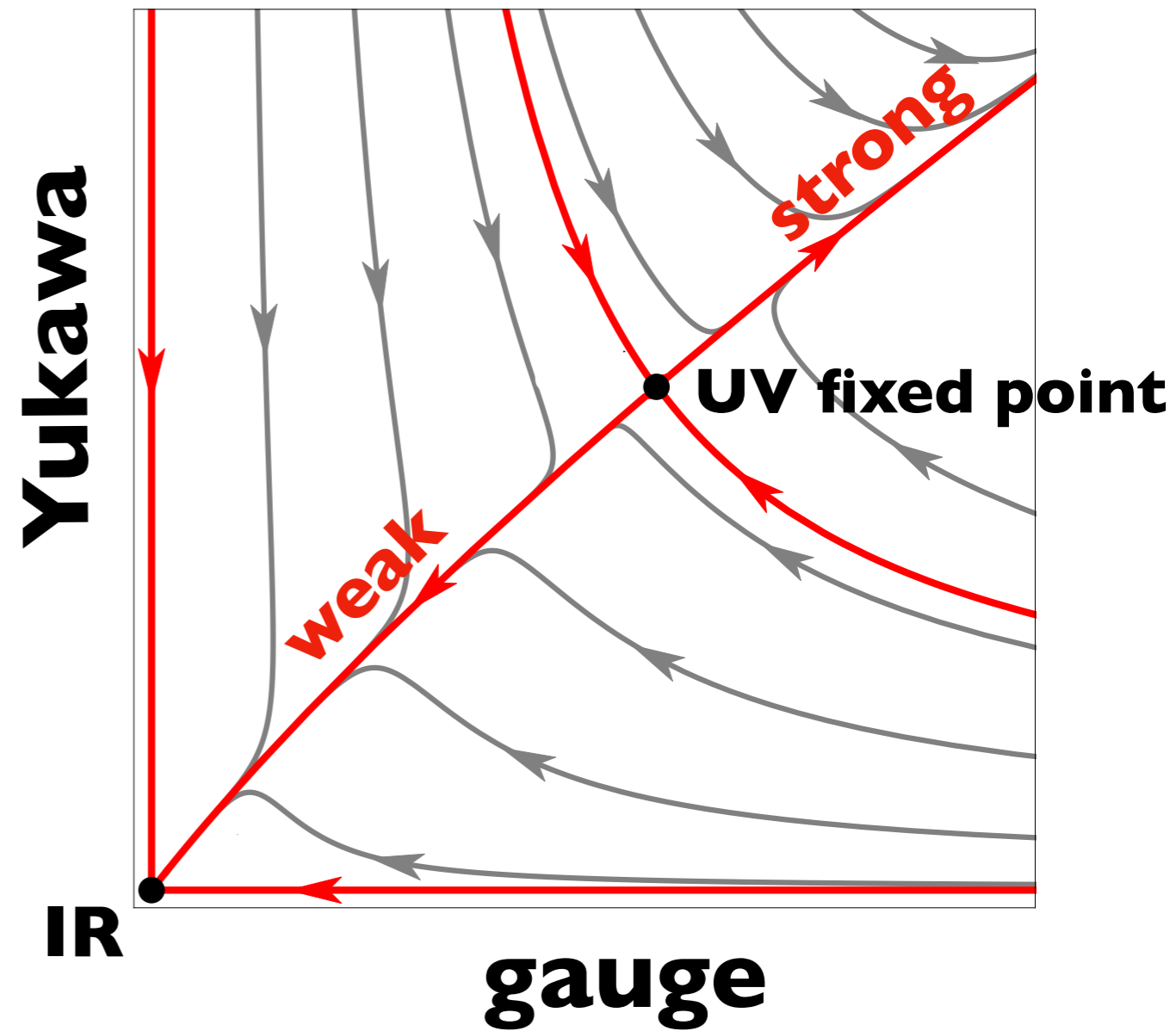
①

**asymptotic
freedom**



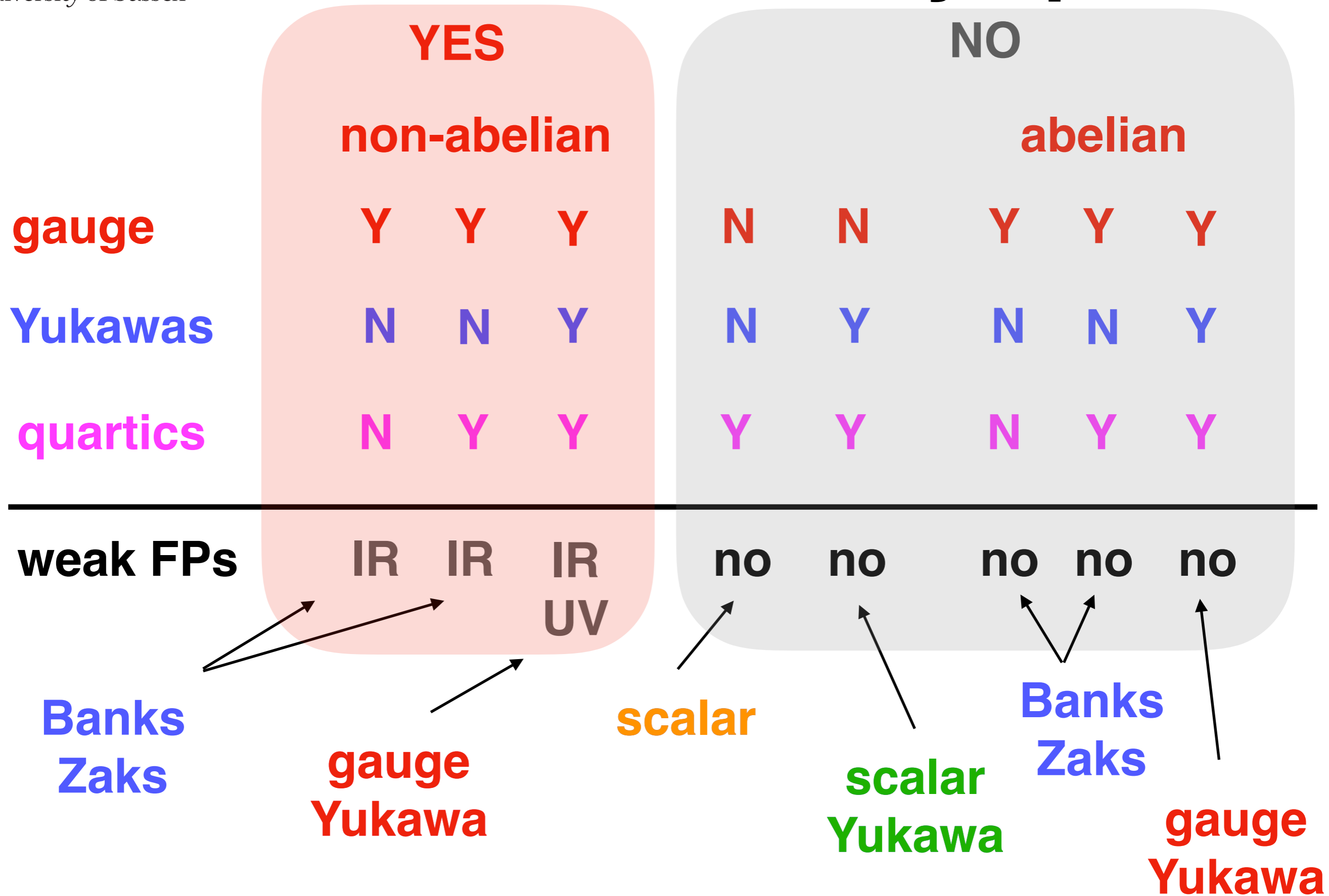
④

**asymptotic
safety**



“enhanced predictivity”

“theory space”



proofs of fixed points & asymptotic safety

general **theorems** for fixed points

AD Bond, DF Litim, **Theorems for Asymptotic Safety of Gauge Theories**, 1608.00519 (EPJC)

AD Bond, DF Litim, **Price of Asymptotic Safety**, 1801.08527 (PRL)

simple gauge theories with matter

DF Litim, F Sannino, **Asymptotic Safety Guaranteed**, 1406.2337 (JHEP)

AD Bond, DF Litim, G Medina Vazquez, T Steudtner, **Conformal window for asymptotic safety**, 1710.07615 (PRD)

semi-simple $SU(N) \times SU(M)$ gauge theories with matter

AD Bond, DF Litim, **More Asymptotic Safety Guaranteed**, 1707.04217 (PRD)

supersymmetric gauge theories with matter

AD Bond, DF Litim, **Asymptotic Safety Guaranteed in Supersymmetry**, 1709.06953 (PRL)

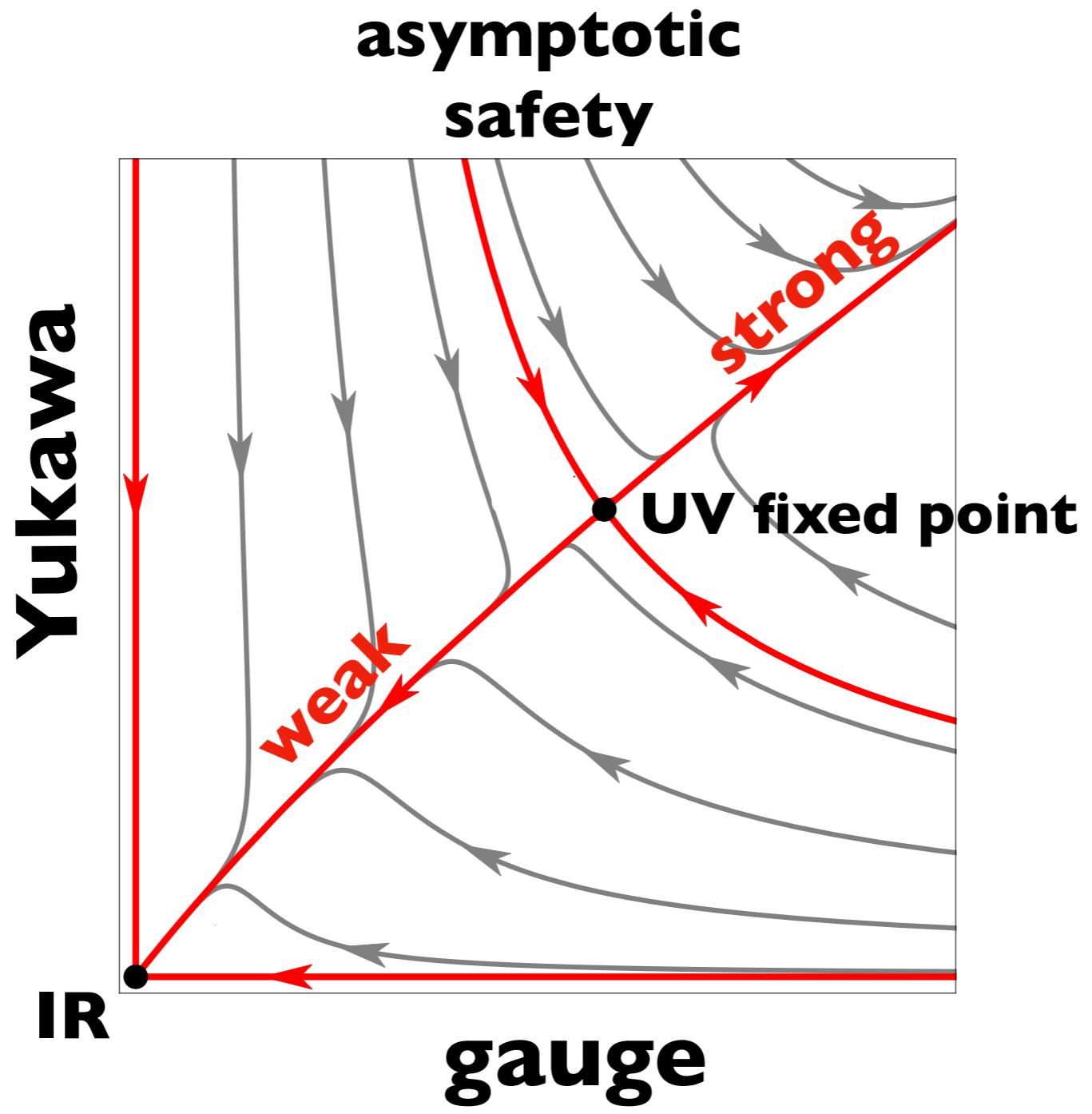
higher order interactions in gauge theories with matter

T Buyukbese, DF Litim, **Asymptotic Safety Beyond Marginal Interactions**, PoS LATTICE2016 (2017) 233

phenomenology and models beyond the Standard Model

A Bond, G Hiller, K Kowalska, DF Litim, **Directions for model building from asymptotic safety**, JHEP1708 (2017) 004

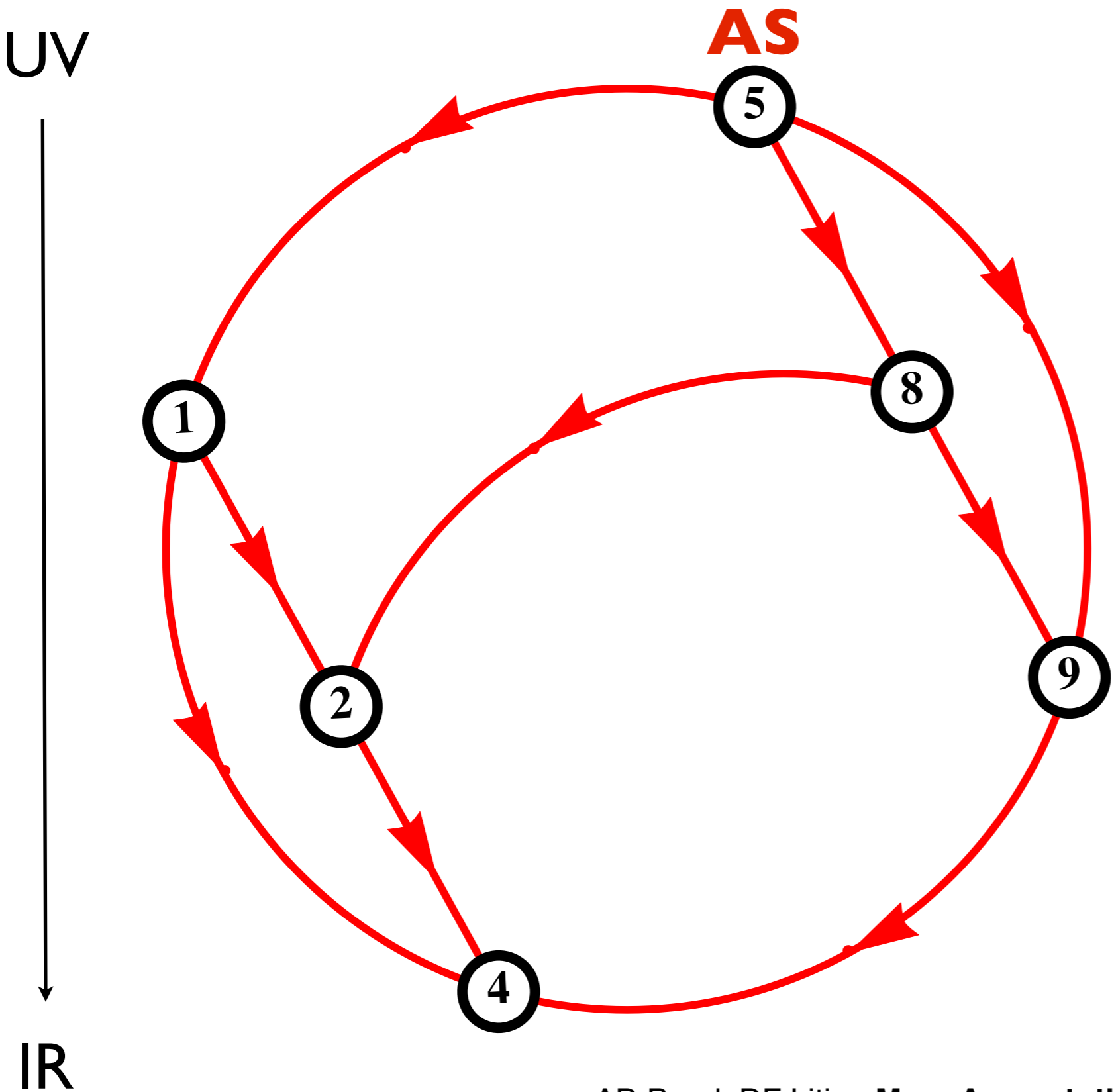
G Hiller, C Hermigos-Feliu, DF Litim, T Steudtner, **Asymptotically safe extensions of the Standard Model and their flavour phenomenology**, Moriond (EW2019) 1905.11020



SU(N)
+1 Yukawa
+2 quartics

(NF fermions
+ mesons)

semi-simple

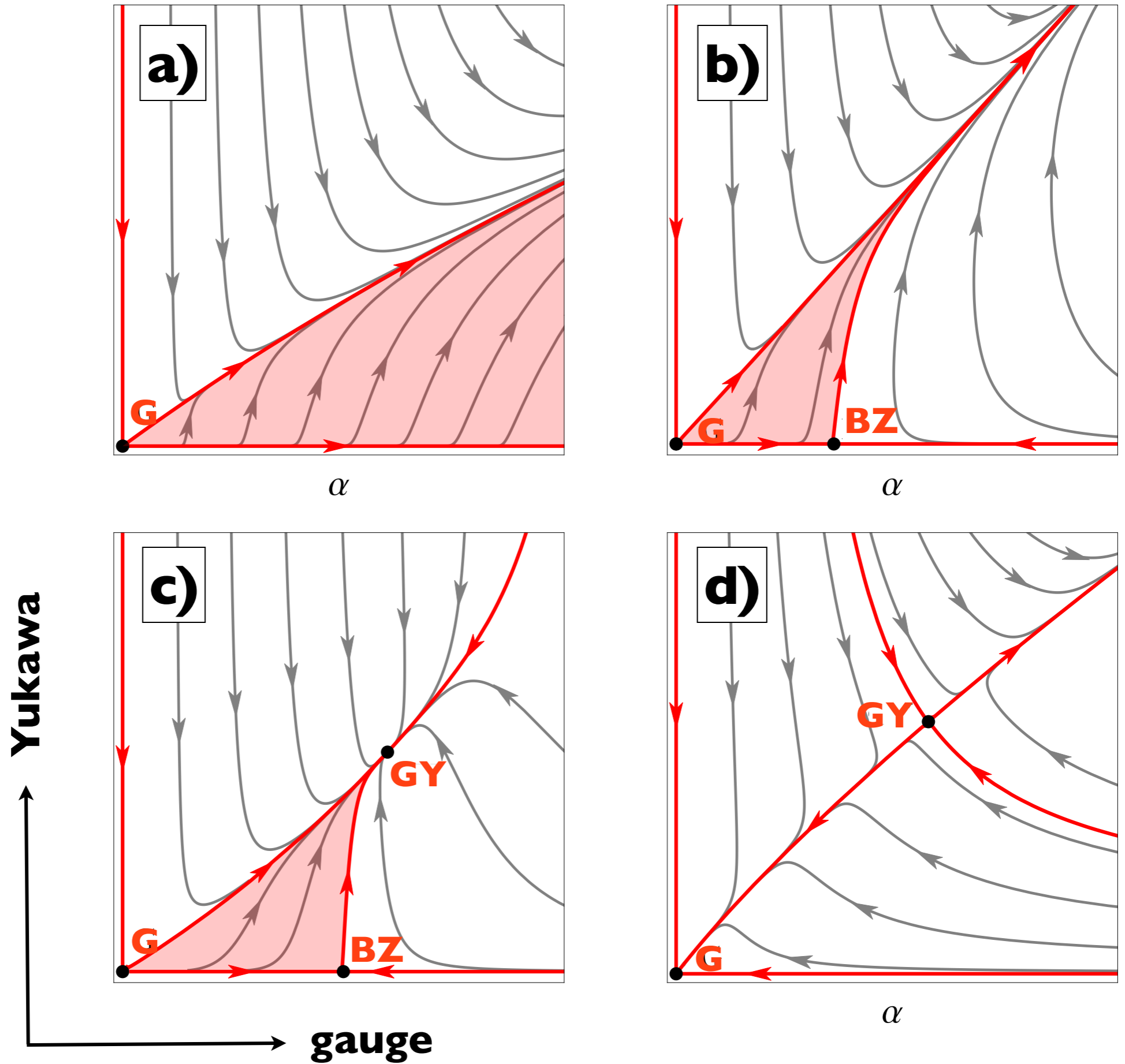


$SU(N) \times SU(M)$
+ 2 Yukawas
+ 5 quartics

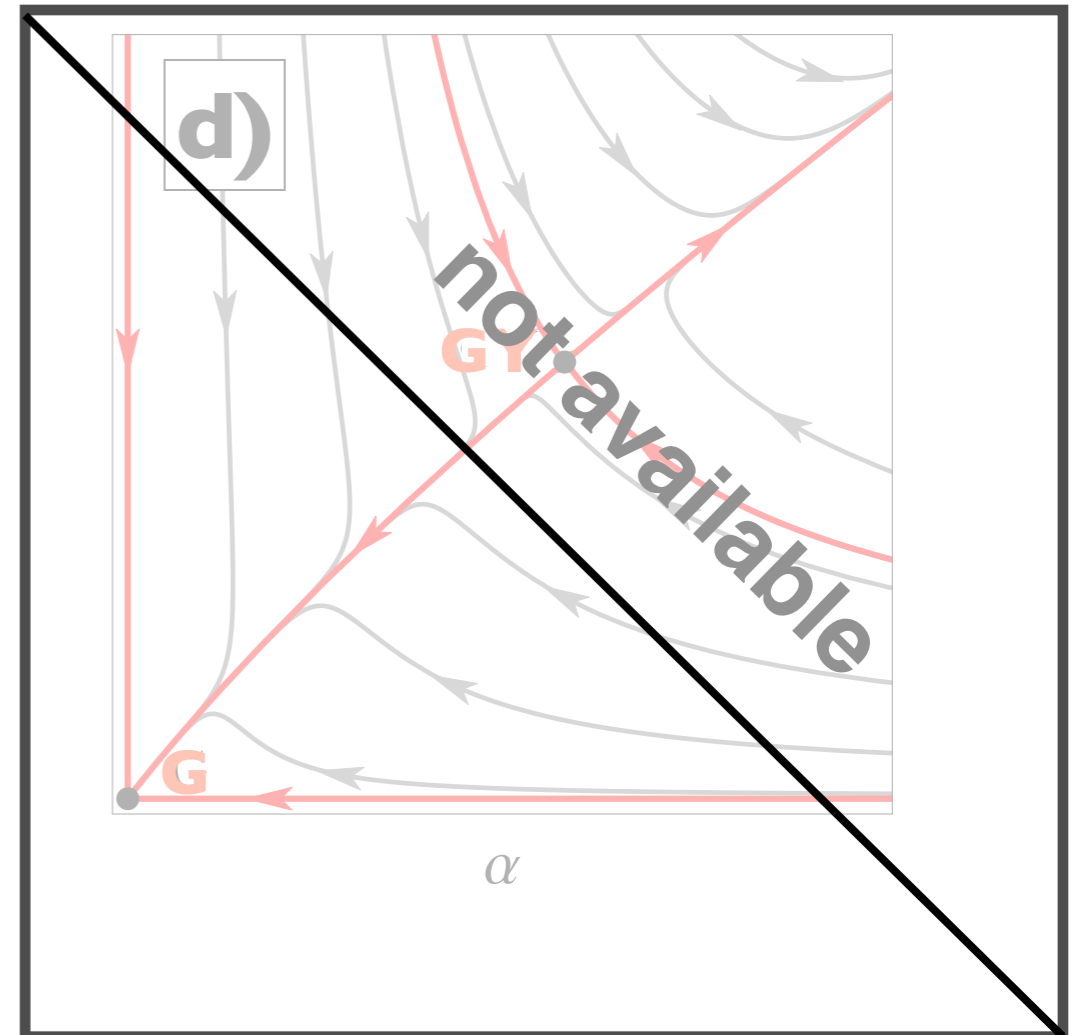
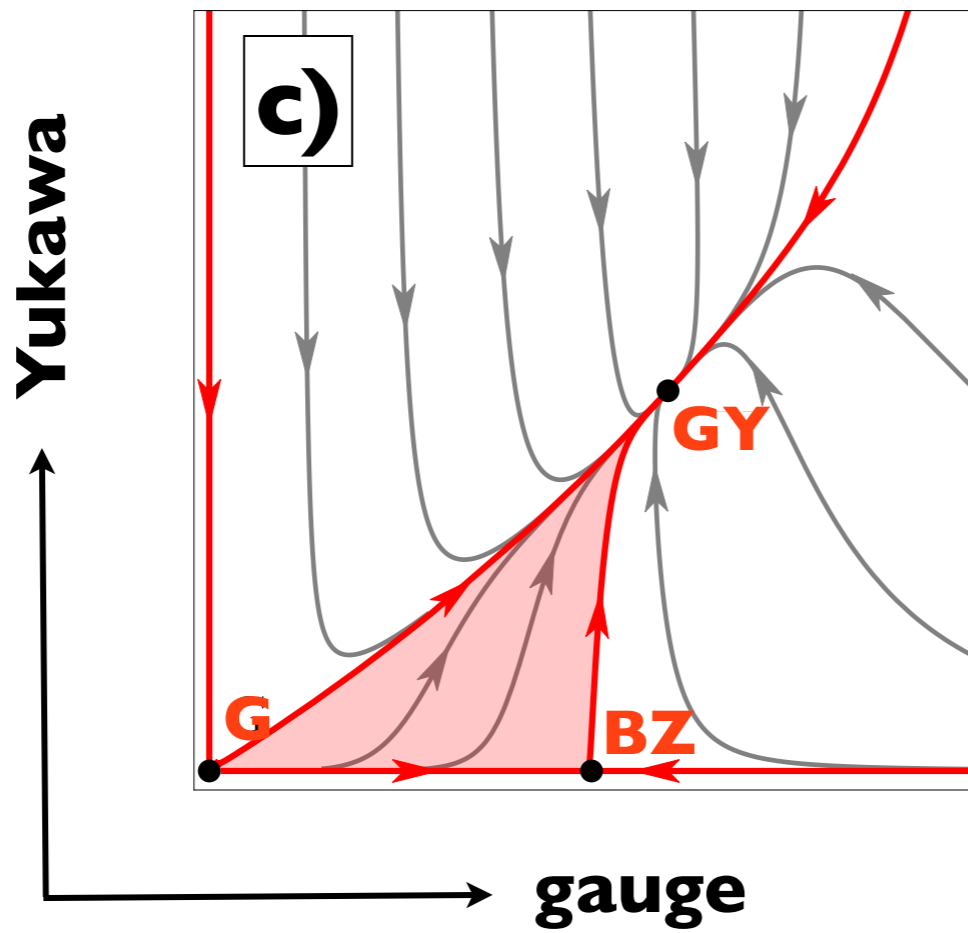
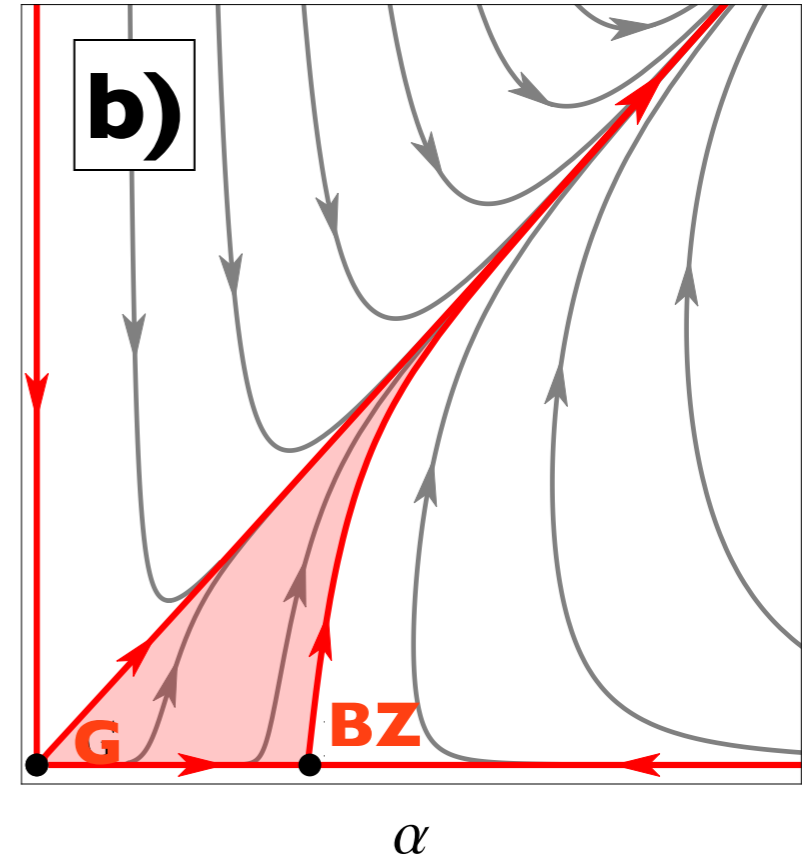
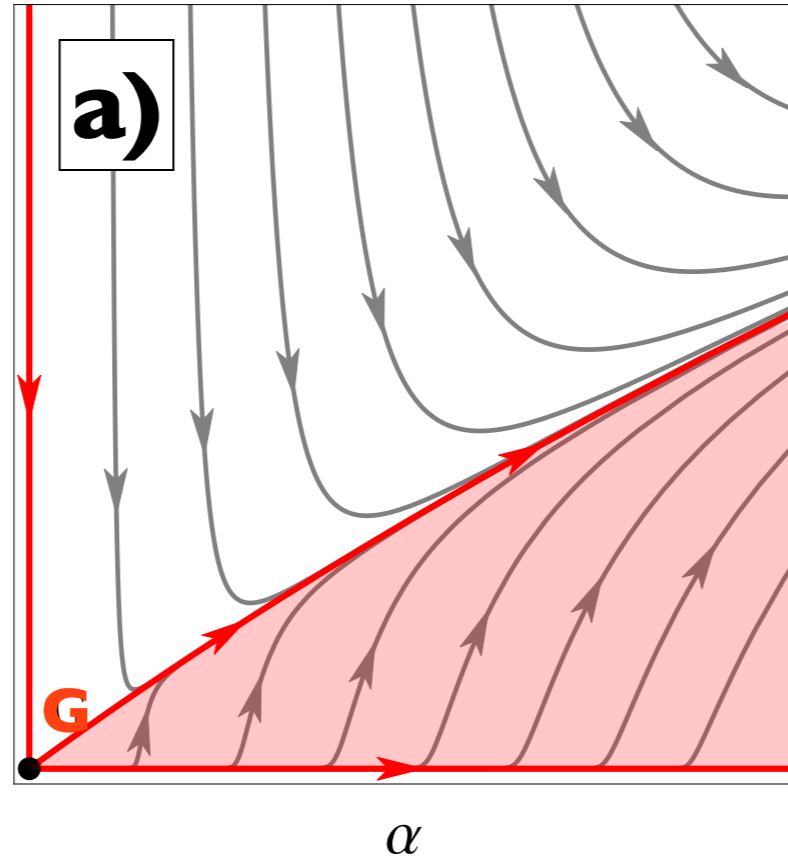
various fermions
and mesons

what about supersymmetry?

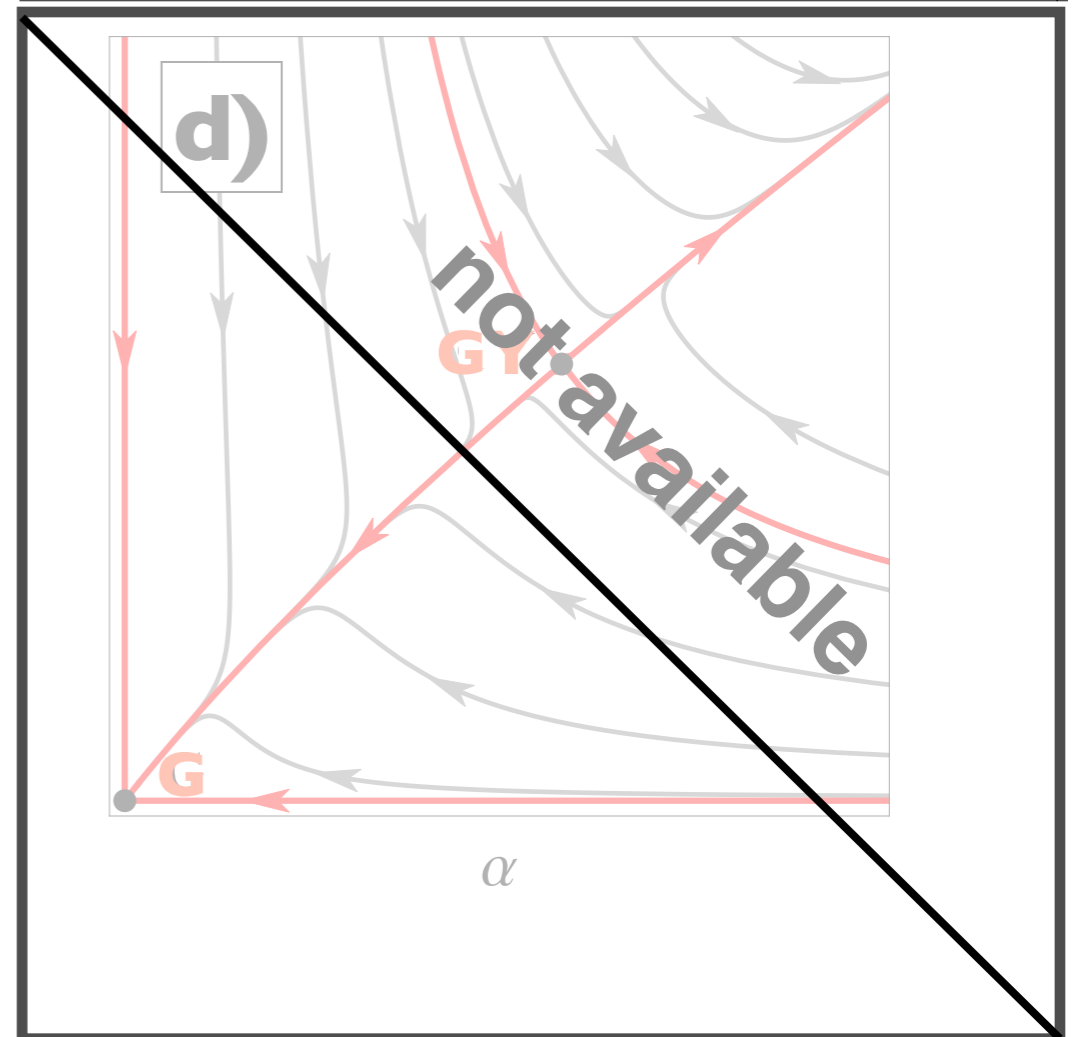
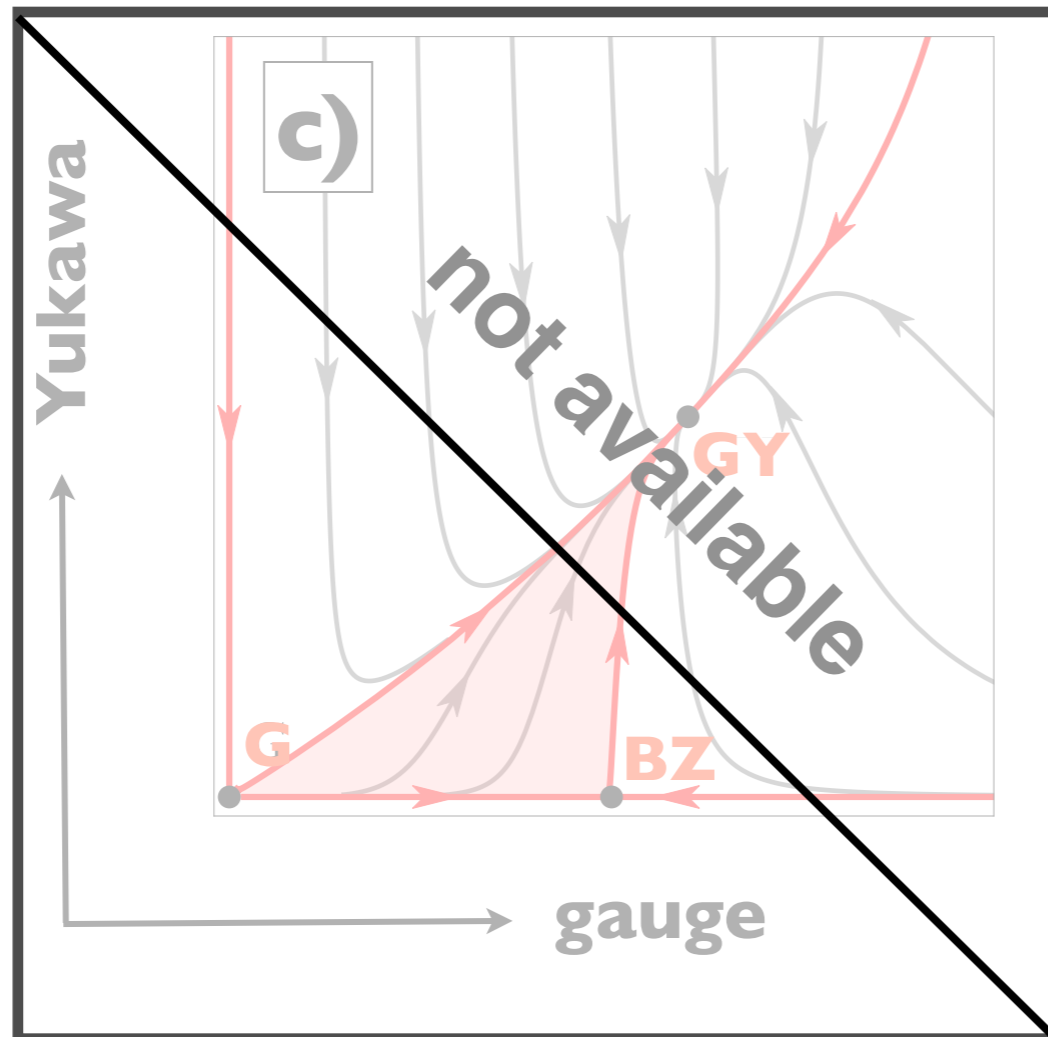
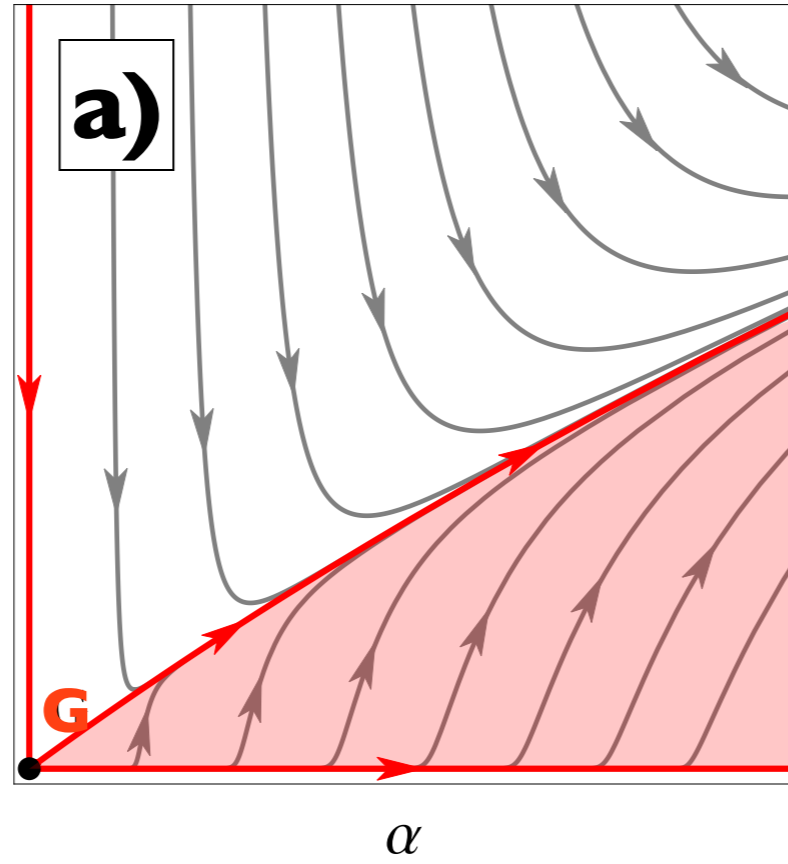
**N=0
SUSY**



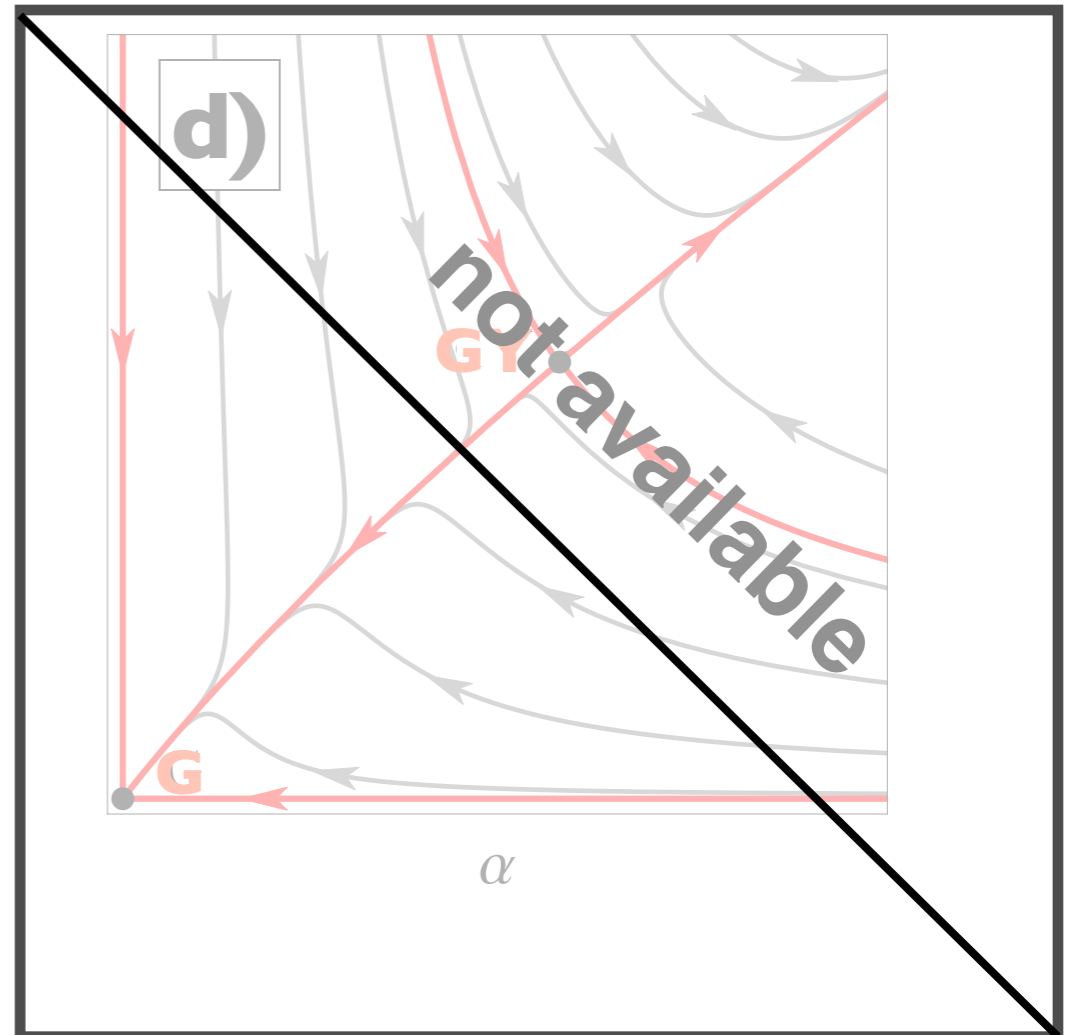
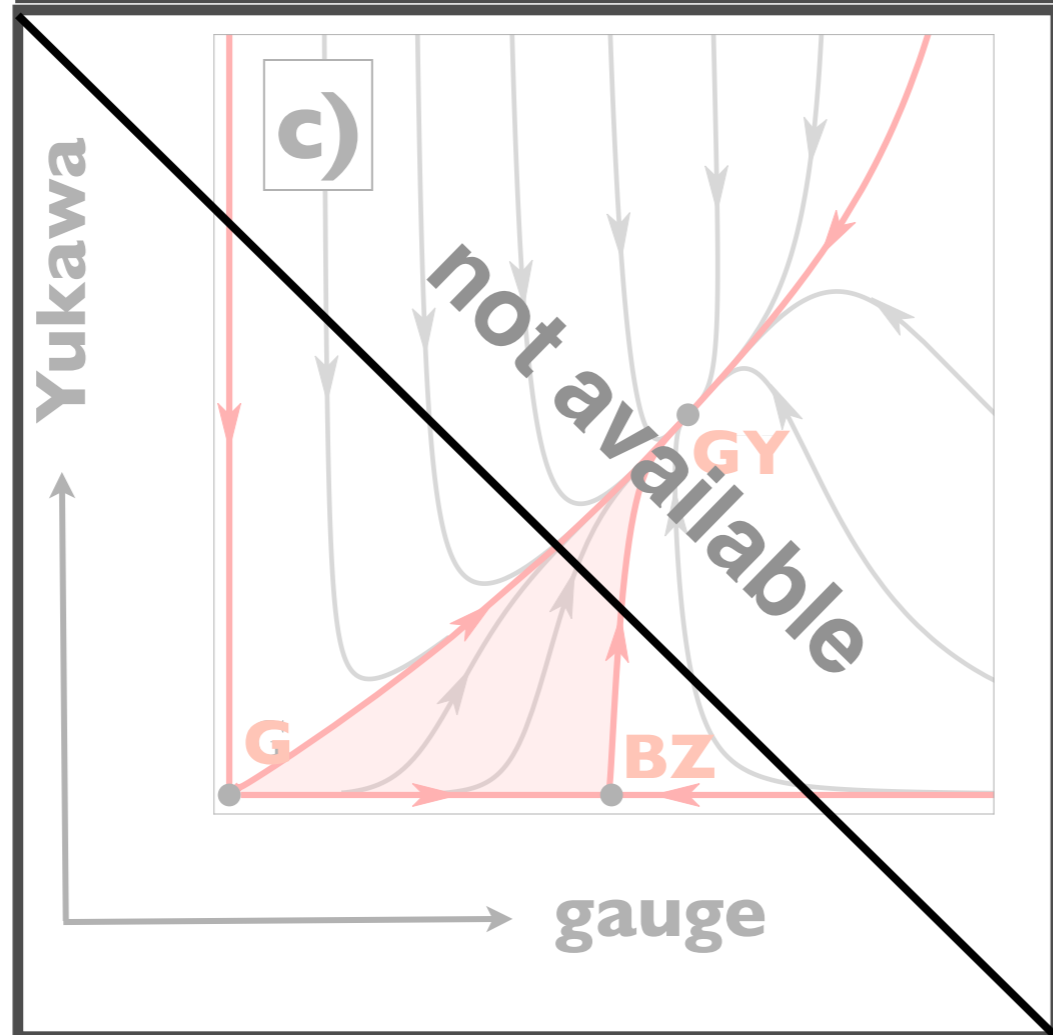
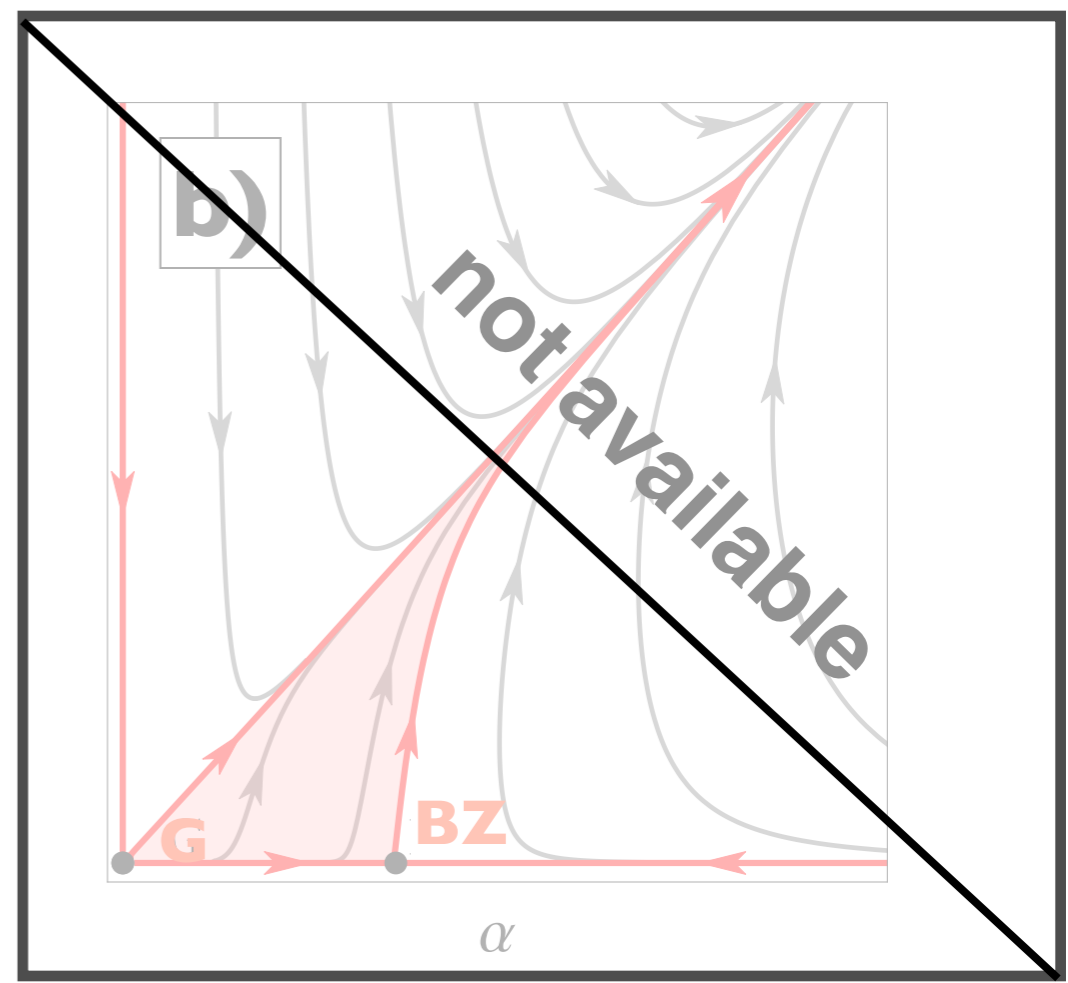
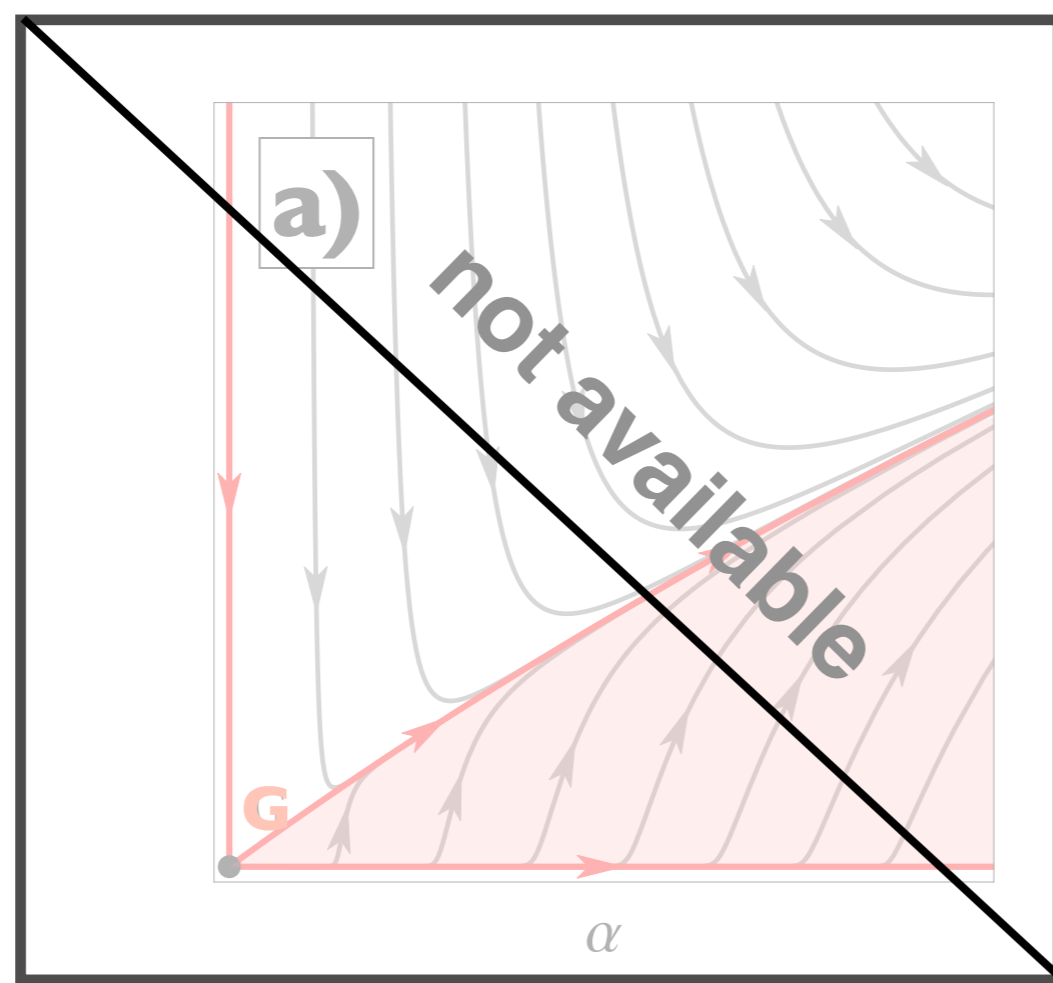
**N=1
SUSY**



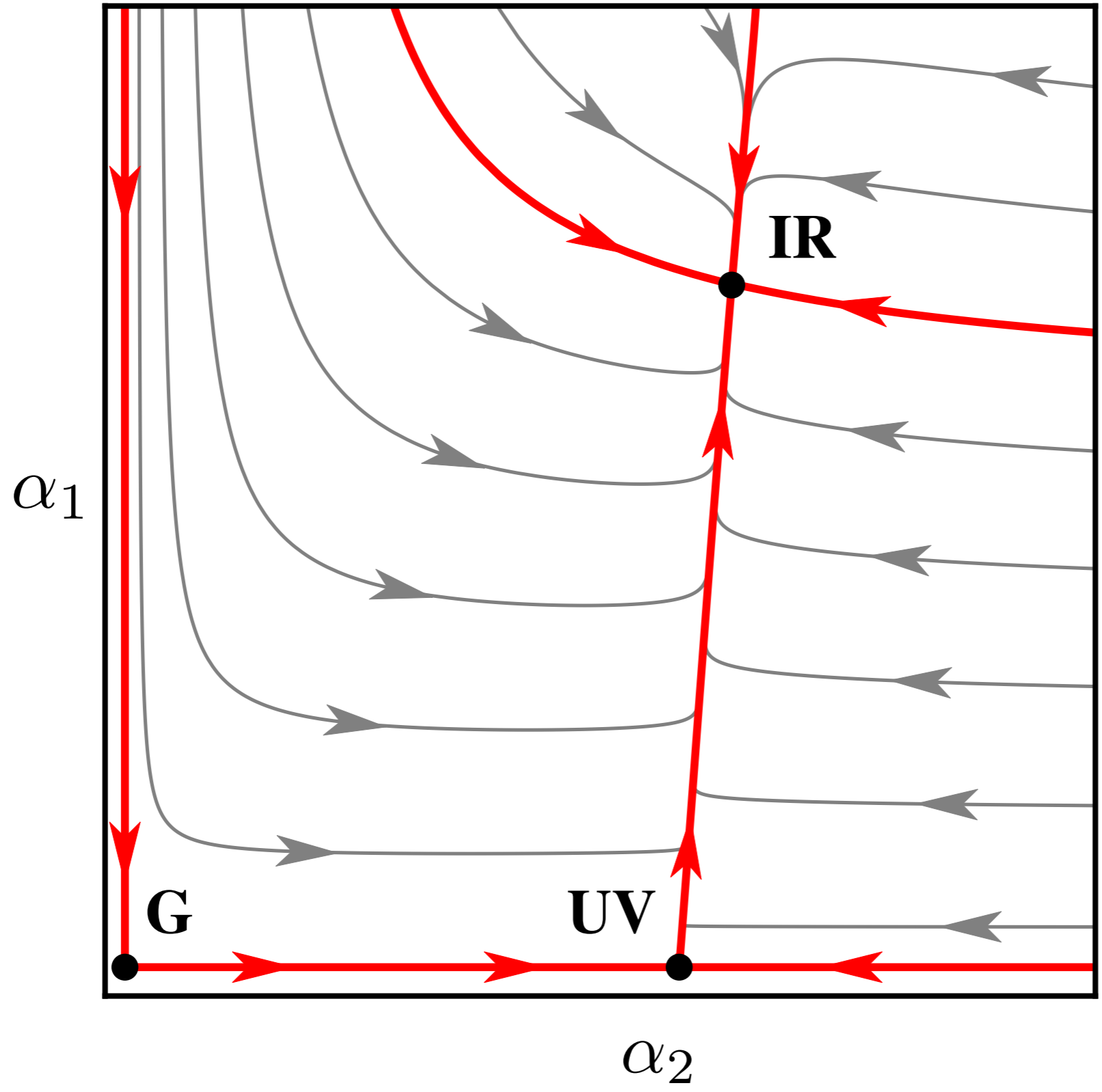
**N=2
SUSY**



**N=4
SUSY**



Susy UV fixed point



SU(N)xSU(M)
+ super-potential

**“Susy
enhances
predictivity”**

gauge coupling

$$\alpha = \frac{g^2}{(4\pi)^2}$$

$$\beta = -B \alpha^2 + C \alpha^3 + \mathcal{O}(\alpha^4)$$

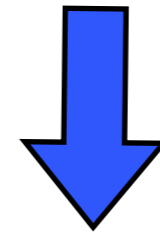
weakly coupled fixed point

$$0 < \alpha^* = B/C \ll 1$$

competition between **matter** and **gauge fields**

why no UV BZ?

$$C = \frac{2}{11} \left[\underbrace{2S_2^F (11C_2^F + 7C_2^G)}_{> 0} + \underbrace{2S_2^S (11C_2^S - C_2^G)}_{> 0} - \underbrace{17B C_2^G}_{> 0} \right]$$

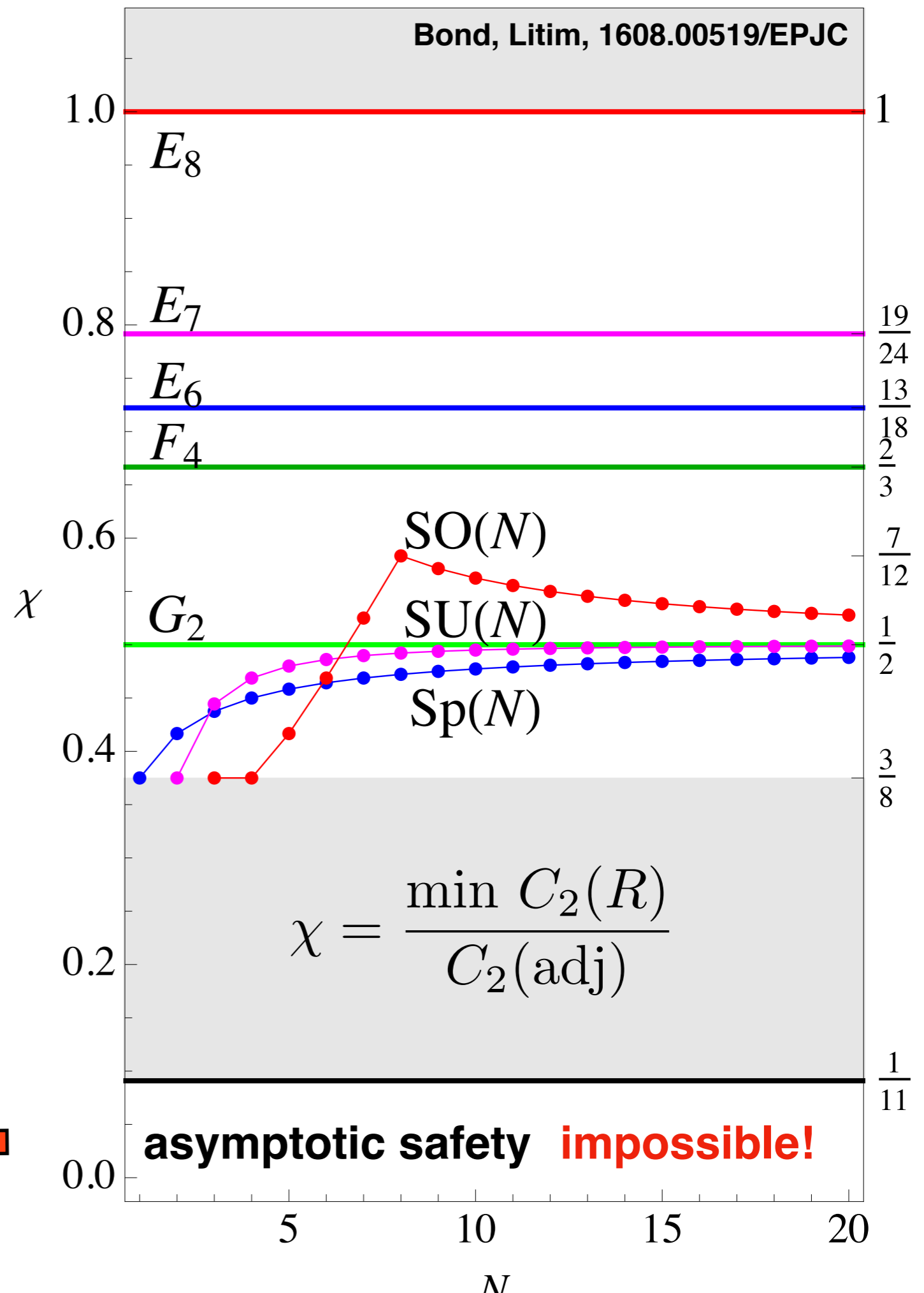
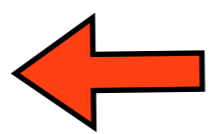


must have

$$C_2^S < \frac{1}{11} C_2^G$$

here's why.

weakly coupled
BZ are never UV



| Case | Condition | Fixed Point |
|--------------|---|--------------|
| <i>i</i>) | $g_i = \mathbf{Y}_{JK}^A = \lambda_{ABCD} = 0$ | Gaussian |
| <i>ii</i>) | some $g_i \neq 0$, all $\mathbf{Y}_{JK}^A = 0$ | Banks-Zaks |
| <i>iii</i>) | some $g_i \neq 0$, some $\mathbf{Y}_{JK}^A \neq 0$ | gauge-Yukawa |

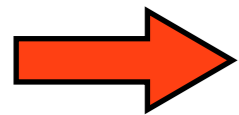
asymptotic safety requires **all types** of matter fields

Yukawa couplings are key

works with **supersymmetry**

implications

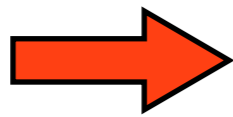
- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields



asymptotic **freedom** and asymptotic **safety**
are **two sides of one and the same medal**

implications

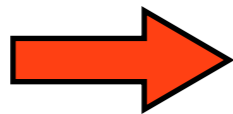
- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields



new BSM territory to UV-complete the Standard Model

implications

- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields
- statistical physics: **universality class** of any weakly coupled 4D critical point contains non-abelian gauge fields



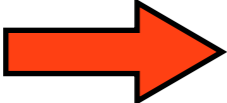
systematic classification of weak critical points in 4D

implications

- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields
- statistical physics: **universality class** of any weakly coupled 4D critical point contains non-abelian gauge fields
- conformal field theory: **any** weakly-coupled **4D CFT** contains non-abelian gauge fields

“any QFT under perturbative control in the deep UV or IR asymptotes to a conformal field theory”

implications

- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields
 - statistical physics: **universality class** of any weakly coupled 4D critical point contains non-abelian gauge fields
 - conformal field theory: **any** weakly-coupled **4D CFT** contains non-abelian gauge fields
-  QFT offers infinitely many 4D CFTs
access to CFT data
complementary to conformal bootstrap

implications

- particle physics: **UV complete 4D theories** (free or safe) require non-abelian gauge fields
- statistical physics: **universality class** of any weakly coupled 4D critical point contains non-abelian gauge fields
- conformal field theory: **any** weakly-coupled **4D CFT** contains non-abelian gauge fields

asymptotic safety BSM

A Bond, G Hiller, K Kowalska, DF Litim, **Directions for model building from asymptotic safety**, JHEP1708 (2017) 004
G Hiller, C Hermigos-Feliu, DF Litim, T Steudtner, **Asymptotically safe extensions of the Standard Model and their flavour phenomenology**, Moriond (EW2019) 1905.11020

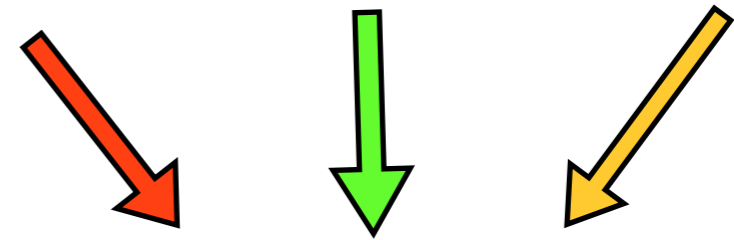
asymptotic safety beyond the SM

minimal framework:

AD Bond, G Hiller, K Kowalska, DF Litim, 1702.01727 (JHEP)

SM gauge symmetry

$$SU(3)_C \times SU(2)_L \times U(1)_Y$$



N_F flavors of BSM fermions

$$\psi_i(R_3, R_2, Y)$$

BSM singlet scalars

$$S_{ij}$$

features: vector-like fermions

global flavor symmetry $U(N_F) \times U(N_F)$

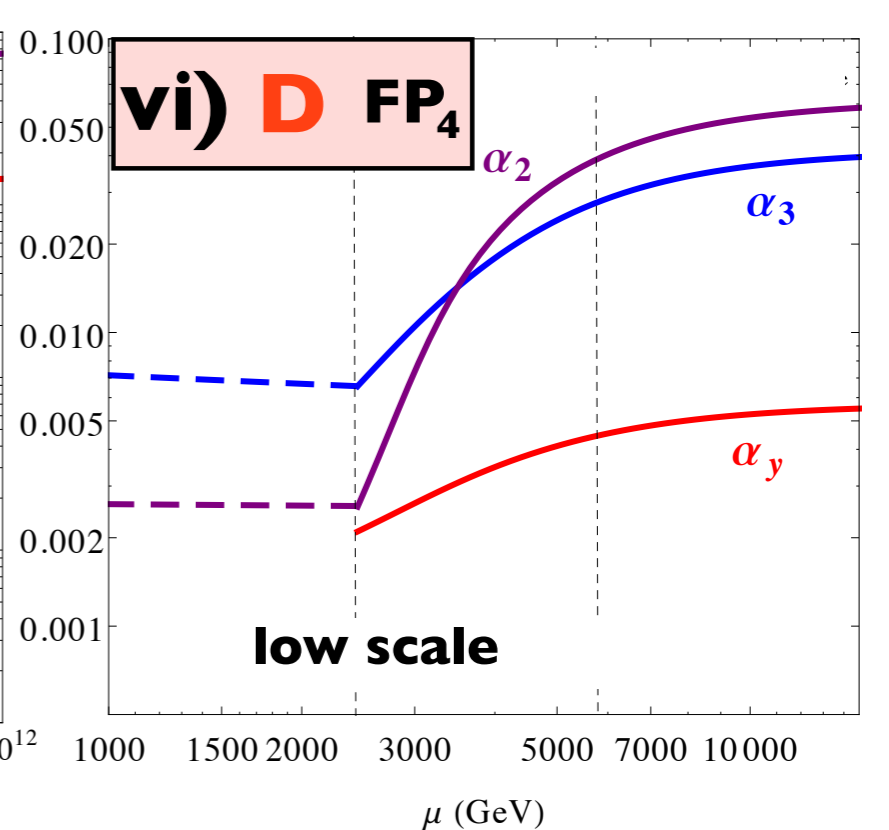
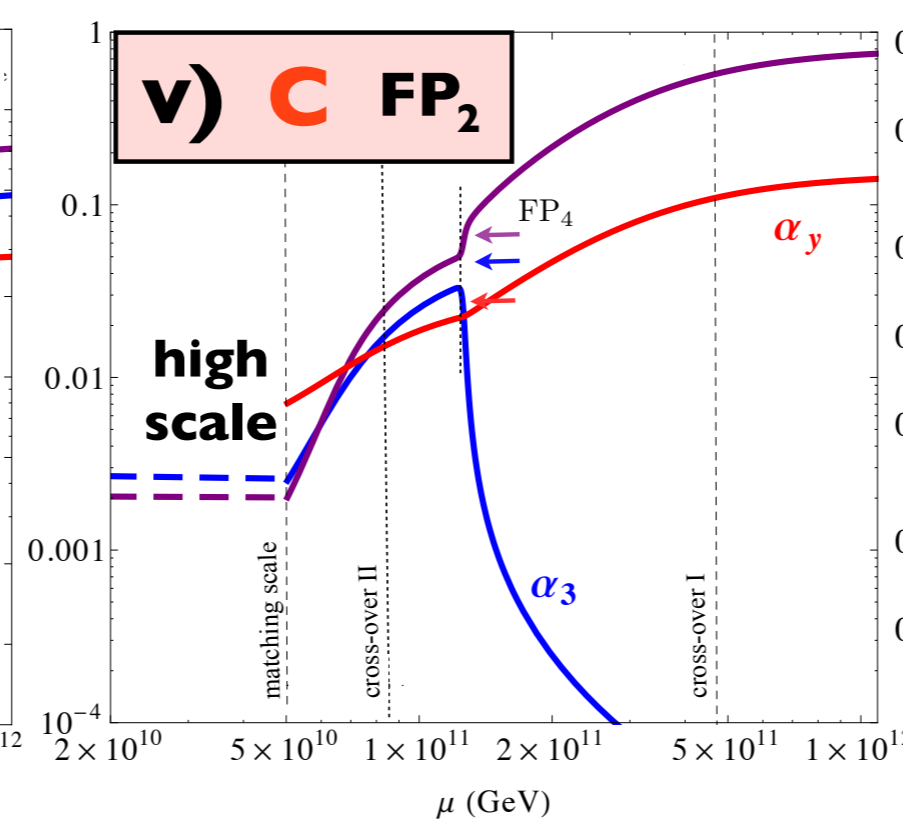
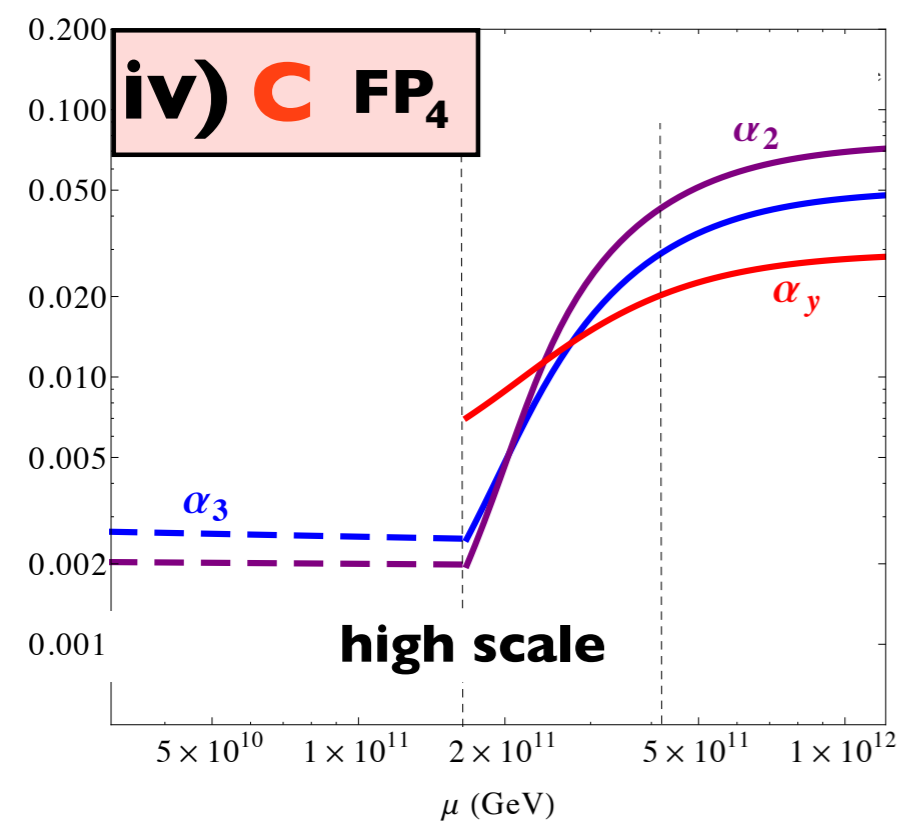
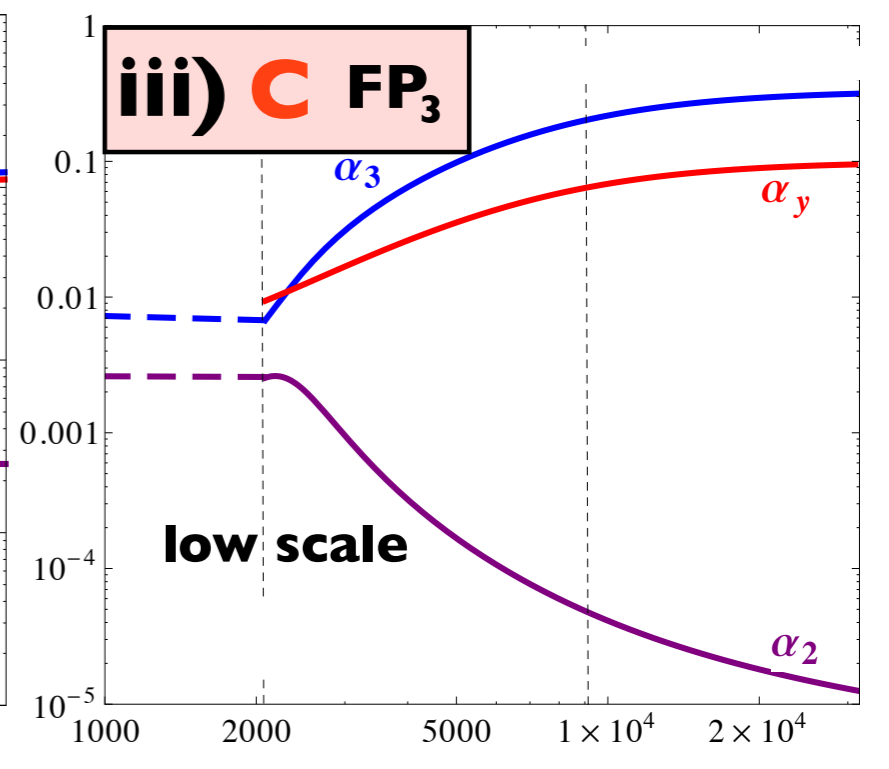
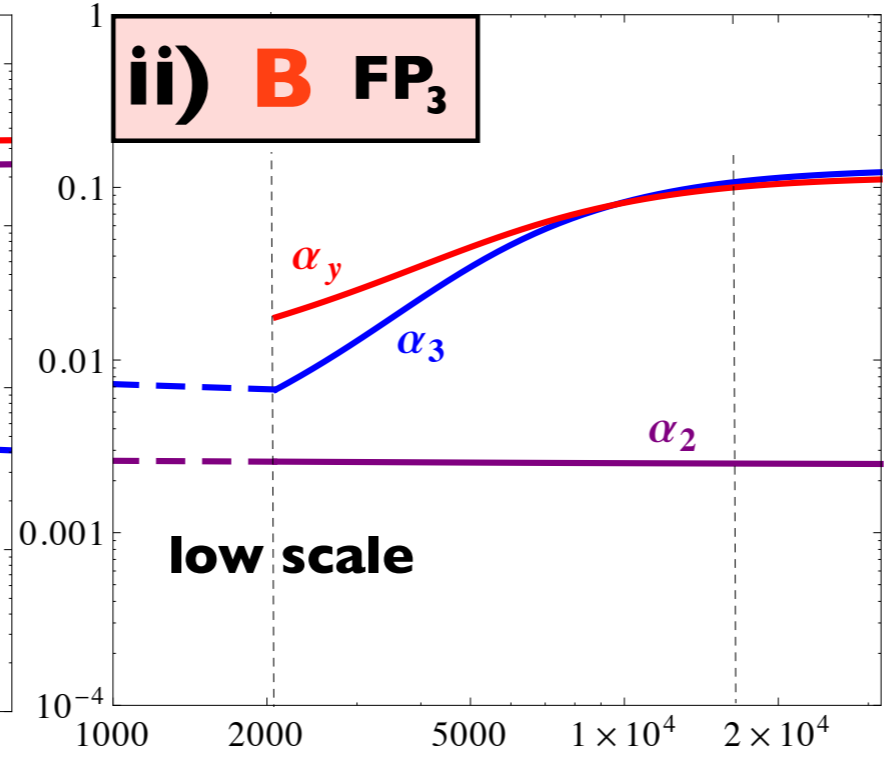
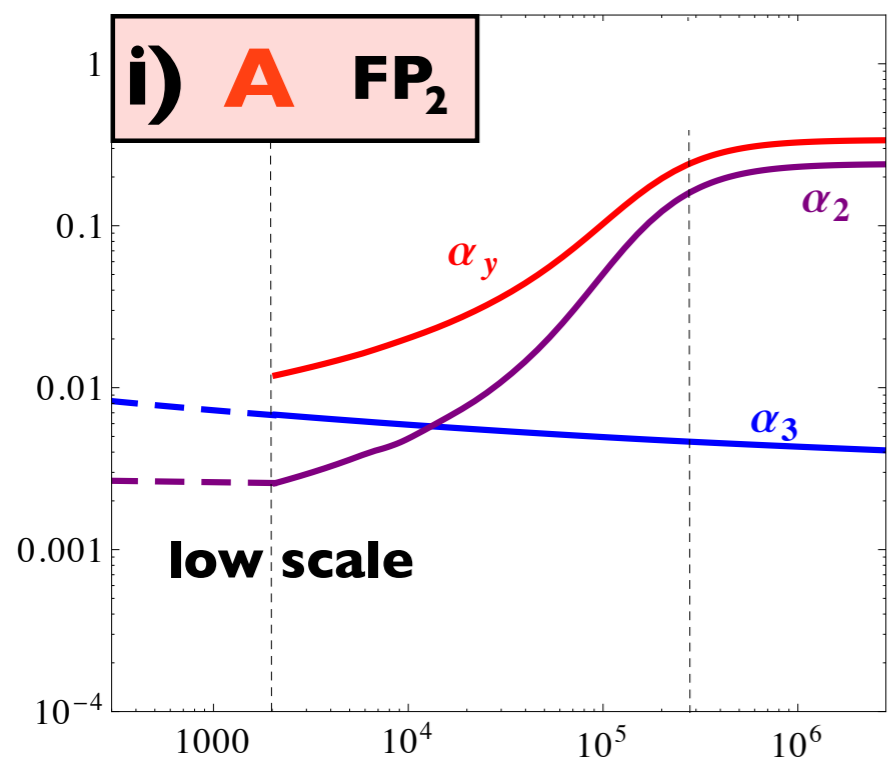
single BSM Yukawa coupling

BSM matter charged under U(1)Y

| model | parameter (R_3, R_2, N_F) | UV fixed points | | | AF for $U(1)_Y$ | info |
|-------|----------------------------------|-----------------|--------------|--------------|--------------------|-------------------|
| | | α_3^* | α_2^* | α_y^* | | |
| A | (1, 4, 12) | 0 | 0.2407 | 0.3385 | $Y > 0.228$ | FP ₂ ● |
| B | (10, 1, 30) | 0.1287 | 0 | 0.1158 | $Y > 0.107$ | FP ₃ ■ |
| | | 0.1292 | 0.2769 | 0.1163 | $Y > 0.114$ | FP ₄ ◆ |
| C | (10, 4, 80) | 0.3317 | 0 | 0.0995 | $Y > 0.024$ | FP ₃ ■ |
| | | 0.0503 | 0.0752 | 0.0292 | $Y > 0.050$ | FP ₄ ◆ |
| D | (3, 4, 290) | 0 | 0.8002 | 0.1500 | $Y > 0.018$ | FP ₂ ● |
| | | 0.0416 | 0.0895 | 0.0066 | $Y > 0.042$ | FP ₂ ● |
| E | (3, 3, 72) | 0.0416 | 0.0615 | 0.0056 | $Y > 0.052$ | FP ₄ ◆ |
| | | 0.1499 | 0.2181 | 0.0471 | $Y > 0.073$ | FP ₄ ◆ |

lower bounds
on hypercharge

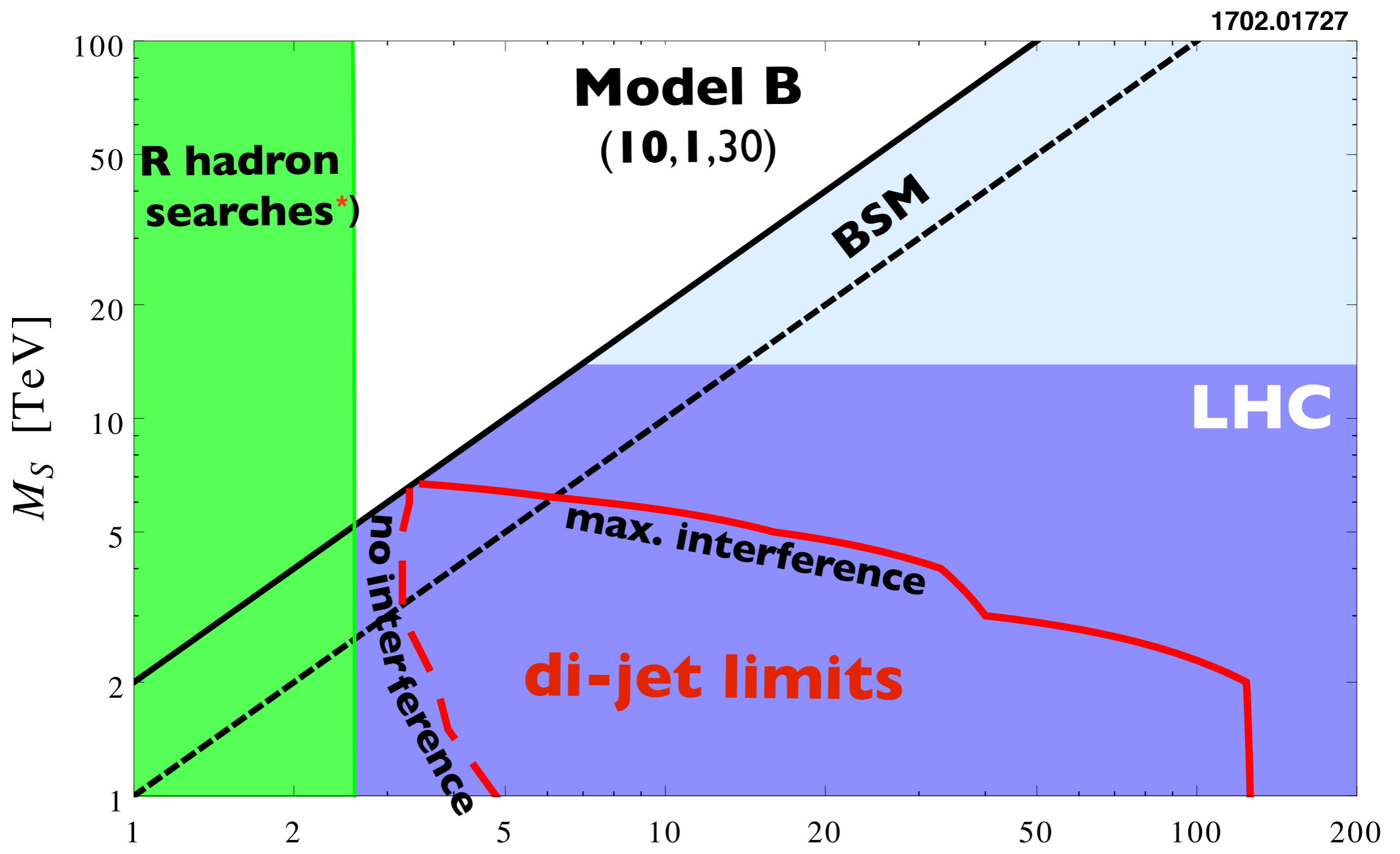
UV fixed points



lower limits on the mass of **long-lived colored particles** that form QCD bound states; exclusion data from ATLAS & CMS (pair production)

| $\psi(R_3, R_2)$ | $R_2 = 1$ | | $R_2 = 2$ | | $R_2 = 3$ | |
|------------------|-------------------|-----------------------|-------------------|-----------------------|-----------|-----------------------|
| R_3 | C_3 | M_ψ^{\min} (TeV) | C_3 | M_ψ^{\min} (TeV) | C_3 | M_ψ^{\min} (TeV) |
| 3 | $5\frac{1}{3}$ | (1.3) | $10\frac{2}{3}$ | (1.4) | 16 | 1.5 |
| 6 | $66\frac{2}{3}$ | 1.7 | $133\frac{1}{3}$ | 1.8 | 200 | 1.9 |
| 8 | 72 | 1.7 | 144 | 1.8 | 216 | 1.9 |
| 10 | 360 | 2.0 | 720 | 2.1 | 1080 | 2.2 |
| 15 | $426\frac{2}{3}$ | 2.0 | $853\frac{1}{3}$ | 2.1 | 1280 | 2.2 |
| 15' | $1306\frac{2}{3}$ | 2.2 | $2313\frac{1}{3}$ | 2.3 | 3920 | 2.4 |

mass exclusion limits



*) fudged from 13 TeV
ATLAS + CMS gluino analysis

M_ψ [TeV]

fixed points in quantum gravity

gravitation

physics of classical gravity

Einstein's theory of general relativity

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -\Lambda g_{\mu\nu} + 8\pi G_N T_{\mu\nu}$$

Newton's coupling

$$G_N = 6.7 \times 10^{-11} \frac{\text{m}^3}{\text{kg s}^3}$$

cosmological constant

$$\Lambda \approx 10^{-35} \text{s}^{-2}$$

what's new with gravity?

degrees of freedom: **spin 2**

perturbatively non-renormalisable

Newton's coupling is **dimensionful** $[G_N] = 2 - D < 0$

asymptotic safety requires **large**

anomalous dimensions

non-perturbative tools mandatory

Einstein-Hilbert

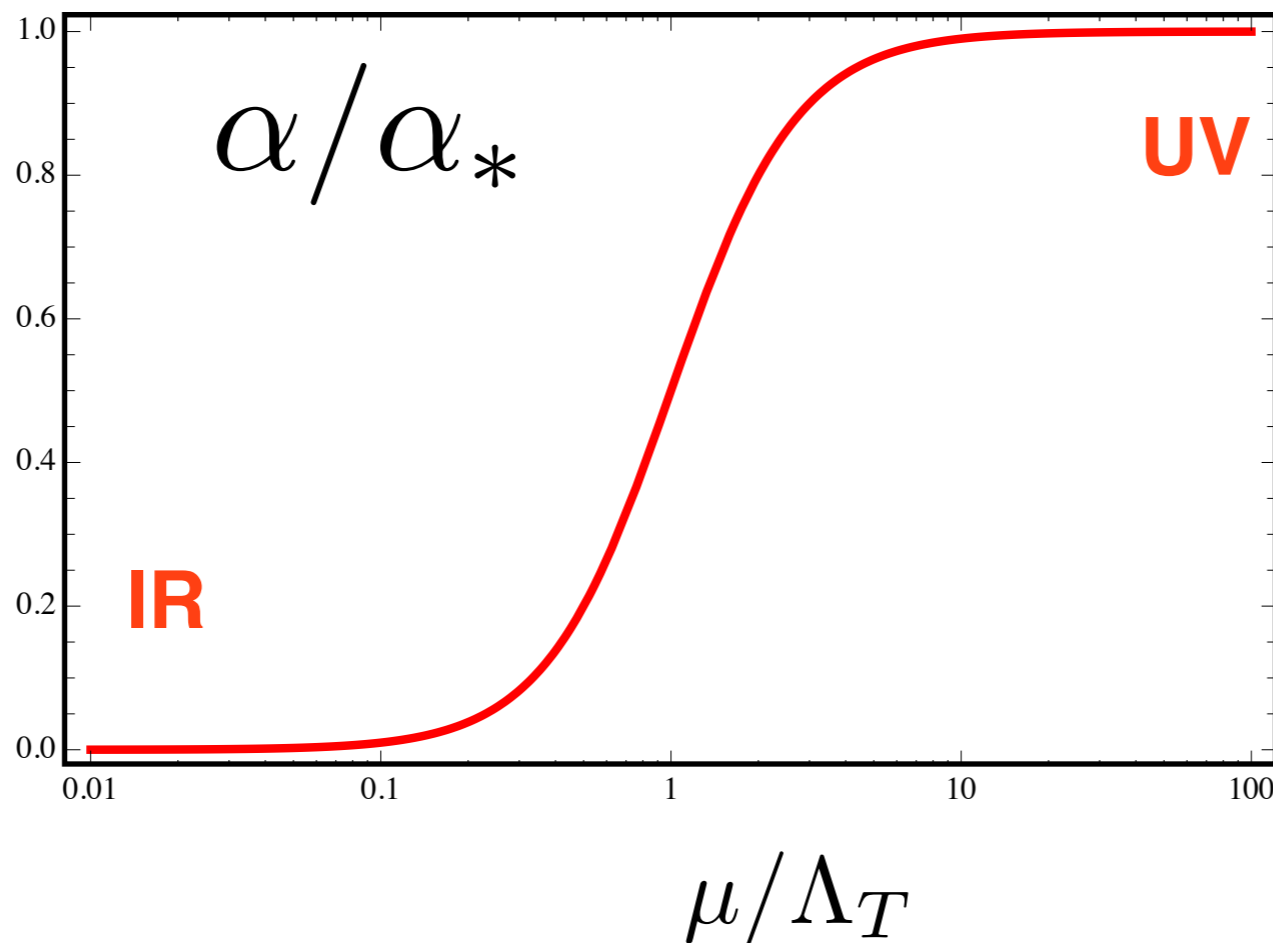
gravitons

dimension

coupling

$$D = 2 + \epsilon : \quad \alpha = G_N(\mu) \mu^{D-2}$$

Gastmans et al '78
Christensen, Duff '78
Weinberg '79
Kawai et al '90



$$G(\mu) \approx \frac{\alpha_*}{\mu^{D-2}}$$

quantum GR

**UV fixed point
implies
weakened gravity**

$G(\mu) \approx G_N$
classical GR

Ricci scalars $\Gamma_k \propto f(R)$

$$\Gamma_k = \int d^4x \sqrt{\det g_{\mu\nu}} \frac{1}{16\pi G} [-R + 2\Lambda] + \sum_{n=2}^{N-1} \lambda_n R^n$$

effective action with
invariants up to mass
dimension $D = 2(N - 1)$

Ricci scalars $\Gamma_k \propto f(R)$

$$\Gamma_k = \int d^4x \sqrt{\det g_{\mu\nu}} \frac{1}{16\pi G} [-R + 2\Lambda] + \sum_{n=2}^{N-1} \lambda_n R^n$$

effective action with
invariants up to mass
dimension $D = 2(N - 1)$

bootstrap search strategy

Falls, DL, Nikolakopoulos, Rahmede '13, '14

- 1 fix **N**, compute RG flow
- 2 deduce fixed point and exponents
- 3 increase **N** to **N+1** and start over at **1**

Ricci scalars $\Gamma_k \propto f(R)$

$$\Gamma_k = \int d^4x \sqrt{\det g_{\mu\nu}} \frac{1}{16\pi G} [-R + 2\Lambda] + \sum_{n=2}^{N-1} \lambda_n R^n$$

effective action with
invariants up to mass
dimension $D = 2(N - 1)$

up to order **$N = 2$** Souma, '99, Reuter, Lauscher '01, Litim '03

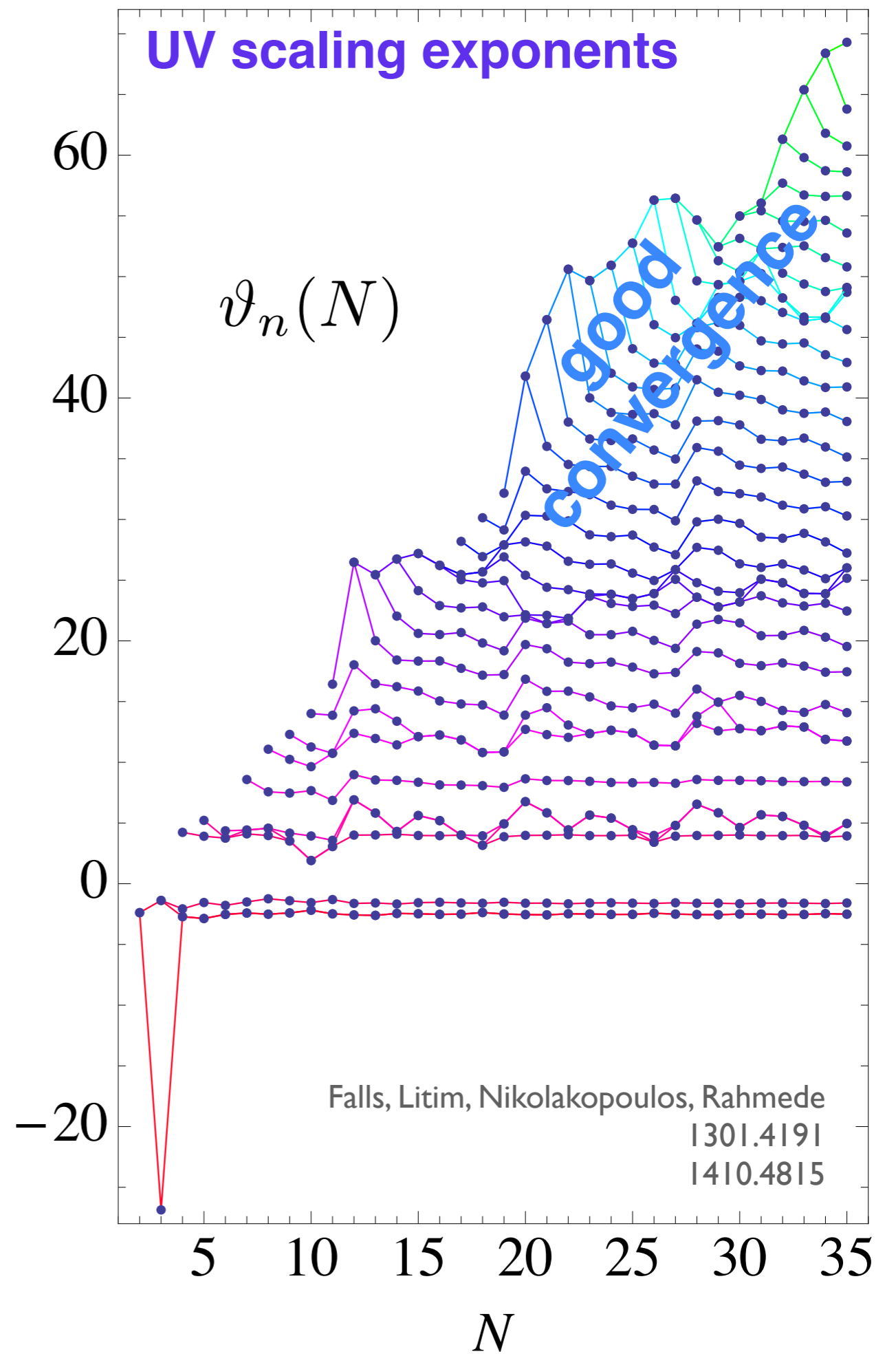
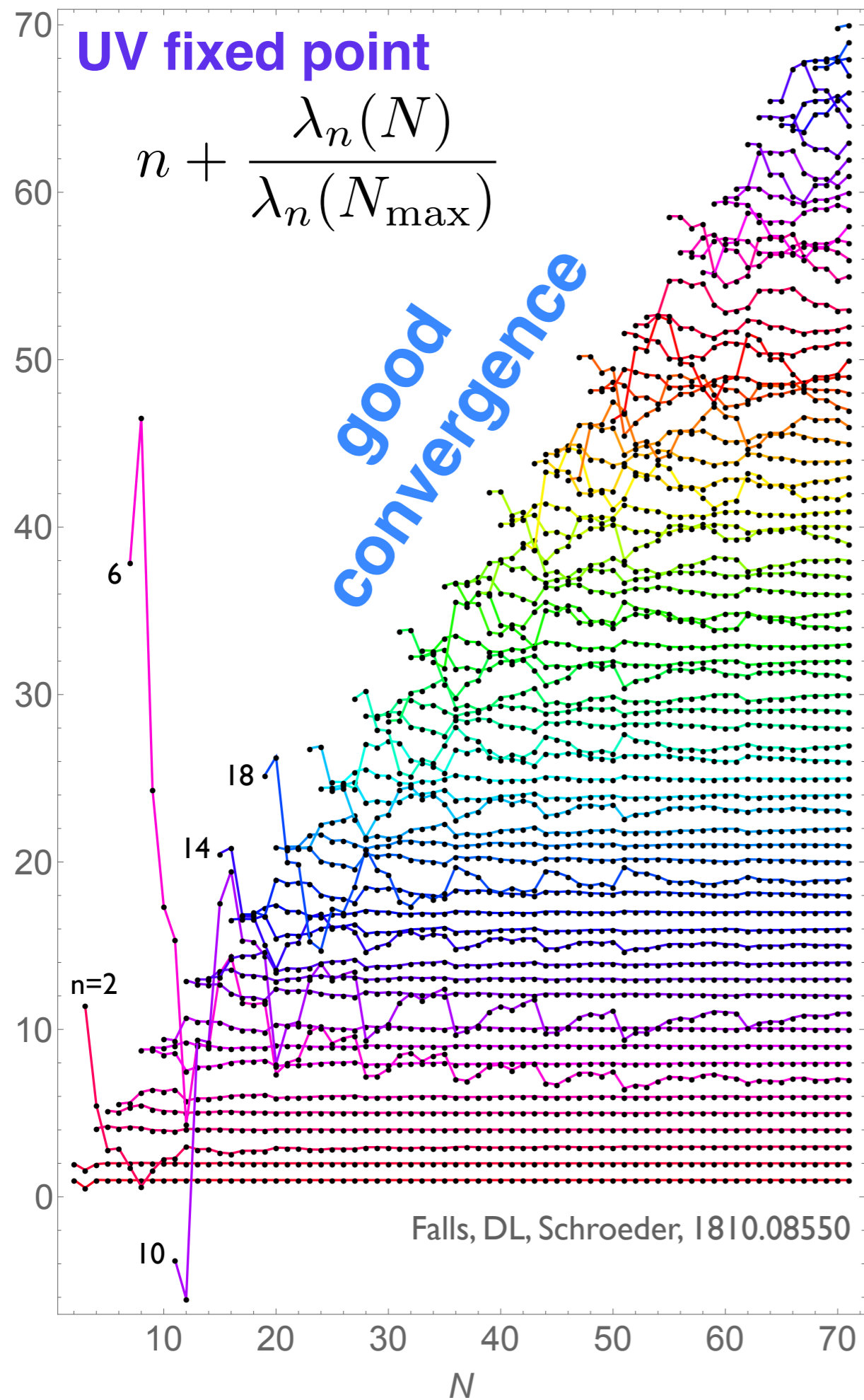
$N = 3$ Reuter, Lauscher '01

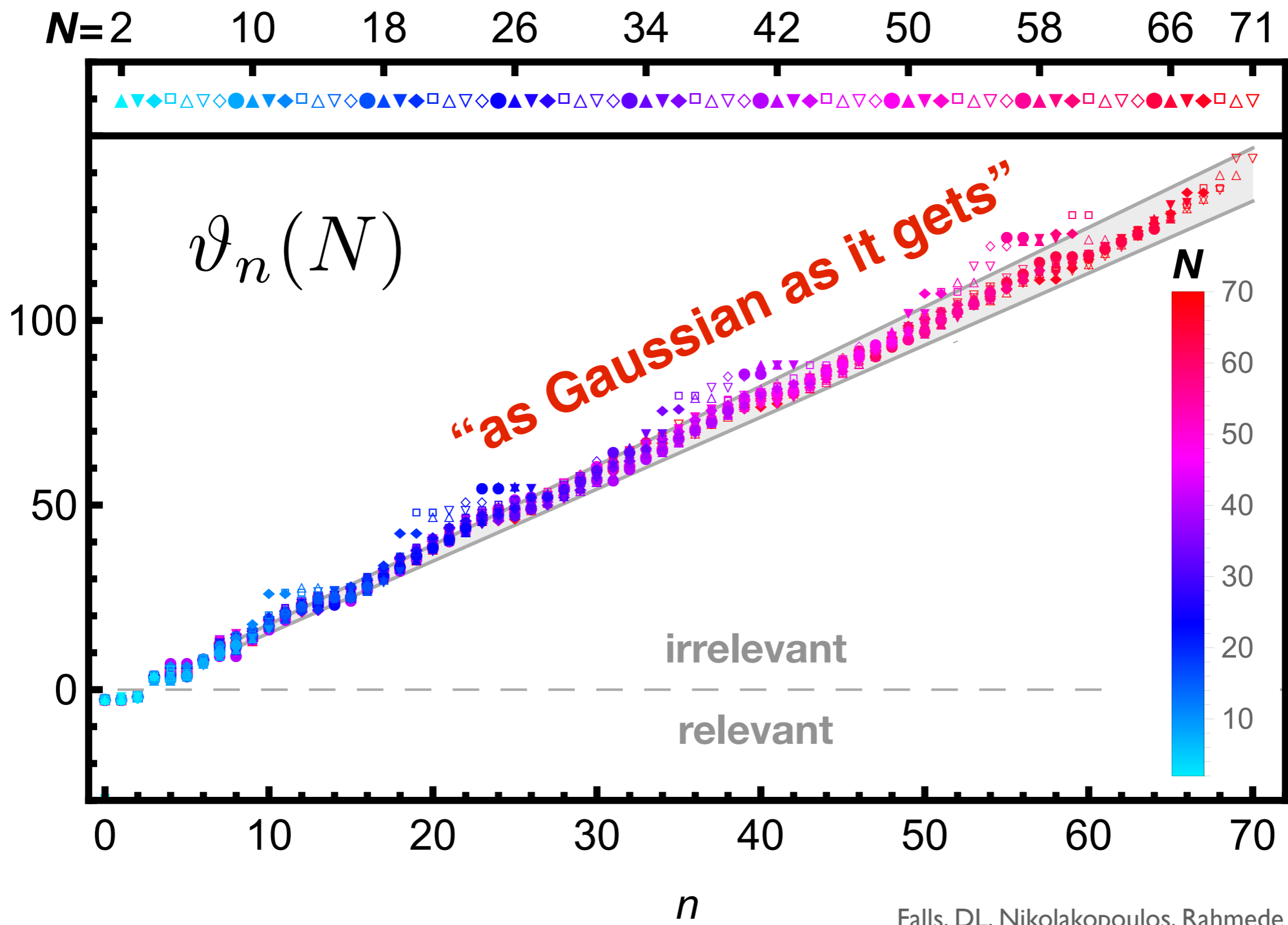
$N = 7$ Codello, Percacci, Rahmede '07

$N = 11$ Bonanno, Contillo, Percacci '10

$N = 35$ Falls, Litim, Nikolakopoulos, Rahmede '13, '14, '16

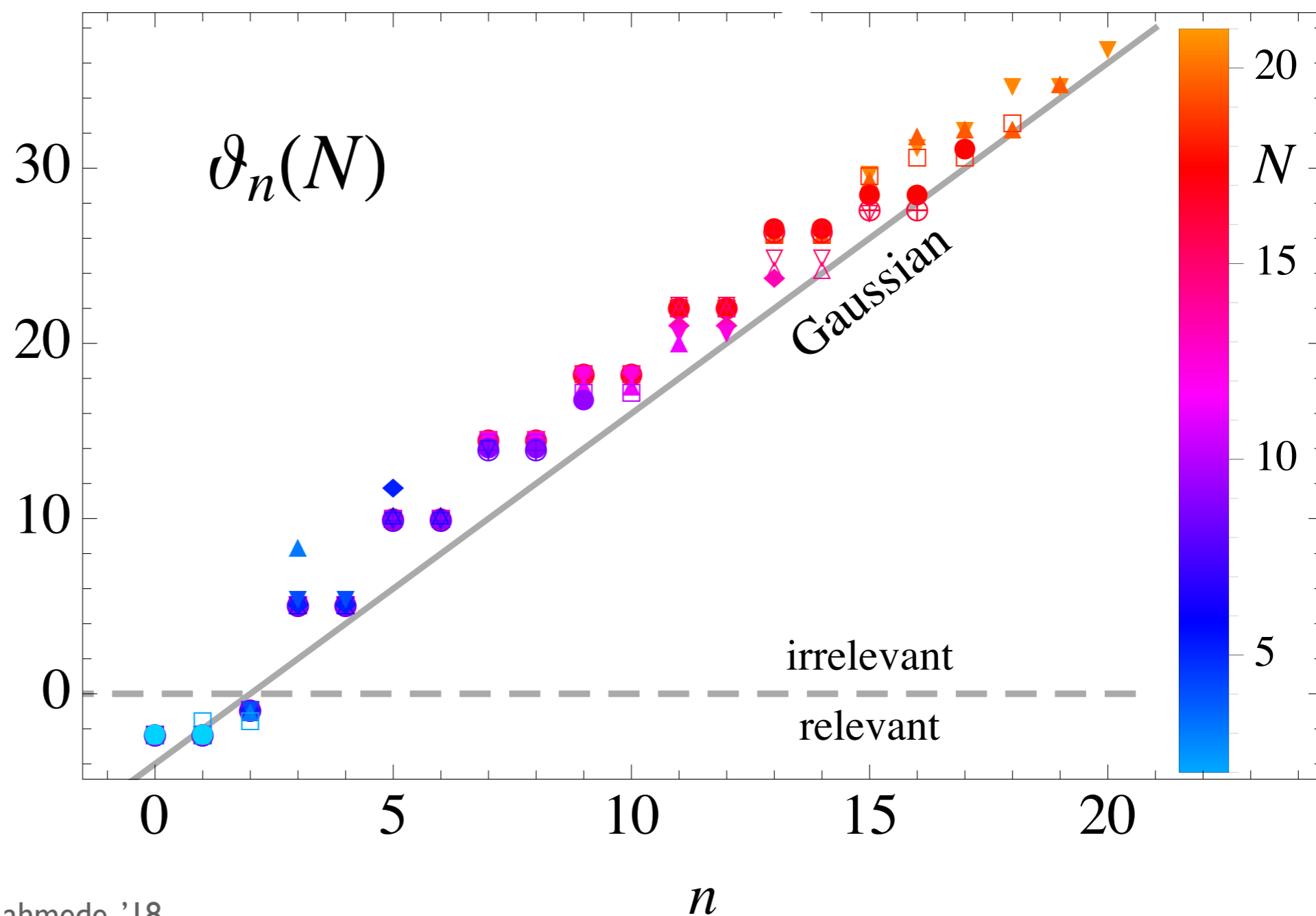
$N = 71$ Falls, Litim, Schroeder '18





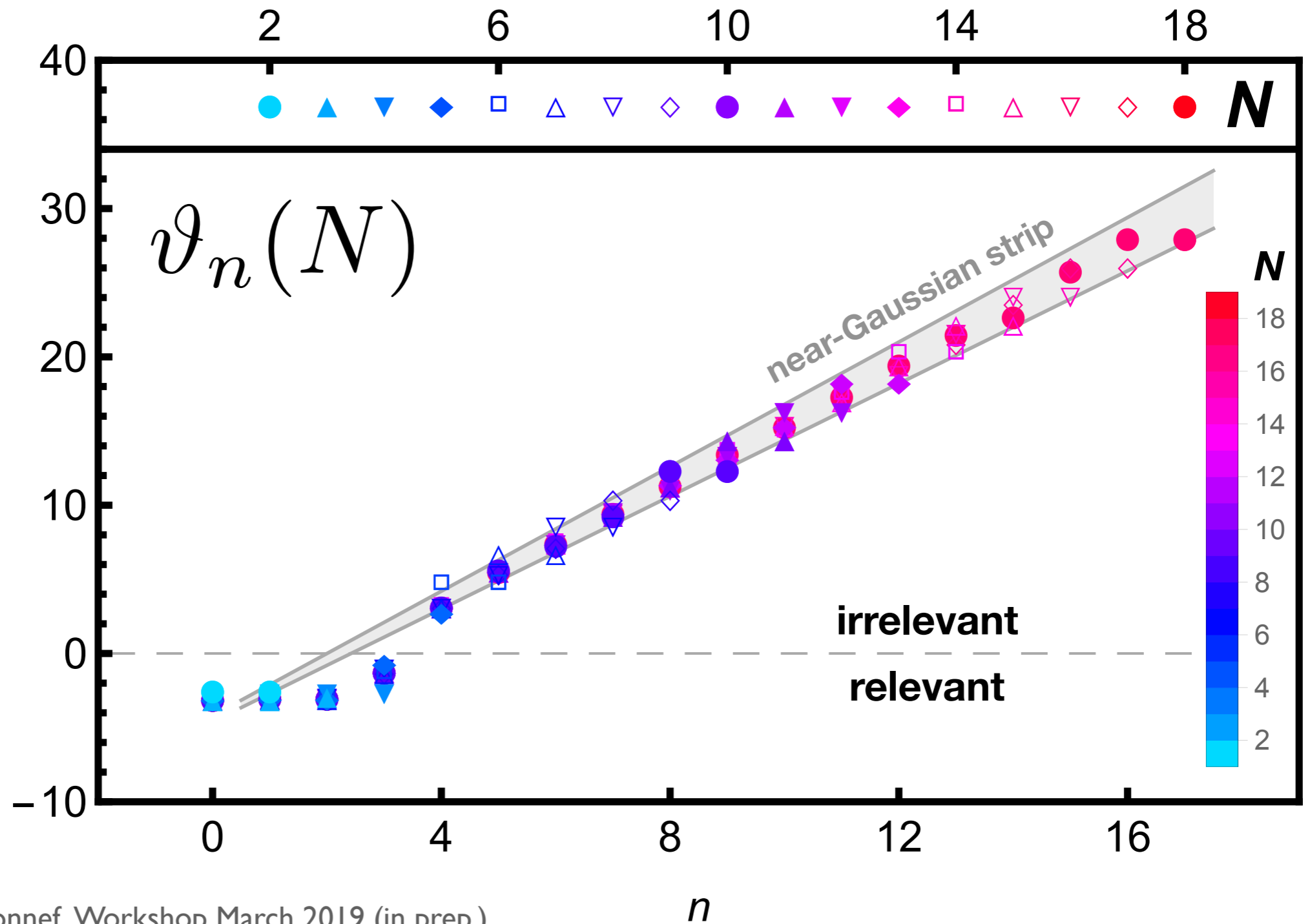
Ricci

$$\Gamma_k = \int d^d x \sqrt{g} [F_k(\text{Ric}^2) + R \cdot Z_k(\text{Ric}^2)]$$



Riemann

$$\Gamma_k = \int d^d x \sqrt{g} [F_k(\text{Riem}^2) + R \cdot Z_k(\text{Riem}^2)]$$

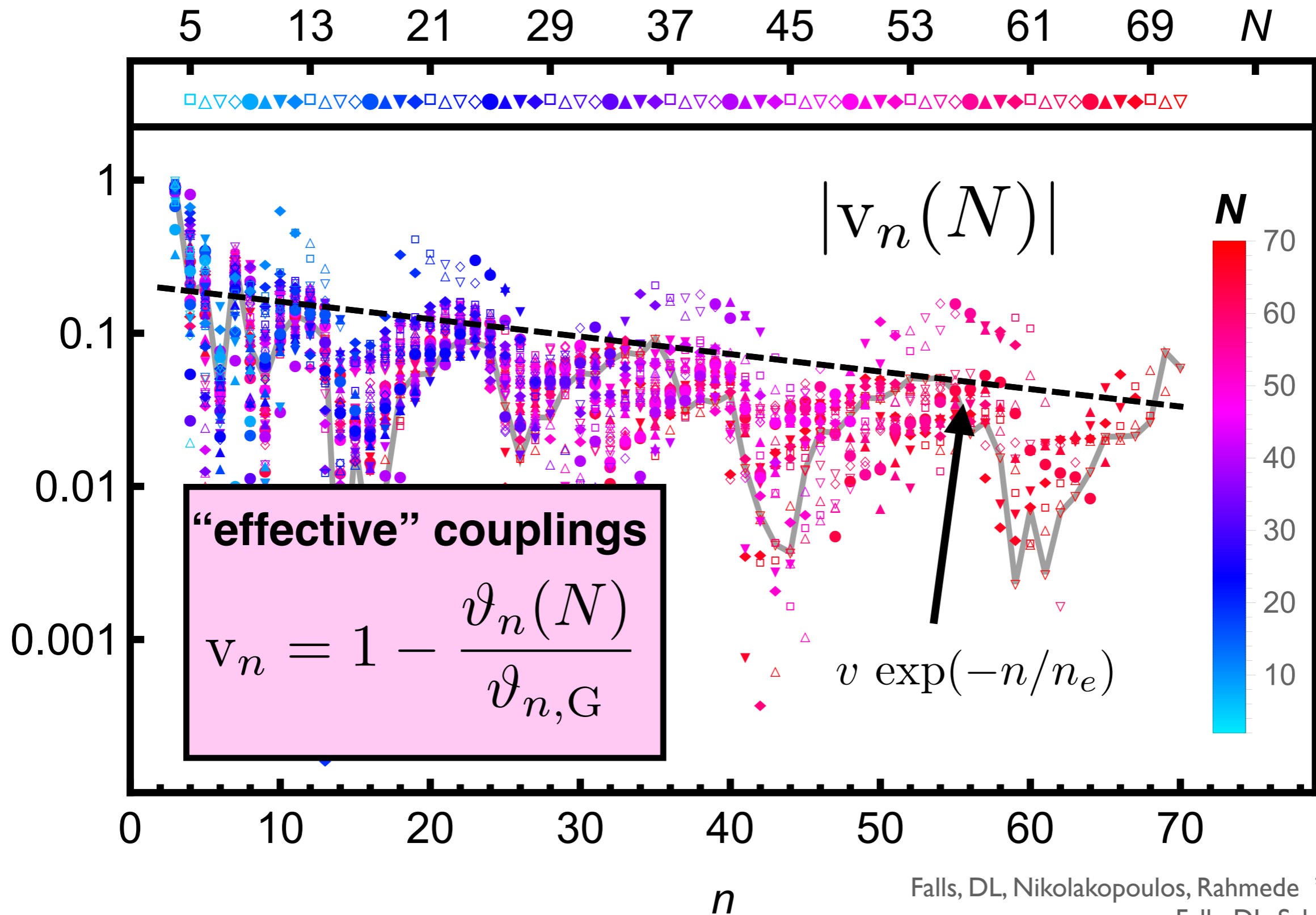


near-Gaussian scaling



signature of weak coupling

weak “effective” coupling



conclusions

complete
understanding of
asymptotic safety
at weak coupling

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works for
simple, semi-simple
and **supersymmetric**
gauge theories,
large variety of
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w or w/o gravity

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