

Displaced Vertices and long-lived charged particles in the NMSSM with Right-Handed Sneutrino

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MSSM: μ Problem

$$W_{\text{MSSM}} \supset -\mu H_1 \cdot H_2$$

After REWSB takes place μ has to be of the order of the EW scale

→ μ Problem *Kim, Nilles '84*

Solution

Introduction of a singlet S

$$W_{\text{NMSSM}} \supset -\lambda S H_1 \cdot H_2 + \frac{1}{3} \kappa S^3$$

After REWSB takes place S acquires a vev v_s and originates $\mu = \lambda v_s$ that is of the order of the EW scale.

Higgs mass in the NMSSM

$$m_h^{2,tree} = m_Z^2 \cos^2 2\beta + \lambda^2 v^2 \sin^2 2\beta,$$

New contribution at the tree-level Higgs mass, so no large loop corrections are needed. Same case as adding triplets. *See Garcia-Pepin's, Nardini's and di Chiara's talks.*

Superpotential

Introduction of a supermultiplet N

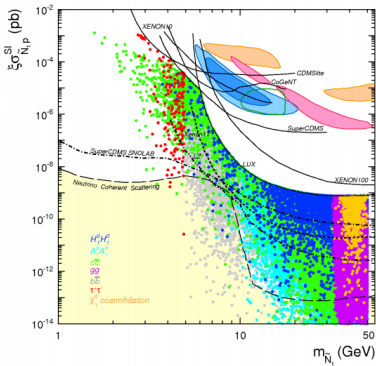
Kitano, Oda '00, Garbrecht, Pallis, Pilaftsis '06 Cerdeno, Munoz, Seto '08 Cerdeno, Seto '09

$$W = W_{\text{NMSSM}} + \lambda_N S N N + y_N L \cdot H_2 N$$

- RH sneutrino: Good dark matter candidate
- RH neutrino: Interesting phenomenology. Low scale see-saw, displaced vertices, long-lived NLSP.

Parameters

5 new parameters: $m_{\tilde{N}}$, λ_N , y_N , A_{λ_N} , A_{y_N}



Sneutrino Dark Matter

- Good Dark Matter candidate
- It is widely extended in the $(\sigma_{\tilde{N}p}^{SI}, m_{\tilde{N}})$ plane.
- It can explain the little excess in gamma rays coming from the Galactic Centre.

Cerdeno, Seto '08

Cerdeno, Seto '09

Cerdeno, Peiro, Seto '11

Cerdeno, Peiro, Robles '14

Mass Matrix

After REWSB takes place the term $\lambda_N S N N$ originates a Majorana mass for N , $M_N = 2\lambda_N v_s$. This term and $y_N L \cdot H_2 N$ provide a neutrino mass matrix:

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & \mathbf{y}_N v_2 \\ \mathbf{y}_N^T v_2 & 2\lambda_N v_s \end{pmatrix}$$

See-saw

In order to obtain a see-saw mechanism we assume $y_N v_2 \ll \lambda_N v_s$, so the masses of the physical states are

$$m_{\nu_1} = \frac{(y_N v_2)^2}{2\lambda_N v_s} \quad m_{\nu_2} = 2\lambda_N v_s$$

Neutrino Yukawa Coupling

To reproduce the tiny neutrino masses $y_N \sim \mathcal{O}(10^{-6})$

Mixing of states

We can write the physical states as a function of the gauge states

$$\nu_i = N_{iL}^\nu \nu_L + N_{iR}^\nu N$$

Neglecting terms of the order $\mathcal{O}(y_N^2)$ and beyond

$$\nu_1 = \nu_L + N_{1R}^\nu N, \quad \nu_2 = N_{2L}^\nu \nu_L + N$$

$$N_{1R}^\nu = -N_{2L}^\nu = \frac{y_N v_2}{2\lambda_N v_s}$$

Heavy Neutrino Decay

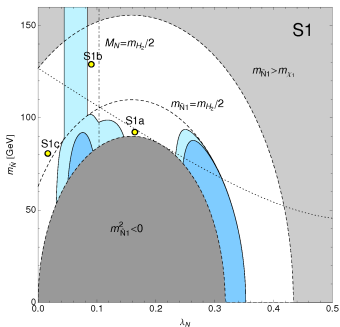
There exists a mixing that makes the heavy neutrino decay.

	S1	S2	S3
$\tan \beta$	2.0	2.5	2.7
(M_1, M_2, M_3)	(500, 650, 1950) GeV	(300, 600, 1800) GeV	(345, 575, 2500) GeV
μ	152 GeV	180 GeV	595 GeV
(λ, κ)	(0.50, 0.27)	(0.60, 0.40)	(0.58, 0.34)
(A_λ, A_κ)	(283, -220) GeV	(265, -50) GeV	(1189, -225) GeV
$m_{H_1^0}$	99.5 GeV	125.7 GeV	125.8 GeV
$m_{H_2^0}$	125.8 GeV	225.7 GeV	656.9 GeV

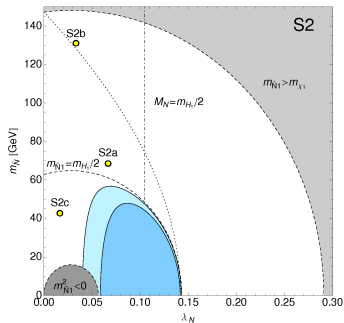
Squarks around 2 TeV

Spectrum computed by `NMSSMTools` → Loop corrections, low energy observables, LEP constraints...

Displaced Vertices: Benchmark points

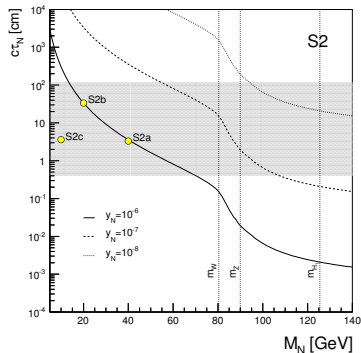
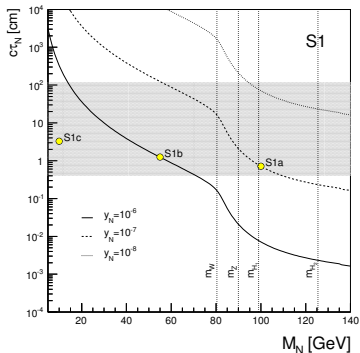


S1a: $m_{\tilde{N}_1} = 20$ GeV, $M_N = 100$ GeV
S1b: $m_{\tilde{N}_1} = 100$ GeV, $M_N = 55$ GeV
S1c: $m_{\tilde{N}_1} = 70$ GeV, $M_N = 10$ GeV



S2a: $m_{\tilde{N}_1} = 40$ GeV, $M_N = 40$ GeV
S2b: $m_{\tilde{N}_1} = 20$ GeV, $M_N = 20$ GeV
S2c: $m_{\tilde{N}_1} = 10$ GeV, $M_N = 10$ GeV

Displaced Vertices: RH Neutrino decays



Process	Signature
$N \rightarrow W^\pm l_i^\mp \rightarrow \nu_j l_j^\pm l_i^\mp$	$2\ell(+\cancel{E}_T)$
$\quad \quad \quad \rightarrow q\bar{q}' l_i^\mp$	ℓjj
$N \rightarrow Z\nu_i \rightarrow \nu_i l_j^\pm l_j^\mp$	$2\ell(+\cancel{E}_T)$
$\quad \quad \quad \rightarrow \nu_i q\bar{q}$	$2j(+\cancel{E}_T)$
$N \rightarrow H_i^0 \nu_i \rightarrow \nu_i l_j^\pm l_j^\mp$	$2\ell(+\cancel{E}_T)$
$\quad \quad \quad \rightarrow \nu_i q\bar{q}$	$2j(+\cancel{E}_T)$
$\quad \quad \quad \rightarrow \nu_i \gamma\gamma$	$2\gamma(+\cancel{E}_T)$

Caution

\cancel{E}_T is not a good variable for displaced vertices

Information about the mass

m_{jjl}^{inv} from the same displaced vertex \rightarrow Peak mass around M_N

m_{ll}^{inv} from the same displaced vertex \rightarrow Endpoint variable:

- If $M_N < m_{W,Z}$: $(m_{ll}^{inv})^2 \simeq m_N^2$

- If $M_N > m_{W,Z}$: $(m_{ll}^{inv})^2 = m_N^2 - m_{W,Z}^2$

MC Event Simulation

We simulate events with CalcHEP 3.4.

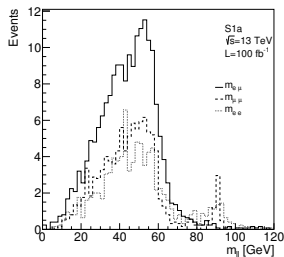
We simulate detector effects

$$\frac{\sigma}{E} = \frac{a}{\sqrt{E/\text{GeV}}} \oplus b$$

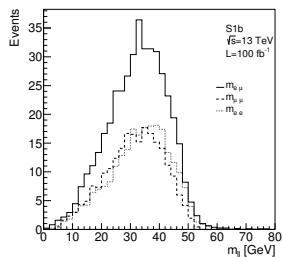
Cuts

- $10 \text{ mm} < c\tau < 100 \text{ mm}$
- Electrons $p_T > 10 \text{ GeV}$ and $|\eta_e| < 2.5$
- Muons $p_T > 6 \text{ GeV}$ and $|\eta_\mu| < 2.5$
- Jet $p_T > 15 \text{ GeV}$ and $|\eta_j| < 2.5$
- Isolated object $\Delta R > 0.4$ ($\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$)

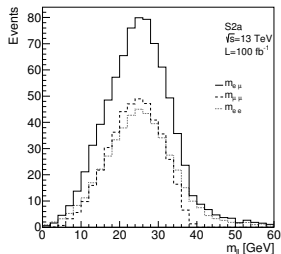
Displaced Vertices: Signatures at the LHC



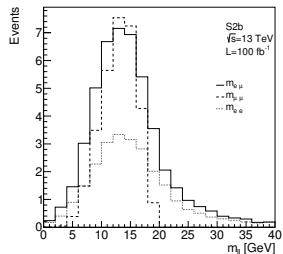
$M_N = 100$ GeV



$M_N = 55$ GeV

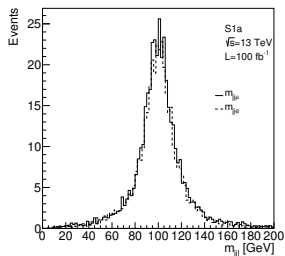


$M_N = 40$ GeV

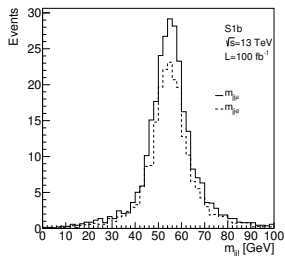


$M_N = 20$ GeV

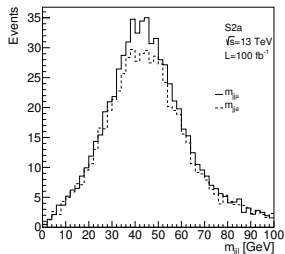
Displaced Vertices: Signals at the LHC



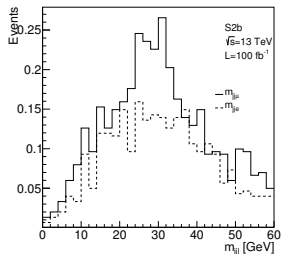
$M_N = 100$ GeV



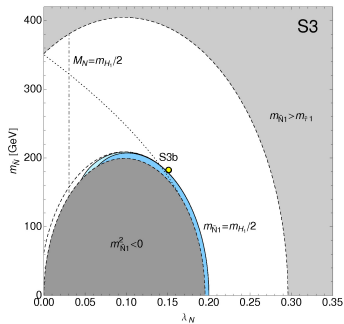
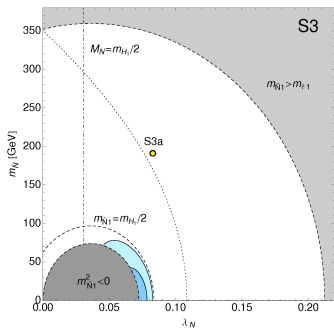
$M_N = 55$ GeV



$M_N = 40$ GeV

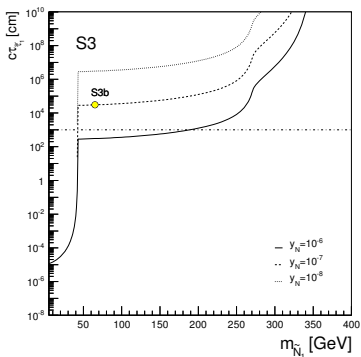
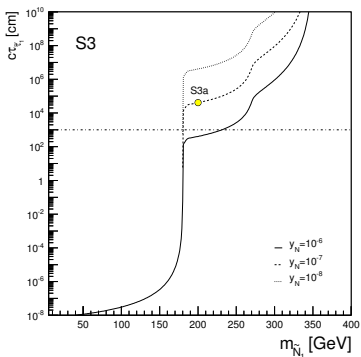


$M_N = 20$ GeV



LHC long-lived staus

LHC constrains the mass of the long-lived staus to be greater than $m_{\tilde{\tau}_1} = 342$ GeV



Cuts

- 2 staus that scape the detector ($c\tau > 10$ m)
- $\beta (\equiv v/c) < 0.95$
- $p_T > 50$ GeV and $|\eta| < 2.5$

$$\sqrt{s} = 13 \text{ TeV and } \mathcal{L} = 100 \text{ fb}^{-1}$$

S3a: 30.3 events
S3b: 28.9 events

- We introduce two singlet superfields to solve the μ problem (S) and the neutrino mass problem (N).
- The introduction of the singlet S adds tree-level contribution to the Higgs mass.
- The introduction of the singlet N adds RH-sneutrinos and RH-neutrinos.
- RH-sneutrinos are well DM candidates.
- RH-neutrinos could produce displaced vertices and the stau to be long-lived.
- Specific signatures at the LHC that are interesting to look for.

Thank you for your attention!