

Boosting stop searches with boosted di-boson

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Tuning in the MSSM

$$\Delta m_h^2 \approx \frac{3 m_t^4}{2 \pi^2 v^2} \left[\ln \left(\frac{M_S^2}{m_t^2} \right) + \frac{X_t^2}{2 M_S^2} \left(1 - \frac{X_t^2}{6 M_S^2} \right) \right] + \dots$$

$$m_h^2 = -2 \left(m_{H_u}^2 + |\mu|^2 \right) \quad (\text{large } \tan \beta)$$

$$X_t = A_t - \frac{\mu}{\tan \beta}, \quad M_S^2 = M_{\tilde{t}_1} M_{\tilde{t}_2}$$

$$\frac{d}{dt} m_{H_u}^2 = \frac{6 |y_t|^2}{(4\pi)^2} \left(m_{\tilde{t}_L}^2 + m_{\tilde{t}_R}^2 + |A_t|^2 \right) + \dots$$

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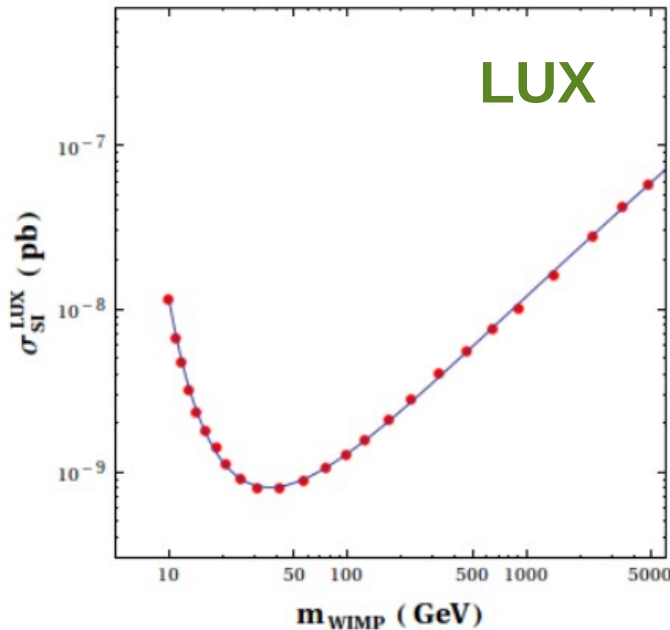
$$\frac{d}{dt} m_{\tilde{t}}^2 = -\frac{8\alpha_s}{3\pi} M_3^2 + \dots$$

gluino sucks effect:
the stop mass itself is attracted to the gluino mass in the IR

Strong LHC bound on gluino mass \dashrightarrow large tuning

pMSSM

Parameter	Scan range (GeV)
M_1	20 – 2000
M_2	100 – 2000
$\tan(\beta)$	2 – 55
μ	100 – 3000
m_A	100 – 2000
A_t	-5000 – 5000
A_b	-5000 – 5000
$m_{\tilde{Q}_3}$	100 – 2000
$m_{\tilde{U}_3}$	100 – 2000
$m_{\tilde{D}_3}$	100 – 2000



Light stop?

Observable	Allowed range
$\text{BR}(B_d \rightarrow X_s \gamma) \times 10^4$	2.78 – 4.08
$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \times 10^9$	1.43 – 4.37
$\text{BR}(B_d \rightarrow \mu^+ \mu^-) \times 10^{10}$	0.79 – 6.80
Ωh^2	≤ 0.1289

+

HiggsBounds+HiggsSignals

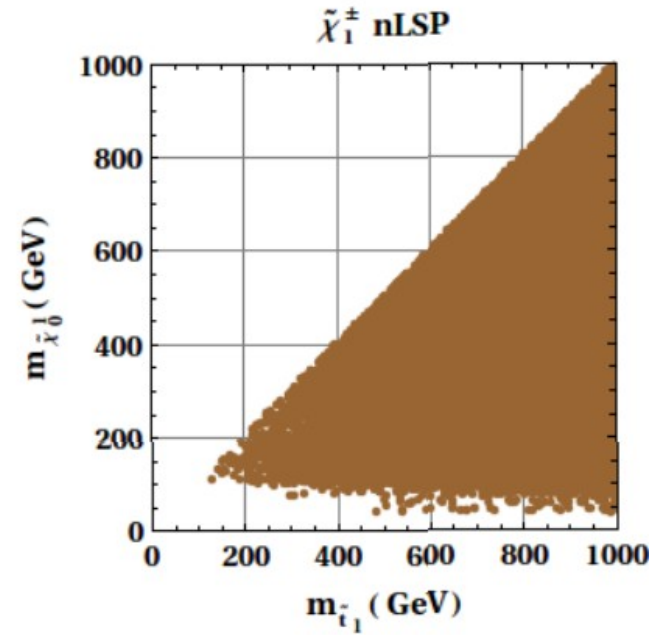
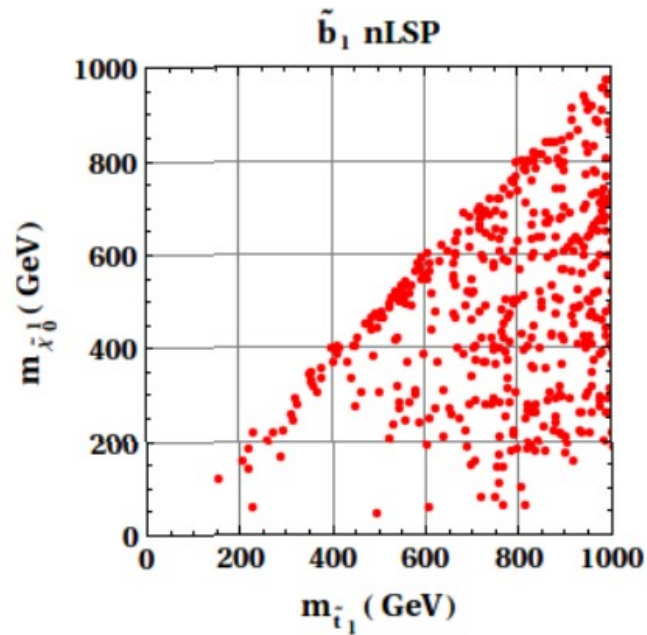
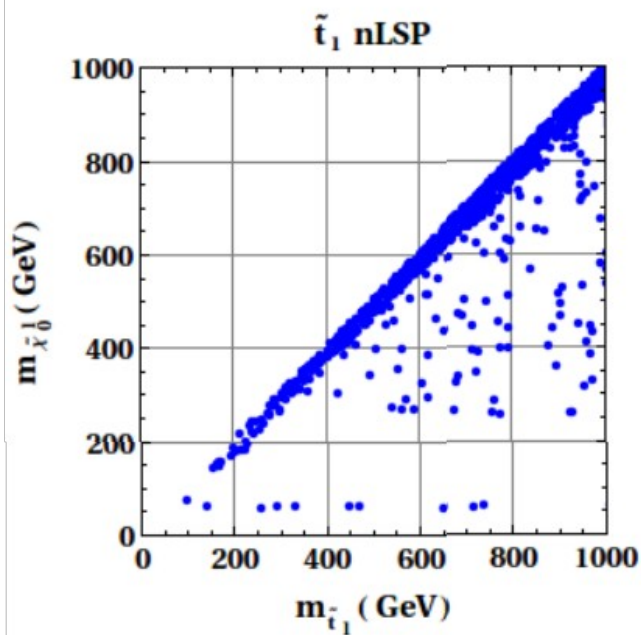
Number of free model parameters $N_p = 10$

$p\text{-value} \geq 0.05$

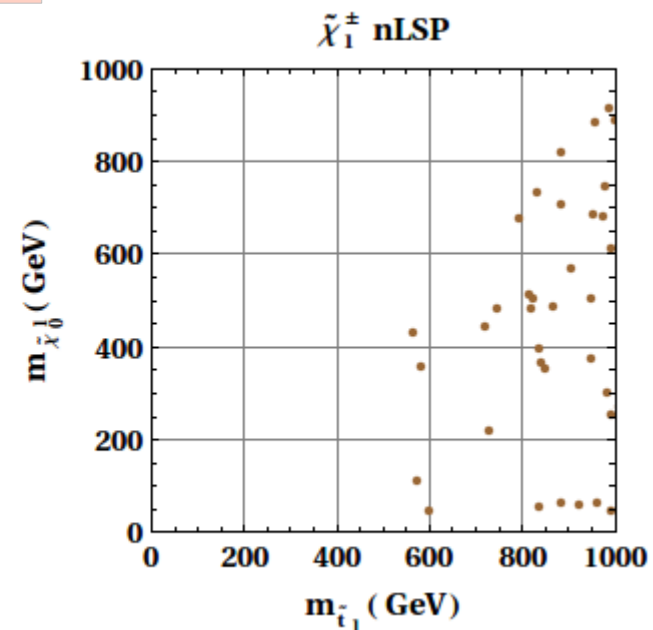
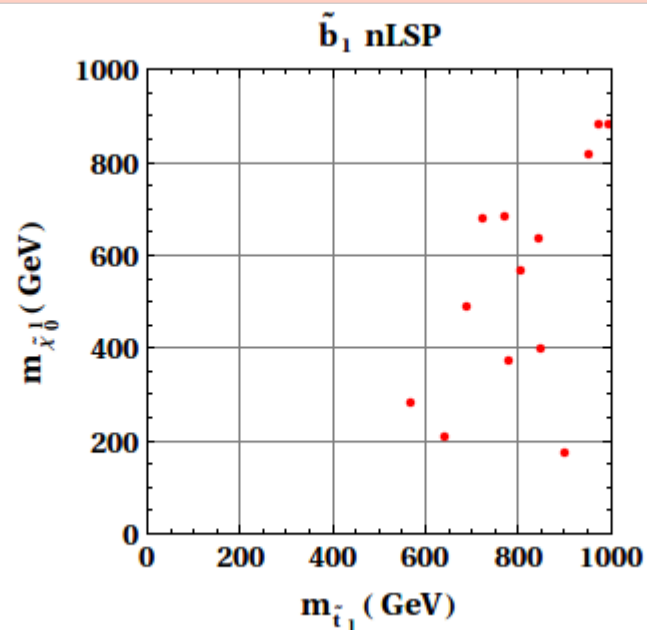
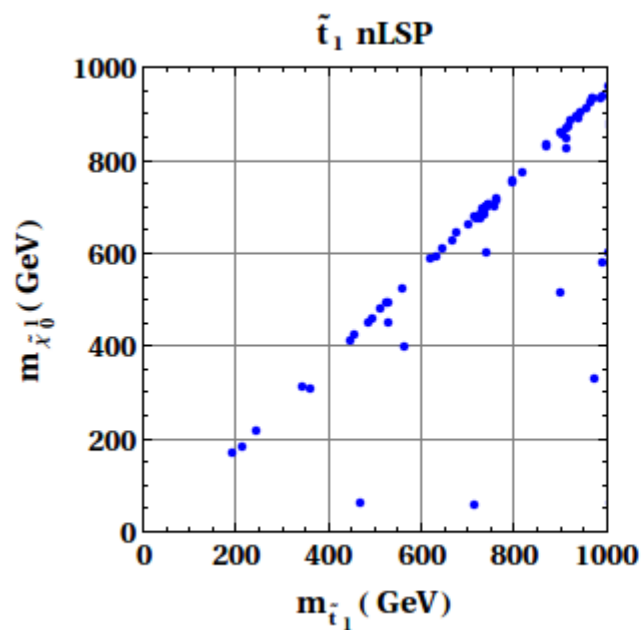
$$\log_{10} \sigma_{\text{SI}}^{\text{LUX}} = -8.569 + \frac{7.029}{(\log_{10} m_{\text{WIMP}})^2} - \frac{7.161}{\log_{10} m_{\text{WIMP}}} + 0.755 \log_{10} m_{\text{WIMP}} - 0.003 (\log_{10} m_{\text{WIMP}})^2$$

$$\sigma_{\text{SI}} < \begin{cases} \sigma_{\text{SI}}^{\text{LUX}} & \text{for } 0.1103 < \Omega h^2 < 0.1289 \\ \frac{0.1196}{\Omega h^2} \sigma_{\text{SI}}^{\text{LUX}} & \text{for } \Omega h^2 < 0.1103 \end{cases}$$

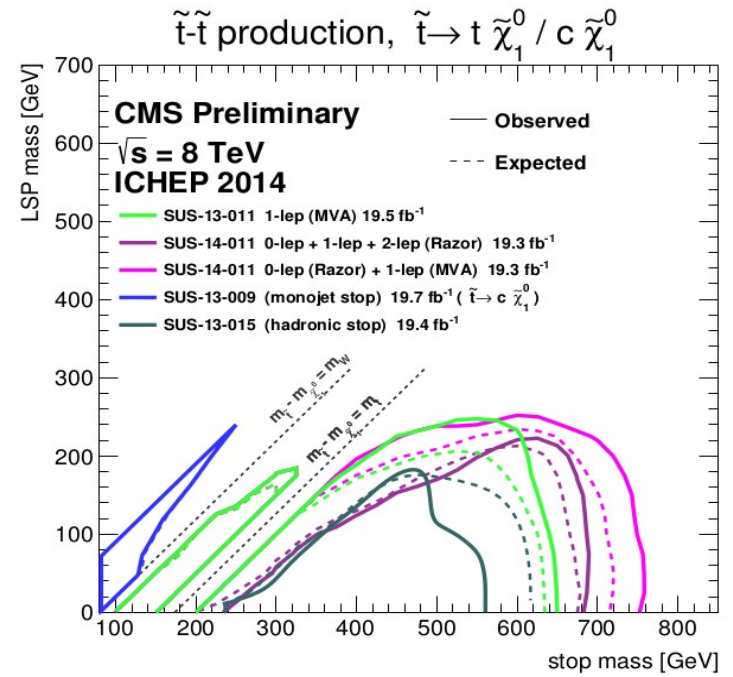
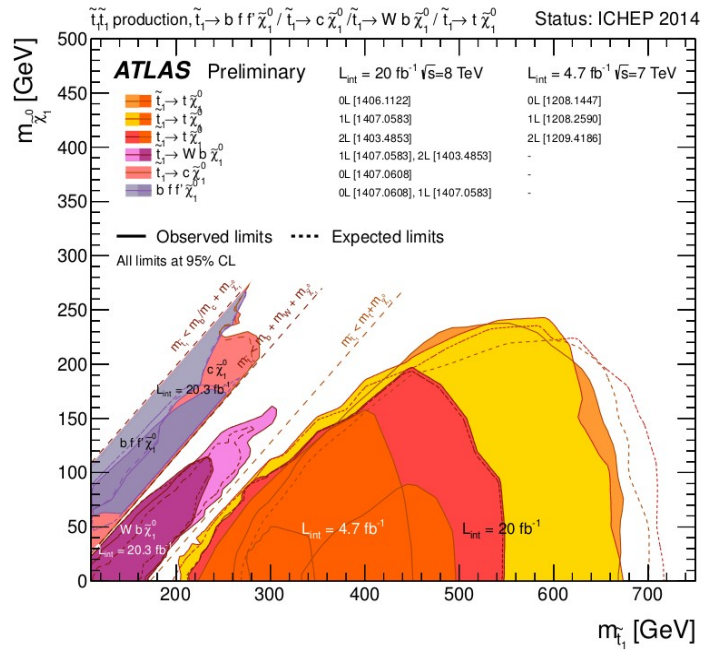
$$\Omega h^2 < 0.1289$$



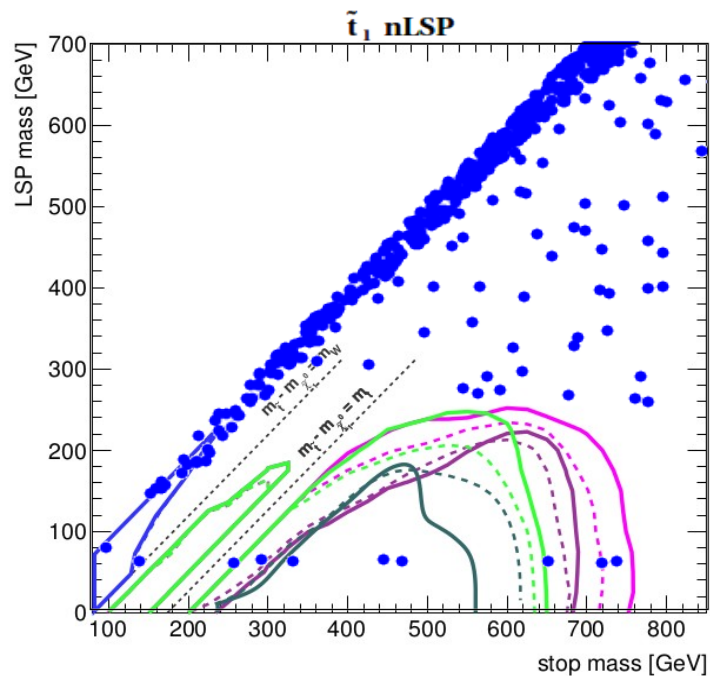
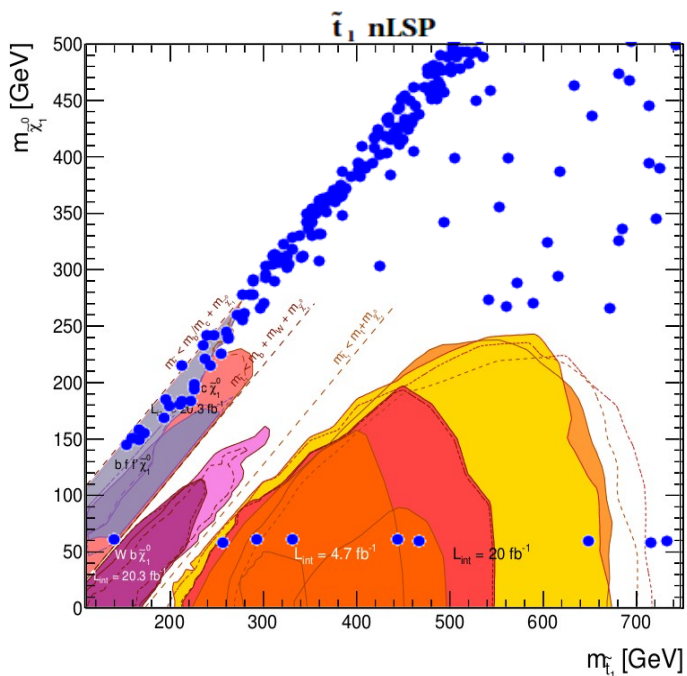
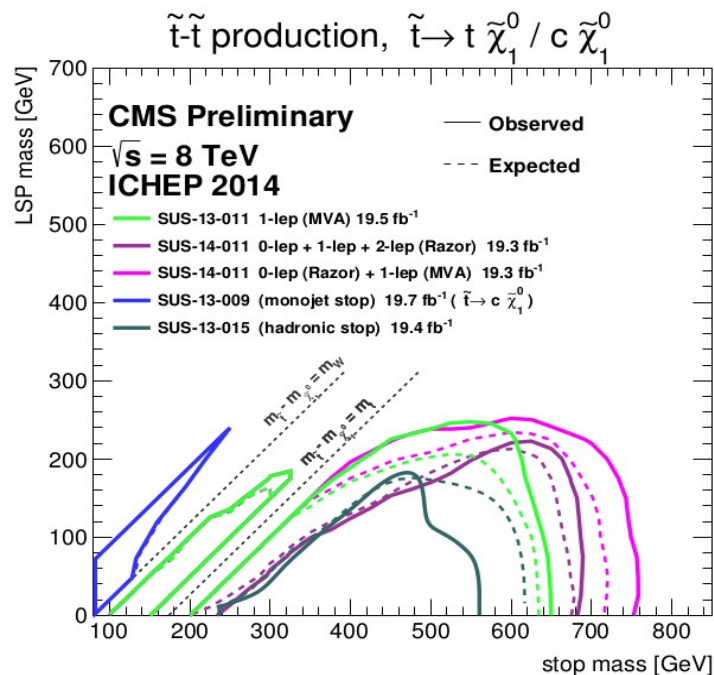
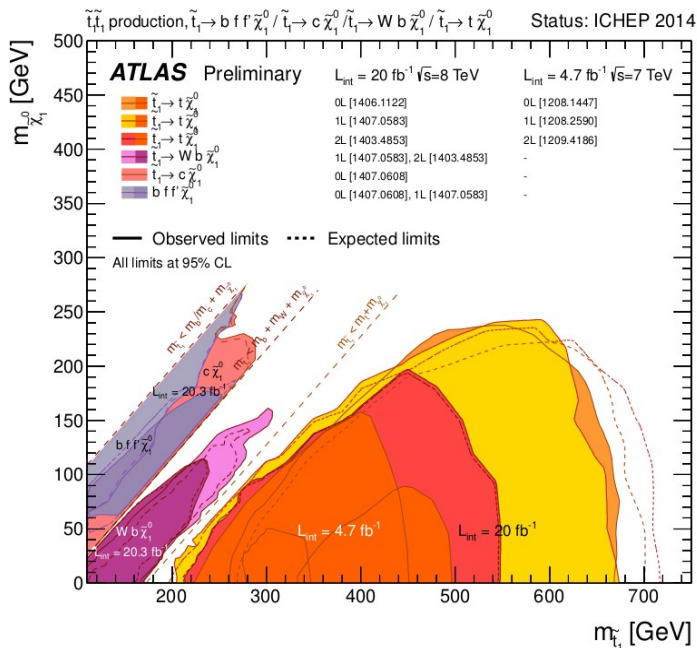
$$0.1103 < \Omega h^2 < 0.1289$$



LHC bounds



LHC bounds



Degenerate \tilde{t}_1 and $\tilde{\chi}_1^0$

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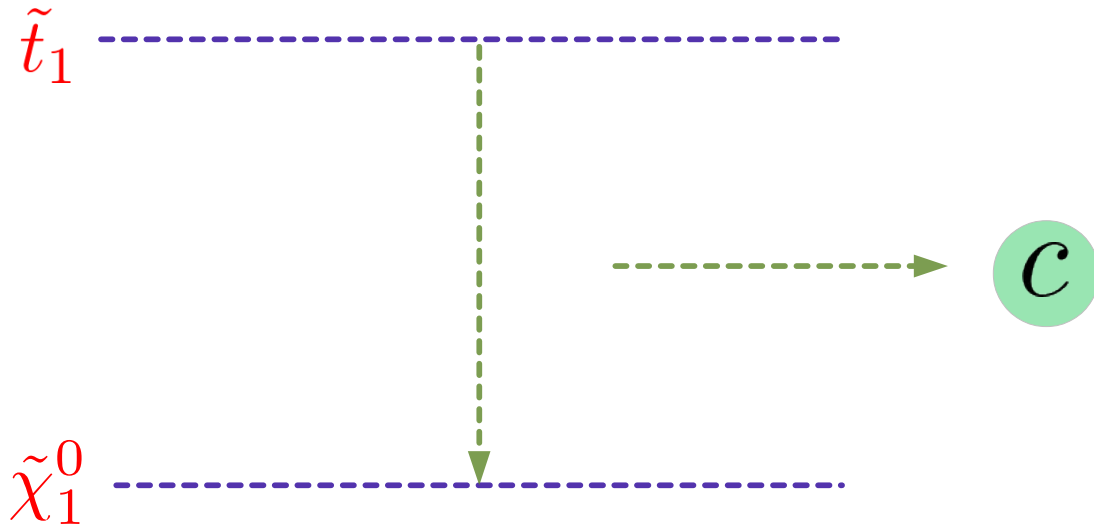
\tilde{t}_1 -----

Degenerate \tilde{t}_1 and $\tilde{\chi}_1^0$

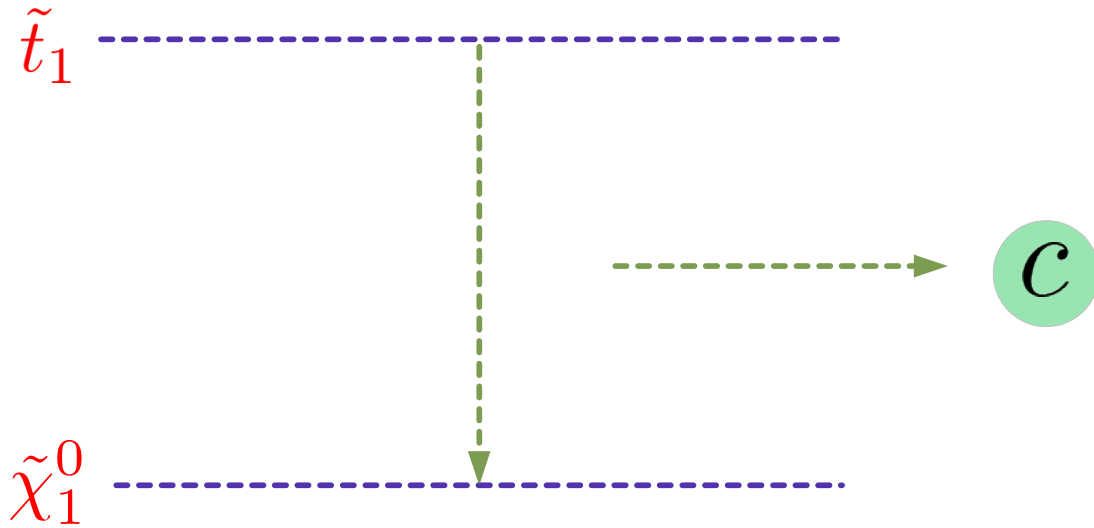
\tilde{t}_1 -----

$\tilde{\chi}_1^0$ -----

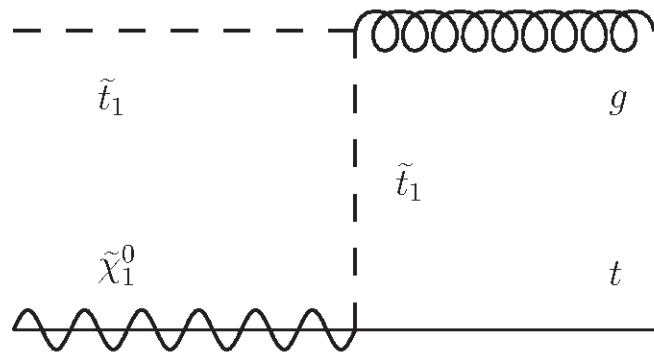
Degenerate \tilde{t}_1 and $\tilde{\chi}_1^0$



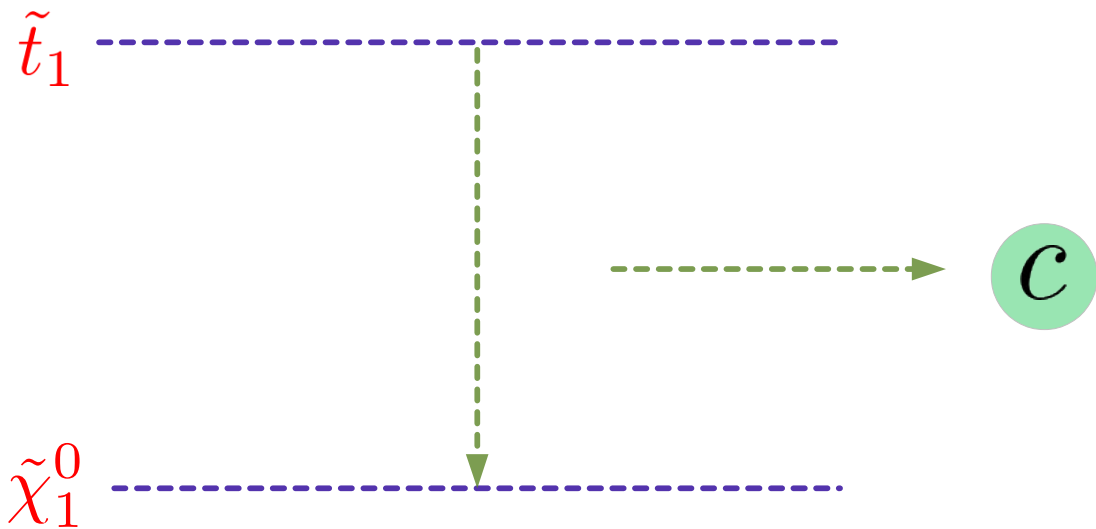
Degenerate \tilde{t}_1 and $\tilde{\chi}_1^0$



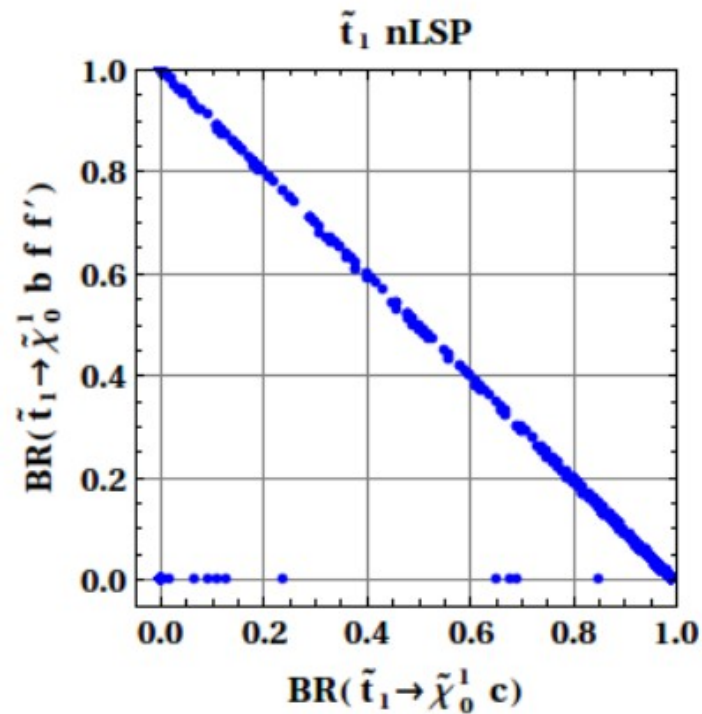
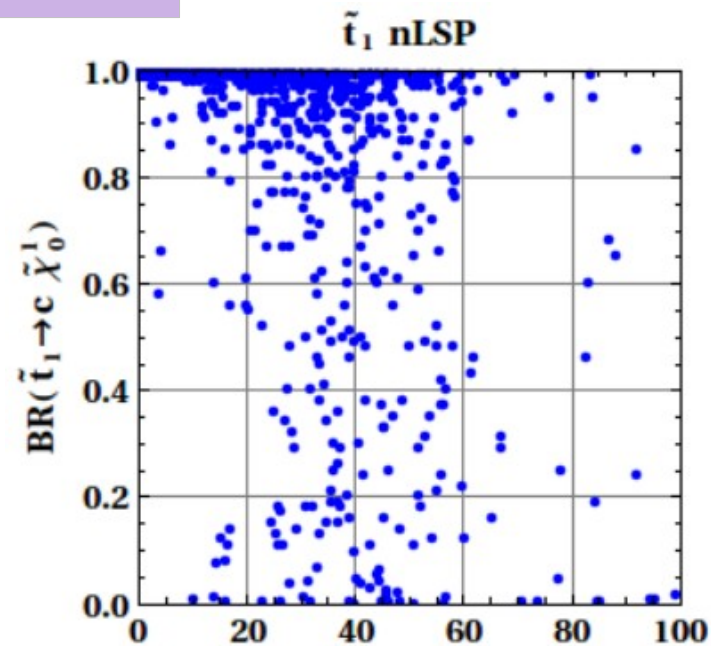
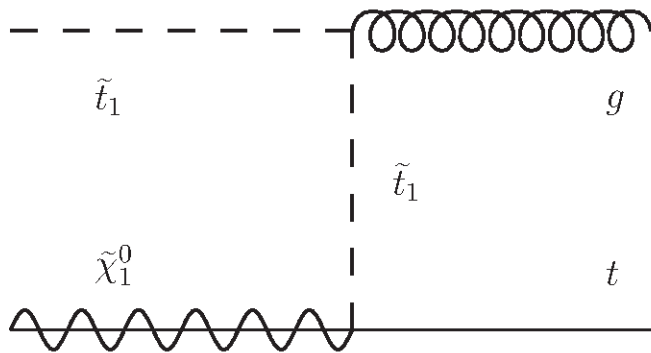
Coannihilation



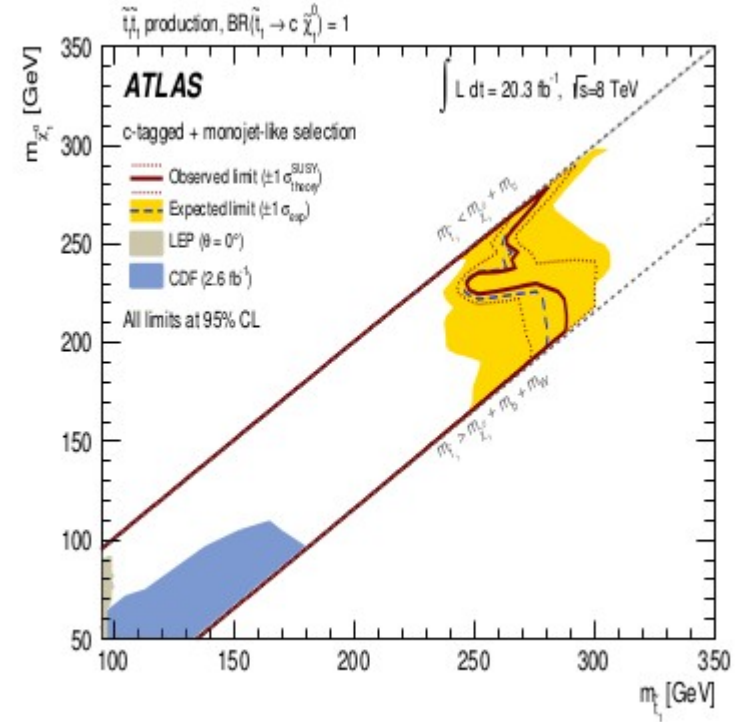
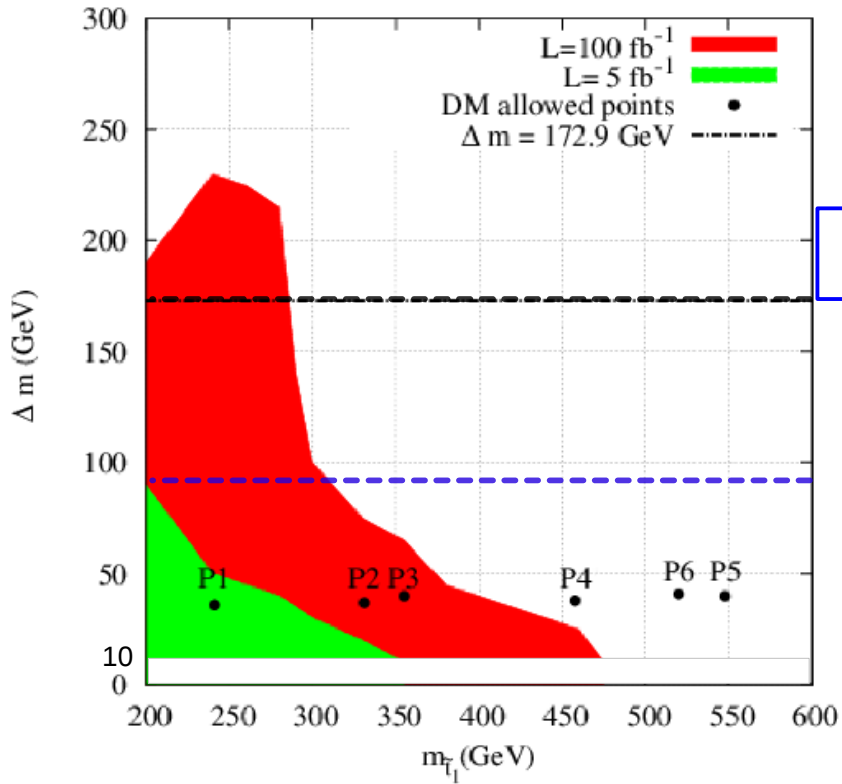
Degenerate \tilde{t}_1 and $\tilde{\chi}_1^0$



Coannihilation

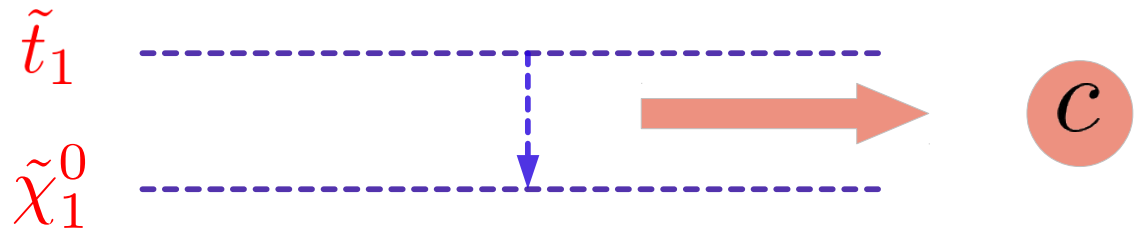
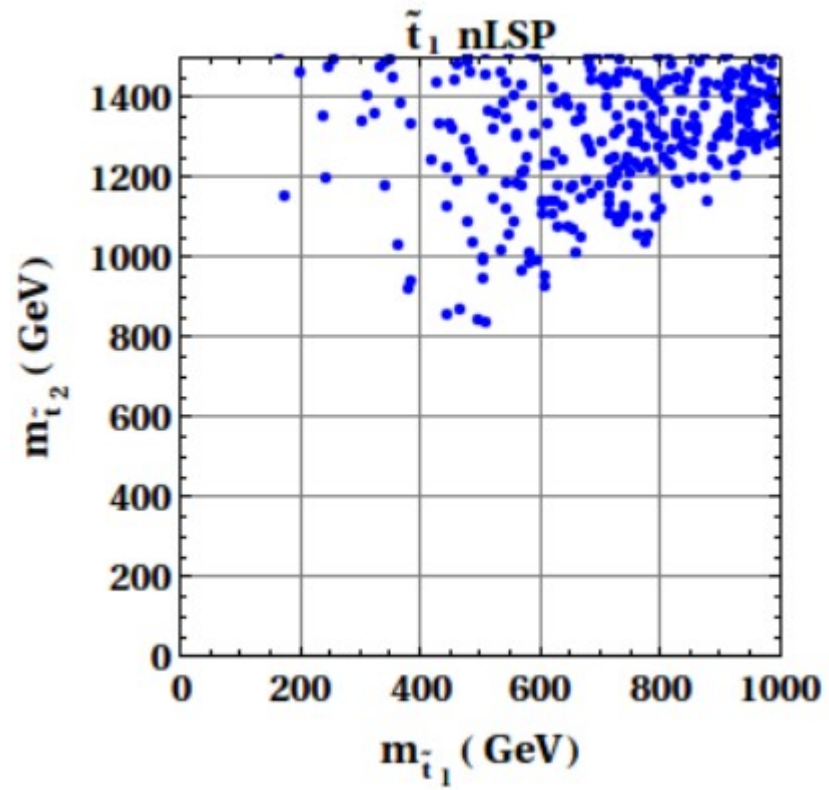


LHC14 projections



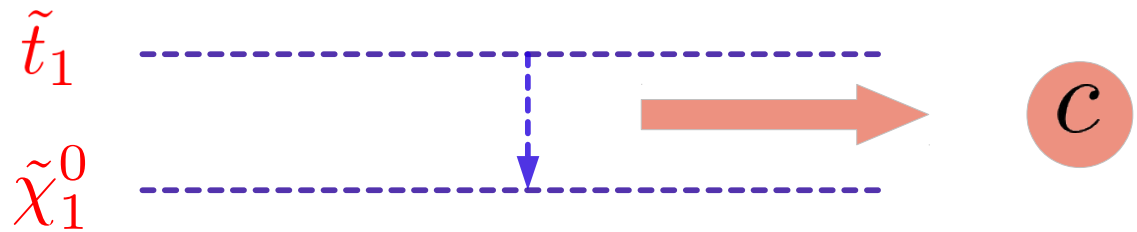
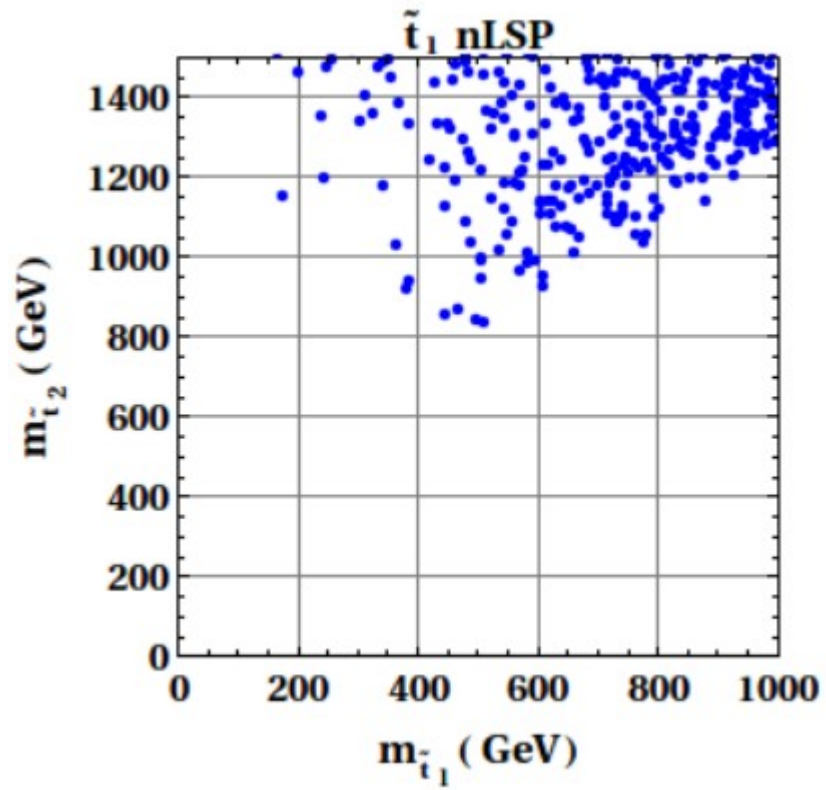
	P1	P2	P3	P4	P5	P6
$m_{\tilde{t}_1}, m_{\tilde{\chi}_1^0}$ (GeV)	241, 205	331, 294	355, 315	458, 420	548, 508	520, 479
Ωh^2	0.119	0.119	0.119	0.119	0.119	0.119
Δm (GeV)	36	37	40	38	40	41

Where is \tilde{t}_2

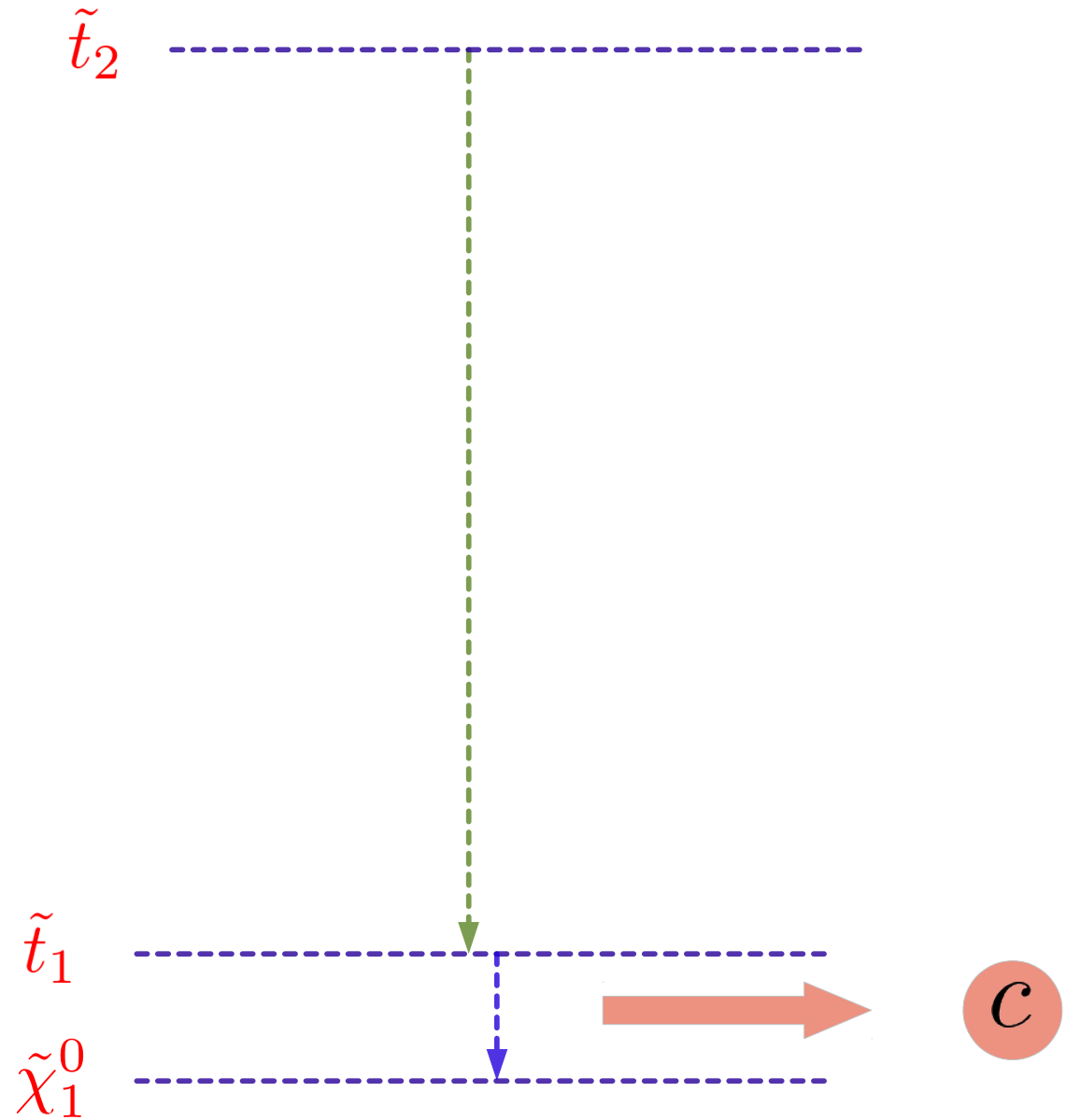
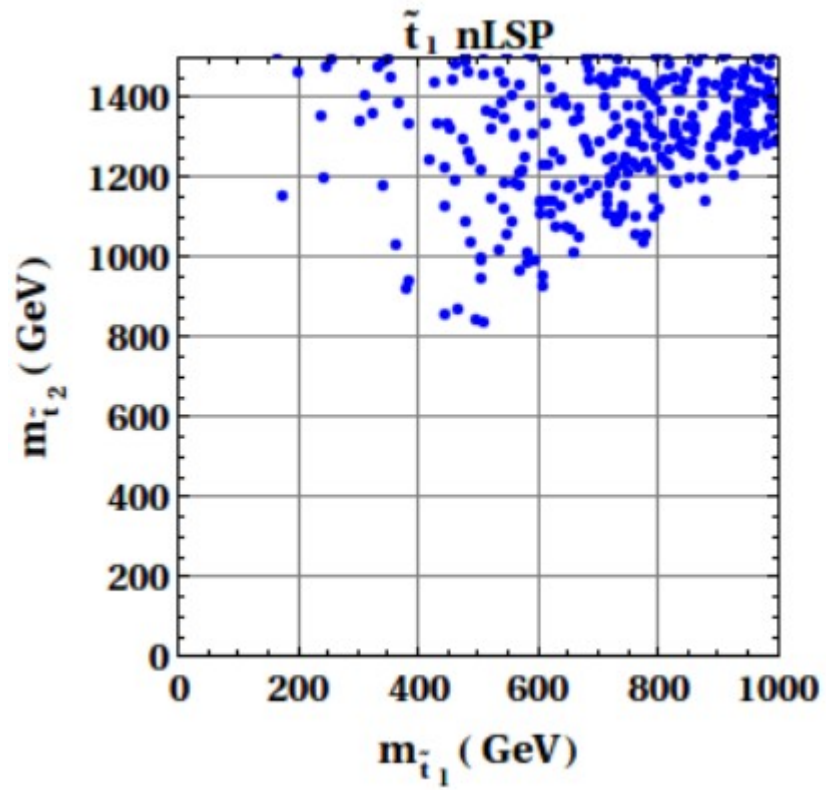


Where is \tilde{t}_2

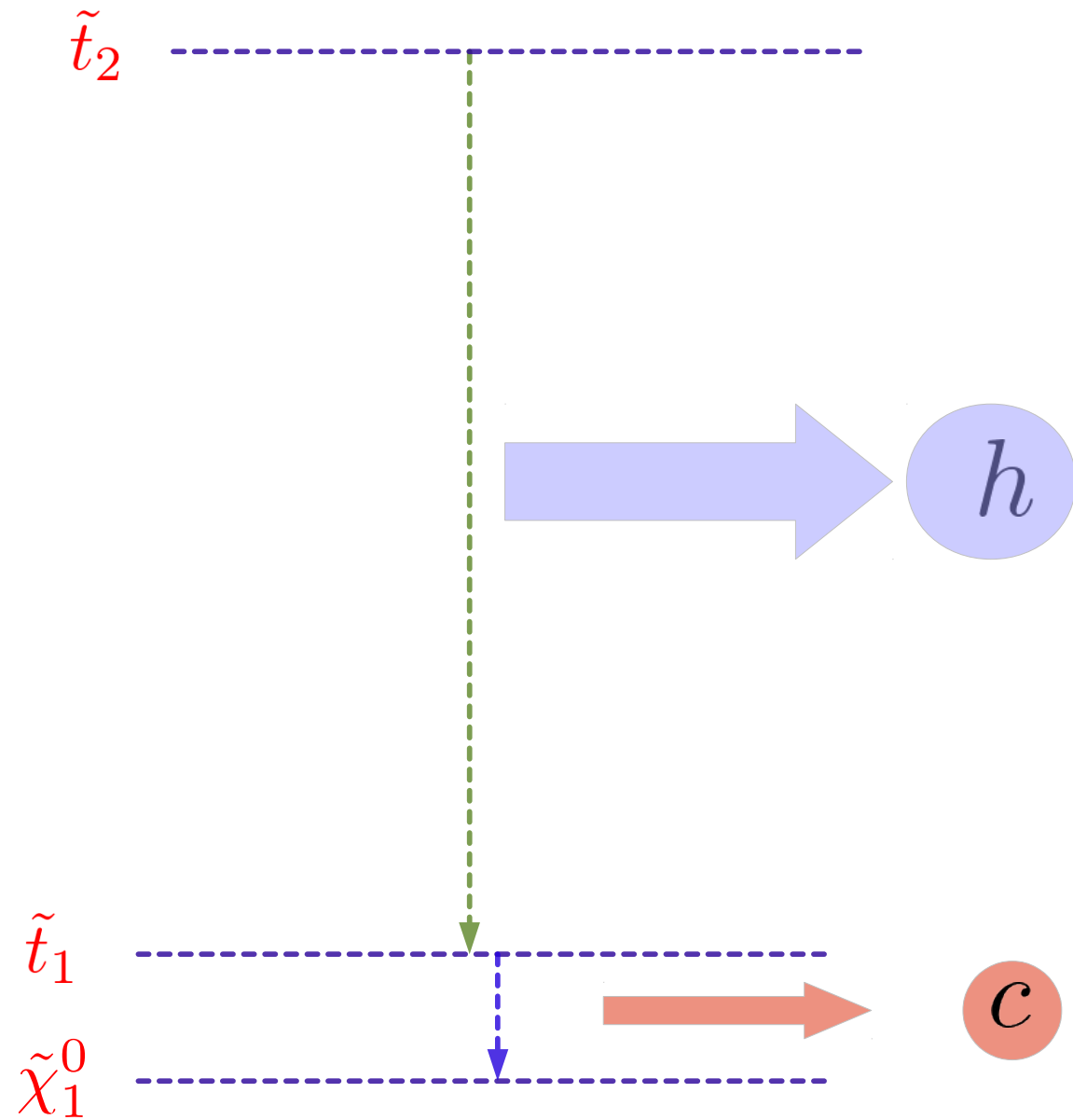
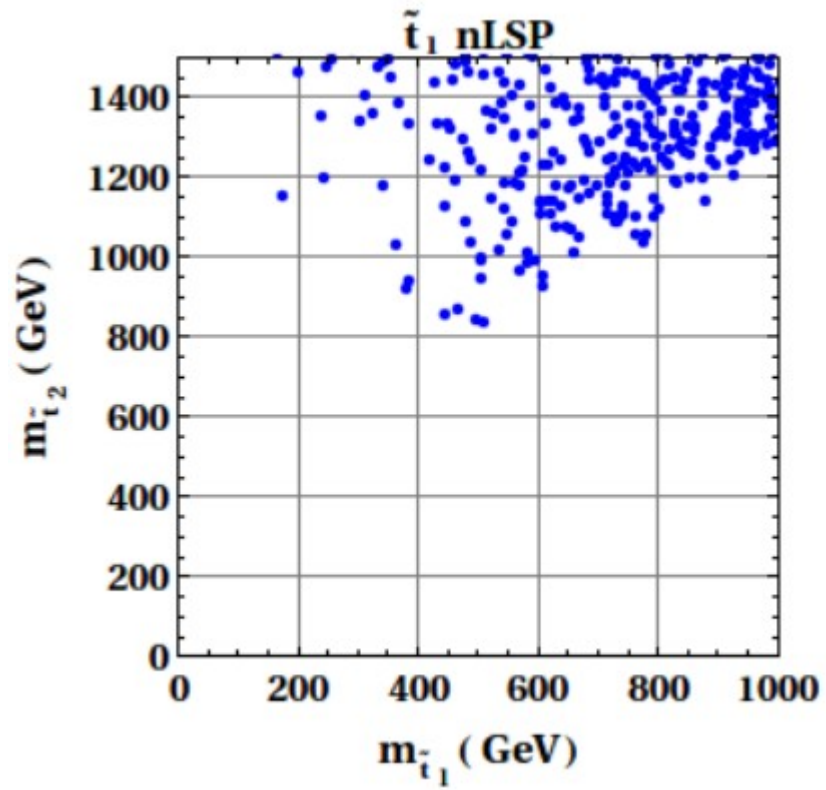
\tilde{t}_2 -----



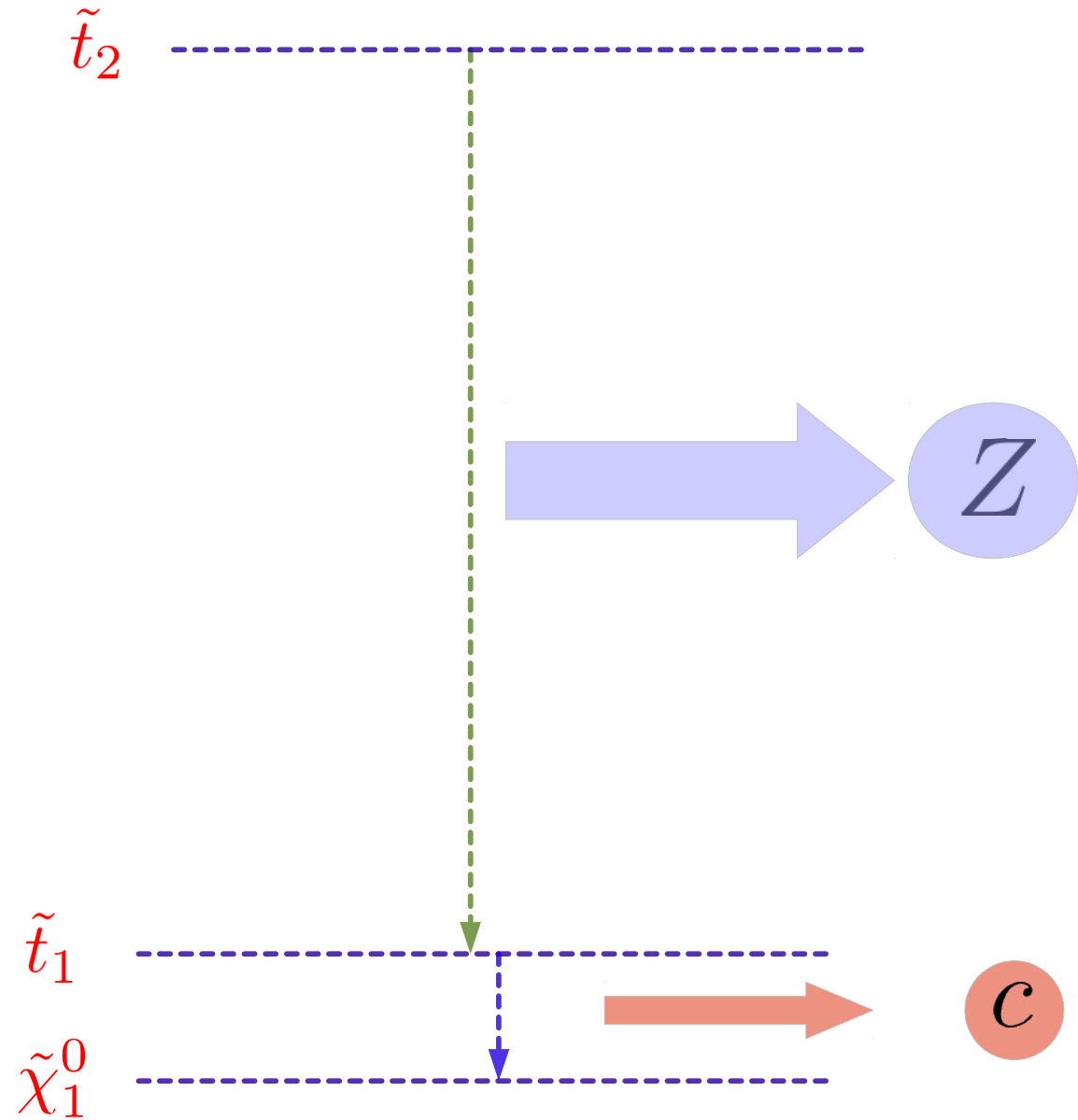
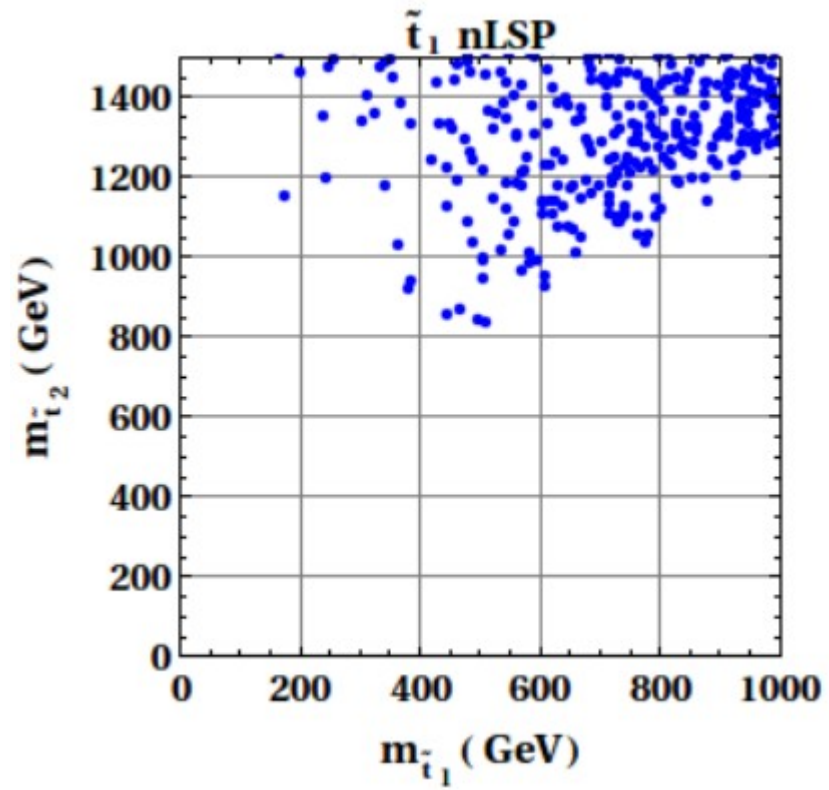
Where is \tilde{t}_2



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Where is \tilde{t}_2



Branching Ratios

DG, PRD 2013

$$\lambda_{\tilde{t}_2\tilde{t}_1 h} \propto 2(\sqrt{2}G_F)^{\frac{1}{2}} M_Z^2 \times$$

$$\left[\left(\frac{2}{3} \sin^2 \theta_W - \frac{1}{4} \right) \cos(2\beta) \sin(2\theta_t) + \frac{1}{2} \frac{m_t}{M_Z^2} \cos(2\theta_t) X_t \right]$$

$$\lambda_{\tilde{t}_2\tilde{t}_1 Z} \approx \frac{g}{2M_W} m_t X_t$$

$$m_{\tilde{t}_{1,2}}^2 = m_t^2 + \frac{1}{2} \left[m_{LL}^2 + m_{RR}^2 \mp \sqrt{(m_{LL}^2 - m_{RR}^2)^2 + 4m_t^2 X_t^2} \right]$$

$$\sin 2\theta_t = \frac{2m_t X_t}{m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2}, \quad \cos 2\theta_t = \frac{m_{LL}^2 - m_{RR}^2}{m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2}$$

$$M_{\tilde{t}}^2 = \begin{pmatrix} m_{LL}^2 & m_t X_t \\ m_t X_t & m_{RR}^2 \end{pmatrix}$$

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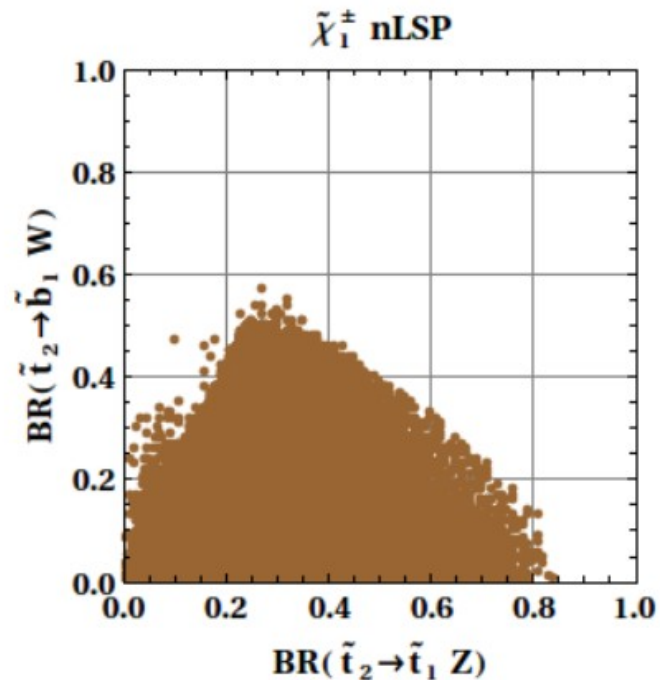
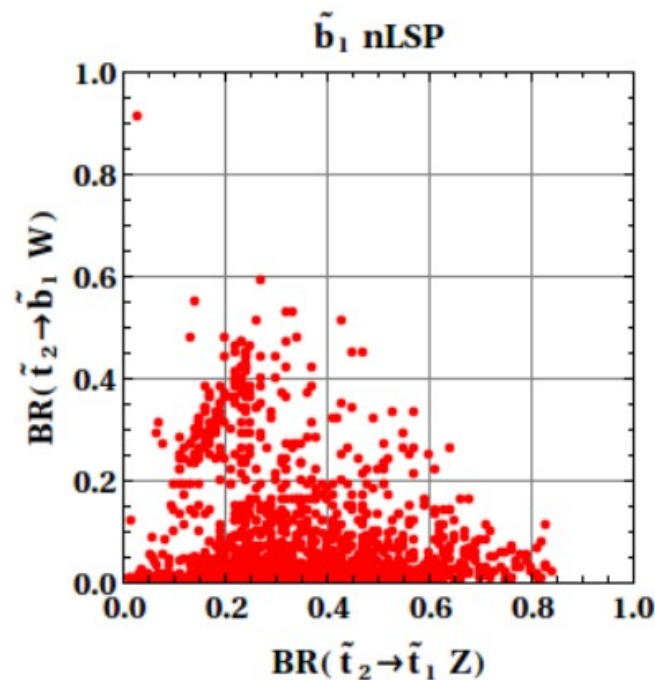
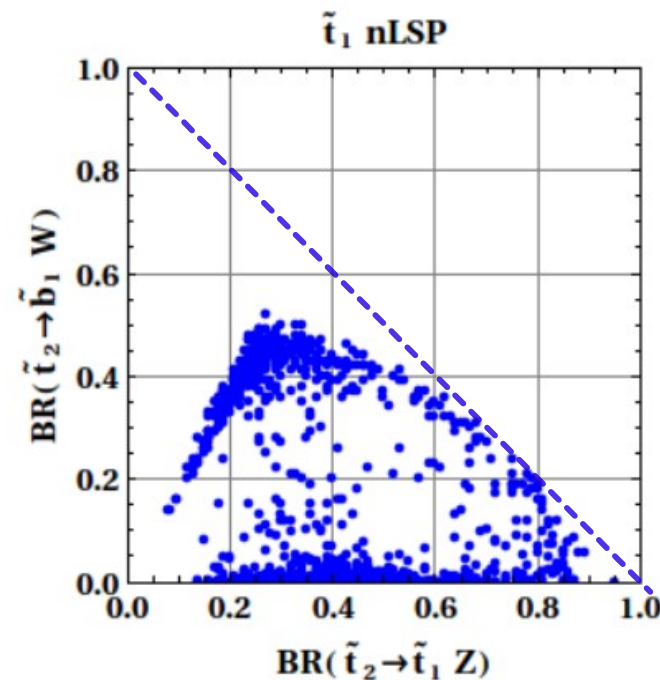
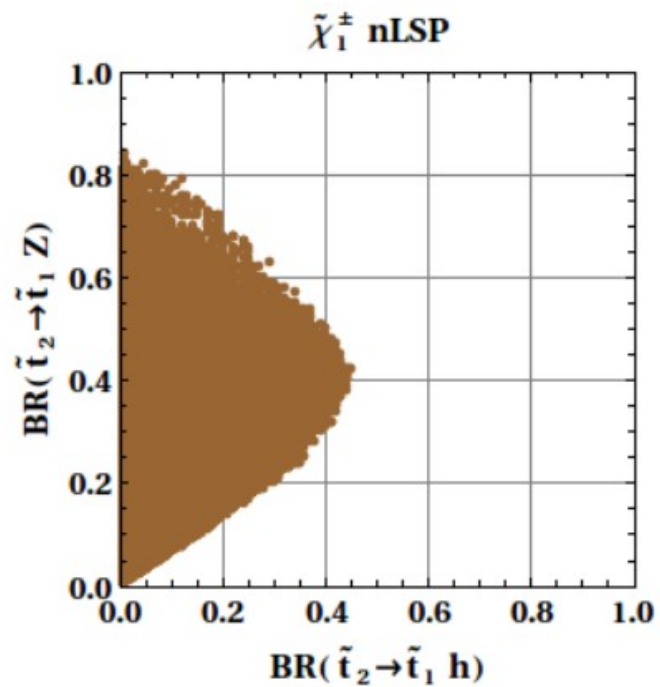
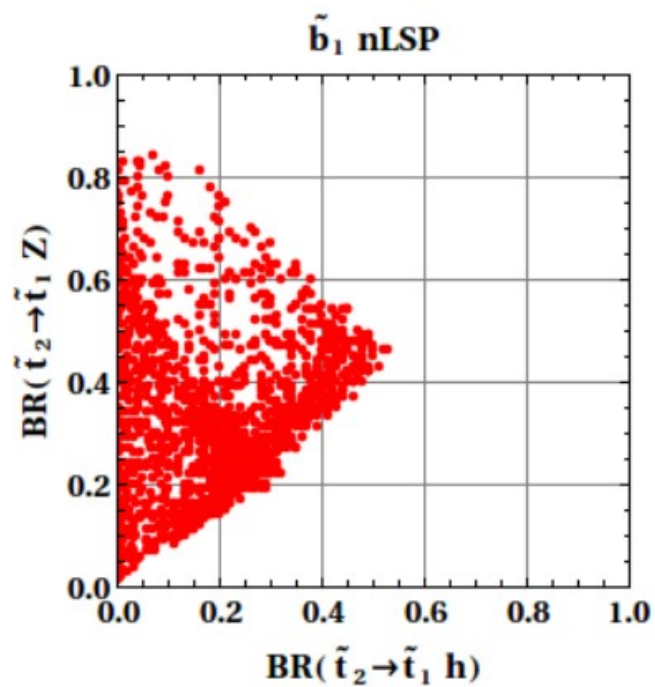
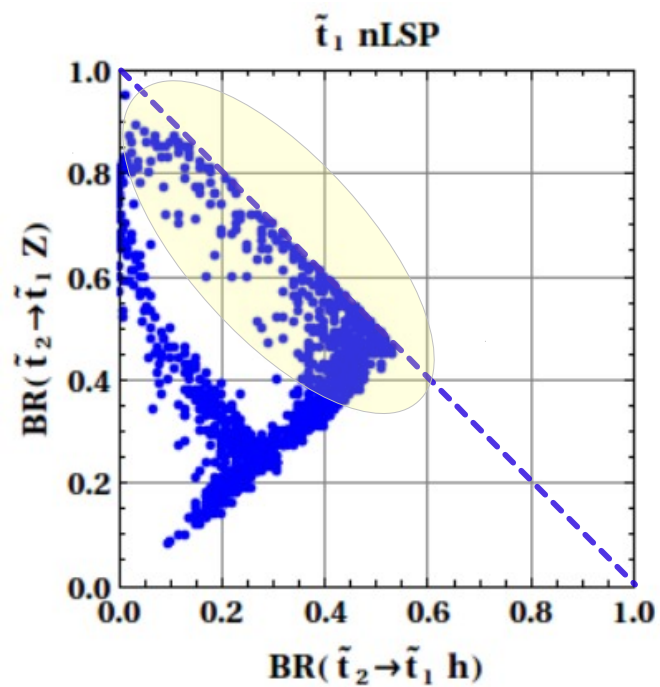
$$\lambda_{\tilde{t}_2\tilde{t}_1 Z} \approx \frac{g}{2M_W} m_t X_t$$

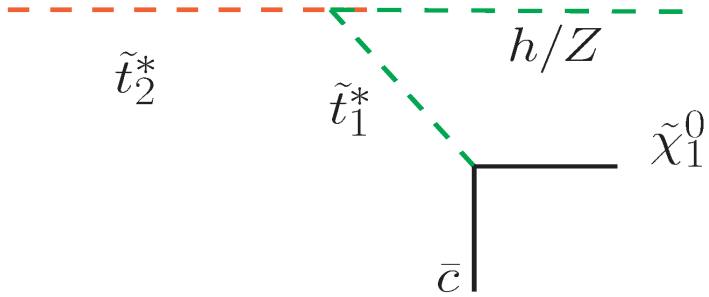
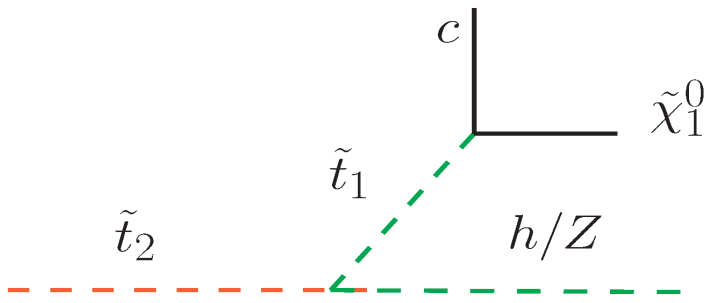
Higgs mass \longrightarrow Large
 $\text{Br}(\tilde{t}_2 \rightarrow \tilde{t}_1 Z/h)$

$$m_{\tilde{t}_{1,2}}^2 = m_t^2 + \frac{1}{2} \left[m_{LL}^2 + m_{RR}^2 \mp \sqrt{(m_{LL}^2 - m_{RR}^2)^2 + 4m_t^2 X_t^2} \right]$$

$$\sin 2\theta_t = \frac{2m_t X_t}{m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2}, \quad \cos 2\theta_t = \frac{m_{LL}^2 - m_{RR}^2}{m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2}$$

$$M_{\tilde{t}}^2 = \begin{pmatrix} m_{LL}^2 & m_t X_t \\ m_t X_t & m_{RR}^2 \end{pmatrix}$$





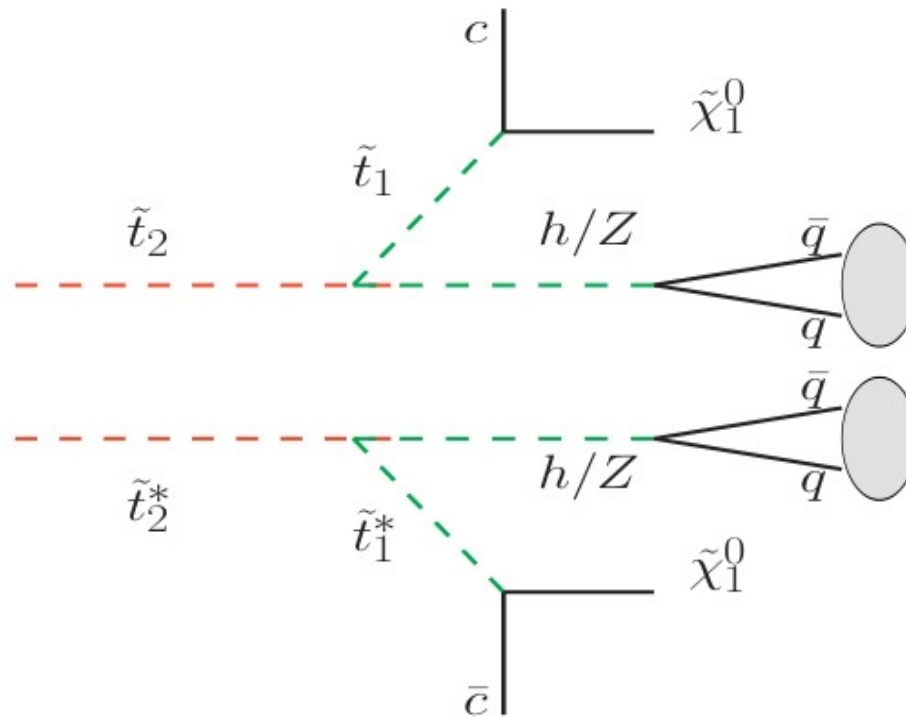
Model:1		Model:2	
pMSSM inputs	Masses	pMSSM inputs	Masses
$M_1 = 300$	$m_{\tilde{t}_2} = 1005$	$M_1 = 410$	$m_{\tilde{t}_2} = 1003$
$M_2 = 650$	$m_{\tilde{t}_1} = 334$	$M_2 = 850$	$m_{\tilde{t}_1} = 434$
$M_3 = 2100$	$m_{\tilde{\chi}_1^0} = 300$	$M_3 = 2600$	$m_{\tilde{\chi}_1^0} = 411$
$\mu = 2000$	$m_{\tilde{\chi}_2^0} = 676$	$\mu = 2000$	$m_{\tilde{\chi}_2^0} = 884$
$m_A = 1500$	$m_{\tilde{\chi}_1^\pm} = 676$	$m_A = 1500$	$m_{\tilde{\chi}_1^\pm} = 884$
$\tan \beta = 10$	$m_h = 125$	$\tan \beta = 7.5$	$m_h = 125$
$m_{Q3} = 1010$		$m_{Q3} = 1050$	
$m_{t_R} = 630$		$m_{t_R} = 770$	
$m_{b_R} = 3000$		$m_{b_R} = 3000$	
$A_t = -1700$		$A_t = -1600$	
$\mathcal{B}(\tilde{t}_2 \rightarrow \tilde{t}_1 Z) = 52\%$ $\mathcal{B}(\tilde{t}_2 \rightarrow \tilde{t}_1 h) = 39\%$ $\mathcal{B}(\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0) = 82\%$		$\mathcal{B}(\tilde{t}_2 \rightarrow \tilde{t}_1 Z) = 56\%$ $\mathcal{B}(\tilde{t}_2 \rightarrow \tilde{t}_1 h) = 41\%$ $\mathcal{B}(\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0) = 90\%$	

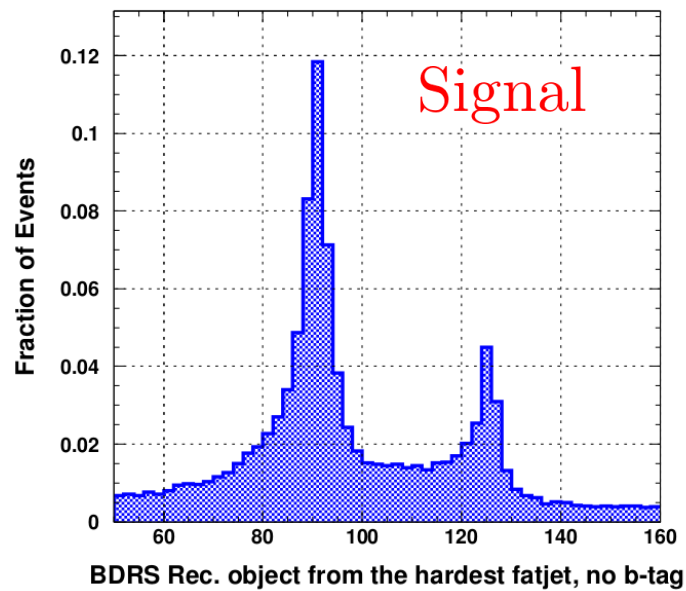
Final State

$$\begin{array}{l}
 \text{pp} \rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* Z Z \rightarrow Z Z \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 \hookrightarrow Z Z + \cancel{p}_T + \text{soft jets} \\
 \text{pp} \rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* Z h \rightarrow Z h \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 \hookrightarrow Z h + \cancel{p}_T + \text{soft jets} \\
 \text{pp} \rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* h h \rightarrow h h \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 \hookrightarrow h h + \cancel{p}_T + \text{soft jets}
 \end{array}$$

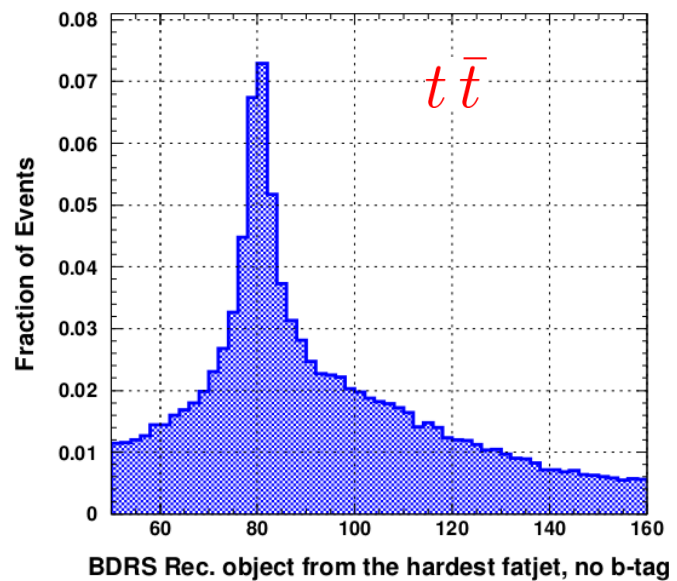
Final State

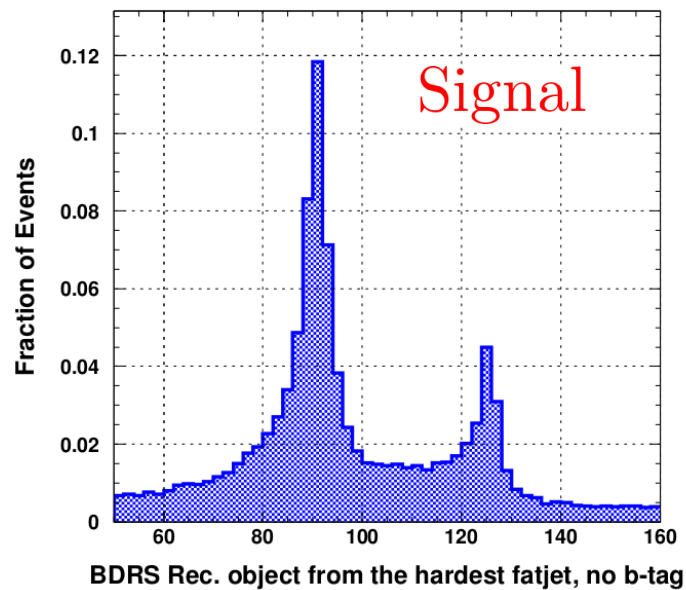
$$\begin{aligned}
 pp &\rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* Z Z \rightarrow Z Z \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 &\hookrightarrow Z Z + \cancel{p}_T + \text{soft jets} \\
 pp &\rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* Z h \rightarrow Z h \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 &\hookrightarrow Z h + \cancel{p}_T + \text{soft jets} \\
 pp &\rightarrow \tilde{t}_2 \tilde{t}_2^* \rightarrow \tilde{t}_1 \tilde{t}_1^* h h \rightarrow h h \tilde{\chi}_1^0 \tilde{\chi}_1^0 c \bar{c} \\
 &\hookrightarrow h h + \cancel{p}_T + \text{soft jets}
 \end{aligned}$$



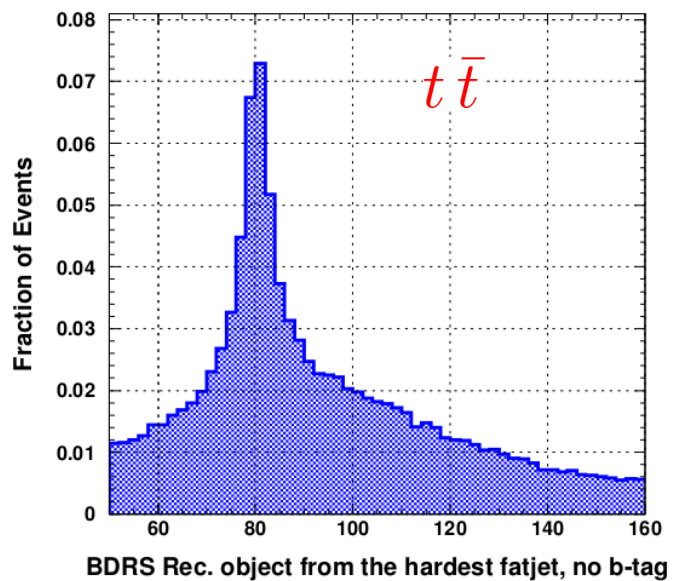
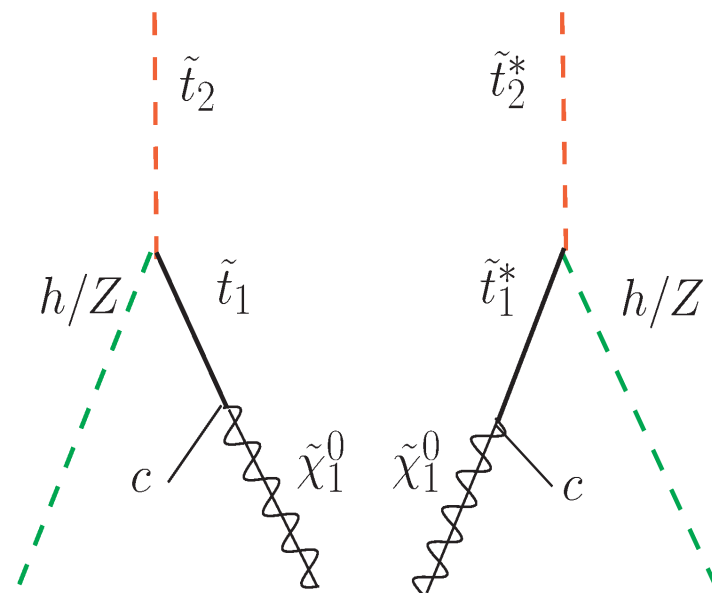


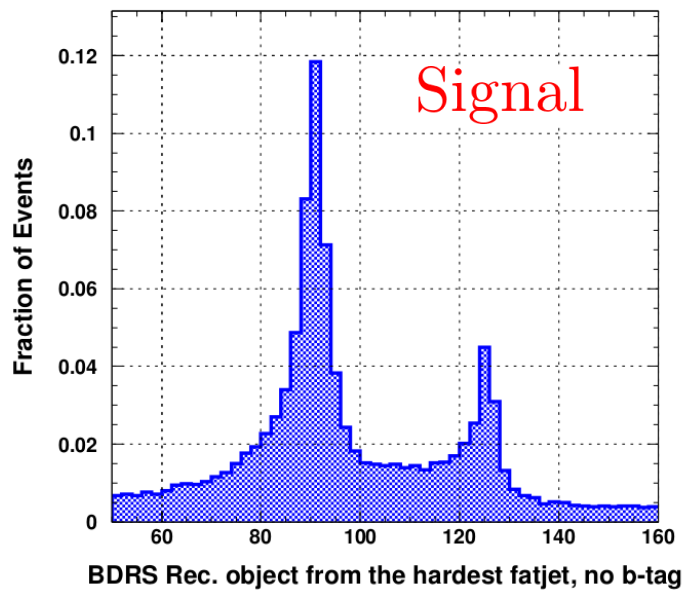
BDRS, R=1.0



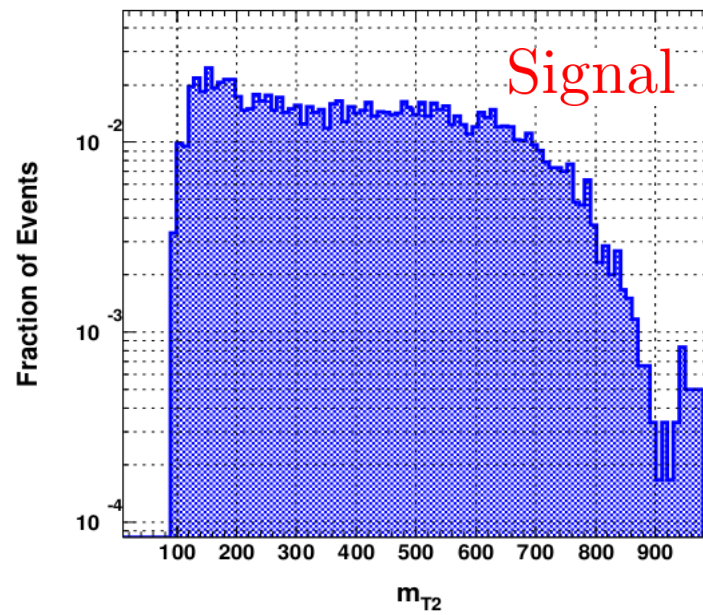
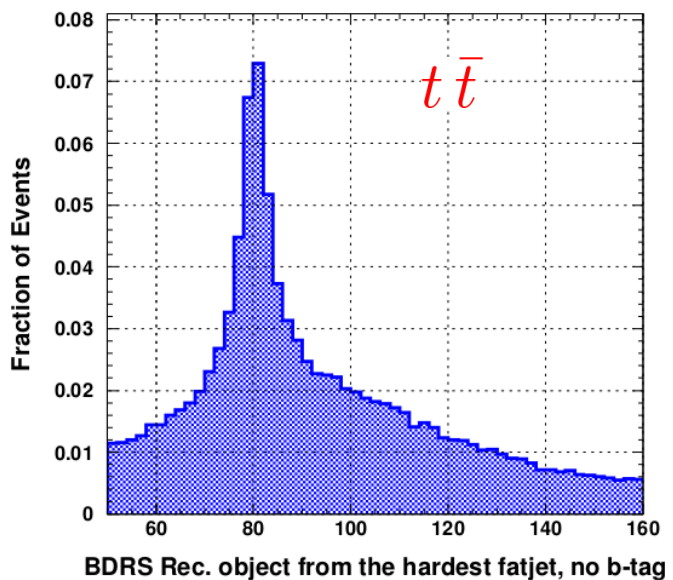
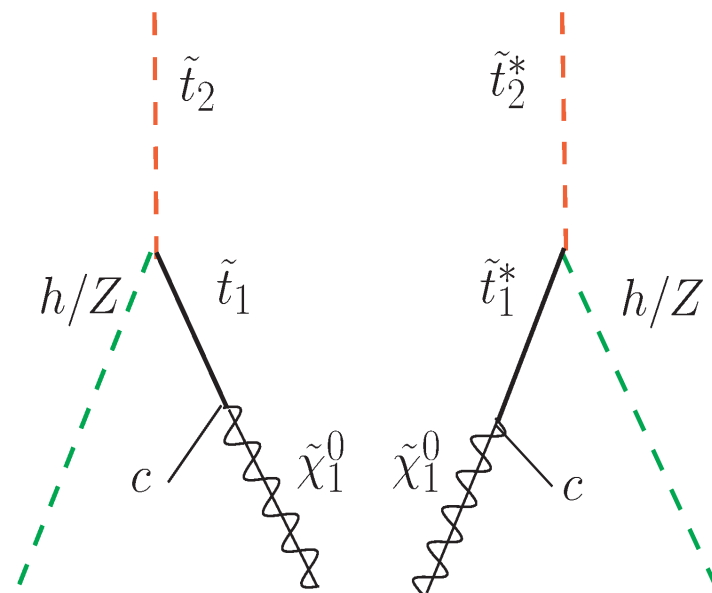


BDRS, R=1.0





BDRS, R=1.0



Results

Process	Production cross-section	Simulated events	No. of events after						Final cross-section (fb)	\mathcal{S} (100 fb ⁻¹)
			S1	S2	S3	S4	S5	S6		
Signal										
Model:1	10 fb	10 ⁵	6012	4902	2736	2359	2143	1718	17.2 × 10 ⁻²	4.3
Model:2	10 fb	10 ⁵	6170	5319	2813	2421	2081	1853	18.5 × 10 ⁻²	4.6
Backgrounds										
$t\bar{t}$	833 pb	10 ⁸	221747	148580	142	41	26	11	9.1 × 10 ⁻²	
$t\bar{t}Z(1j)$	1.12 pb	226110	2484	1444	8	7	1	1	0.5 × 10 ⁻²	
$t\bar{t}W^\pm(1j)$	770 fb	276807	1365	787	5	3	3	2	0.5 × 10 ⁻²	
$t\bar{t}h(1j)$	700 fb	231064	1893	1027	2	2	2	2	0.6 × 10 ⁻²	
$t/\bar{t}W^\pm(1j)$	64 pb	6518431	7596	5801	13	9	3	3	2.9 × 10 ⁻²	
$P_1 P_2 P_3(1j)$ ($P_i \in W Z h$)	500 fb	313350	1475	1093	10	5	4	2	0.3 × 10 ⁻²	
$P_1 P_2 + 1j/2j$ ($P_i \in Z h$)	5.5 pb	738779	2927	2646	3	3	3	3	2.2 × 10 ⁻²	
Total Background									16.1 × 10 ⁻²	

$hh/hZ/ZZ$

\cancel{p}_T M_{T2}

DG, PRD 2013

Boosted dibosons from mixed heavy top squarks

Diptimoy Ghosh*

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(Received 14 August 2013; published 19 December 2013)

The lighter mass eigenstate (\tilde{t}_1) of the two top squarks, the scalar superpartners of the top quark, is extremely difficult to discover if it is almost degenerate with the lightest neutralino ($\tilde{\chi}_1^0$), the lightest stable supersymmetric particle in the R-parity conserving supersymmetry. The current experimental bound on \tilde{t}_1 mass in this scenario stands only around 200 GeV. For such a light \tilde{t}_1 , the heavier top squark (\tilde{t}_2) can also be around the TeV scale. Moreover, the high value of the Higgs (h) mass prefers the left- and right-handed top squarks to be highly mixed, allowing the possibility of a considerable branching ratio for $\tilde{t}_2 \rightarrow \tilde{t}_1 h$ and $\tilde{t}_2 \rightarrow \tilde{t}_1 Z$. In this paper, we explore the above possibility together with the pair production of $\tilde{t}_2 \tilde{t}_2^*$, giving rise to the spectacular diboson + missing transverse energy final state. For an approximately 1 TeV \tilde{t}_2 and a few hundred GeV \tilde{t}_1 the final state particles can be moderately boosted, which encourages us to propose a novel search strategy employing the jet substructure technique to tag the boosted h and Z . The reconstruction of the h and Z momenta also allows us to construct the transverse mass M_{T2} , providing an additional efficient handle to fight the backgrounds. We show that a 4–5 σ signal can be observed at the 14 TeV LHC for ~ 1 TeV \tilde{t}_2 with 100 fb⁻¹ integrated luminosity.

Thank you very much!

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Phys.Rev. D88 (2013) 115013, arXiv:1308.0320

Thank you very much!

Additional Material

M_{T2}

Let us define m_T first : Consider the decay $A \rightarrow B \tilde{\chi}$

$$m_A^2 = m_B^2 + m_{\tilde{\chi}}^2 + 2(E_{TA}E_{T\tilde{\chi}} \cosh(\Delta\eta) - \mathbf{p}_{TB} \cdot \mathbf{p}_{T\tilde{\chi}})$$

$$\cosh(\Delta\eta) \geq 1$$



$$\eta = \frac{1}{2} \ln[(E + p_z)/(E - p_z)]$$

$$E_T = \sqrt{\mathbf{p}_T^2 + m^2}$$

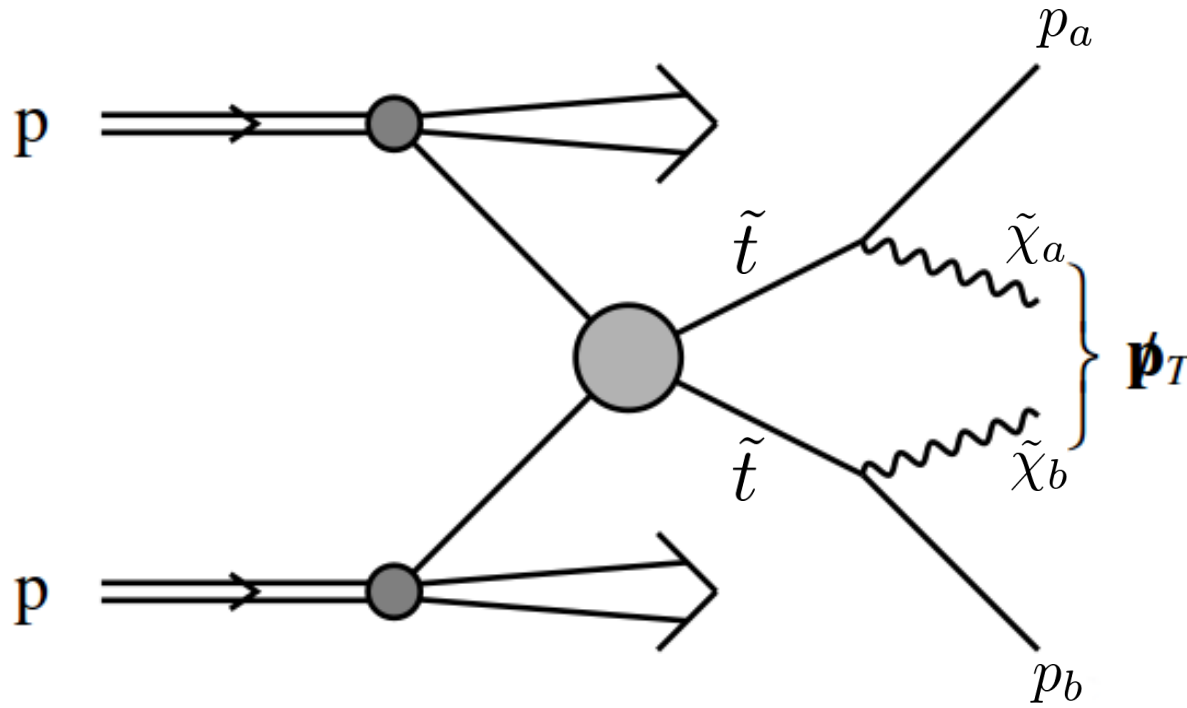
$$m_A^2 \geq m_B^2 + m_{\tilde{\chi}}^2 + 2(E_{TA}E_{T\tilde{\chi}} - \mathbf{p}_{TB} \cdot \mathbf{p}_{T\tilde{\chi}})$$



$$m_A \geq \sqrt{m_B^2 + m_{\tilde{\chi}}^2 + 2(E_{TA}E_{T\tilde{\chi}} - \mathbf{p}_{TB} \cdot \mathbf{p}_{T\tilde{\chi}})}$$

$$m_T(\mathbf{p}_{TB}, \mathbf{p}_{T\tilde{\chi}}) \quad \text{for given } m_B, m_{\tilde{\chi}}$$

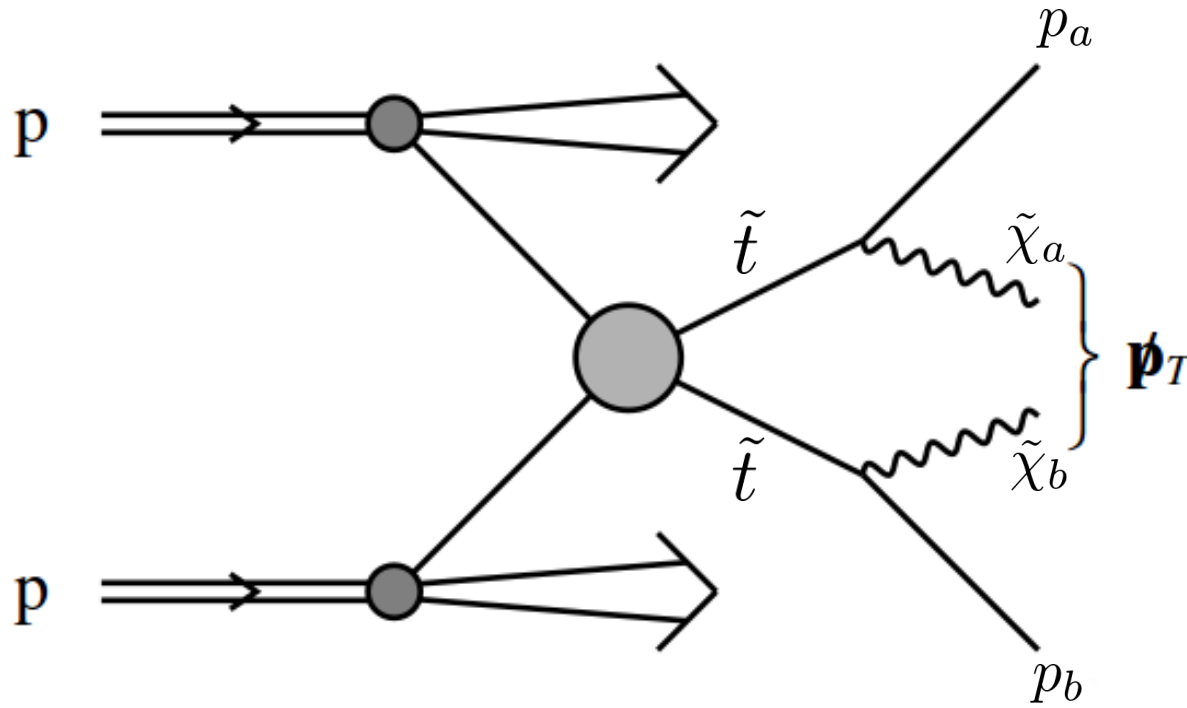
Now consider a decay with two sources of missing energy :



C. G. Lester and D. J. Summers, 1999

$$\mathbf{p}_T = \mathbf{p}_{T\tilde{\chi}_a} + \mathbf{p}_{T\tilde{\chi}_b}$$

Now consider a decay with two sources of missing energy :

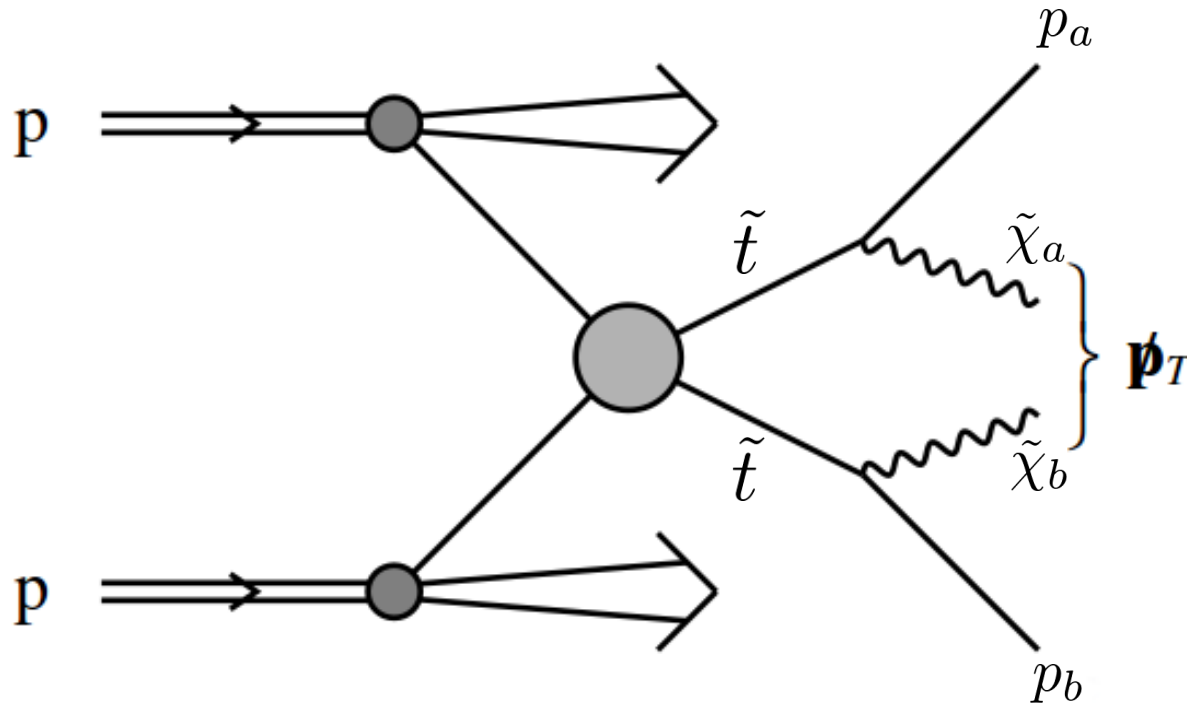


C. G. Lester and D. J. Summers, 1999

$$\mathbf{p}_T = \mathbf{p}_{T\tilde{\chi}_a} + \mathbf{p}_{T\tilde{\chi}_b}$$

$$m_{\tilde{t}} \geq m_T(\mathbf{p}_{T a}, \mathbf{p}_{T\tilde{\chi}_a})$$

Now consider a decay with two sources of missing energy :

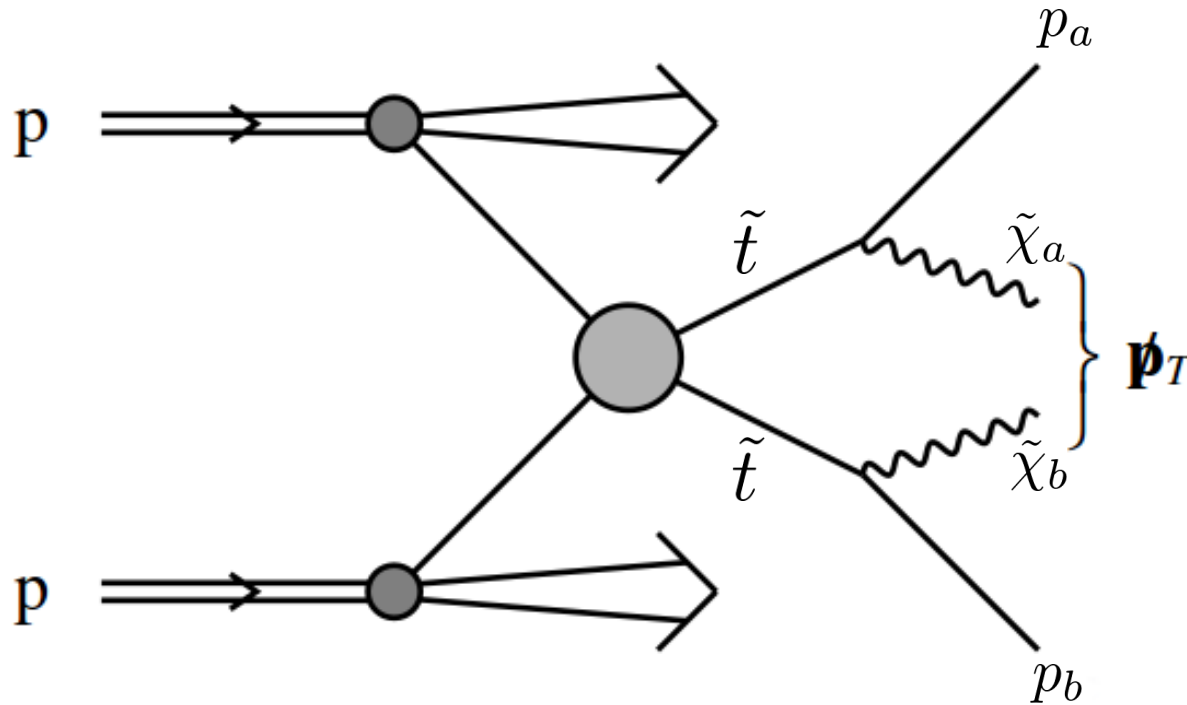


C. G. Lester and D. J. Summers, 1999

$$\mathbf{p}_T = \mathbf{p}_{T\tilde{\chi}_a} + \mathbf{p}_{T\tilde{\chi}_b}$$

$$m_{\tilde{t}} \geq \{m_T(\mathbf{p}_{T a}, \mathbf{p}_{T\tilde{\chi}_a}), m_T(\mathbf{p}_{T b}, \mathbf{p}_{T\tilde{\chi}_b})\}$$

Now consider a decay with two sources of missing energy :

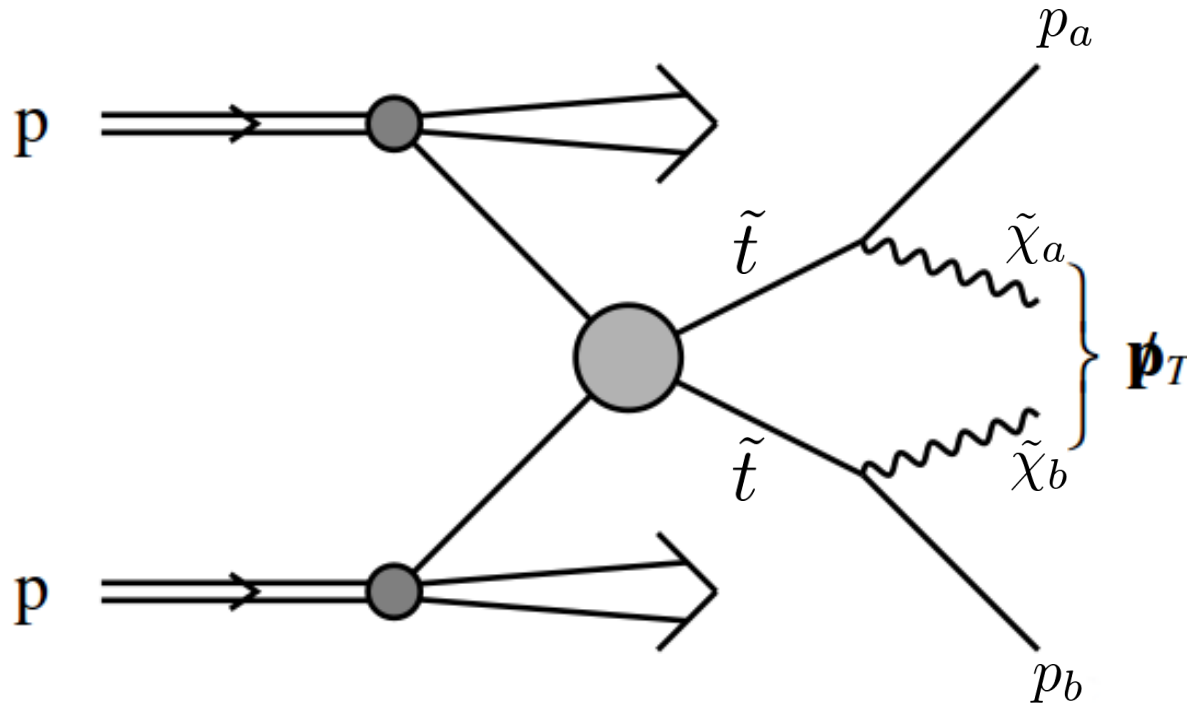


C. G. Lester and D. J. Summers, 1999

$$\mathbf{p}_T = \mathbf{p}_{T\tilde{\chi}_a} + \mathbf{p}_{T\tilde{\chi}_b}$$

$$m_{\tilde{t}} \geq \max\{m_T(\mathbf{p}_{T a}, \mathbf{p}_{T\tilde{\chi}_a}), m_T(\mathbf{p}_{T b}, \mathbf{p}_{T\tilde{\chi}_b})\}$$

Now consider a decay with two sources of missing energy :



C. G. Lester and D. J. Summers, 1999

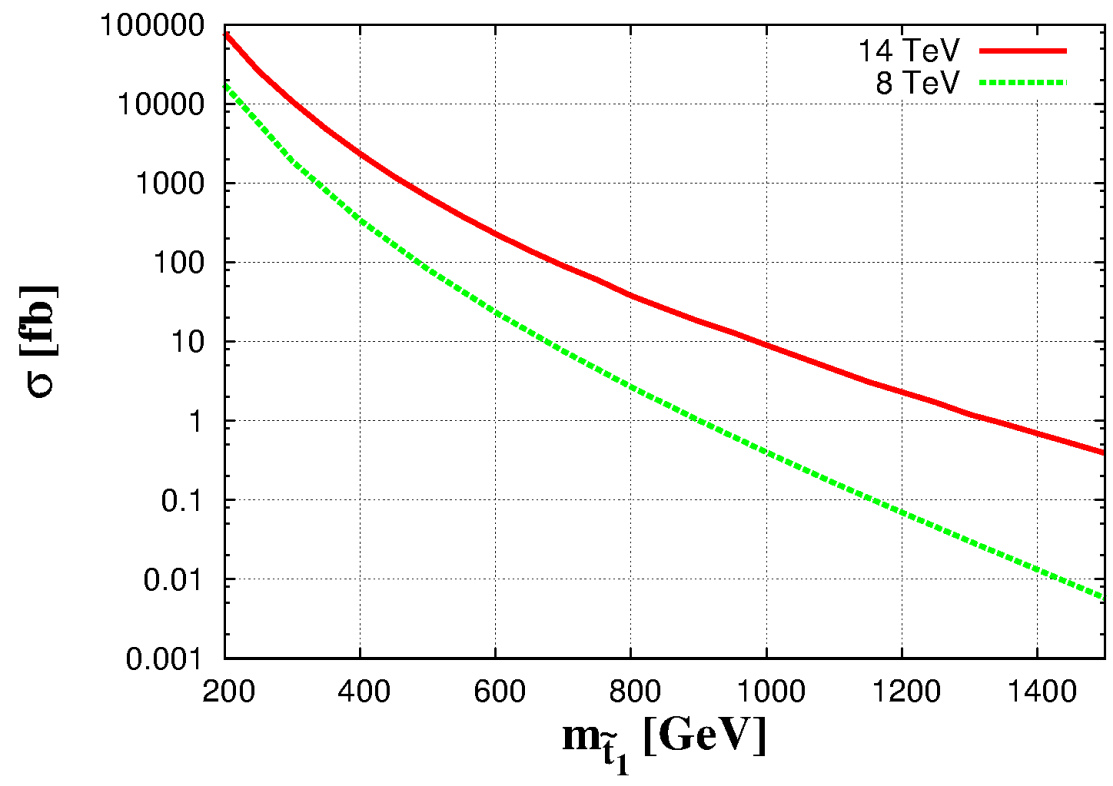
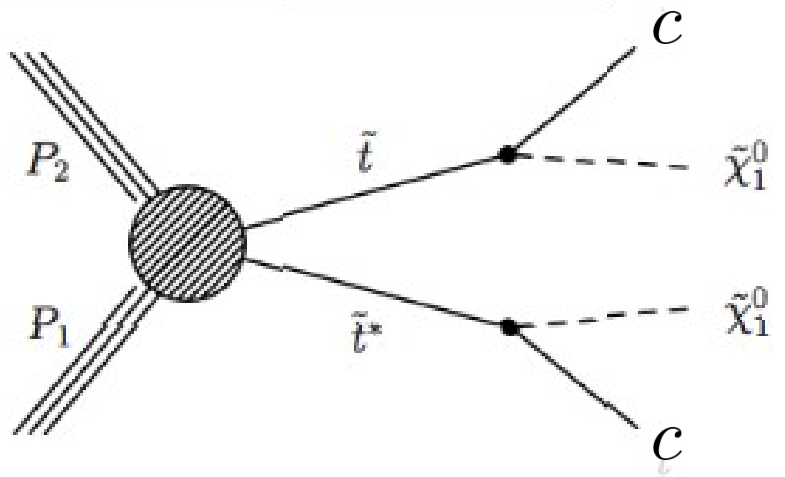
$$\not{p}_T = \mathbf{p}_{T\tilde{\chi}_a} + \mathbf{p}_{T\tilde{\chi}_b}$$

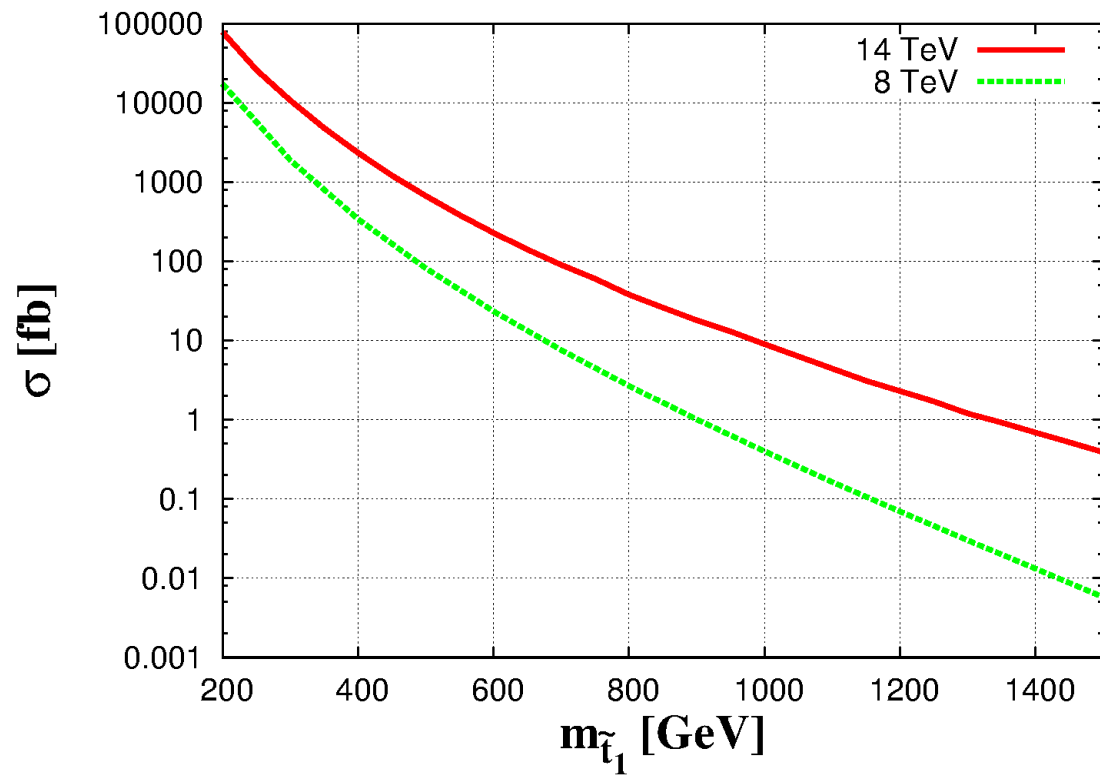
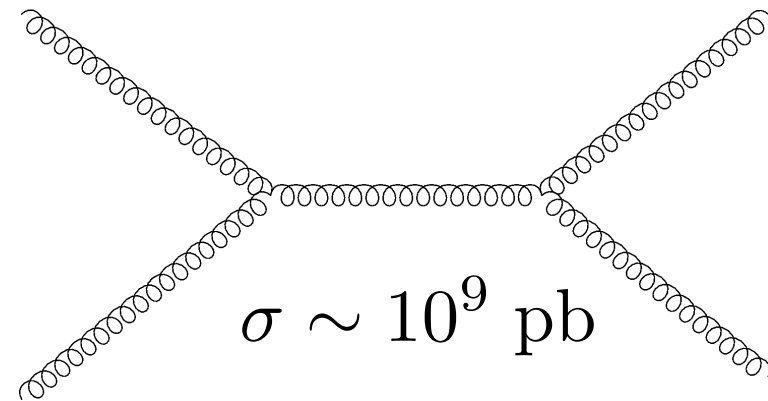
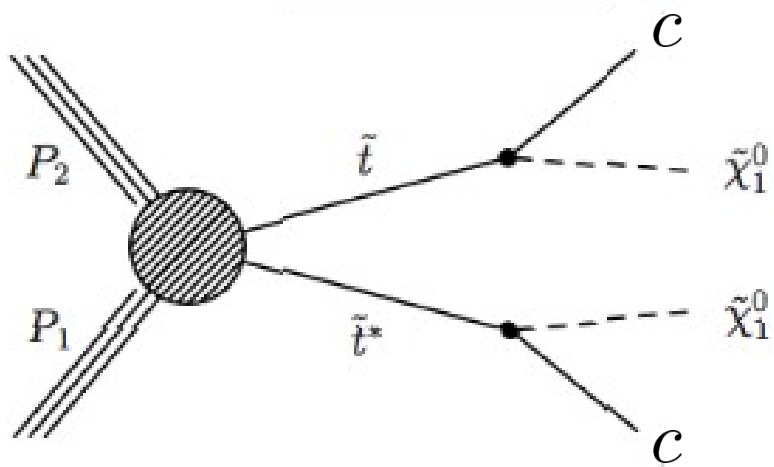
$$m_{\tilde{t}} \geq \max\{m_T(\mathbf{p}_{Ta}, \mathbf{p}_{T\tilde{\chi}_a}), m_T(\mathbf{p}_{Tb}, \mathbf{p}_{T\tilde{\chi}_b})\}$$

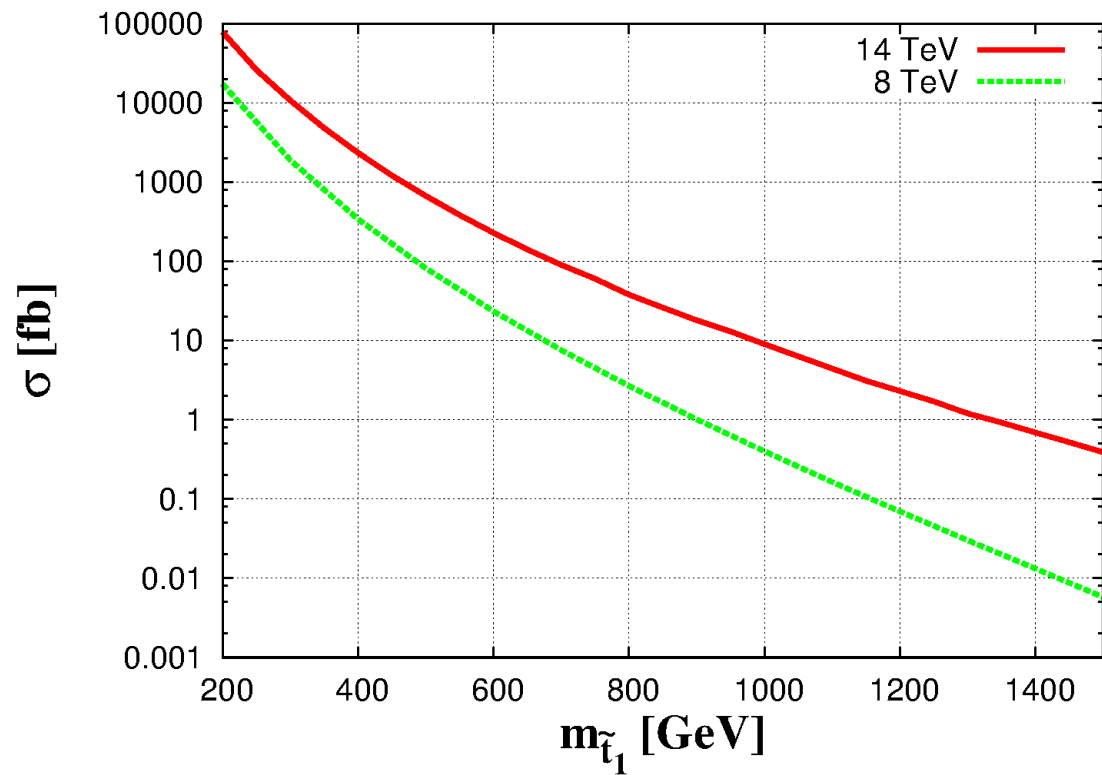
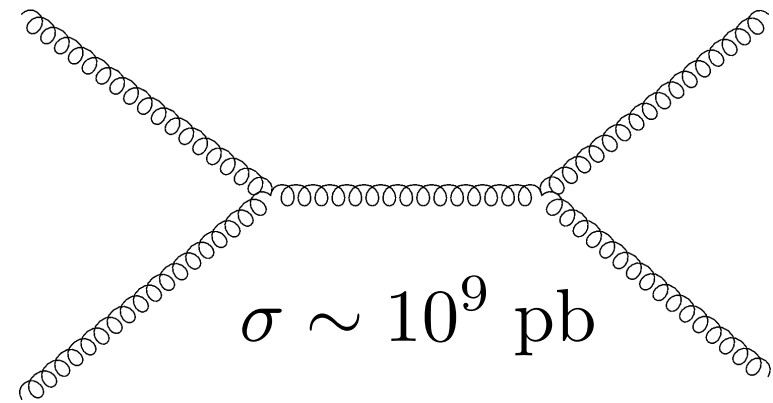
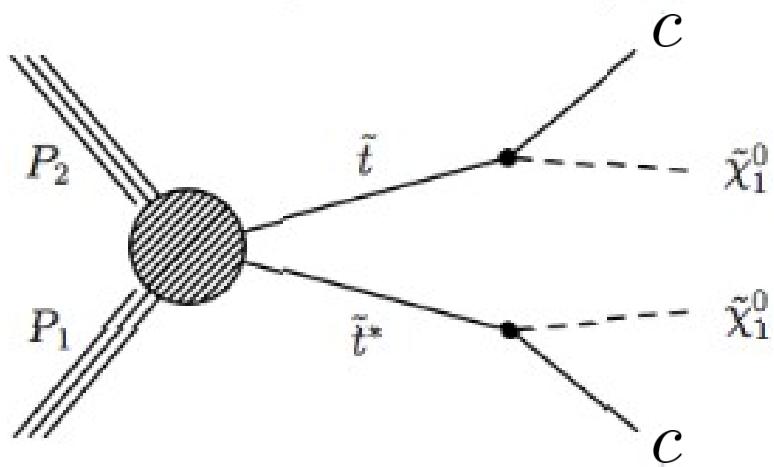


$$m_{\tilde{t}} \geq \min_{\not{p}_T = \not{p}_{T1} + \not{p}_{T2}} \left[\max\{m_T(\mathbf{p}_{Ta}, \not{p}_{T1}), m_T(\mathbf{p}_{Tb}, \not{p}_{T2})\} \right]$$

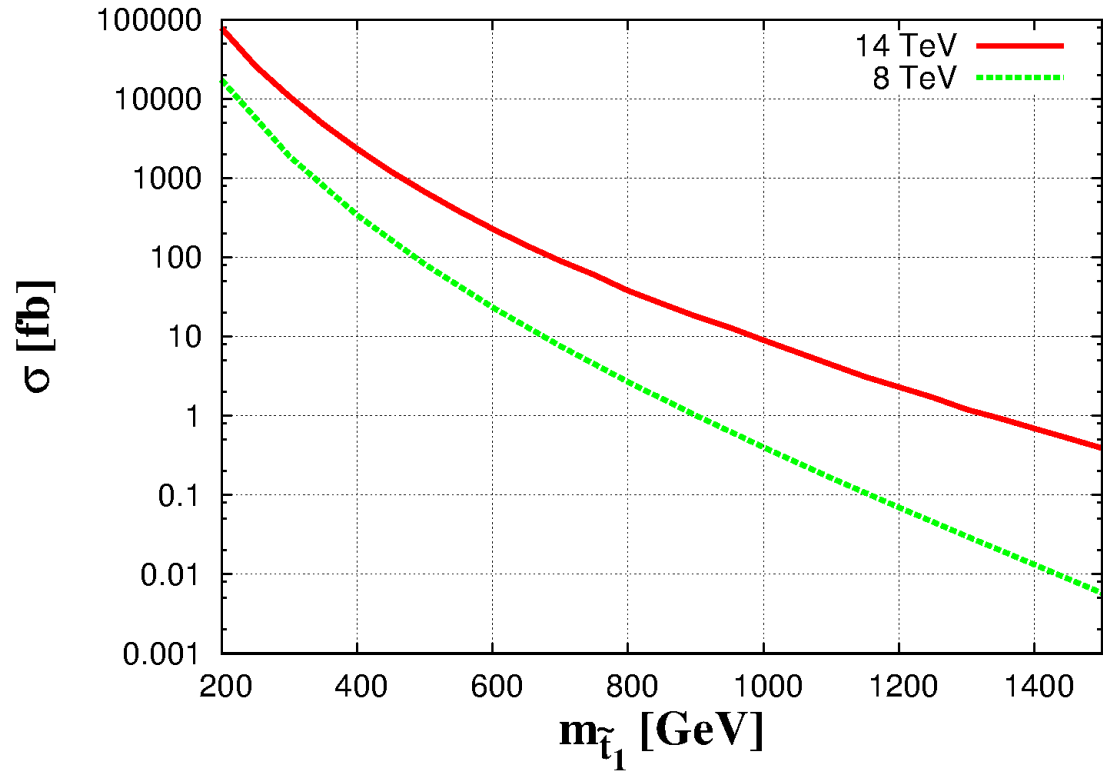
$M_{T2}(\mathbf{p}_{Ta}, \mathbf{p}_{Tb}, \not{p}_T)$



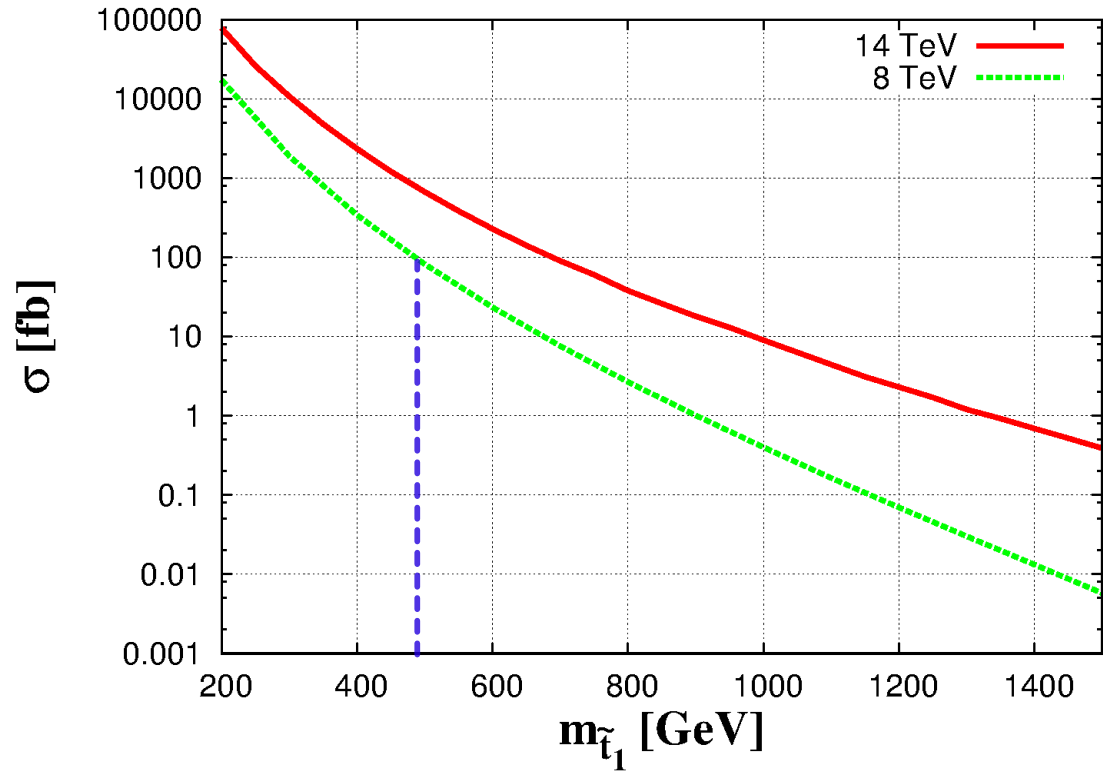




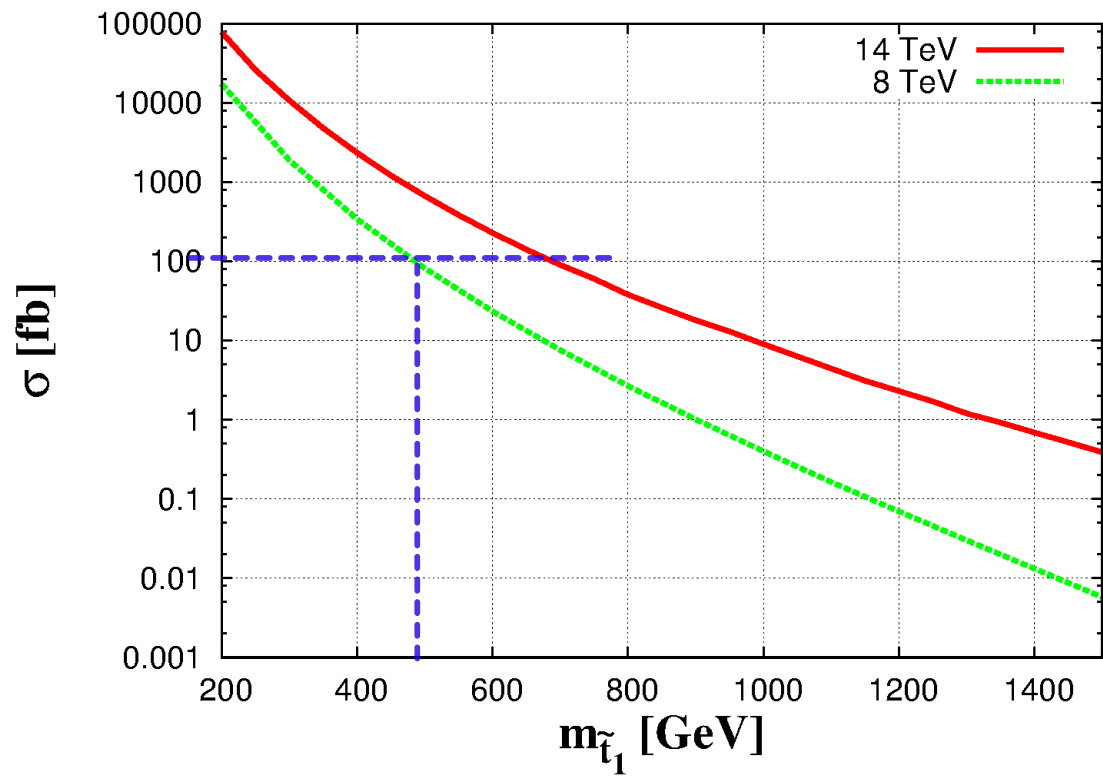
8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



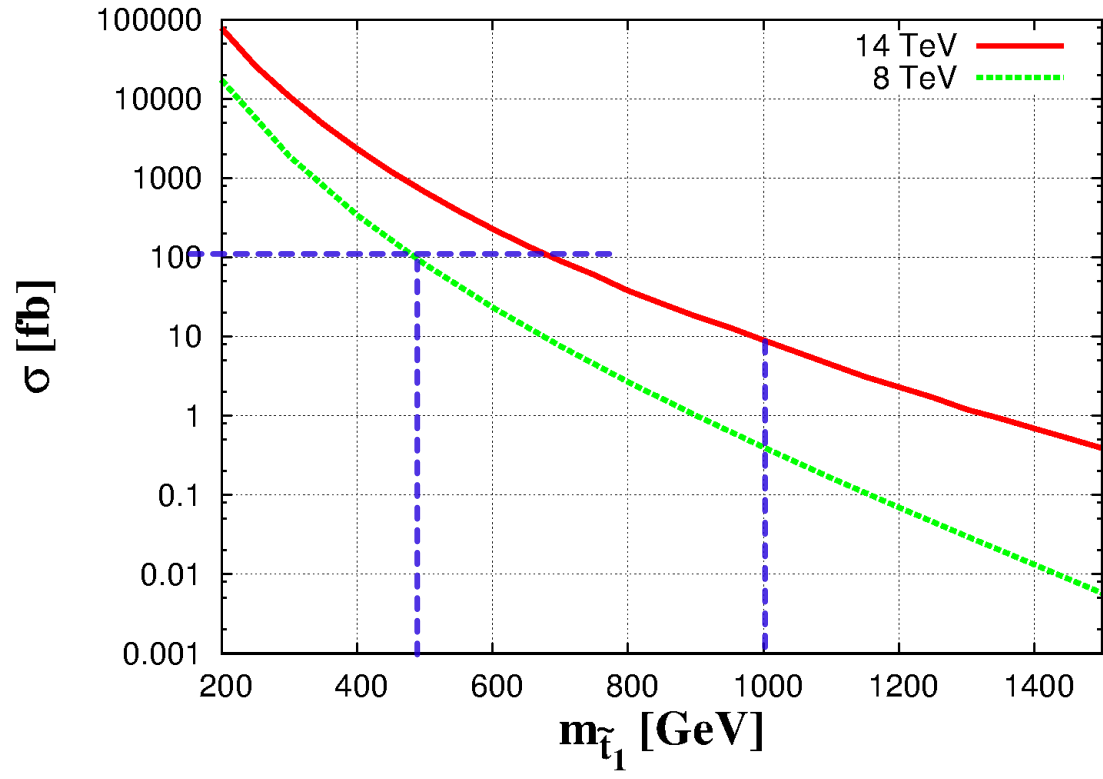
8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



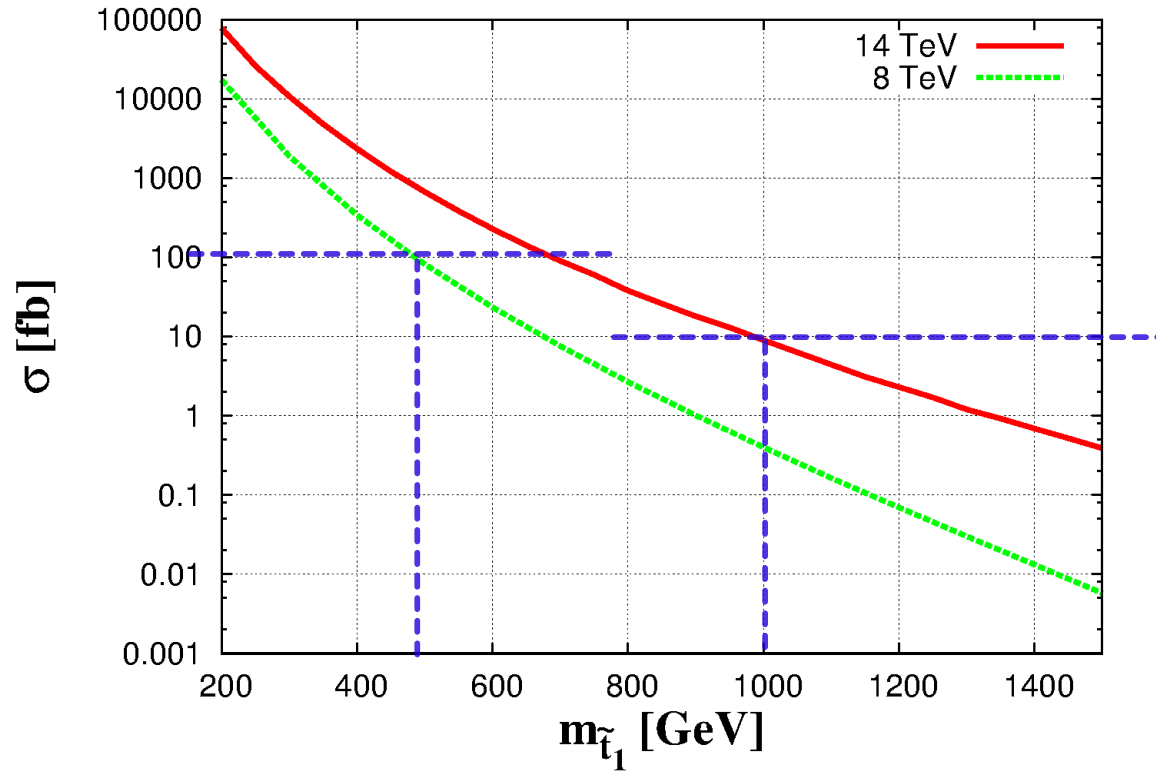
8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



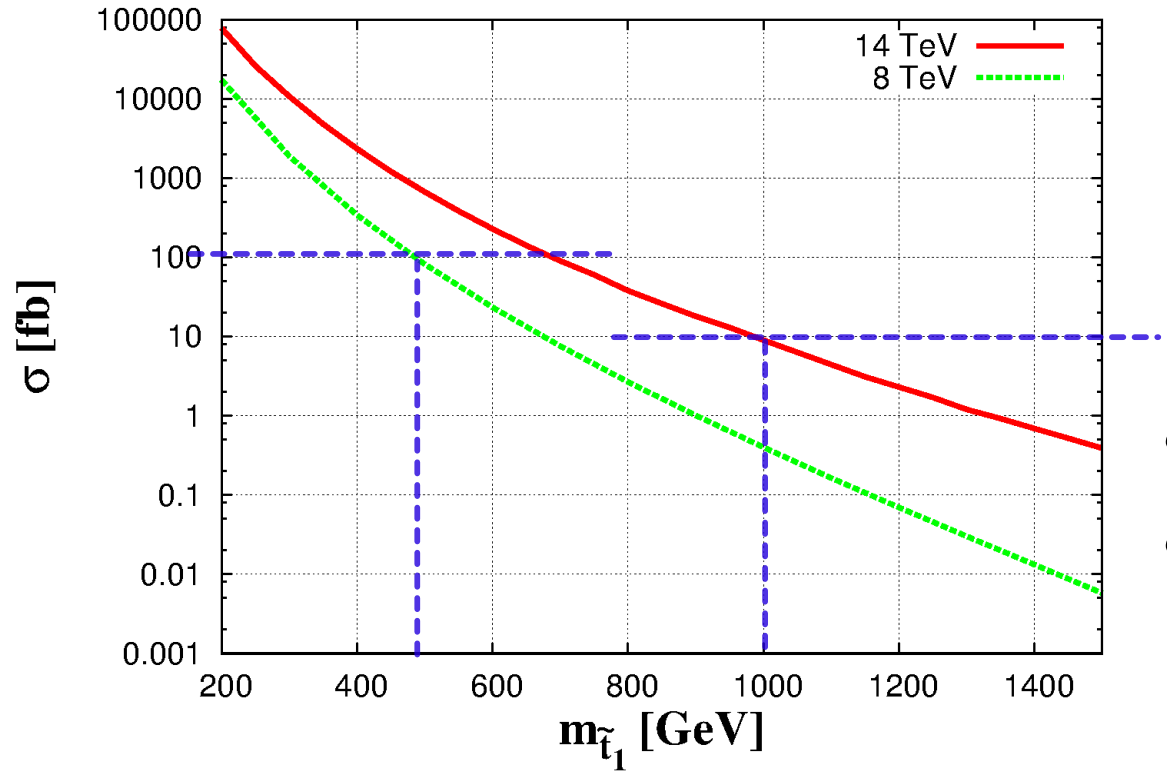
8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$

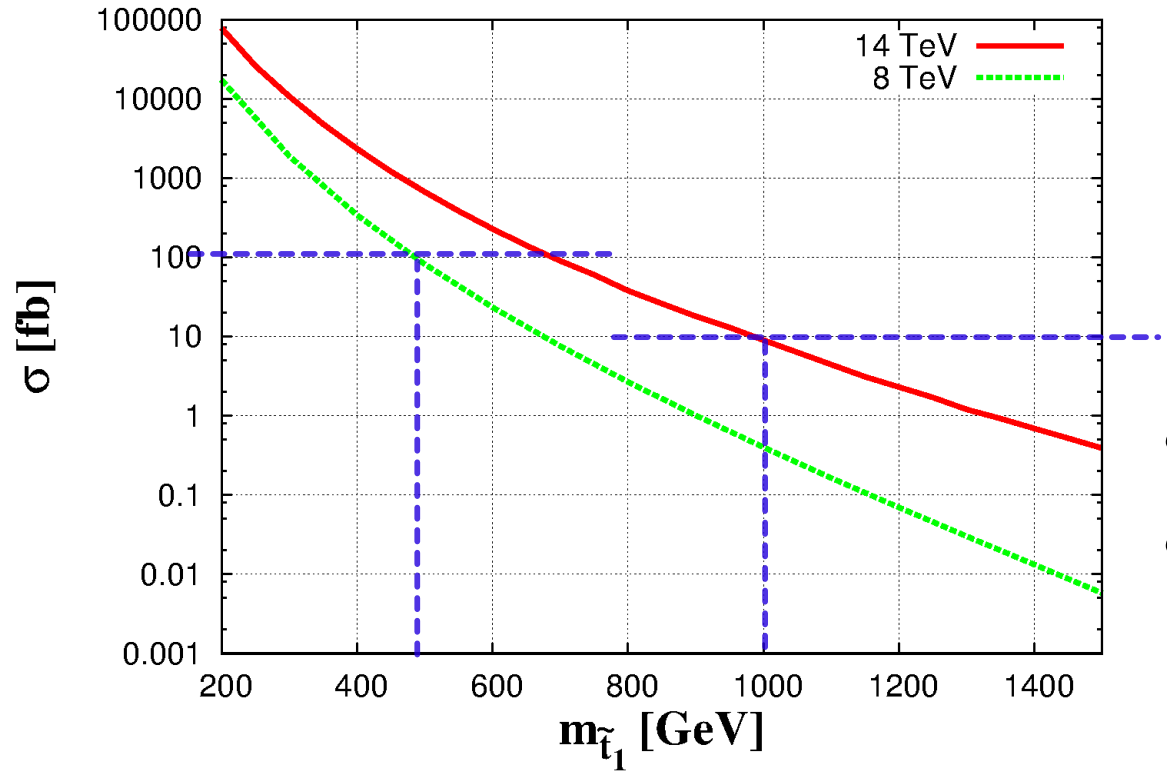


8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



$\sigma_{t\bar{t}} \sim 225 \text{ pb}$ 8TeV
 $\sigma_{t\bar{t}} \sim 925 \text{ pb}$ 14TeV

8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



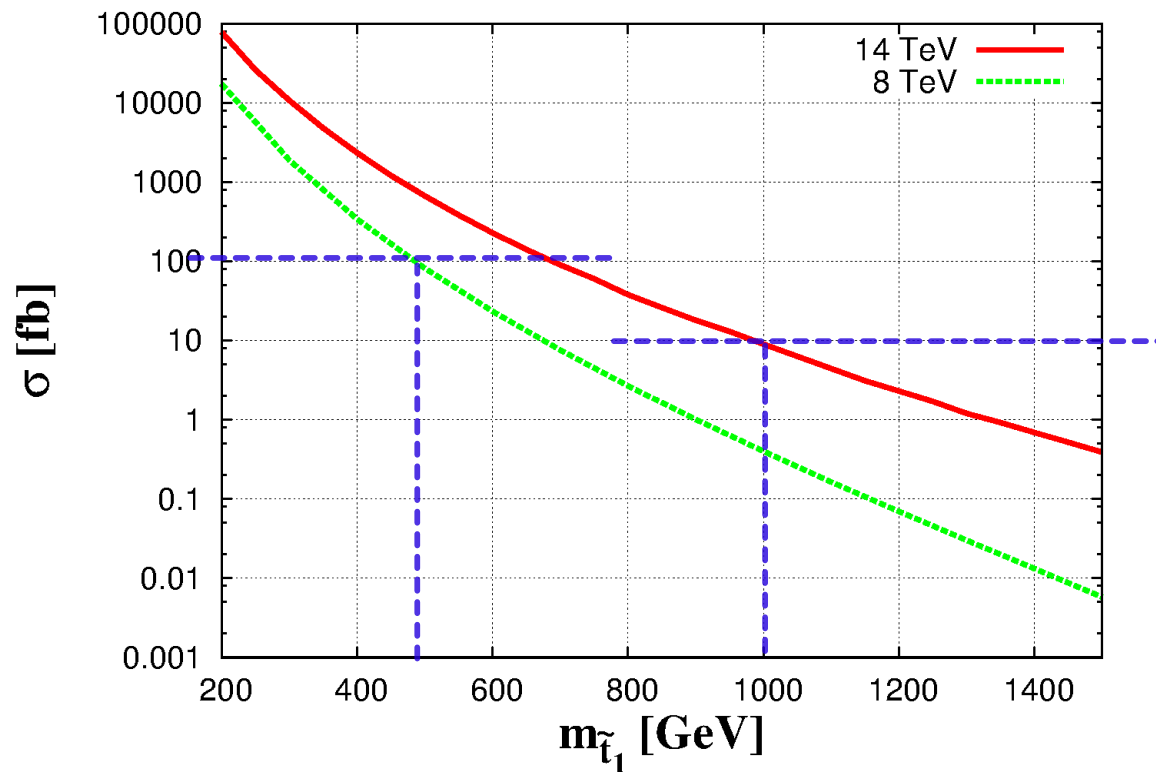
$\sigma_{t\bar{t}} \sim 225 \text{ pb} \quad 8\text{TeV}$
 $\sigma_{t\bar{t}} \sim 925 \text{ pb} \quad 14\text{TeV}$

$$S \quad \dashrightarrow \quad \frac{S}{10}$$

8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$

vs.

14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



$\sigma_{t\bar{t}} \sim 225 \text{ pb}$ 8TeV

$\sigma_{t\bar{t}} \sim 925 \text{ pb}$ 14TeV

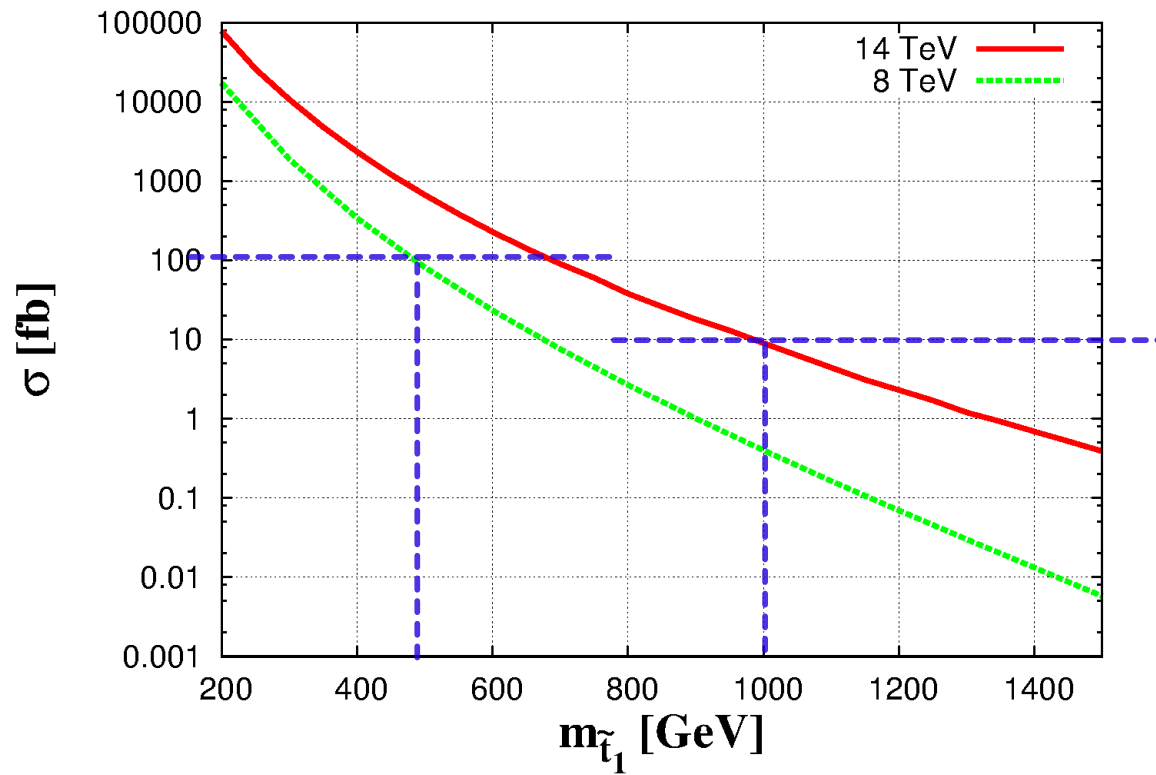
$$S \dashrightarrow \frac{S}{10}$$

$$B \dashrightarrow B \times 4$$

8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$

vs.

14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



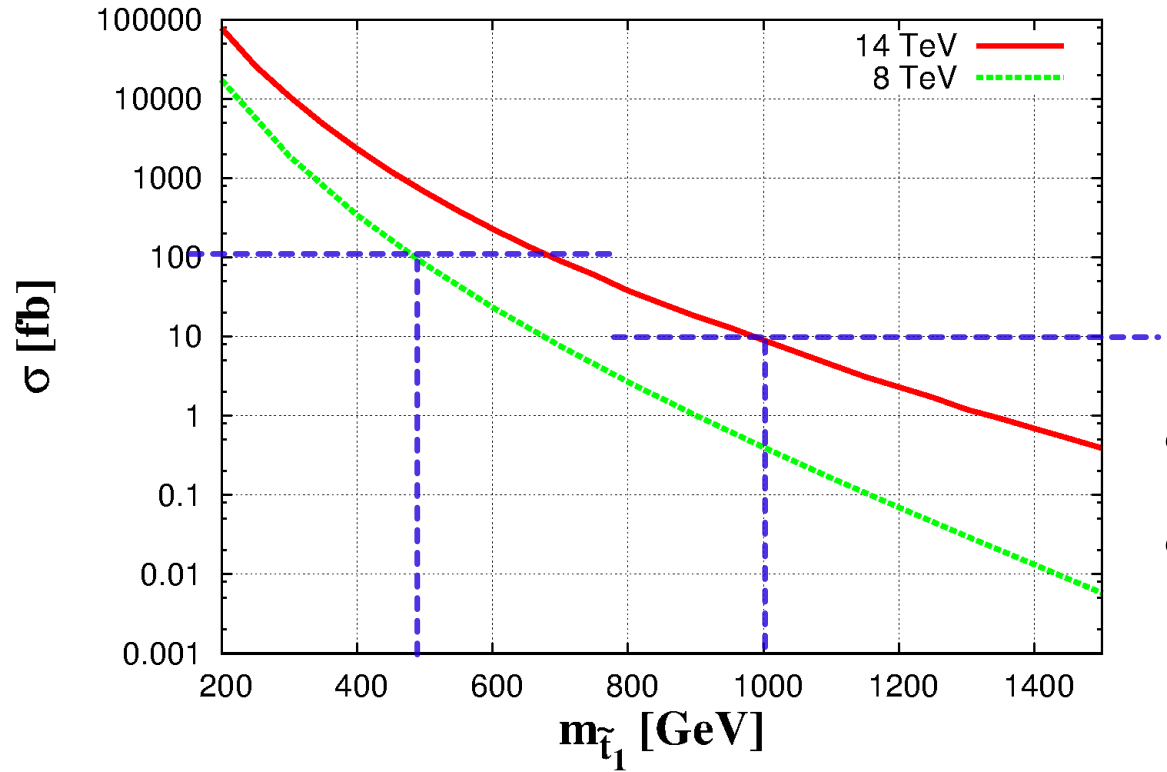
$\sigma_{t\bar{t}} \sim 225 \text{ pb}$ 8TeV

$\sigma_{t\bar{t}} \sim 925 \text{ pb}$ 14TeV

$$S \dashrightarrow \frac{S}{10} \times 4$$

$$B \dashrightarrow B \times 4$$

8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$ **vs.** 14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



$\sigma_{t\bar{t}} \sim 225 \text{ pb}$ 8TeV
 $\sigma_{t\bar{t}} \sim 925 \text{ pb}$ 14TeV

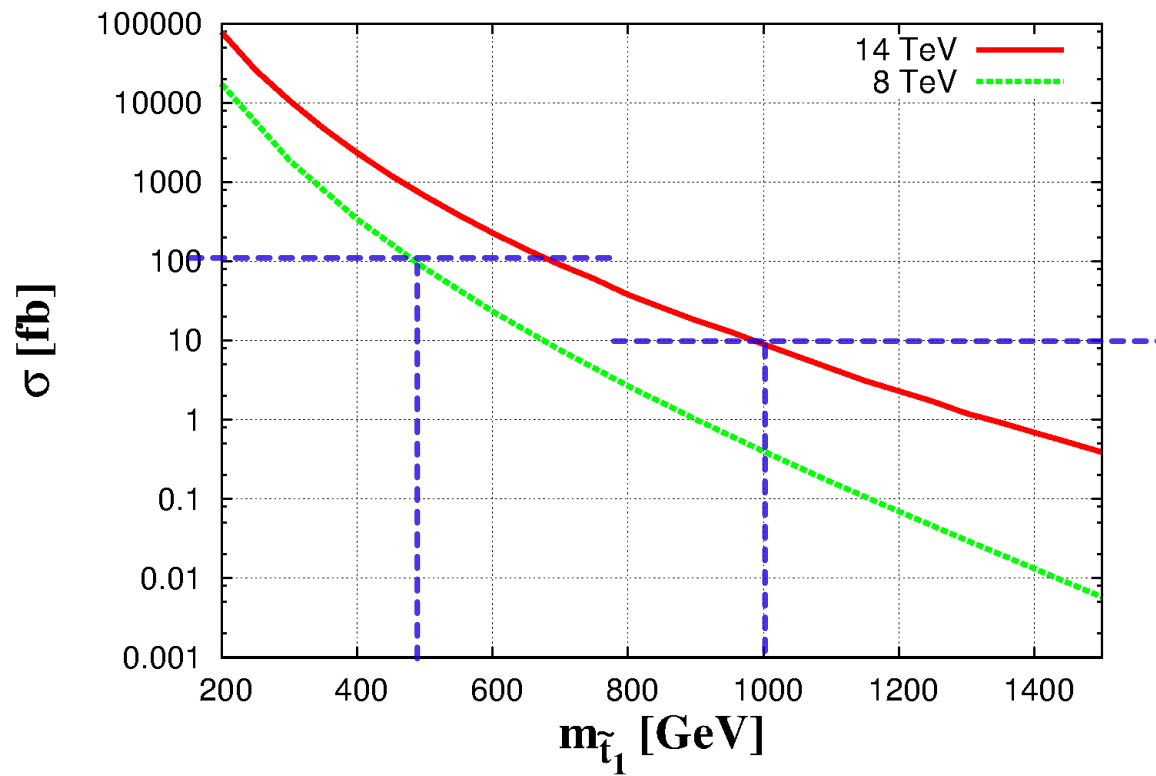
$$S \dashrightarrow \frac{S}{10} \times 4$$

$$B \dashrightarrow B \times 4 \times 4$$

8TeV, 25fb^{-1} , $m_{\tilde{t}_1} = 500\text{GeV}$

vs.

14TeV, 100fb^{-1} , $m_{\tilde{t}_1} = 1000\text{GeV}$



$\sigma_{t\bar{t}} \sim 225 \text{ pb}$ 8TeV

$\sigma_{t\bar{t}} \sim 925 \text{ pb}$ 14TeV

$$S \dashrightarrow \frac{S}{10} \times 4$$

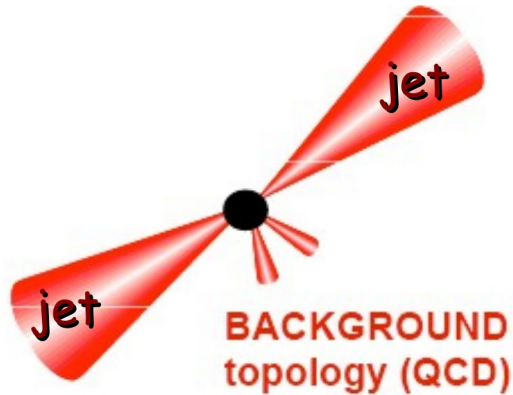
$$B \dashrightarrow B \times 4 \times 4$$

$$\frac{S}{\sqrt{B}} \dashrightarrow \frac{1}{10} \times \frac{S}{\sqrt{B}}$$

$$\alpha_T = \frac{p_T^{j2}}{\sqrt{H_T^2 - \cancel{H}_T^2}}$$

$$H_T = \sum_j |\vec{p}_T^j|$$

$$\cancel{H}_T = \left| \sum_j \vec{p}_T^j \right|$$



Jets are back-to-back in ϕ

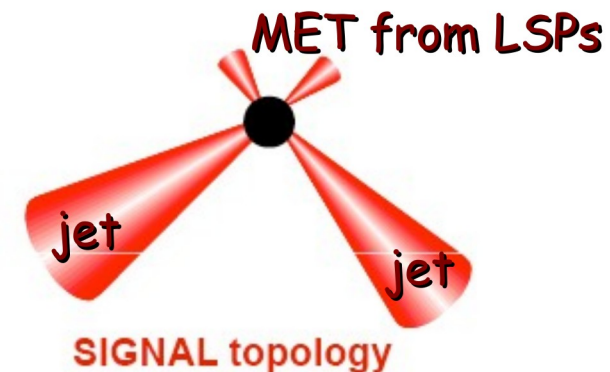
$$\alpha_T = 0.5$$

in the case of an imbalance in the measured p_T s of back-to-back jets

$$\alpha_T < 0.5$$

when the two jets are not back-to-back and balancing genuine MET

$$\alpha_T > 0.5$$



Dark matter relic density

