

Recent accelerator long baseline neutrino oscillation results

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& the T2K Collaboration

1. Introduction
2. Recent T2K & MINOS results
3. Prospects for the next few years
4. Longer term future (10⁺ years)

Neutrino physics – a constant source of surprise & mystery!

Compared with the quarks and charged leptons, why are neutrino masses so small?

Neutrino mixing angles are generally much larger than quark mixing angles; large CP violation effects are **allowed** in the neutrino sector

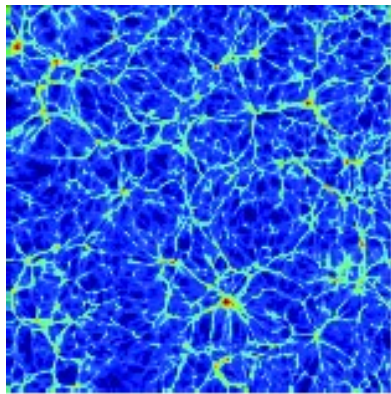
Quarks/CKM

$$V_q = \begin{pmatrix} 0.976 & 0.22 & 0.004 \\ -0.22 & 0.98 & 0.04 \\ 0.007 & -0.04 & 1 \end{pmatrix}$$

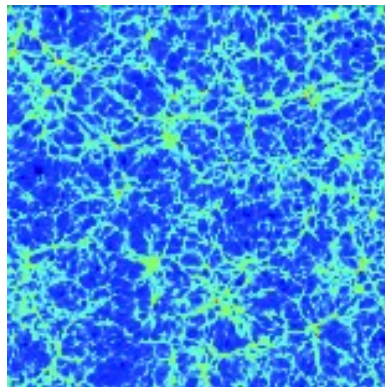
Neutrinos/PMNS

$$U_\ell = \begin{pmatrix} 0.85 & -0.54 & 0.16 \\ 0.33 & 0.62 & -0.72 \\ -0.40 & -0.59 & -0.70 \end{pmatrix}$$

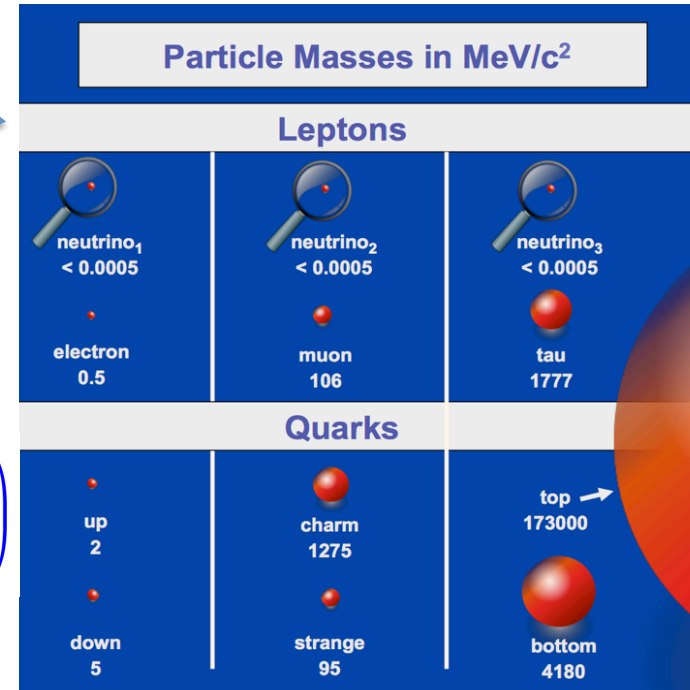
massless ν



$\Sigma \text{Mass} = 1.9 \text{eV}/c^2$



Baryon density



Neutrinos play a fundamental role in the evolution of the Universe

Can neutrinos explain the matter-antimatter asymmetry of the Universe? e.g. via CP violation and Leptogenesis

Agarwal, Feldman 2010

The 3 Neutrino Flavour Paradigm[§]

Pontecorvo-Maki-Nakagawa-Sakata (PMNS) mixing matrix

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Flavour states Mass states

$$s_{ij} = \sin \theta_{ij}$$

δ = CP violating phase

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} = \begin{matrix} \text{“Atmospheric”} \\ \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \end{matrix} \times \begin{matrix} \text{“Reactor”} \\ \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \end{matrix} \times \begin{matrix} \text{“Solar”} \\ \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{matrix}$$

Oscillation phenomena have been convincingly observed using **solar, atmospheric, reactor and accelerator neutrinos**, establishing the 3 neutrino flavour SM paradigm

(§ ignoring Majorana phases)

Neutrino Properties

Known:

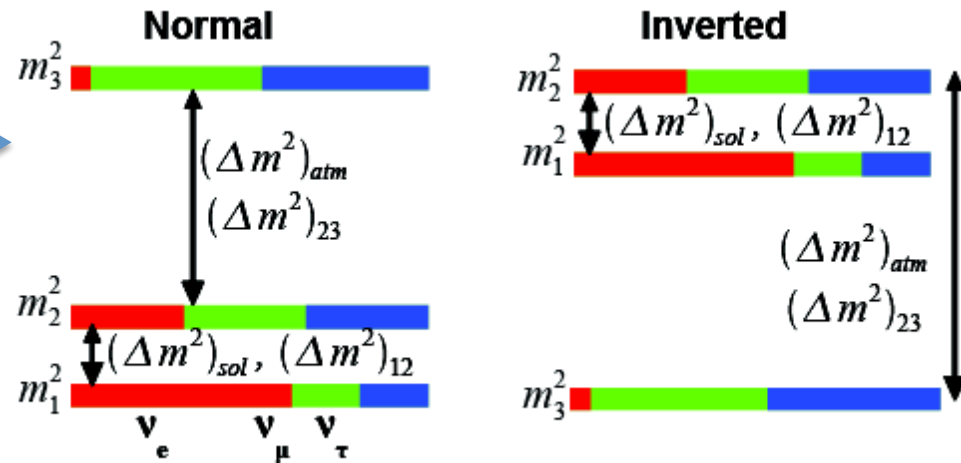
- have mass
- oscillate between flavours
- measured values for
 - $\theta_{12}, \theta_{23}, \theta_{13}$
 - $\Delta m_{21}^2, |\Delta m_{32}^2|$

Parameter	Value	Precision (%)
Δm_{21}^2	$7.5 \cdot 10^{-5} \text{ eV}^2$	2.6
θ_{12}	34°	5.4
Δm_{32}^2	$2.4 \cdot 10^{-3} \text{ eV}^2$	2.6
θ_{23}	42°	~ 10
θ_{13}	9°	8.5

Capozzi et al.
ArXiv:1312.2878
Talk by M. Tortola

Unknown:

- order of the mass states (hierarchy) →
- absolute mass of the neutrino
- Is θ_{23} maximal ($=45^\circ$)? Or which octant?
- are there more than 3 neutrino types?
- Majorana or Dirac?
- the value of δ_{CP}



Neutrino Mass Hierarchy

Neutrino Oscillation Measurements at Accelerators

Muon neutrino disappearance

(in the absence of matter effects)

$$P(\nu_\mu \rightarrow \nu_\mu) \simeq 1 - 4 \cos^2(\theta_{13}) \sin^2(\theta_{23}) \left[1 - \cos^2(\theta_{13}) \right. \\ \left. \times \sin^2(\theta_{23}) \right] \sin^2 \frac{\Delta m_{31}^2 L}{4E}$$

Survival probability

Electron neutrino appearance

$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \\ - \frac{\sin 2\theta_{12} \sin 2\theta_{23}}{2 \sin \theta_{13}} \sin \frac{\Delta m_{21}^2 L}{4E} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \sin \delta_{\text{CP}} \\ + (\text{CP even term, solar term, matter effect term})$$

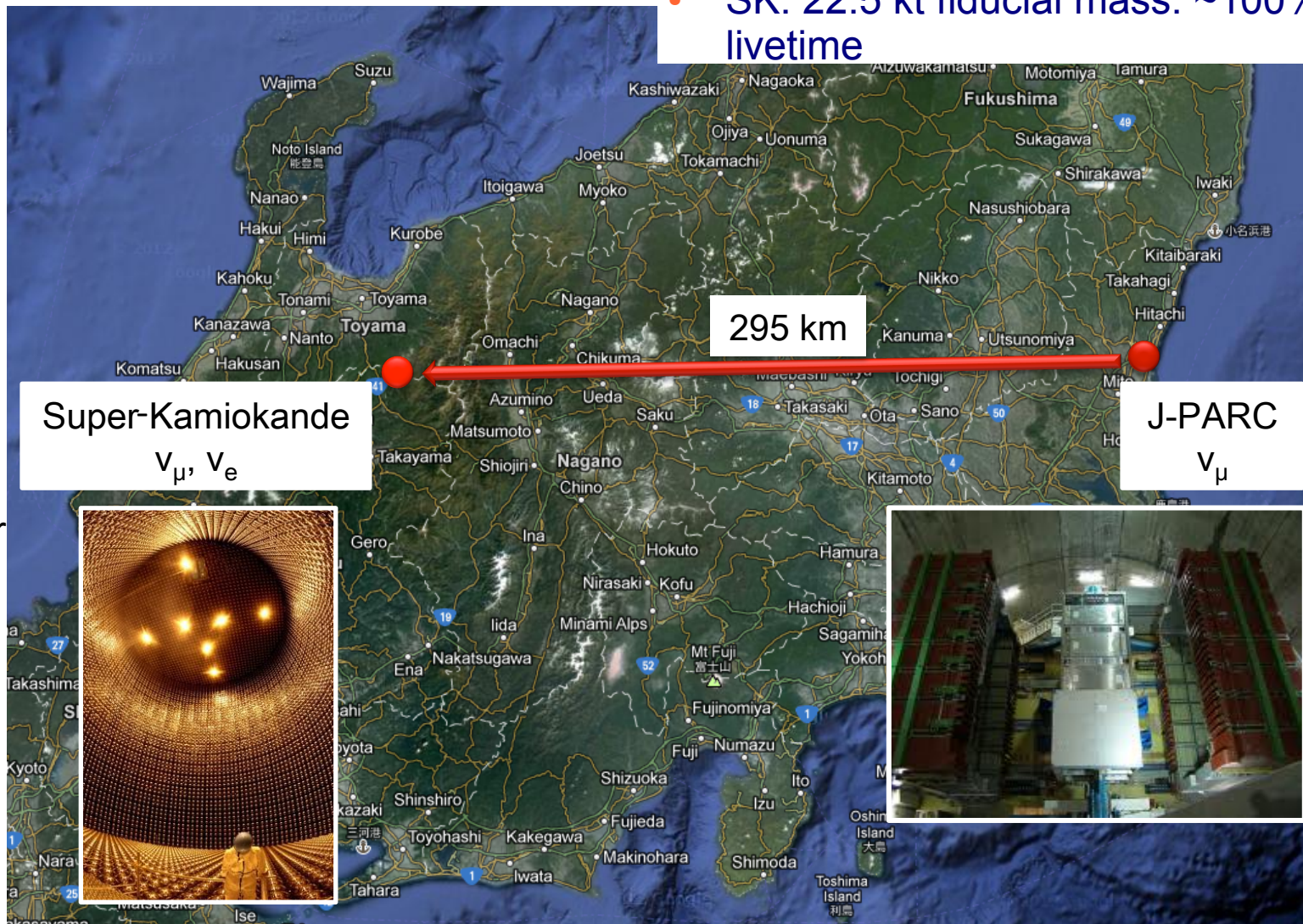
Disappearance channel : sensitivity to θ_{23} and (subleading) to the octant

Appearance channel : sensitivity to θ_{13} and (subleading) to the CP phase

T2K experiment

T2K= Tokai to Kamioka

- Primary proton beam: 30 GeV/c, 235 kW (RUN4) $6.57 \cdot 10^{20}$ Proton On Target (8% of the final design exposure)
- SK: 22.5 kt fiducial mass. $\sim 100\%$ livetime



Super-Kamiokande

ν_μ, ν_e

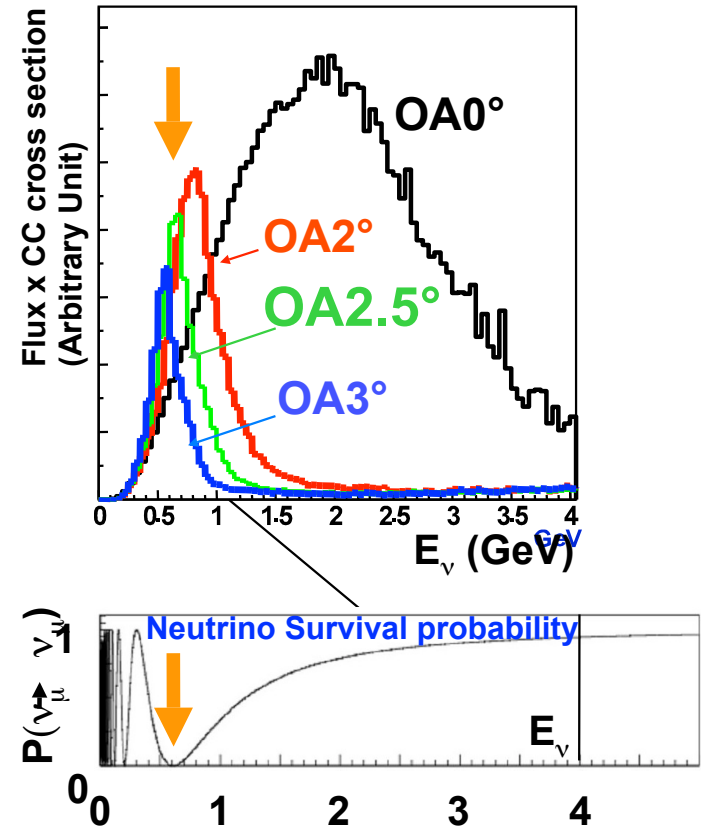
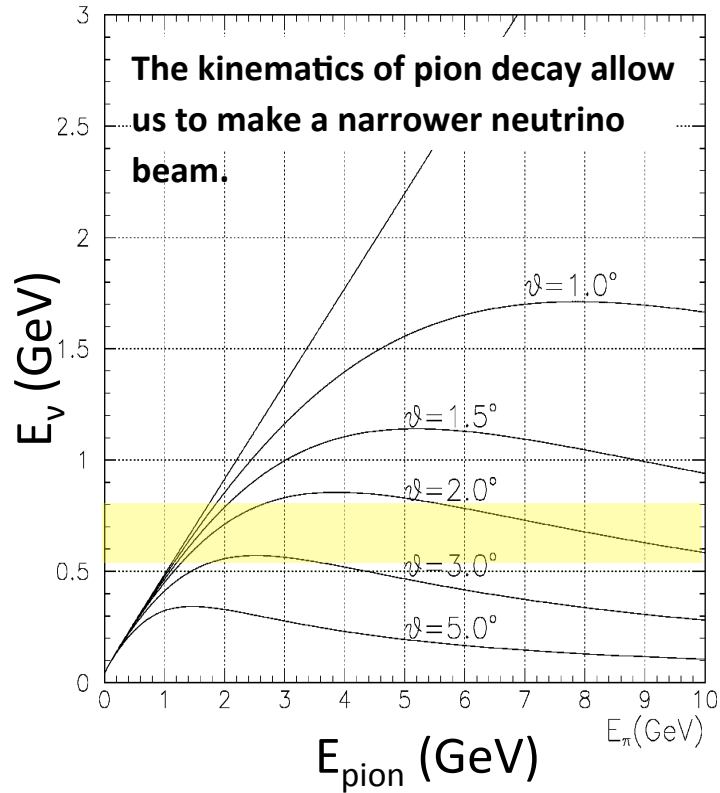
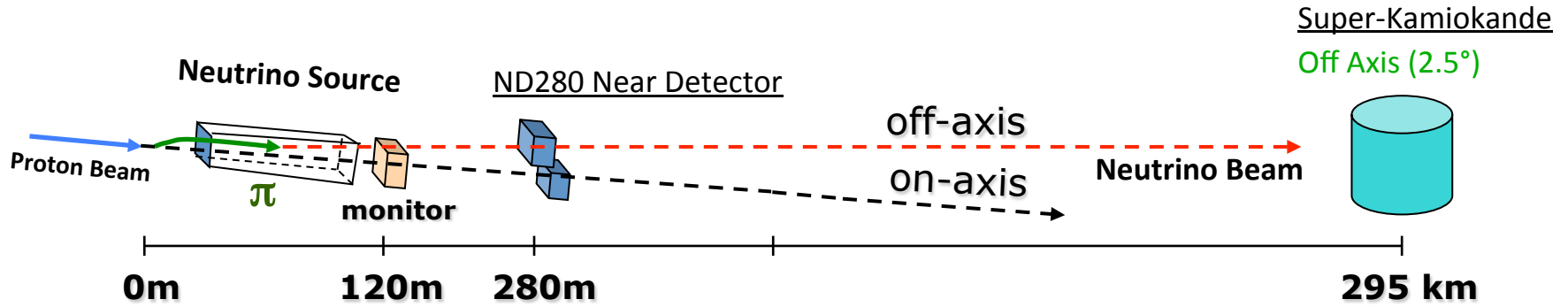
J-PARC

ν_μ

Far
Detector

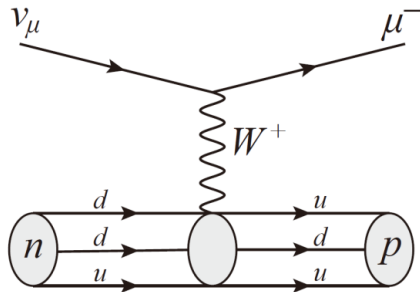
Near
Detector

The T2K Off-Axis Beam



T2K Muon Neutrino (ν_μ) Disappearance Measurement

Phys. Rev. Lett. 112, 181801 (2014)

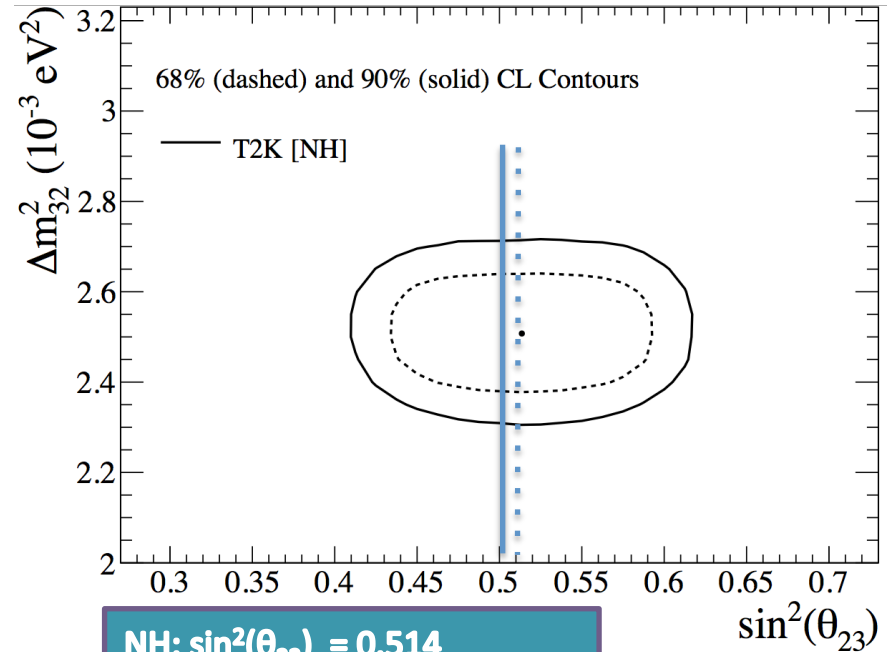
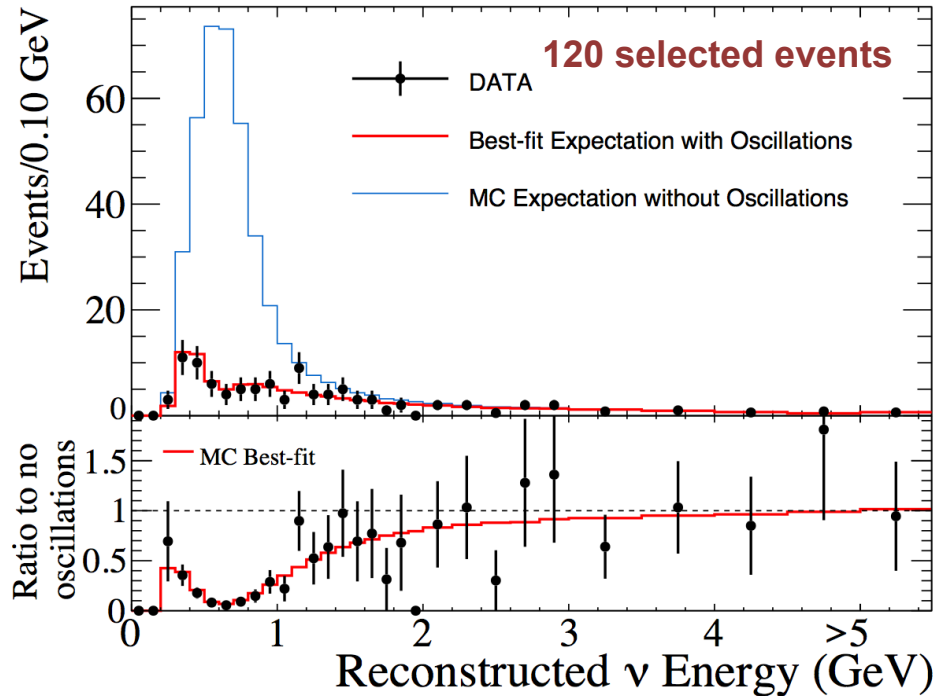


CCQE

Deficit of ν_μ CCQE events at Super-Kamiokande

Maximal mixing is not the same as maximum disappearance if θ_{13} is not zero!

Shows the power of the off-axis technique!



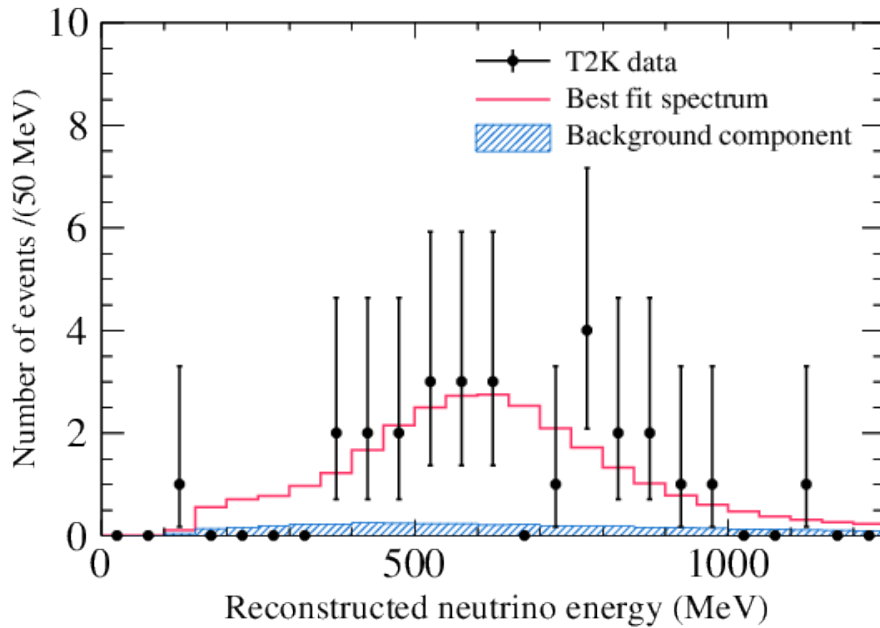
NH: $\sin^2(\theta_{23}) = 0.514$
NH: $\Delta m_{32}^2 = 2.51 \times 10^{-3} \text{ eV}^2$

For θ_{13} given by reactor experiments

T2K Electron Neutrino (ν_e) Appearance Observation

ν_e CCQE @ Super-K

Phys. Rev. Lett. 112, 061802 (2014)



4.92 ± 0.55 events expected background

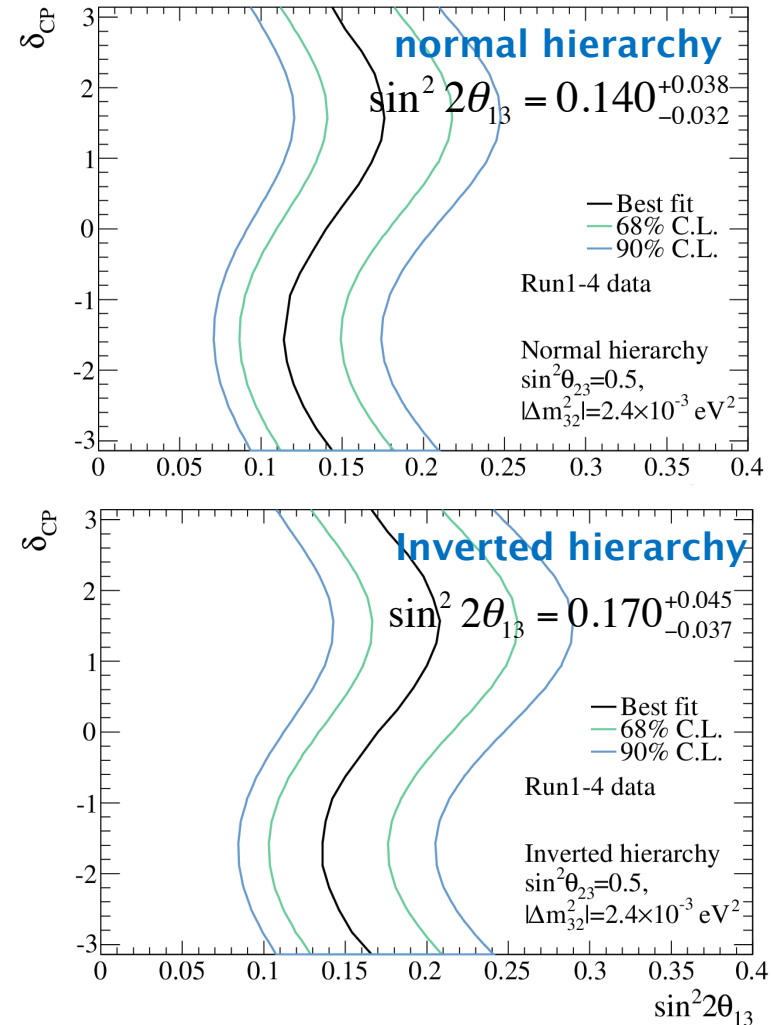
28 events observed

21.6 events expected @ $\sin^2 2\theta_{13} = 0.1$

$$\delta_{CP} = 0, \sin^2 \theta_{23} = 0.5$$

7.3 σ significance for non-zero θ_{13}

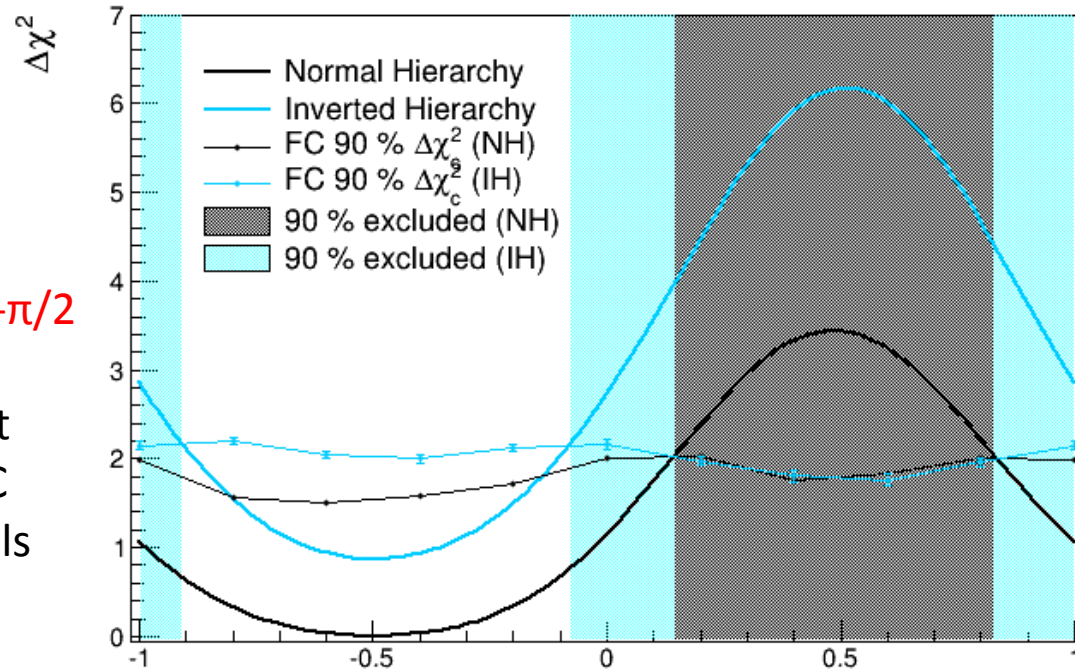
First ever observation ($>5\sigma$) of an explicit ν appearance channel



T2K Joint $\nu_e + \nu_\mu$ Analysis : Constraints on δ_{CP}

PRELIMINARY

- Combined likelihood ratio fit to the ν_e and ν_μ CCQE samples
- Using PDG 2013 θ_{13} value, T2K finds a (slight) preference for $\delta_{CP} = -\pi/2$ and the **Normal mass hierarchy**
- Similar results from an independent T2K analysis using Markov chain MC producing Bayesian credible intervals



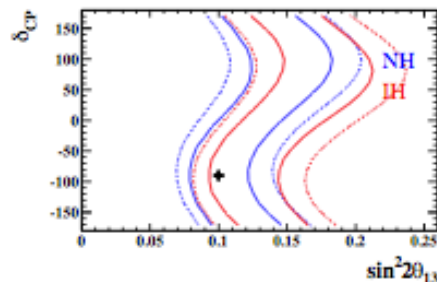
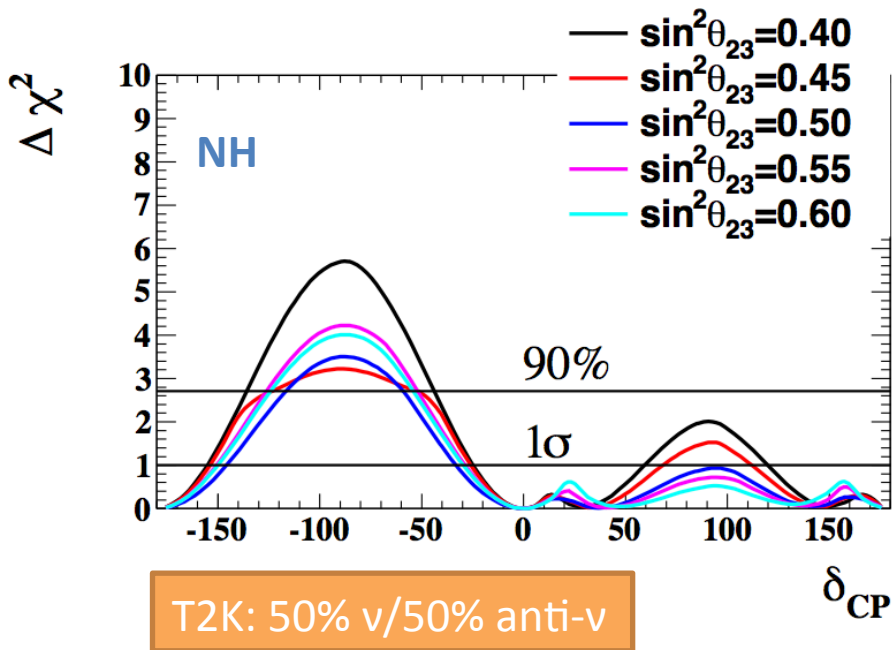
If Nature has chosen $\delta_{CP} \approx -\pi/2$, this will be very helpful experimentally
AND this solution satisfies the **leptogenesis bound** with no additional CP violation

$$|\sin \theta_{13} \sin \delta_{CP}| \geq 0.11$$

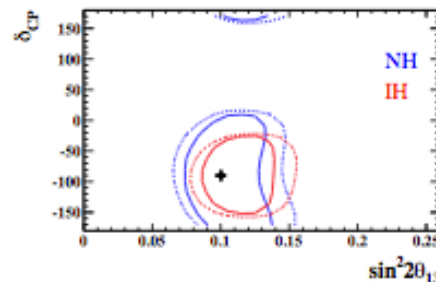
Pascoli, Petcov, Riotta 2007

Future Sensitivity to CPV using T2K

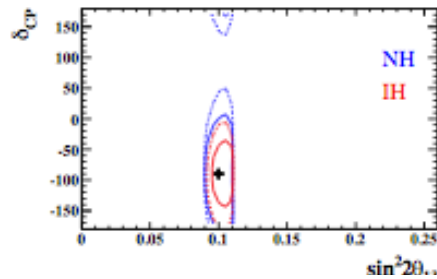
$$\sin^2 2\theta_{13} = 0.1, \delta_{CP} = -\pi/2$$



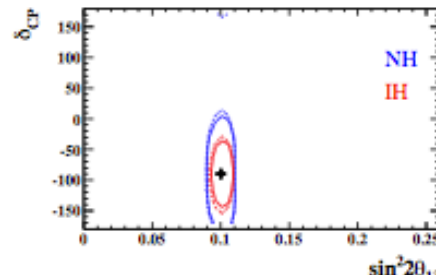
100% ν (true NH)



50% ν /50% anti- ν (true NH)



100% ν (true NH)
w/ Reactor constraint



50% ν /50% anti- ν (true NH)
w/ Reactor constraint

No systematics

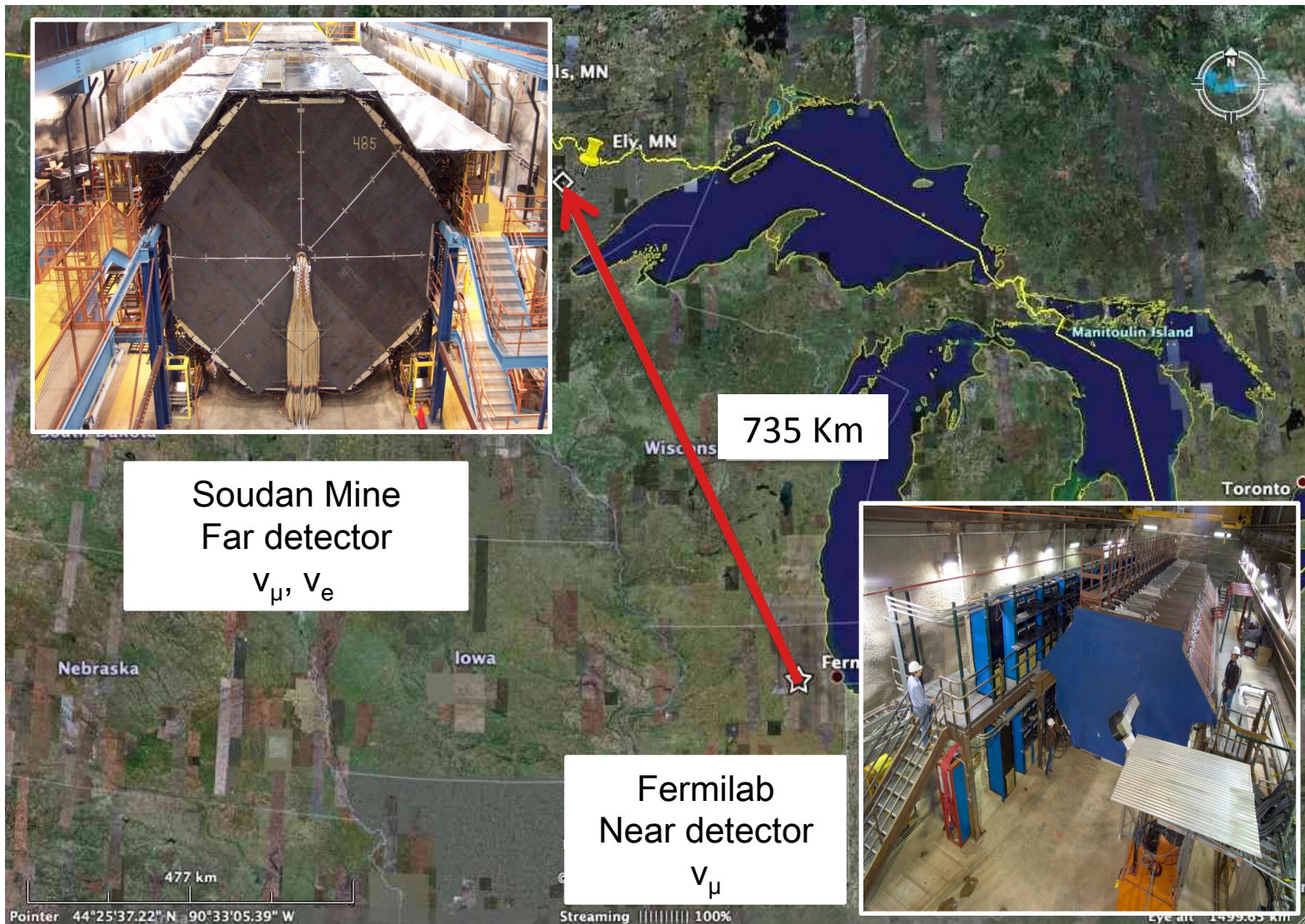
5% error on signal, 10% on background

T2K studies indicate our best sensitivity will be for 50% ν /50% anti- ν running. Anti- ν running also opens a large new physics program.

MINOS

High energy on-axis beam

MINOS= **M**ain **I**njector **N**eutrino **O**scillation **S**earch



MINOS 3 Flavor Oscillation Analysis

► Analysis combines:

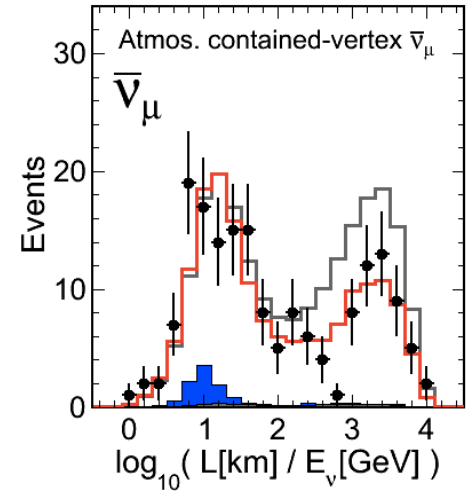
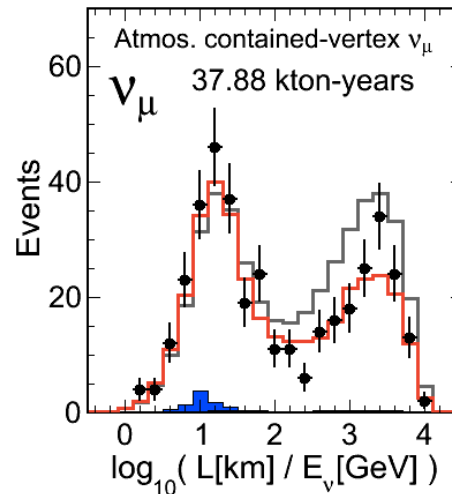
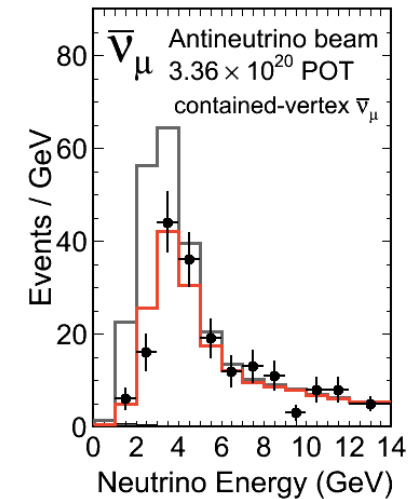
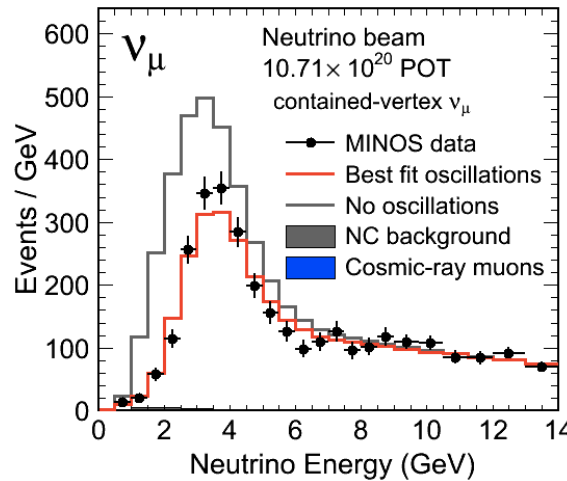
- Full MINOS ν_μ -CC and $\bar{\nu}_\mu$ -CC disappearance sample - *PRL* **112**, 191801 (2014)
- Full ν_e -CC, $\bar{\nu}_e$ -CC appearance sample, described in *PRL* **110** 171801 (2013)
- Full MINOS and **new MINOS+ atmospheric neutrino samples**

► Sensitive to θ_{13} , θ_{23} octant, mass hierarchy, and δ_{CP} from ν_e sample

► Sensitivity enhanced by atmospheric neutrinos:

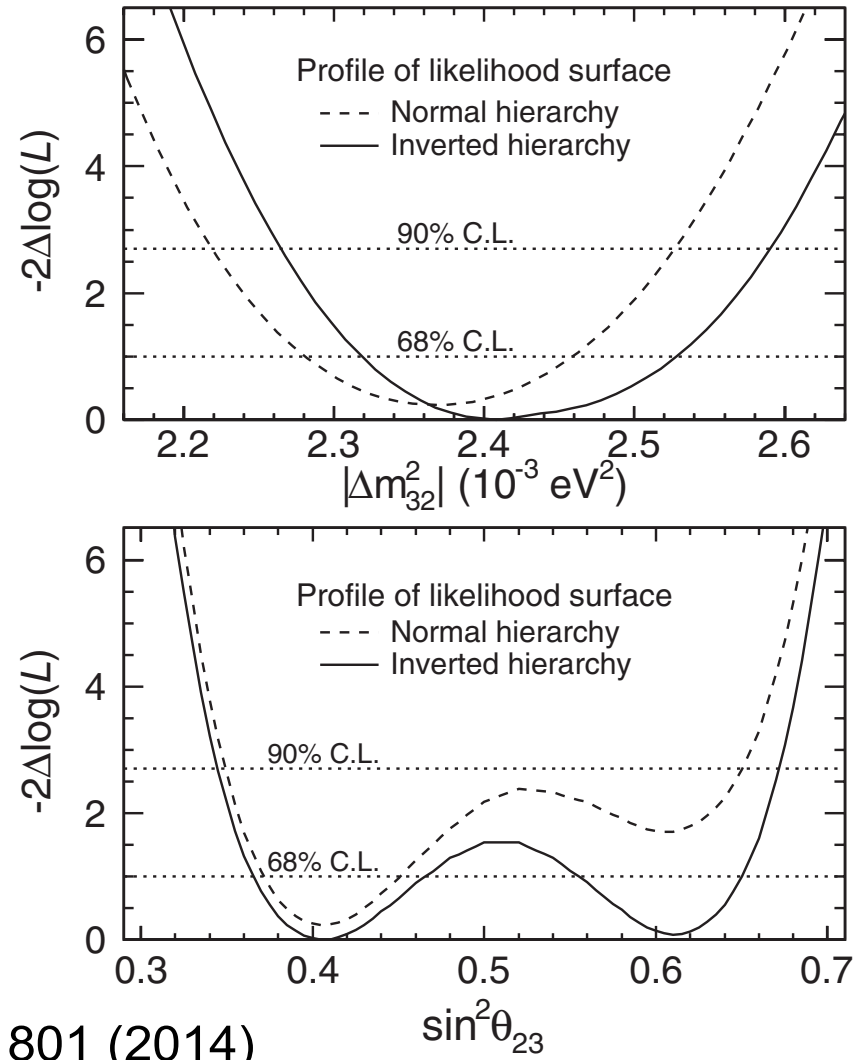
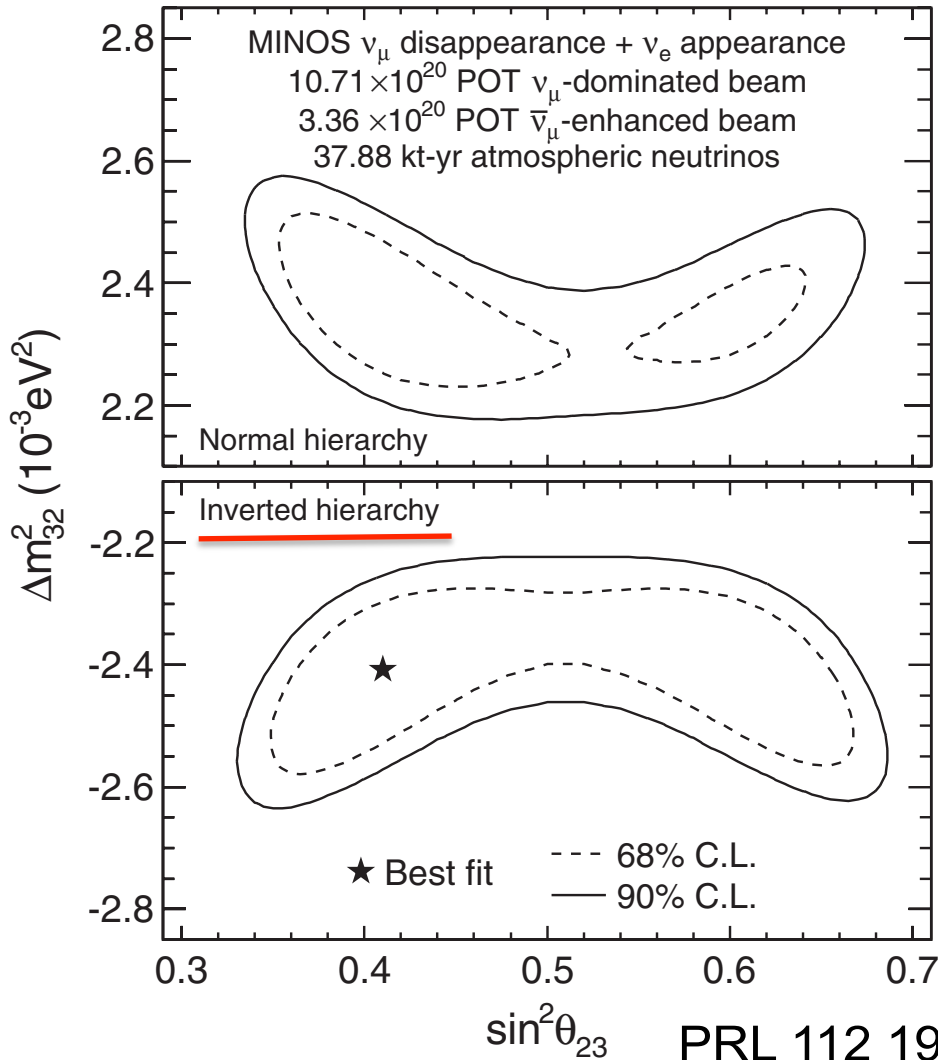
- Matter effects in multi-GeV, upward-going events
- Effect seen in neutrinos or antineutrinos, depending on hierarchy

► MINOS first to probe effect with event-by-event charge separation



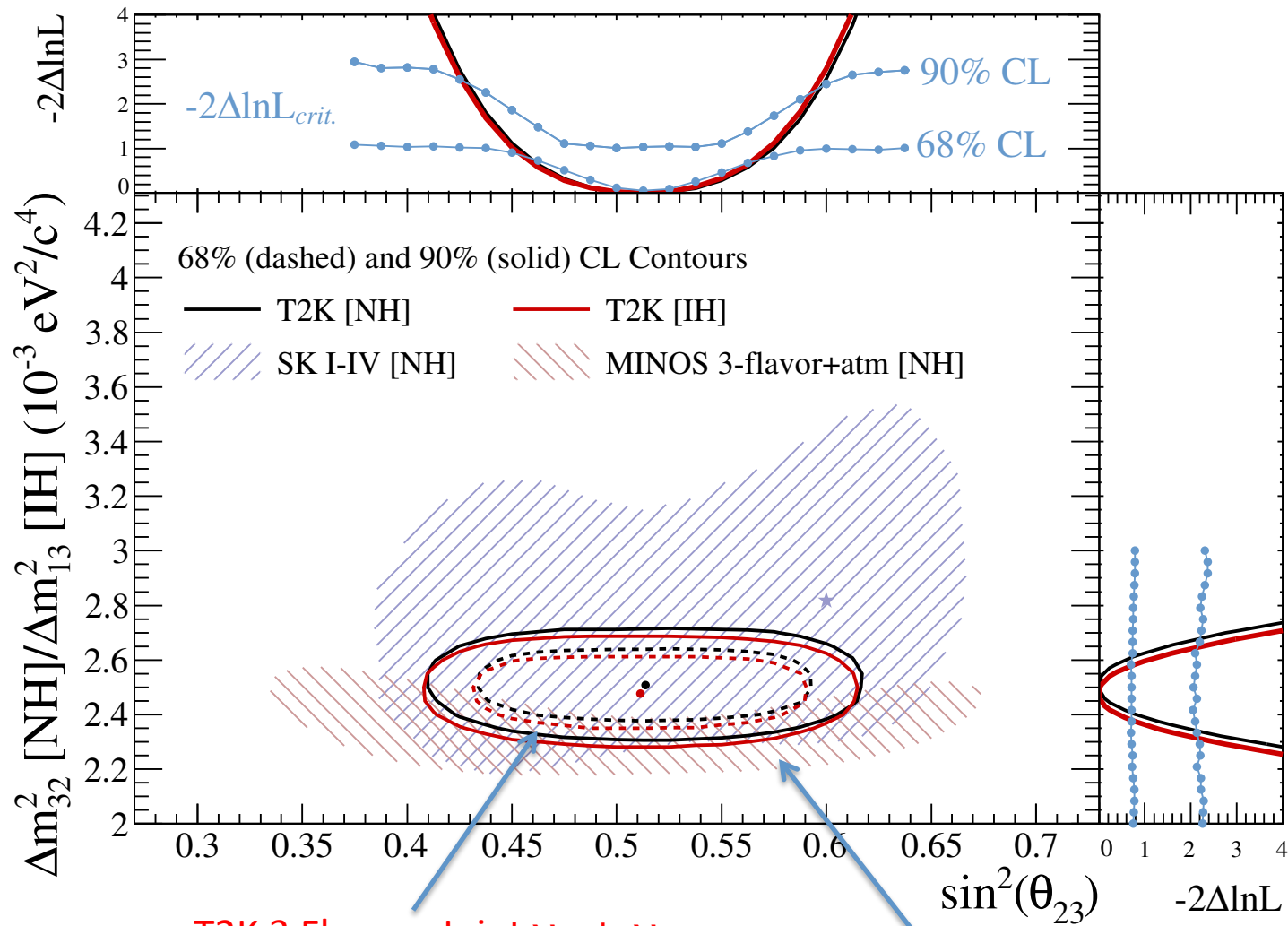
MINOS 3 Flavor Oscillation Analysis

Most precise measurement of $|\Delta m_{32}^2|$ - consistent with Maximal Mixing



PRL 112 191801 (2014)

Comparison of T2K with MINOS Disappearance/Atmospheric Results



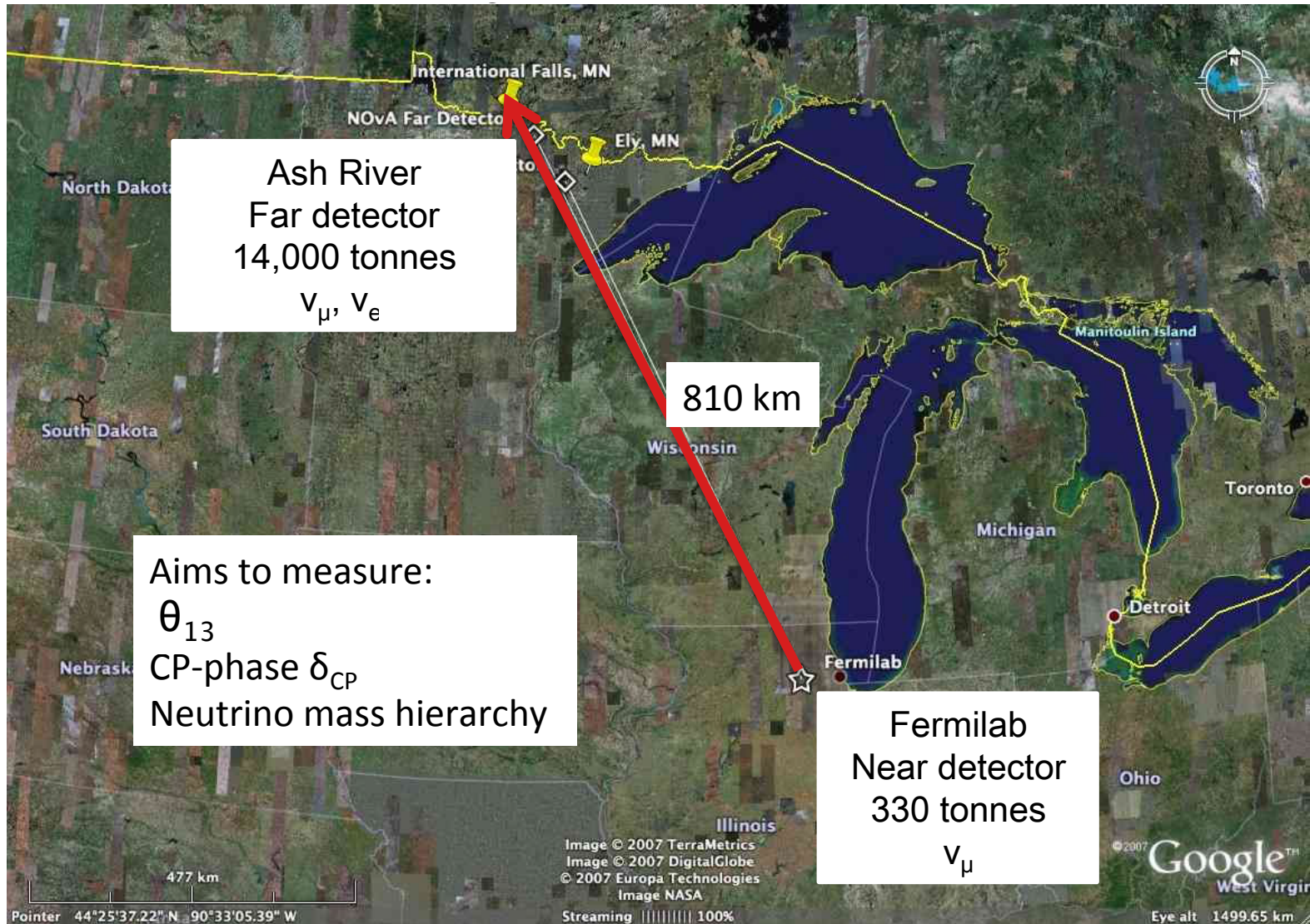
T2K 3 Flavour Joint $\nu_e + \nu_\mu$

MINOS 3 Flavour + Atm.

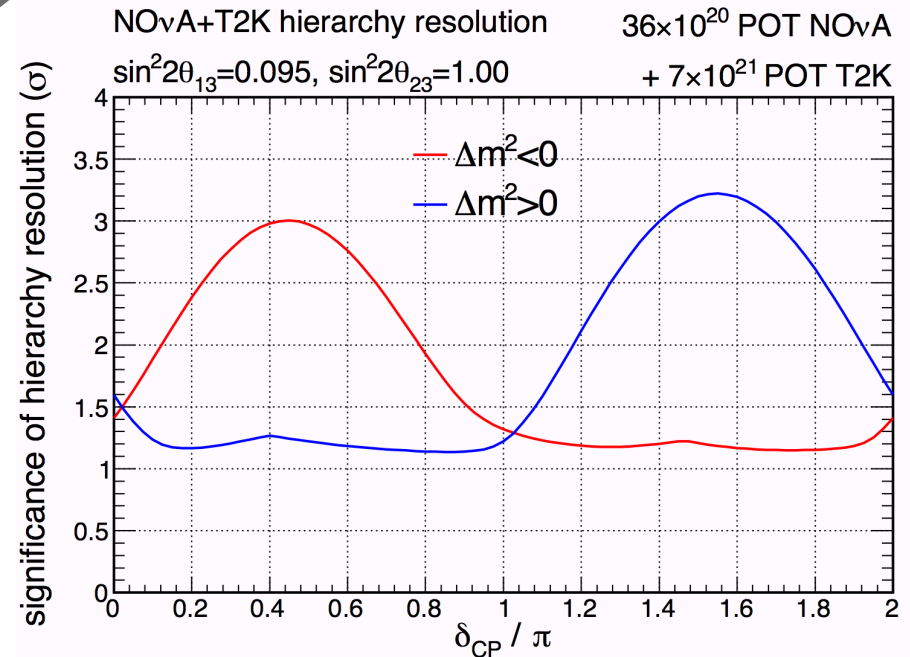
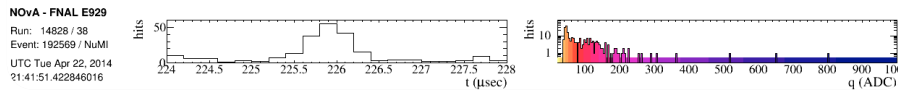
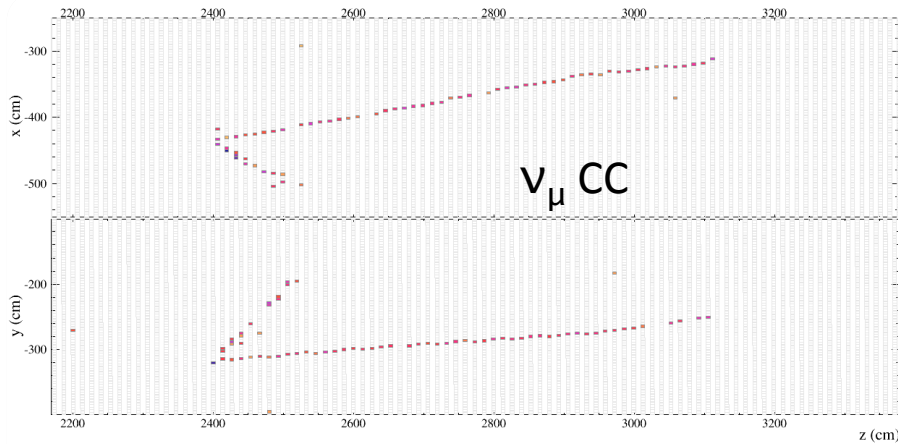
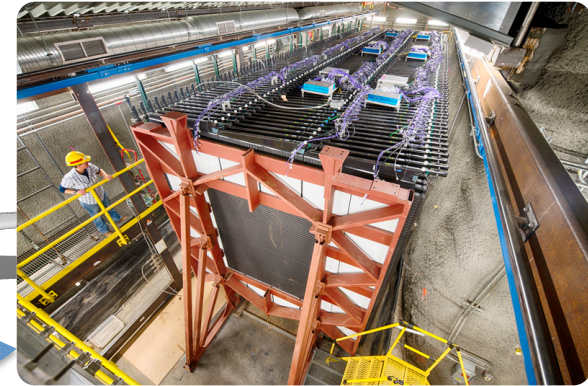
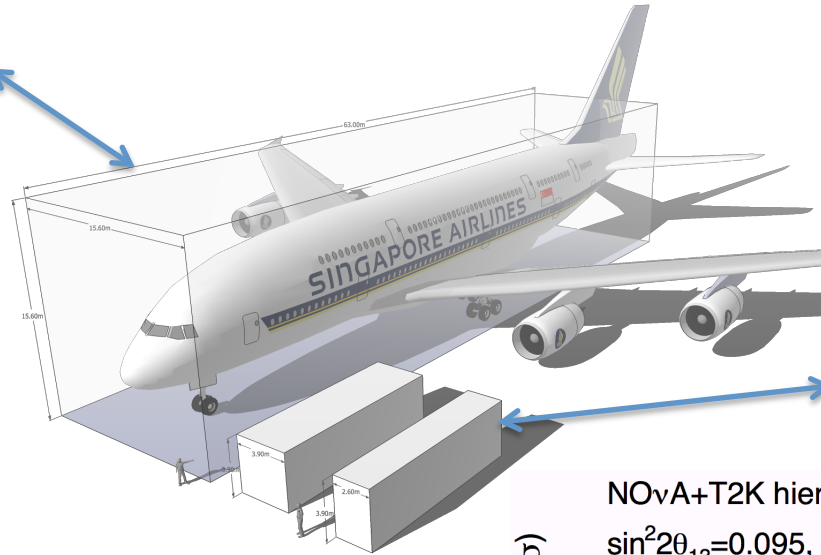
NOvA

High energy off-axis beam

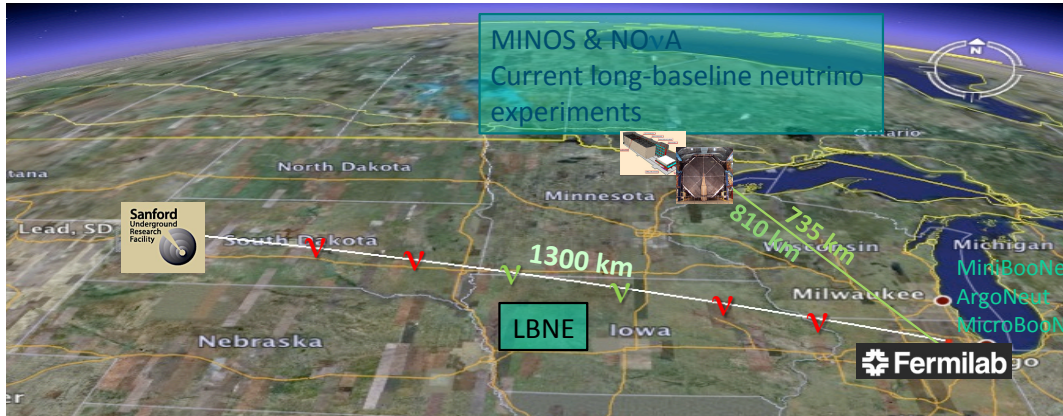
NOvA = NuMI Off-Axis ν_e Appearance



First Neutrino Interactions Observed in NOvA!

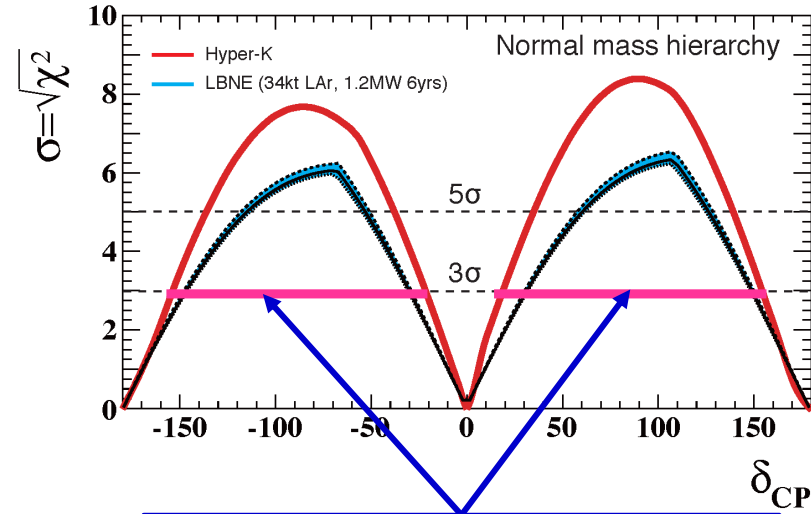


LBN(E/F) What next (mid-2020's)?

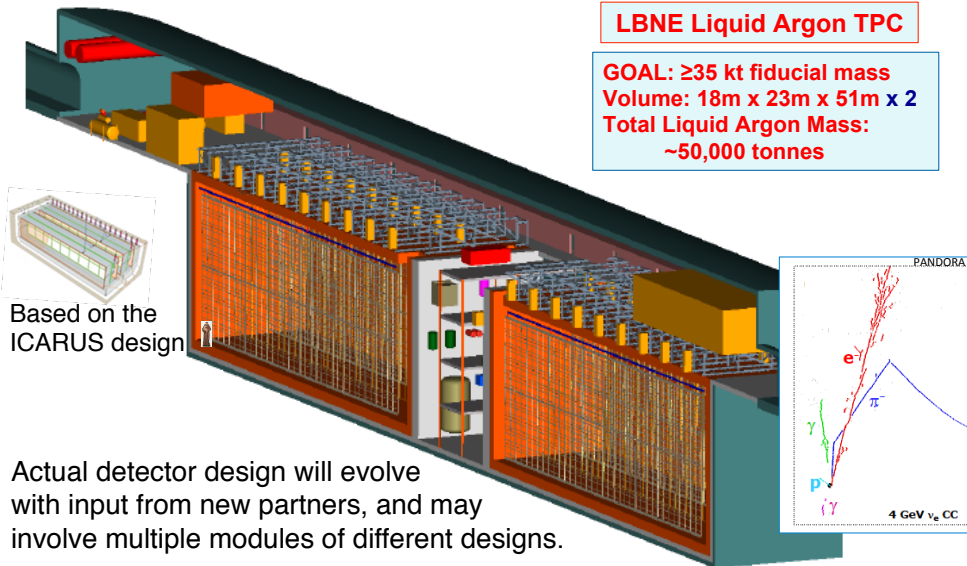


Comprehensive CP Violation, Mass Hierarchy, Non-Standard Interactions and Matter Effects Need **Longer Baseline** and High Intensity **Broadband** Neutrino/Anti-Neutrino Beam

Sensitivity of δ



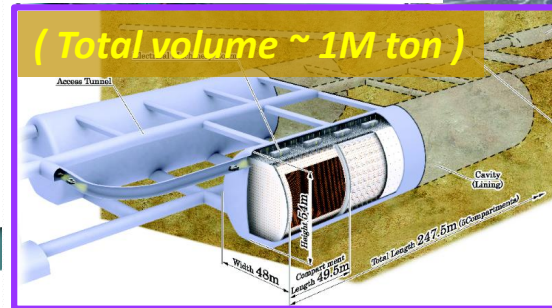
76% (3σ) of the δ is covered by J-PARC + Hyper-Kamiokande



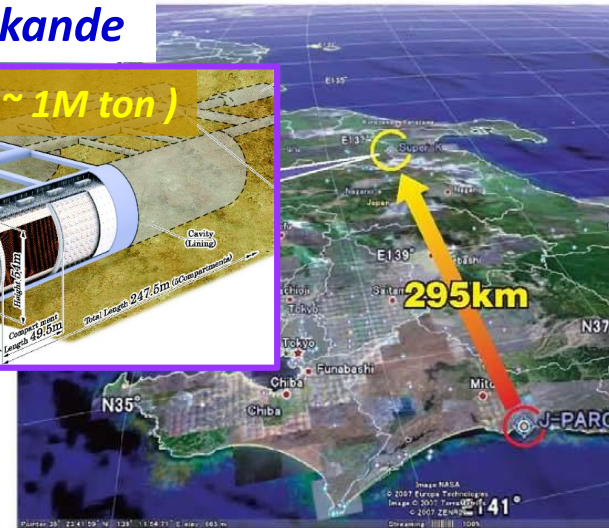
Based on the ICARUS design

Actual detector design will evolve with input from new partners, and may involve multiple modules of different designs.

Hyper-Kamiokande



T2HK



Summary & Conclusions

- During the last 1.5 decades the study of neutrino interactions has revealed many unexpected discoveries, establishing the **3 neutrino flavour mixing paradigm** (PMNS) which implies new physics beyond the Standard Model.
- The field is entering the **few % precision** era due to dedicated experimental efforts on 4 continents.
- The experiments are becoming sensitive to **CPV effects** via the interplay of reactor and accelerator observables.
- Major efforts are ongoing towards answering the remaining open questions and providing precise tests of the current model.

Back-up Slides

Neutrino Oscillations

In a simplified 2 flavour model...

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

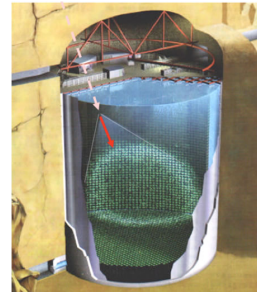
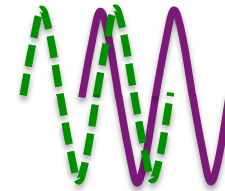
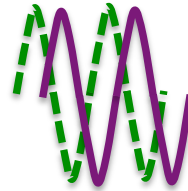
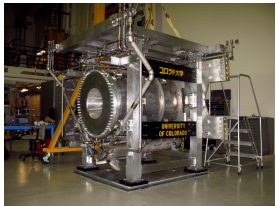
If neutrino **flavour eigenstates** are different to neutrino **mass eigenstates**, propagation over a baseline L of a neutrino of given flavour induces a phase shift, with the possibility of the appearance of the other neutrino flavour

$$\nu_\mu = -\sin \theta \nu_1 + \cos \theta \nu_2$$

$$\nu_1 \rightarrow \exp(-ip_1 x) \nu_1$$

$$\nu_2 \rightarrow \exp(-ip_2 x) \nu_2$$

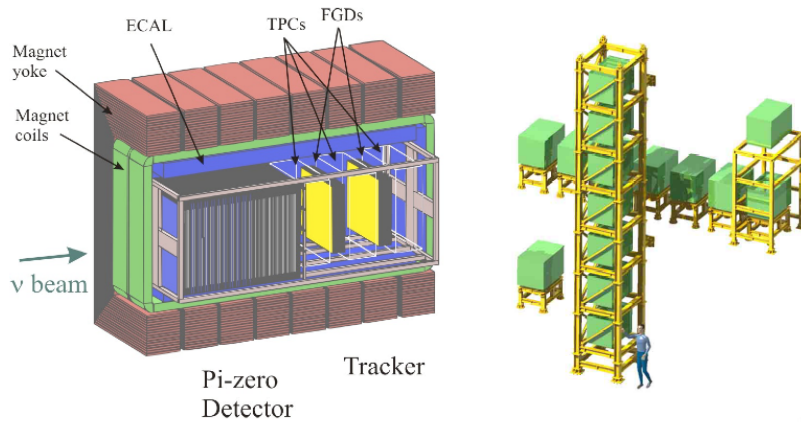
$$\Delta\Phi = \Delta m^2 L / (4E)$$



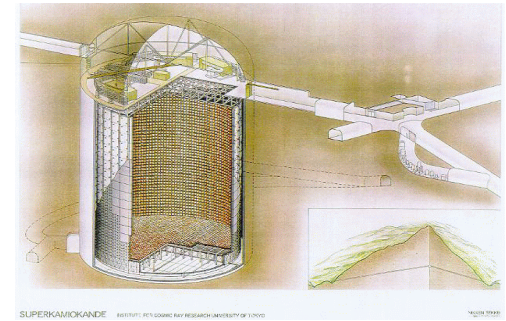
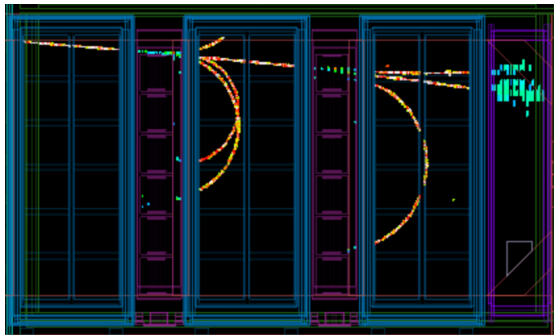
$$\text{Prob}(\nu_\mu \rightarrow \nu_e) = \sin^2(2\theta) \sin^2(\Delta m^2 L / 4E)$$

The T2K Experiment Overview

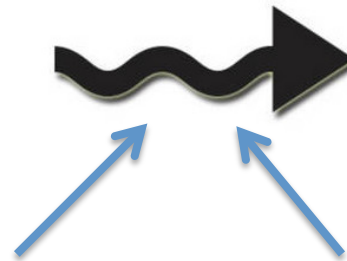
Uses the **J-PARC accelerator complex** for the beam.



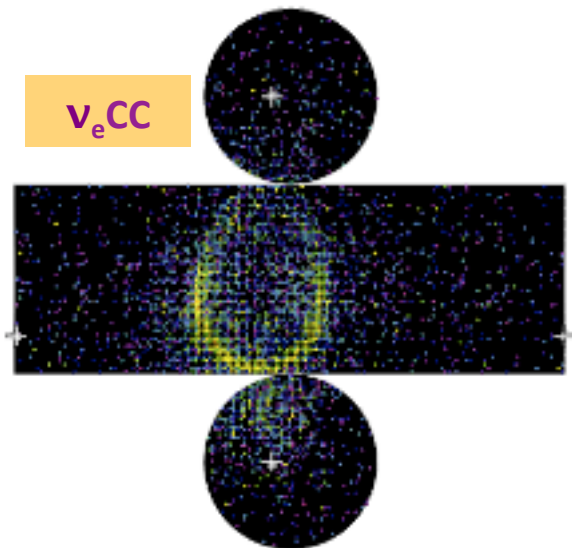
On and off-axis hybrid **near detectors** at 280m



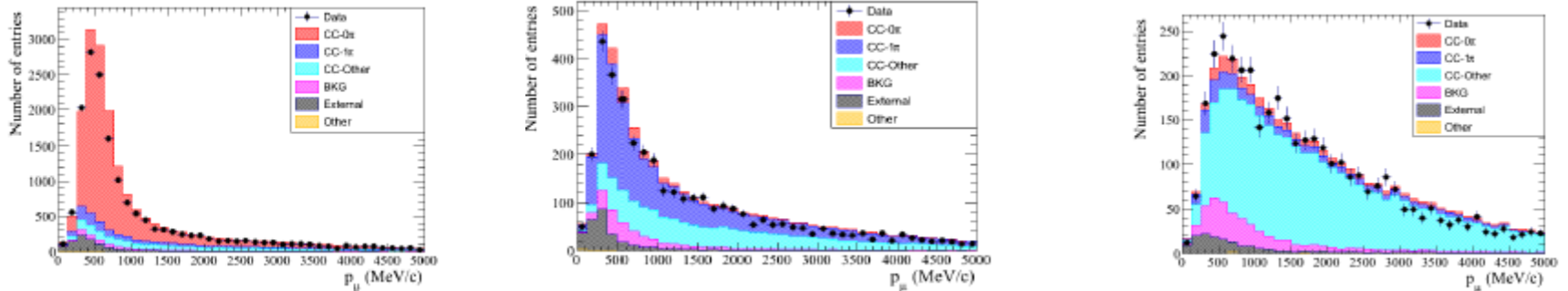
Far detector: Super-K at 295 km.



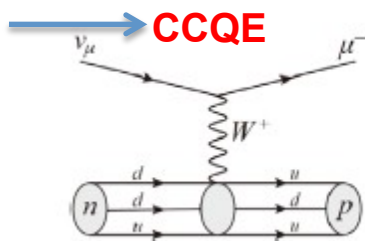
Use information from hadron production experiment NA61



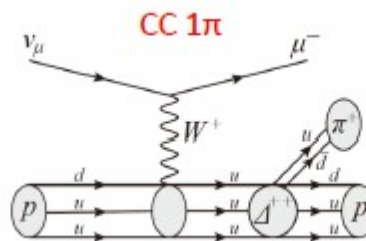
T2K Near Detector Constraint



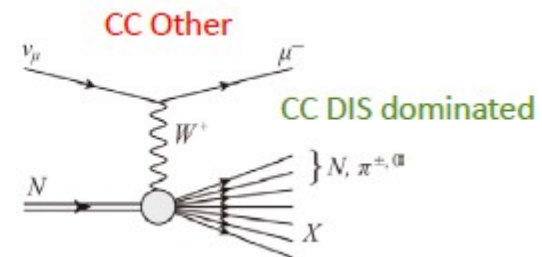
Signal channel for oscillation analysis



CCQE



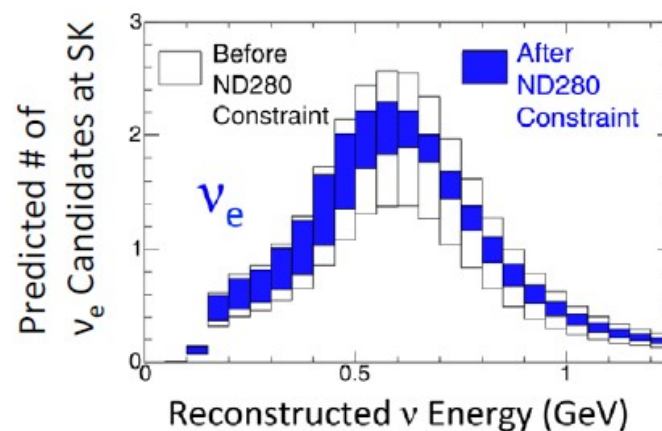
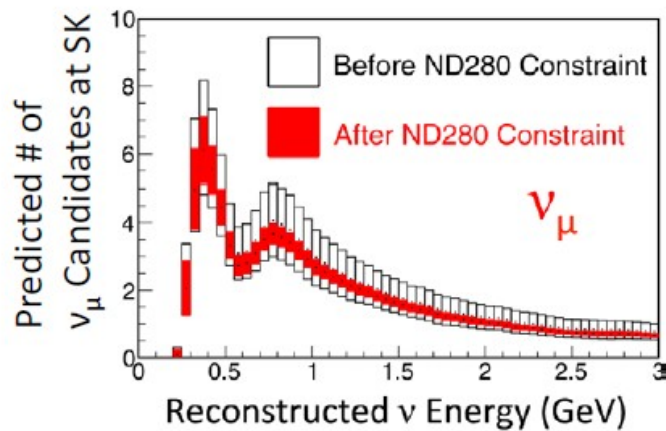
CC 1 π



CC Other

CC DIS dominated

Flux and cross-section systematic uncertainty on N_{SK} significantly reduced to $\sim 7\%$

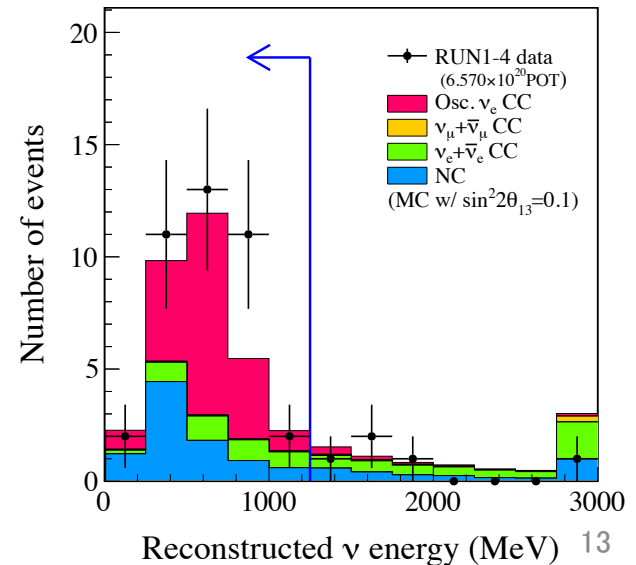
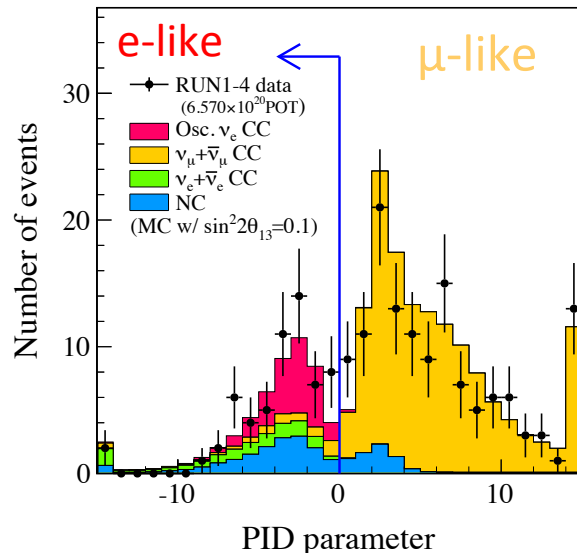
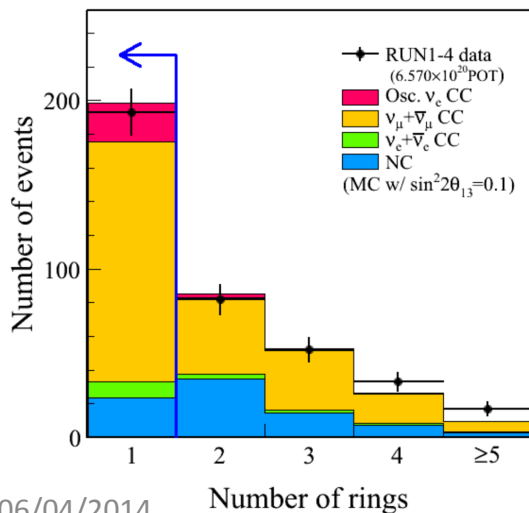


ν_e event selection

T2K has made improvements in background and error reduction.

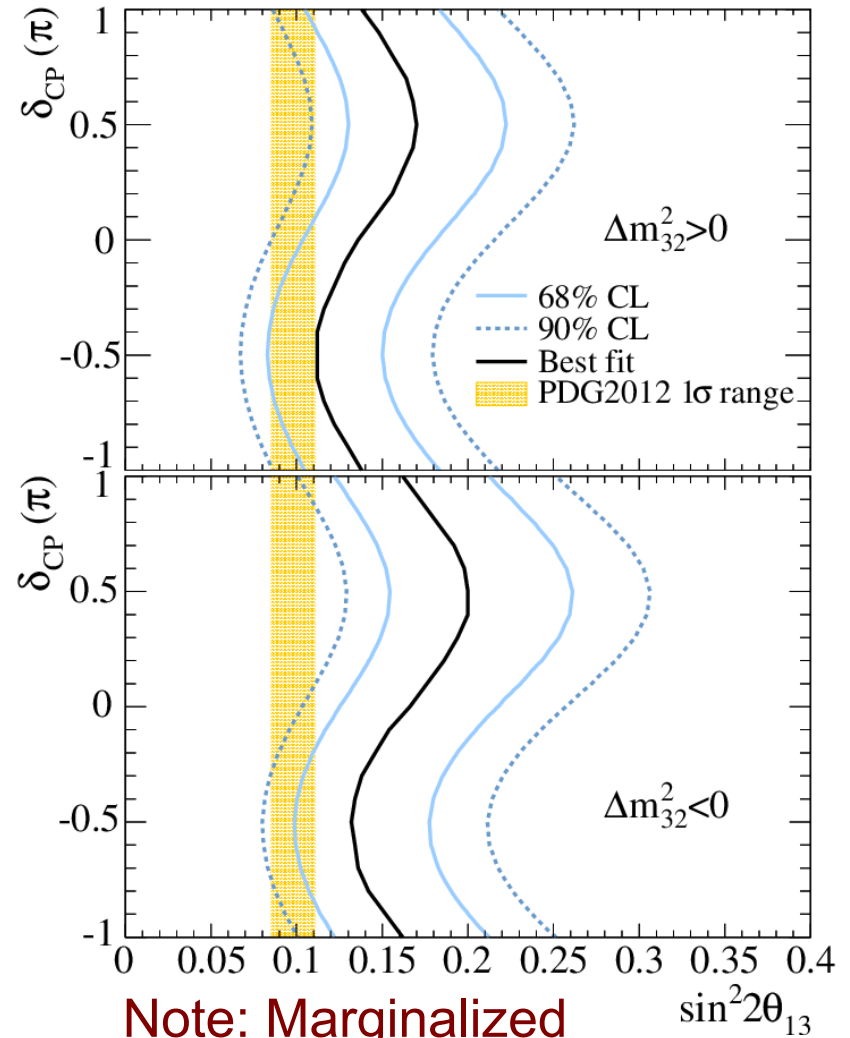
- Fully-contained fiducial volume (FCFV) event
- Single-ring e-like event
- $E_{\text{visible}} > 100$ MeV
- # of decay electron = 0
- $0 < E_{\nu}^{\text{rec}} < 1250$ MeV
- π^0 cut

28 events
in 6.57×10^{20} POT



T2K Let's think about these regions!

- Comparing with the external reactor constraint the best overlap is for the normal hierarchy with $\delta_{\text{CP}} = -\pi/2$.
- This is a **lucky point!**
- You also need to increase the θ_{23} mixing angle to account for the number of observed events.



Note: Marginalized
over θ_{23} and Δm_{32}^2

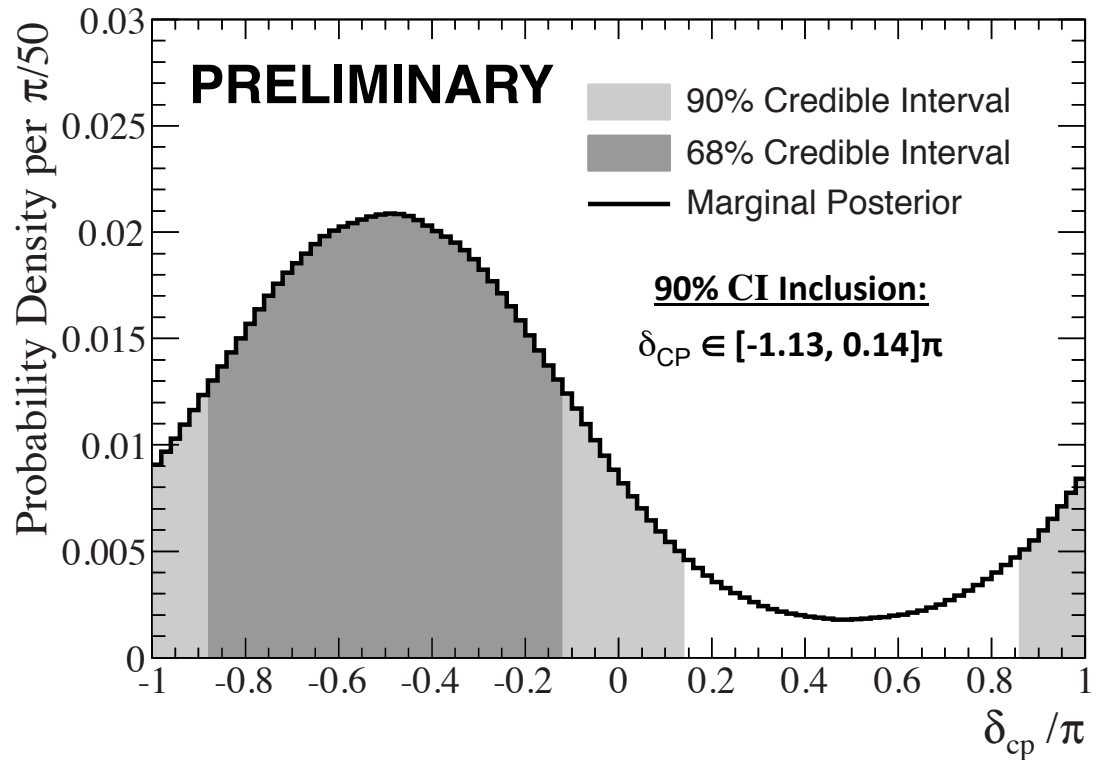
T2K Joint $\nu_\mu + \nu_e$ Bayesian Analysis

Assuming flat priors for $\sin^2\theta_{23}$, $|\Delta m^2_{32}|$; $P(\text{NH}) = P(\text{IH}) = 0.5$

New Result:

Use a Markov Chain Monte Carlo (MCMC) with both T2K-SK $\nu_\mu + \nu_e$ and ND280 samples.

Note: Marginalized over hierarchy.

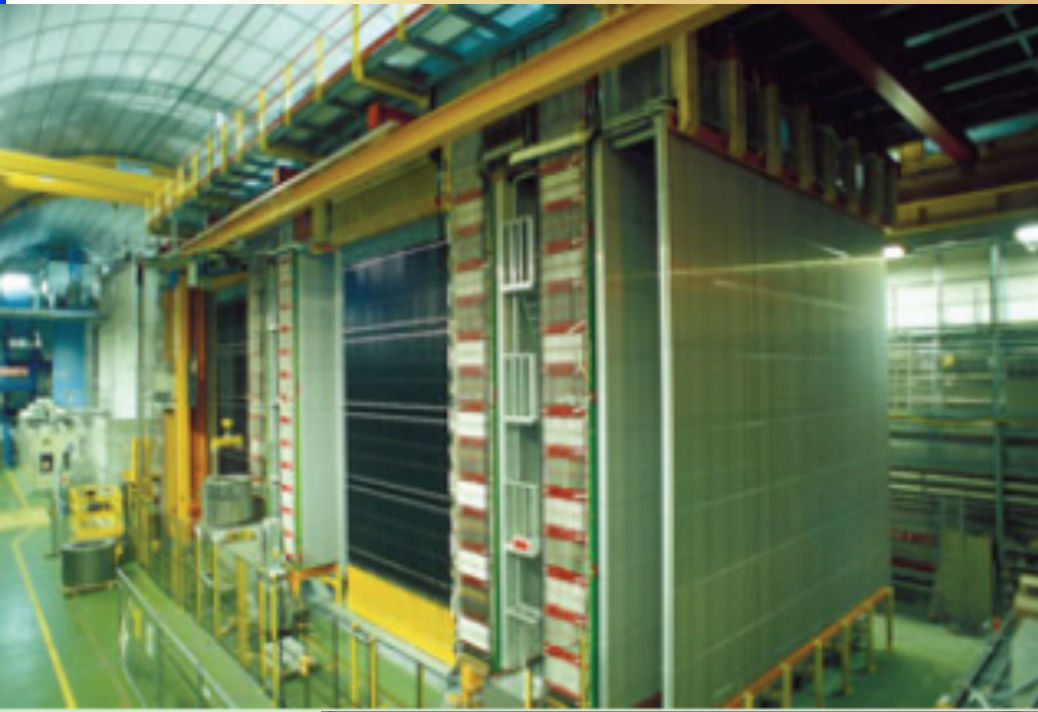


Can compare the probabilities for each MH and θ_{23} octant combination in the posterior probabilities.

(%)	NH	IH	Sum
$\sin^2\theta_{23} \leq 0.5$	18	8	26%
$\sin^2\theta_{23} > 0.5$	50	24	74%
Sum	68%	32%	

PRELIMINARY

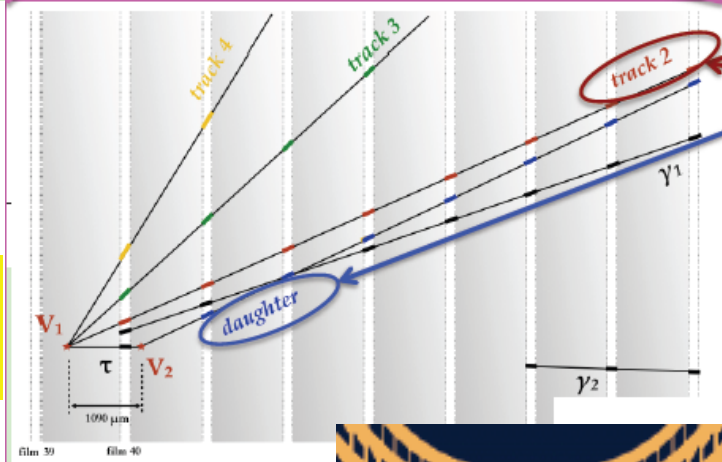
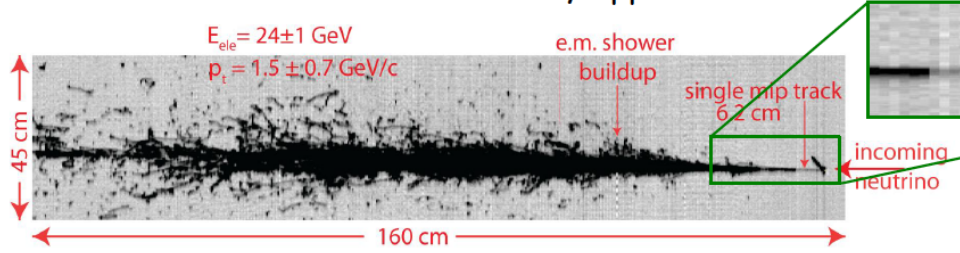
Long Baseline ν Oscillations



Decay channel	Expected signal $\Delta m_{23}^2 = 2.32 \text{ meV}^2$	Total background	Observed
$\tau \rightarrow h$	0.4 ± 0.08	0.033 ± 0.006	2
$\tau \rightarrow 3h$	0.57 ± 0.11	0.155 ± 0.03	1
$\tau \rightarrow \mu$	0.52 ± 0.1	0.018 ± 0.007	1
$\tau \rightarrow e$	0.61 ± 0.12	0.027 ± 0.005	0
Total	2.1 ± 0.42	0.23 ± 0.04	4

no oscillation excluded at 4.2σ CL

- The evolution of the actual dE/dx from a single track to an e.m. shower for the electron shower is clearly apparent from individual wires.



ICARUS



Leptogenesis and low energy CP-violation in neutrino physics

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$$|Y_B| \cong 2.8 \times 10^{-13} |\sin \delta| \left(\frac{s_{13}}{0.2} \right) \left(\frac{M_1}{10^9 \text{ GeV}} \right). \quad (91)$$

The asymmetry of interest is predominantly in the lepton number $L_e + L_\mu$. Thus, in order to reproduce the observed baryon asymmetry, taken to lie in the interval $8.0 \times 10^{-11} \lesssim |Y_B| \lesssim 9.2 \times 10^{-11}$, $s_{13} |\sin \delta|$ and M_1 , in the case analyzed, should satisfy

$$2.9 \lesssim |\sin \delta| \left(\frac{s_{13}}{0.2} \right) \left(\frac{M_1}{10^{11} \text{ GeV}} \right) \lesssim 3.3. \quad (92)$$

Given that $s_{13} |\sin \delta| \lesssim 0.2$, the lower bound in this inequality can be satisfied only for $M_1 \gtrsim 2.9 \times 10^{11} \text{ GeV}$. Recalling that the flavor effects in leptogenesis of interest are fully developed for $M_1 \lesssim 5 \times 10^{11} \text{ GeV}$, we obtain a *lower bound on the values of $|s_{13} \sin \delta|$ and s_{13} for which we can have successful leptogenesis in the case considered:*

$$\underline{|\sin \theta_{13} \sin \delta| \gtrsim 0.11, \quad \sin \theta_{13} \gtrsim 0.11.} \quad (93)$$