



Radiative Decays at LHCb

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on behalf of the LHCb Collaboration



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Overview

The LHCb experiment.

Theory introduction:

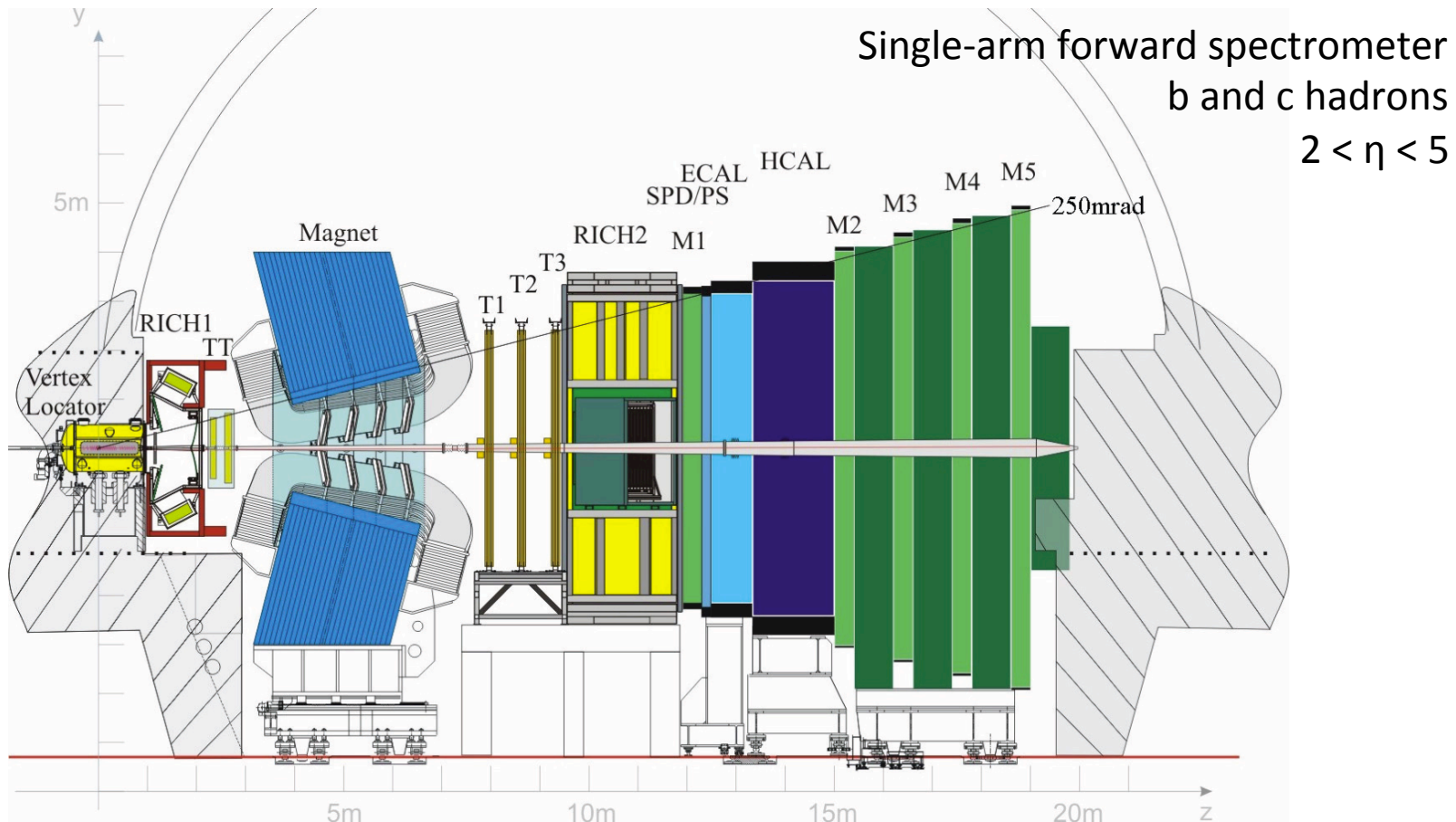
- Radiative decays within the SM and New Physics scenarios.

Measurement of the photon polarization with $B \rightarrow K\pi\pi\gamma$:

- Photon polarization in $B \rightarrow K_{\text{res}}\gamma$ decays.
- Up-down asymmetry. [\[PRL 112, 161801 \(2014\)\]](#)
- Angular fit.

Summary and Future.

The LHCb Experiment

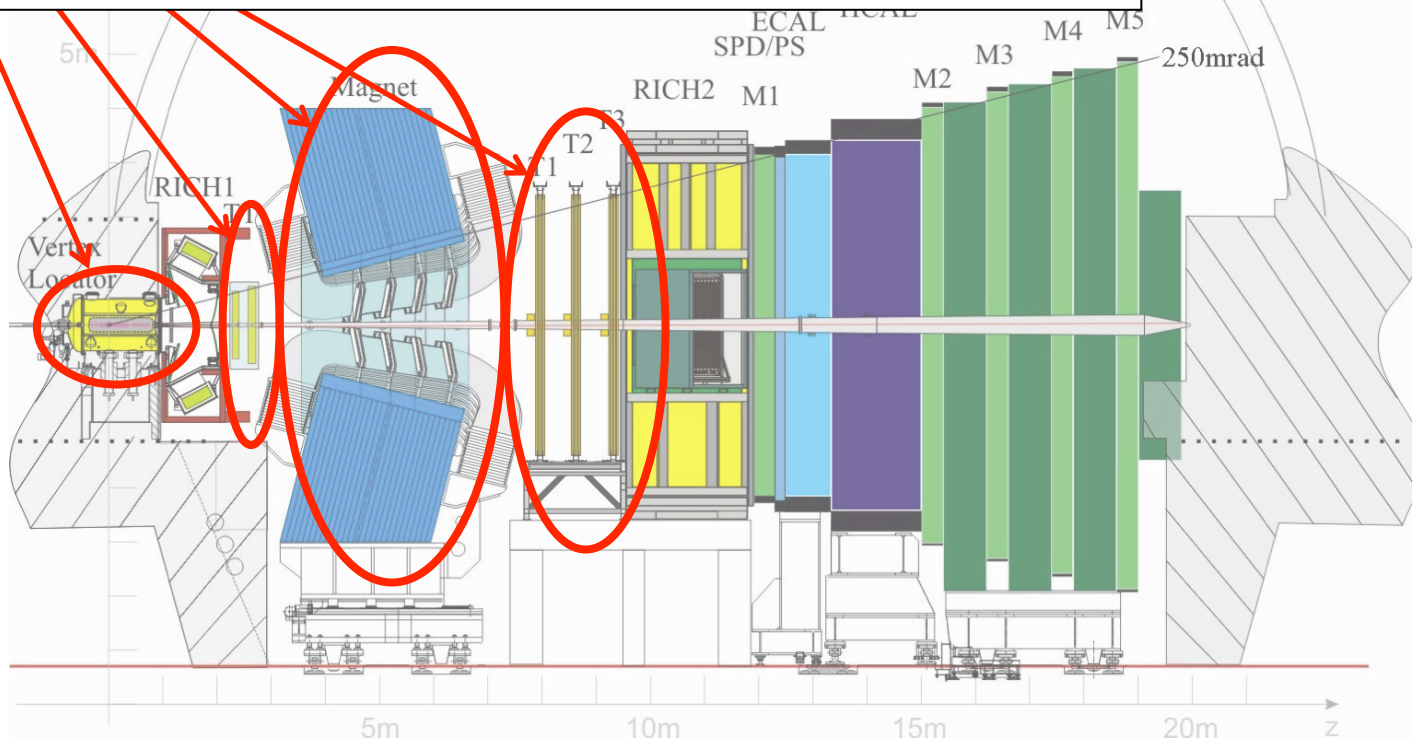


The LHCb Experiment

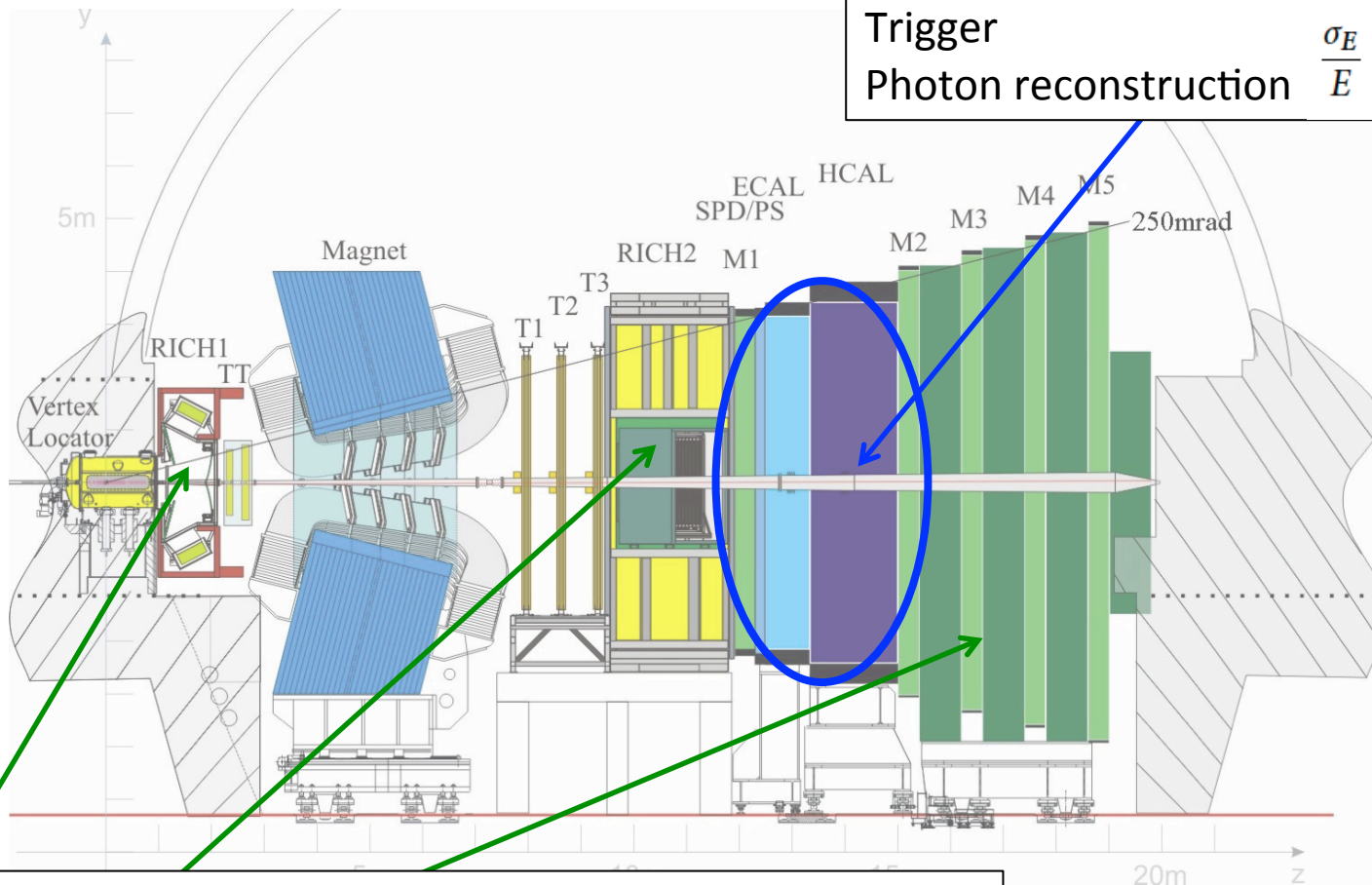
Precise tracking:

Good momentum ($\Delta p/p \sim 0.4\%$ at 5GeV) resolution.

Good IP ($20\mu\text{m}$ for high- p_T tracks) and decay time ($\sim 45\text{fs}$) resolution.



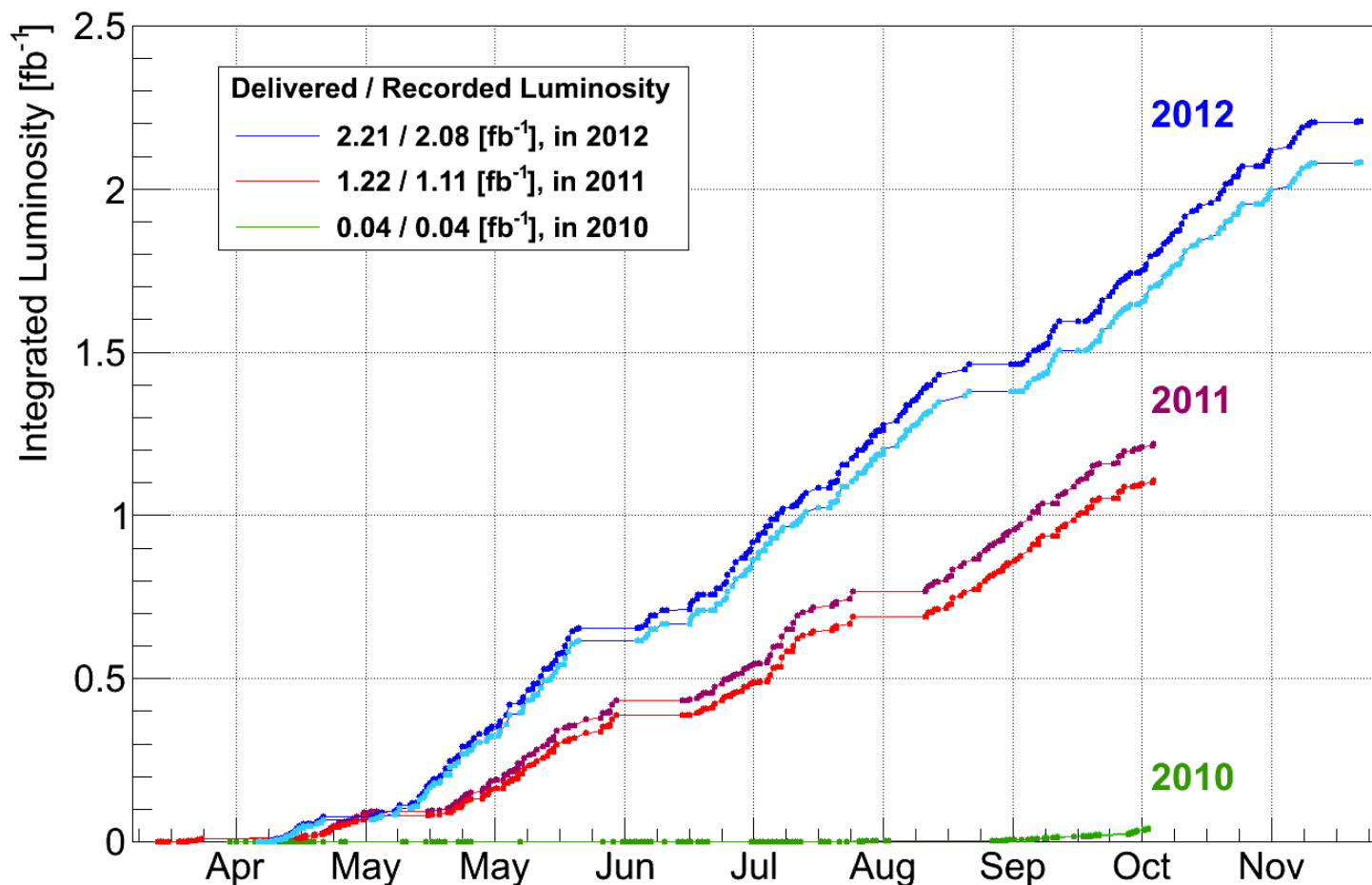
The LHCb Experiment



Calorimeter system:
 Trigger $\frac{\sigma_E}{E} = \frac{10\%}{\sqrt{E}} \oplus 1\%$
 Photon reconstruction

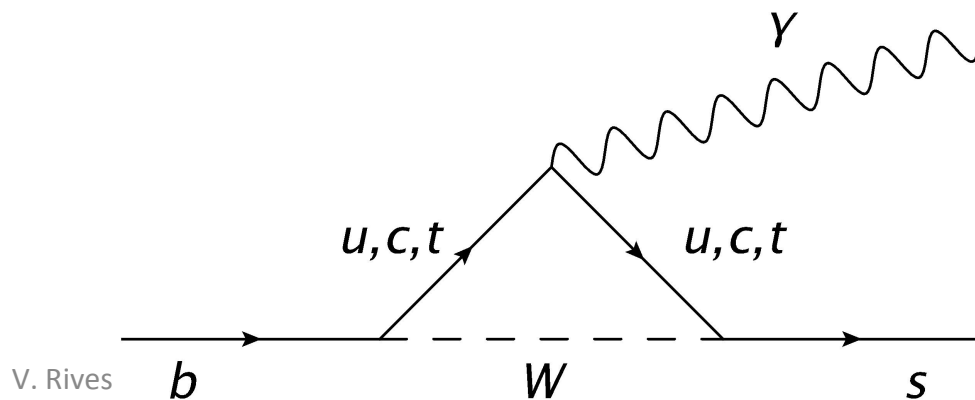
Excellent particle identification:
 π/K separation over 2-100 GeV ($\epsilon_K \sim 90\%$ for $\sim 5\% \pi \rightarrow K$ mis-ID)
 Powerful muon ID ($\epsilon_\mu \sim 97\%$ for $\sim 1-3\% \pi \rightarrow \mu$ mis-ID)

LHCb Run-I Performance



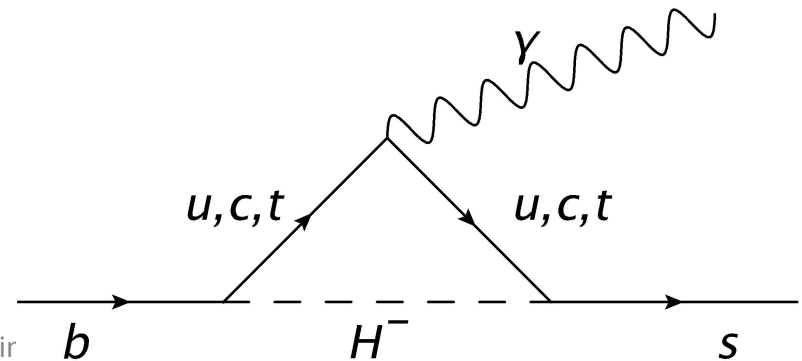
Radiative B Decays within the SM

- Radiative decays are FCNC processes ($b \rightarrow s\gamma$), not allowed at tree level:
 - They proceed through penguin transitions.
- Many possible observables reachable at the LHC:
 - Branching ratios, photon polarization (as a null-test of the SM) and different asymmetries (CP, isospin).
- Exclusive decays difficult from theoretical point of view due to form factor:
 - Need to find form-factor free observables (use of ratios that cancel out the form factors..).



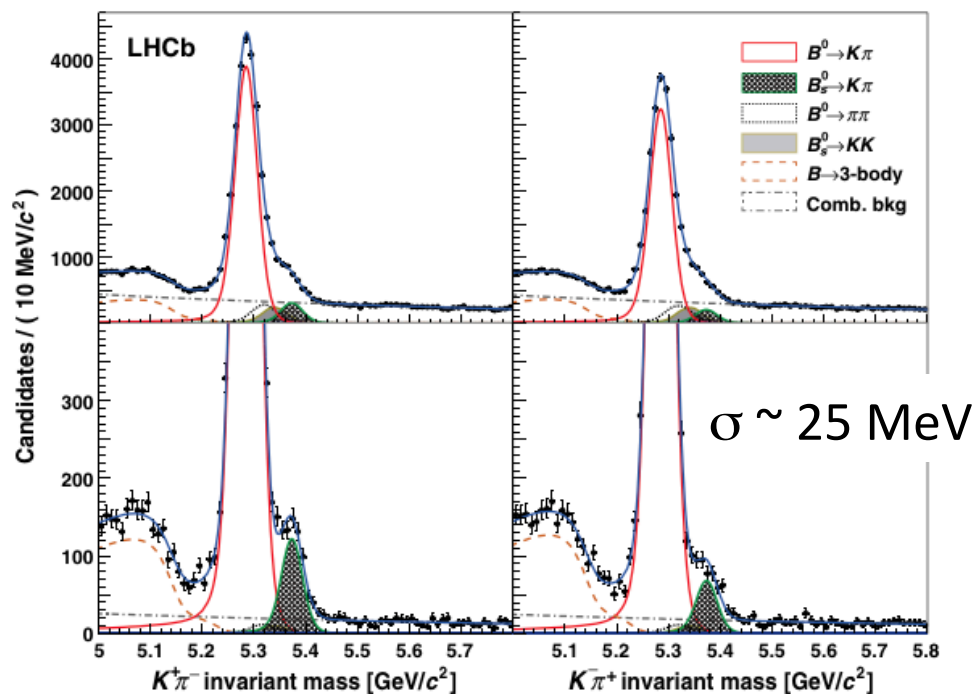
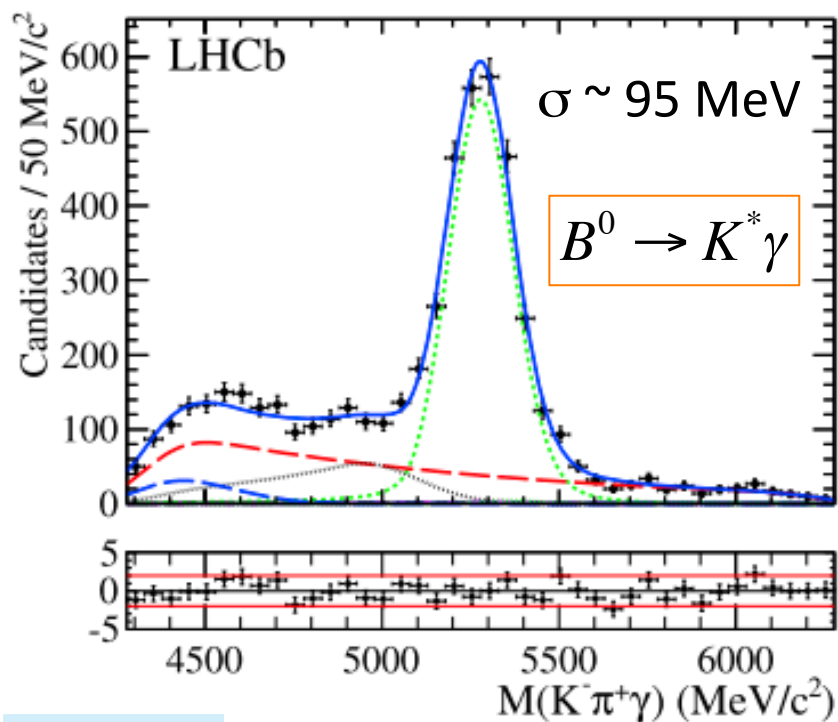
Radiative B Decays within New Physics

- The photon in the final state allows studies not reachable through other analyses.
- The SM predicts the photon to be (almost completely) left-handed polarized (corrections of order m_s/m_b).
- Several NP models introduce right-handed currents. New particles could change the chirality inside the loop, producing chiral enhancement:
 - m_t/m_b from LRSM [Babu *et al.*, [Phys.Lett.B333:196-201,1994](#)].
 - m_{SUSY}/m_b in SUSY with δ_{RL} mass insertions [Gabbiani *et al.*, [Nuclear Physics B 477 \(1996\) 321-352](#)].



Challenges for Radiative Decays at LHCb

- Distinct experimental signature with a high E_T photon:
 - Large levels of background are expected in a pp machine.
- Mass resolution dominated by photon reconstruction.



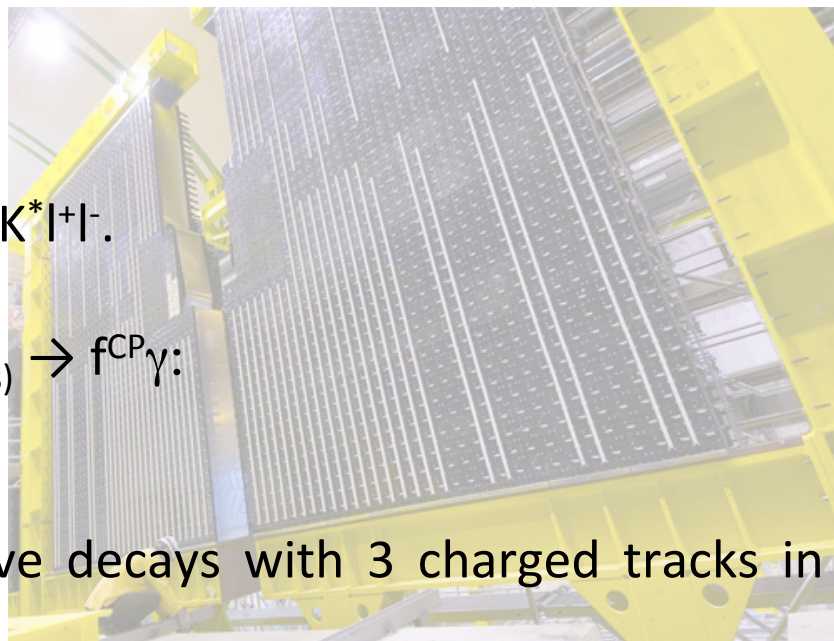
[Nucl. Phys. B 867 (2012) 1-18]

[PRL 110(2013) 221601]

Measuring the γ Polarization

Different approaches:

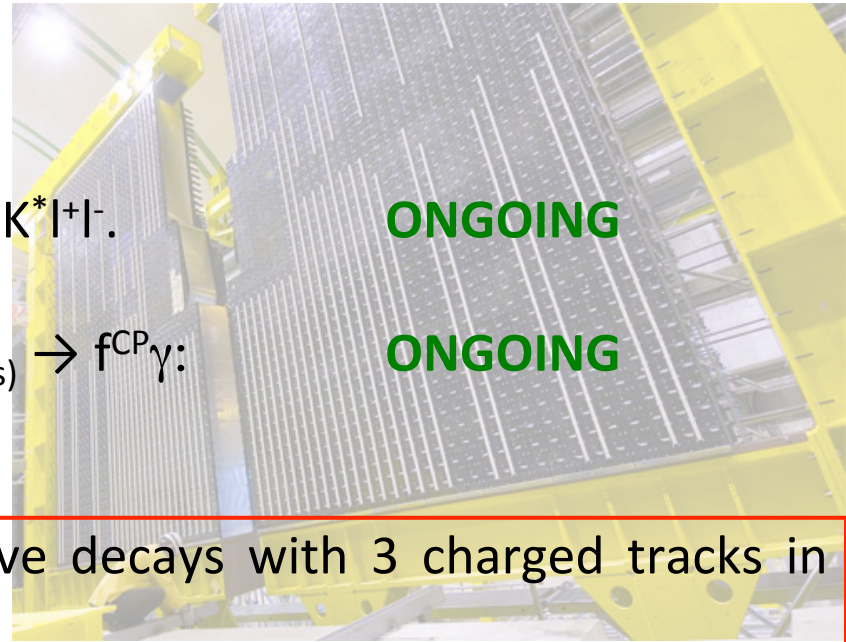
- Transverse asymmetry in $B^0 \rightarrow K^{*+}l^-$.
- Time-dependent analyses of $B_{(s)} \rightarrow f^{CP}\gamma$:
 - $B_s \rightarrow \phi\gamma$ and $B^0 \rightarrow K_s\pi\gamma$
- Angular distribution of radiative decays with 3 charged tracks in the final state:
 - $B \rightarrow K\pi\pi\gamma$, $B \rightarrow \phi K\gamma$, $B \rightarrow \pi\pi\pi\gamma$
- b-baryons decays: $\Lambda_b \rightarrow \Lambda^{(*)}\gamma$



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ONGOING

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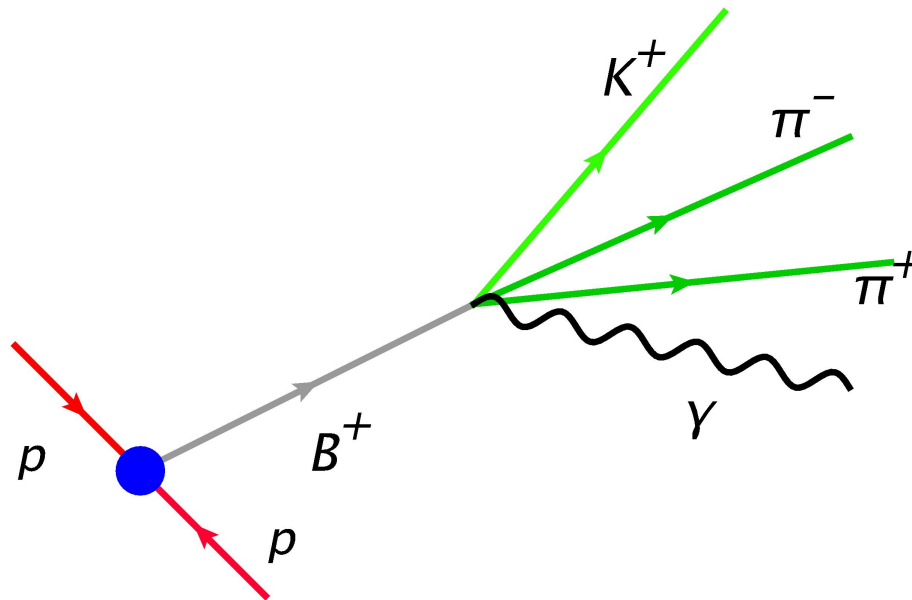
Today's talk, with $B \rightarrow K\pi\pi\gamma$
 [PRL 112, 161801 (2014)]

$B \rightarrow K\pi\pi\gamma$ at LHCb

Inclusive study of the $K\pi\pi$ system with mass range of [1.1, 1.9] GeV.

Use of the whole Run-I dataset, corresponding to 3fb^{-1} .

Study of angular distributions to search for the photon polarization.

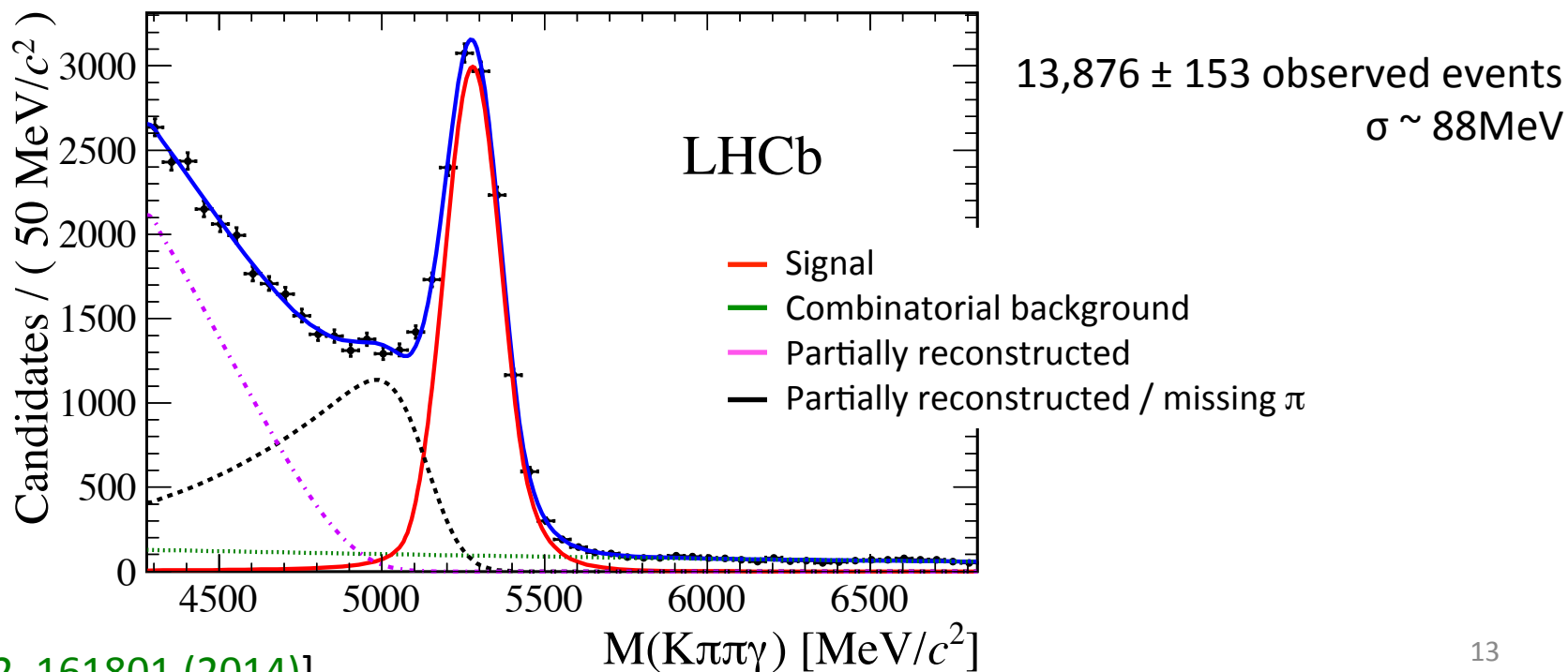


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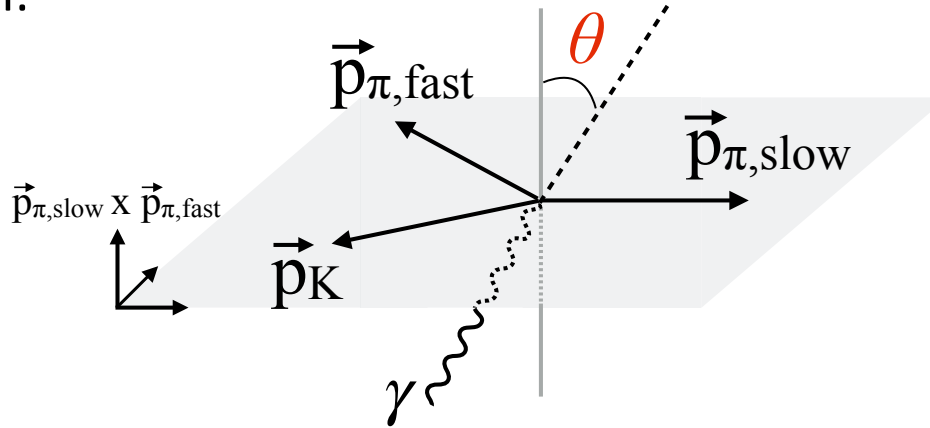
Use of the whole Run-I dataset, corresponding to 3fb^{-1} .

Study of angular distributions to search for the photon polarization.



Towards the photon polarization

The angular structure of the decay can be used to find the photon polarization:



$$A_{UD} \equiv \frac{\int_0^1 d\cos\theta \frac{d\Gamma}{d\cos\theta} - \int_{-1}^0 d\cos\theta \frac{d\Gamma}{d\cos\theta}}{\int_{-1}^1 d\cos\theta \frac{d\Gamma}{d\cos\theta}} = C\lambda_\gamma$$

where C accounts for the integral over the Dalitz plot and the angular distribution [[PRL 112, 161801 \(2014\)](#)].

Photon Polarization in $B \rightarrow K_{res}\gamma$

Photon polarization given by:

$$\lambda_\gamma = \frac{|c_R|^2 - |c_L|^2}{|c_R|^2 + |c_L|^2}$$

Weak amplitudes

Gronau *et al.* [[PRD 66 054008](#)] show that the γ polarization is independent of the K resonance:

$$\frac{|c_R|}{|c_L|} = \frac{|C_{7R}|}{|C_{7L}|} \Rightarrow \lambda_\gamma = \frac{|C_{7R}|^2 - |C_{7L}|^2}{|C_{7R}|^2 + |C_{7L}|^2} \quad \begin{array}{l} +1 \text{ for } \bar{b} \\ -1 \text{ for } b \end{array}$$

The amplitude of a K resonance decay can be written in terms of the helicity amplitude J_μ :

$$A_{L(R)}(s, s_{13}, s_{23}, \cos\theta) = \varepsilon_{K,L(R)}^\mu J_\mu$$

Polarization vector

Contains all the amplitude info

“easy” if only **one** (1^+) intermediate resonance [Kou *et al.*, [PhysRevD.83.094007](#)] [Gronau *et al.*, [PhysRevD.66.054008](#)]:

γ polarization goes with odd powers of $\cos\theta$

$$\frac{d\Gamma(B \rightarrow K_{res}\gamma \rightarrow K\pi\pi\gamma)}{ds ds_{13} ds_{23} d\cos\theta} \propto \frac{1}{4} |\vec{J}|^2 (1 + \cos^2\theta) + \lambda_\gamma \frac{1}{2} \cos\theta \text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]$$

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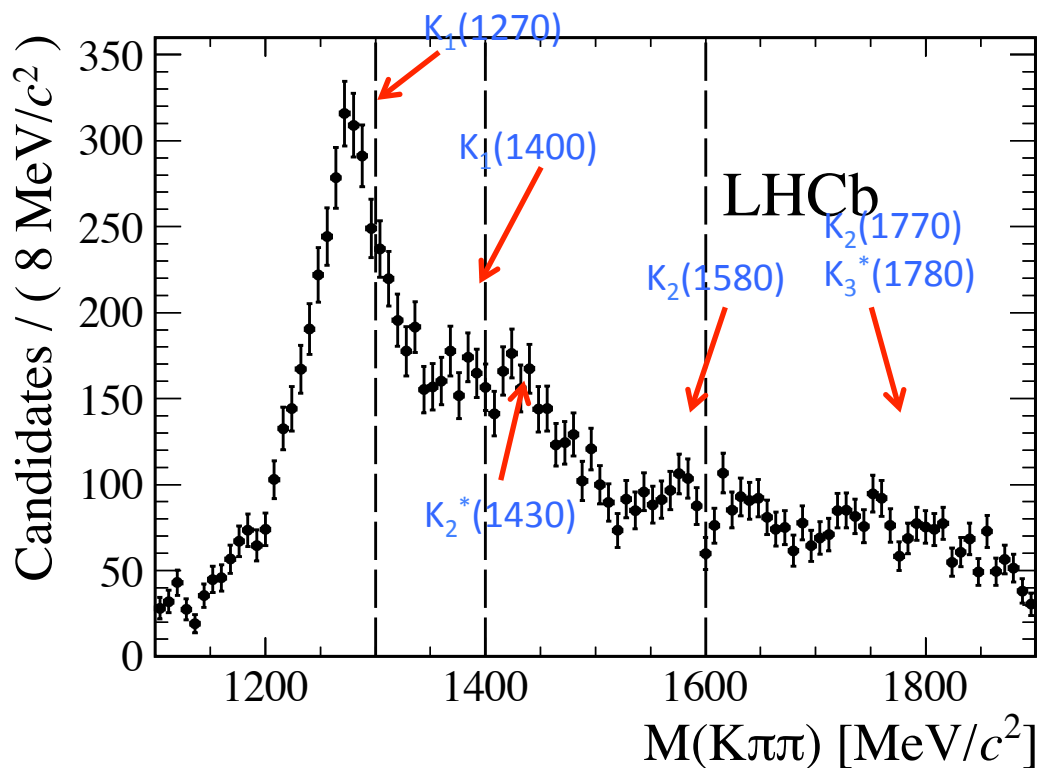
“easy” if only **one** (1^+) intermediate resonance [Kou *et al.*, [PhysRevD.83.094007](#)] [Gronau *et al.*, [PhysRevD.66.](#)] does with odd

Unfortunately, not the case

$$\frac{d\Gamma(B \rightarrow K_{res}\gamma \rightarrow K\pi\pi\gamma)}{ds ds_{13} ds_{23} d\cos\theta} \propto \frac{1}{4} |\vec{J}|^2 (1 - \dots)$$



B → Kππγ



Many contributions, interference between 1^+ , 1^- , 2^+ resonances [Gronau *et al*, [PhysRevD.66.054008](#)] impossible to separate without full amplitude analysis.

Inclusive analysis with theoretically motivated binning (chosen beforehand). The angular analysis is performed region by region.

$$\frac{d\Gamma(\Sigma B \rightarrow K_{res}\gamma \rightarrow P_1 P_2 P_3 \gamma)}{ds ds_{13} ds_{23} \cos\theta} \propto \sum_{j=even} a_j(s_{13}, s_{23}) \cos^j \theta + \lambda_\gamma \sum_{j=odd} a_j(s_{13}, s_{23}) \cos^j \theta$$

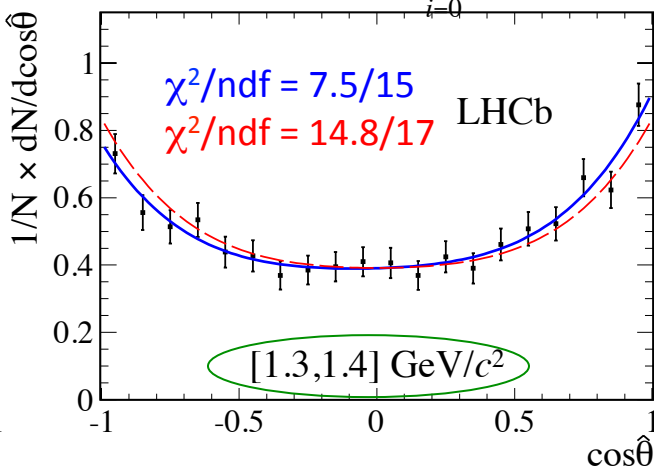
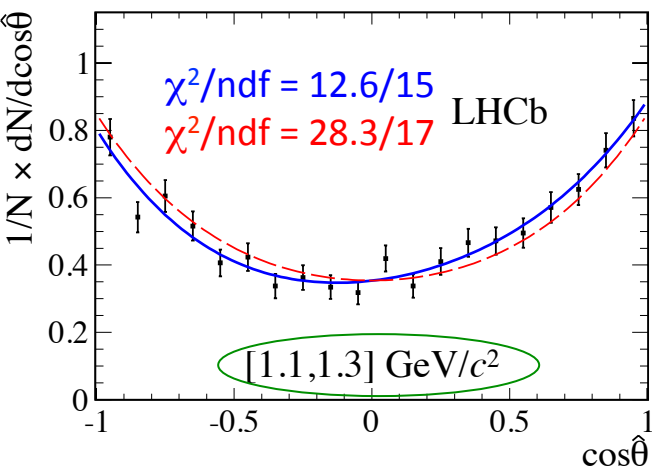
λ_γ goes with odd powers of $\cos\theta$



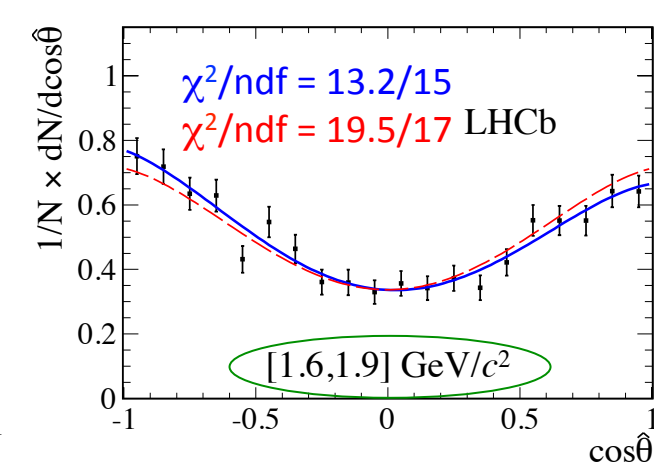
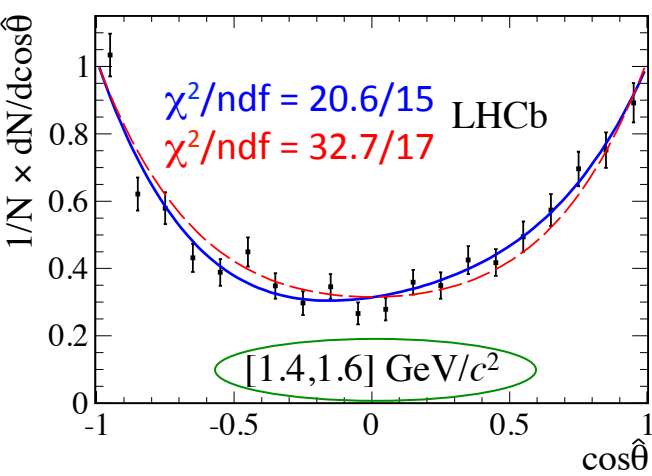
Angular fit coefficients

Angular distributions for each $K\pi\pi$ mass region fitted with combination of Legendre polynomials (up to order 4):

$$f(\cos\hat{\theta}; c_0 = 0.5, c_1, c_2, c_3, c_4) = \sum_{i=0}^4 c_i L_i(\cos\hat{\theta})$$



— Nominal Fit
 - - No polarization



[PRL 112, 161801 (2014)]

Angular fit analysis

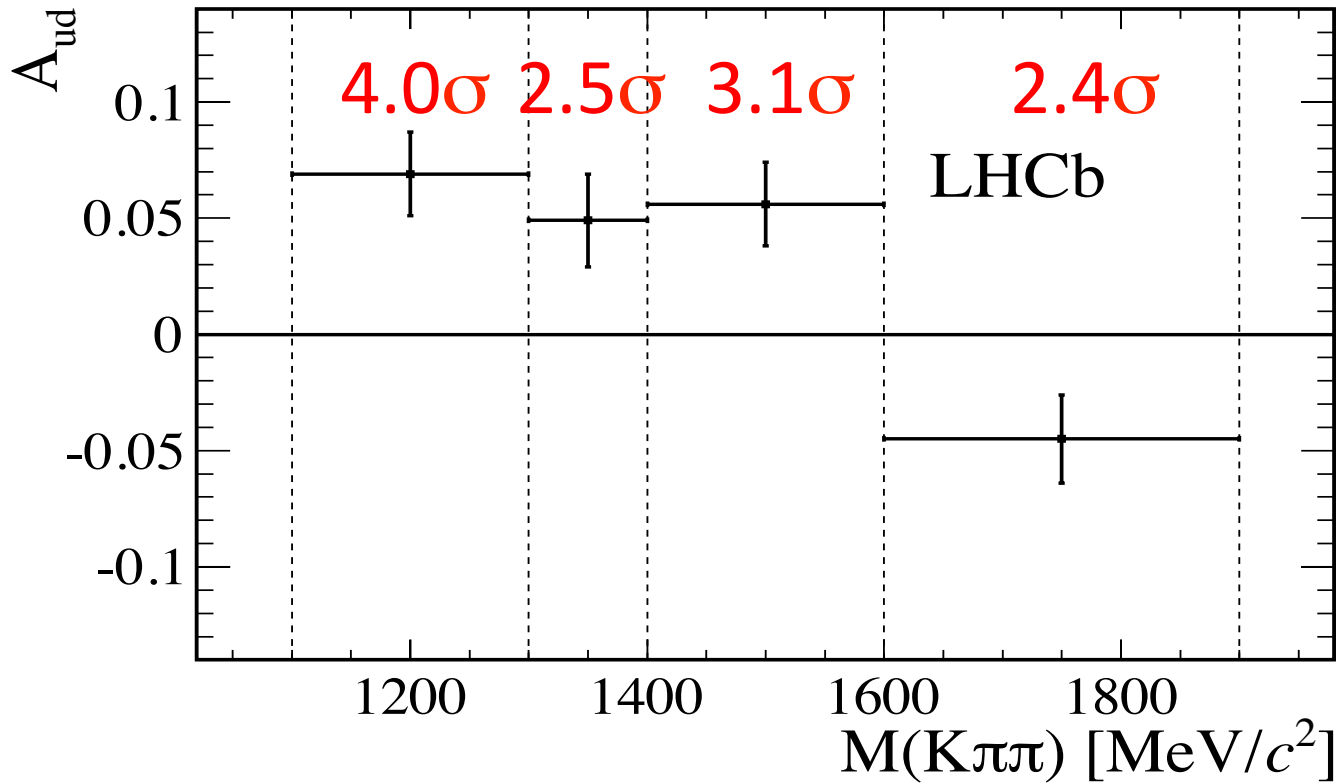
From the different values for the coefficients we can extract a value for the A_{UD} for each of the $K\pi\pi$ invariant mass region.

The up-down asymmetry is determined by: $A_{UD} = c_1 - \frac{c_3}{4}$

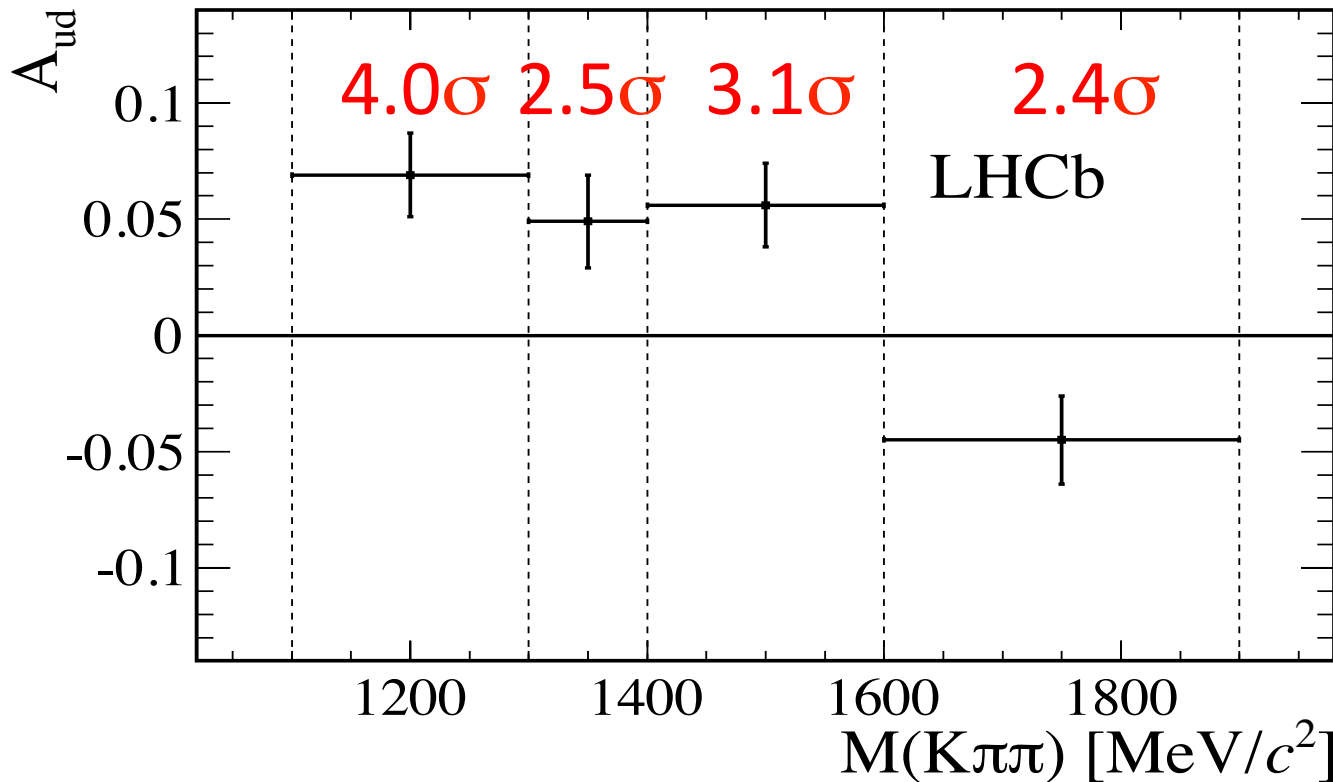
	[1.1, 1.3]	[1.3, 1.4]	[1.4, 1.6]	[1.6, 1.9]
c_1	6.3 ± 1.7	5.4 ± 2.0	4.3 ± 1.9	-4.6 ± 1.8
c_2	31.6 ± 2.2	27.0 ± 2.6	43.1 ± 2.3	28.0 ± 2.3
c_3	-2.1 ± 2.6	2.0 ± 3.1	-5.2 ± 2.8	-0.6 ± 2.7
c_4	3.0 ± 3.0	6.8 ± 3.6	8.1 ± 3.1	-6.2 ± 3.2
A_{UD}	6.9 ± 1.7	4.9 ± 2.0	5.6 ± 1.8	-4.5 ± 1.9

This is a statistically-limited analysis.

Up-down asymmetry results



Up-down asymmetry significance



Combining the four bins, the significance of the A_{UD} being different from zero is of 5.2σ . This can't be translated into a measurement of the γ polarization due to theoretical limitations.

First observation of photon polarization in $b \rightarrow s\gamma$ transitions

Summary and Future

A study of $B \rightarrow K\pi\pi\gamma$ decay is performed on 3fb^{-1} data sample.

We have reported the first observation of photon polarization in $b \rightarrow s\gamma$ transitions, with a 5.2σ significance.

More theoretical input is needed to translate measured up-down asymmetry into a measurement of the photon polarization.

The measurement of the photon polarization at LHCb is also promising with:

- Proper time distribution of $B_s \rightarrow \phi\gamma$.
- Transverse asymmetry in $B \rightarrow K^*ee$.
- Angular distribution in $B \rightarrow \phi K\gamma$.
- Radiative b-baryon decays: $\Lambda_b \rightarrow \Lambda^{(*)}\gamma$, $\Xi_b \rightarrow \Xi^{(*)}\gamma$.

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THANK YOU!

Back-up

Up-down asymmetry

$$A_{UD} \equiv \frac{\int_0^1 d\cos\theta \frac{d\Gamma}{d\cos\theta} - \int_{-1}^0 d\cos\theta \frac{d\Gamma}{d\cos\theta}}{\int_{-1}^1 d\cos\theta \frac{d\Gamma}{d\cos\theta}} = \frac{3}{4} \lambda_\gamma \frac{\int ds ds_{13} ds_{23} \text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{\int ds ds_{13} ds_{23} |\vec{J}|^2}$$

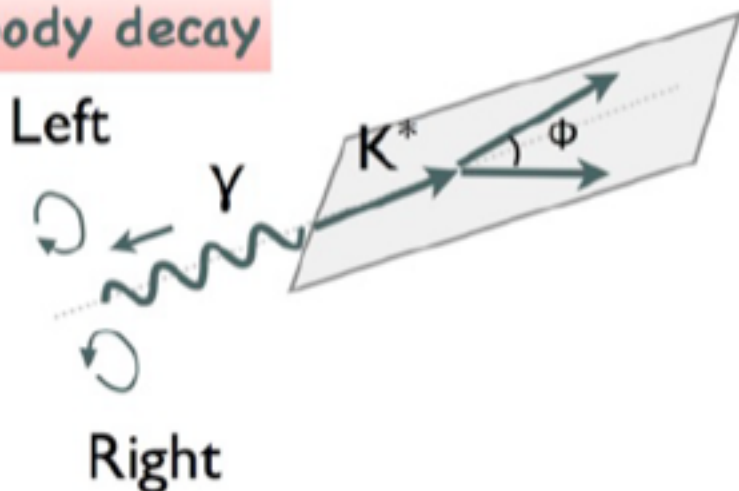
In the case of a single resonance

If J is known, the up-down asymmetry allows the computation of the photon polarization

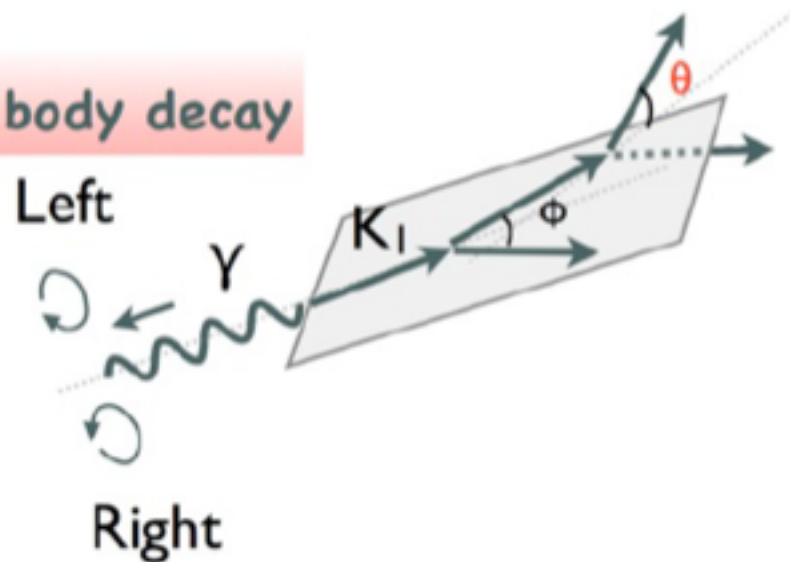
Why three charged particles?

Three tracks is the minimum needed to build a P-odd triple product proportional to the photon polarization using the final state momenta

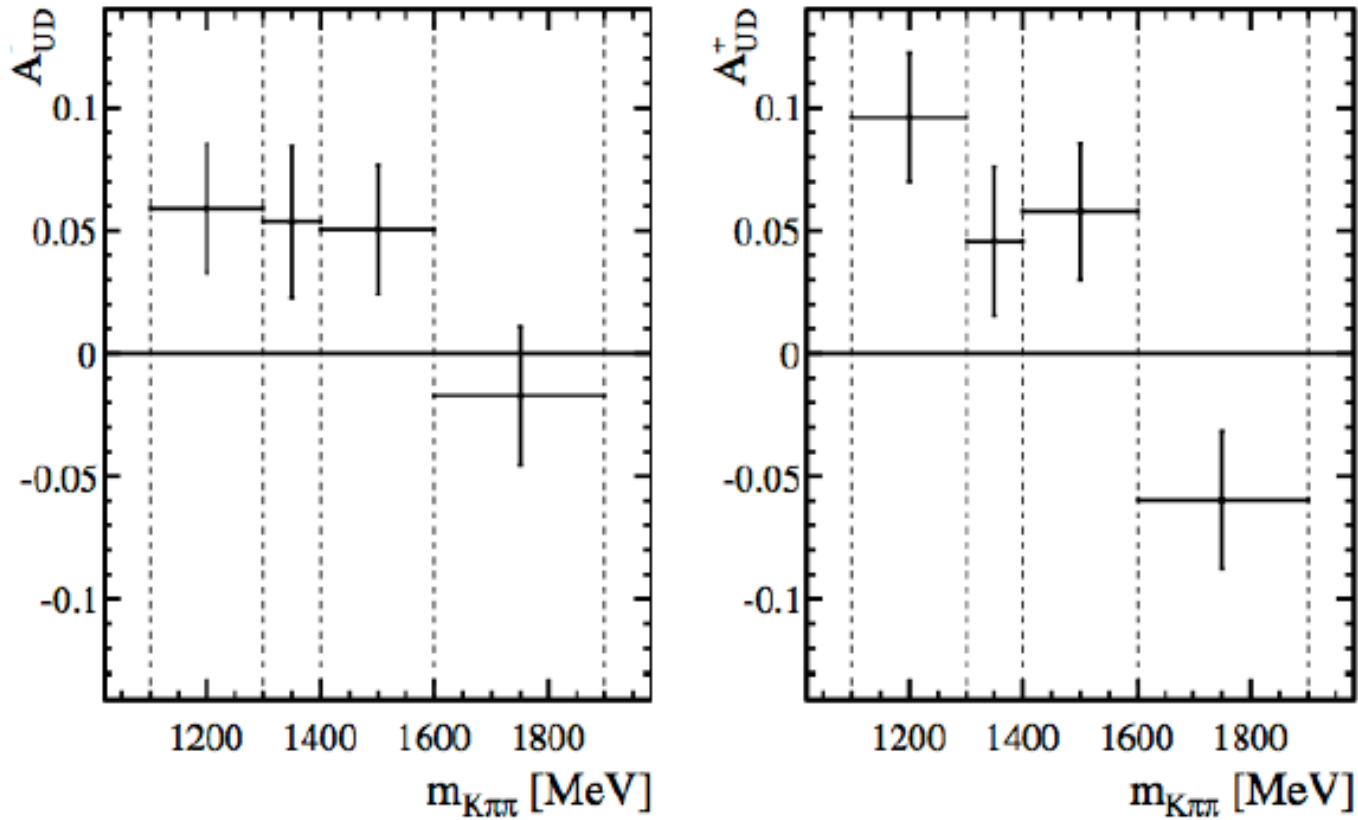
2 body decay



3 body decay



Up-down asymmetry with counting method



The counting method gives compatible results