
Improved τ -weapons for Higgs hunting

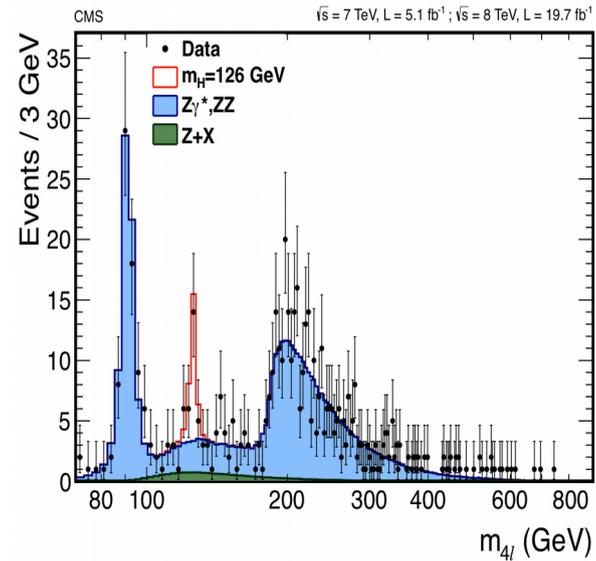
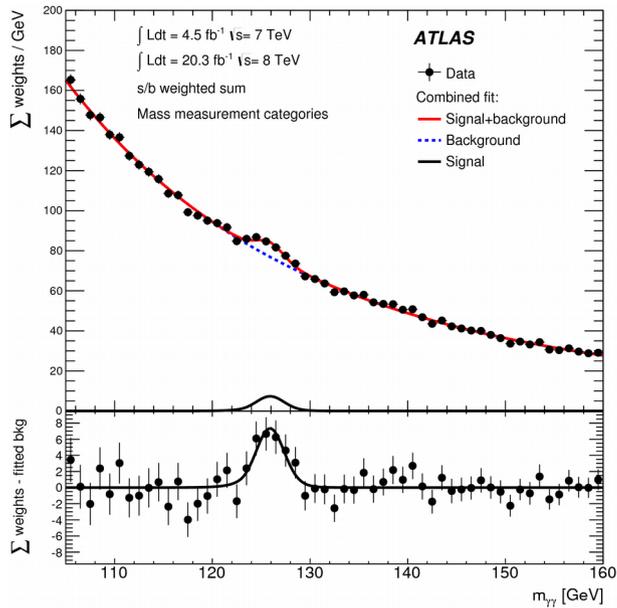
Oscar Vives



SUSY 2014

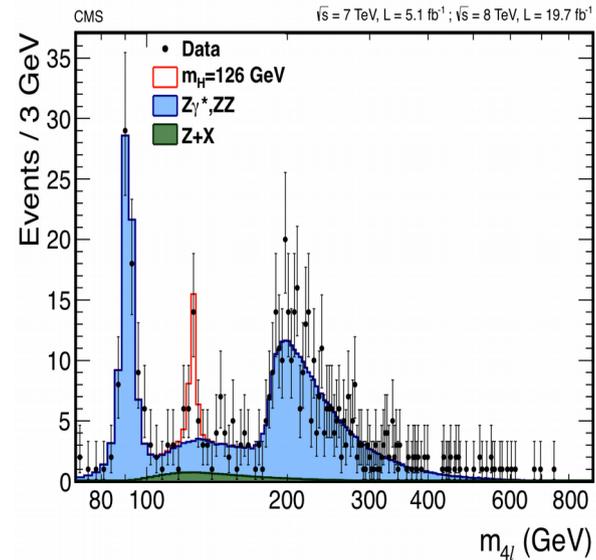
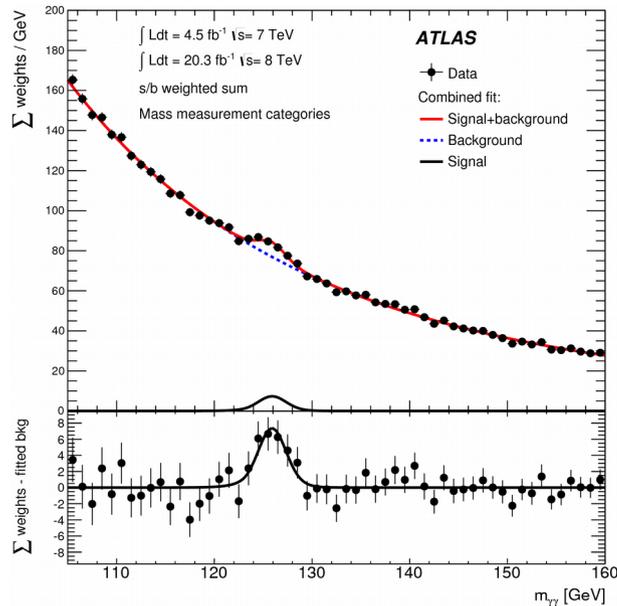
Manchester 21-26 July

4th July 2012



"Higgs" boson discovered...

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But... what is it??, SM Higgs??,
MSSM??, other ??

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Generic 2HdM in the presence of SUSY.

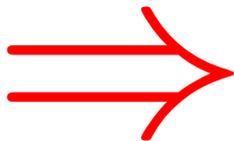
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Analysis in **CMSSM**, **PMSSM**...

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Generic **2HdM** in the presence of **SUSY**.



New couplings, extra scalar states...

- G. Barenboim, C. Bosch, M. L. Lopez-Ibañez and O. Vives, Phys. Rev. D90 (2014) 015003, arXiv:hep-ph/1311.7321.
- G. Barenboim, C. Bosch, M. L. Lopez-Ibañez and O. Vives, JHEP 1311, 051(2013), arXiv:hep-ph/1307.5973.

SCENARIO:

2HdM in generic MSSM at M_w with CP violation.

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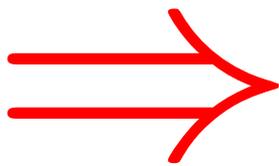
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- Lightest Higgs corresponds to the observed scalar state at LHC
- Heavier neutral and charged Higgs masses (all close in the decoupling limit) and Higgs mixings free. Only constrained by LHC (and indirect) observables.
- MSSM parameters independent and free, subject to experimental constraints.
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General model including all possible MSSMs

2 Higgs doublets with CP violating phase

$$\Phi_1 = \begin{pmatrix} \frac{1}{\sqrt{2}} (v_1 + \phi_1 + ia_1) \\ \phi_1^- \end{pmatrix}; \quad \Phi_2 = e^{i\xi} \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (v_2 + \phi_2 + ia_2) \end{pmatrix}$$

Generic 3 x 3 Higgs mass matrix.

$$M_H^2 = (\phi_1, \phi_2, a) \begin{pmatrix} M_{S_{2 \times 2}}^2 & M_{SP_{1 \times 2}}^2 \\ M_{PS_{2 \times 1}}^2 & M_P^2 \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ a \end{pmatrix}$$

Diagonalized by an unknown unitary matrix \mathcal{U} :

$$\mathcal{U} \cdot M_H^2 \cdot \mathcal{U}^T = \text{Diag} (M_{H_1}^2, M_{H_2}^2, M_{H_3}^2) .$$

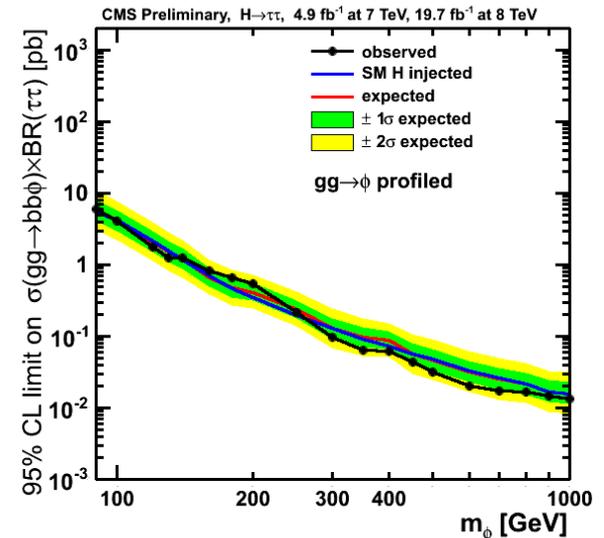
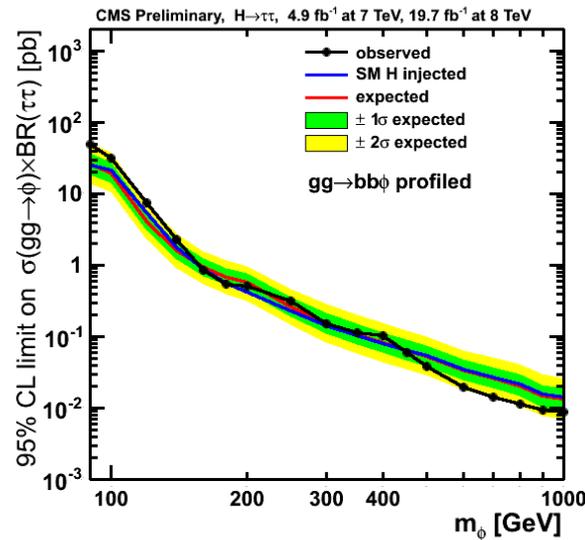
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- LHC constraints:

At 95% C.L. $0.75 \leq \mu_\gamma = \frac{\sigma(pp \rightarrow H) \times \text{BR}(H \rightarrow \gamma\gamma)}{\sigma(pp \rightarrow H)_{\text{SM}} \times \text{BR}(H \rightarrow \gamma\gamma)_{\text{SM}}} \leq 1.55$

$H \rightarrow \tau\tau$
Searches of
extra Higgs

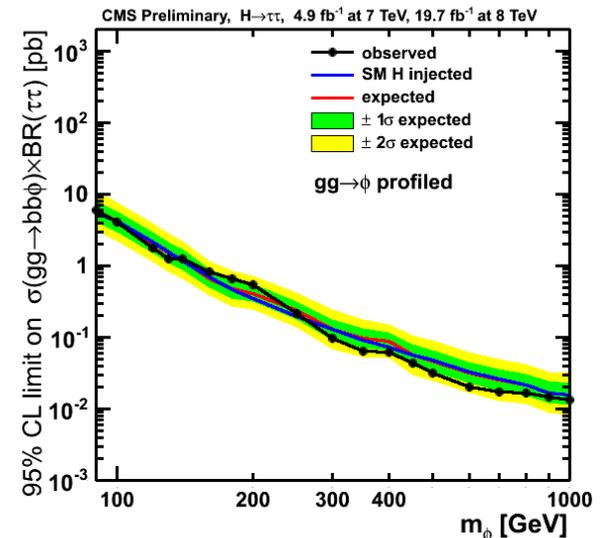
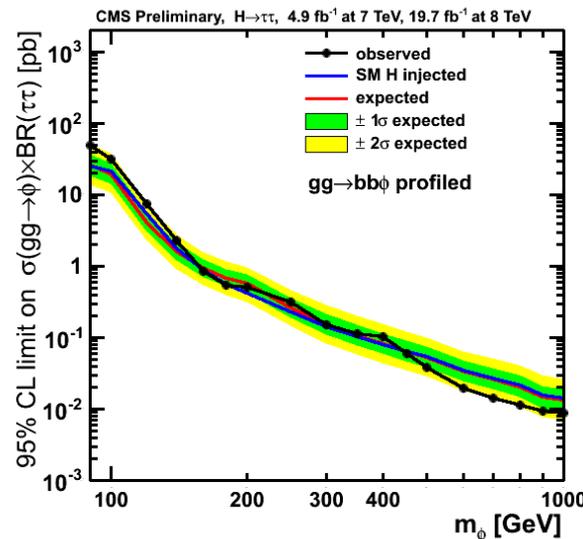


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- Indirect constraints:

At 95% C.L. $2.99 \times 10^{-4} \leq \text{BR}(B \rightarrow X_s \gamma) \leq 3.87 \times 10^{-4}$

also $B_s \rightarrow \mu^+ \mu^-$, and at low masses, $t \rightarrow H^+ b$

Model Analysis

Higgs to $\gamma\gamma$ at LHC

$$\sigma(pp \rightarrow H \rightarrow \gamma\gamma) = \sigma(pp \rightarrow H_1) \times \Gamma(H_1 \rightarrow \gamma\gamma)/\Gamma_{H_1}$$

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1.- Decay width

$$\Gamma(H_a \rightarrow \gamma\gamma) = \frac{M_{H_a}^3 \alpha^2}{256\pi^3 v^2} \left[|S_a^\gamma(M_{H_a})|^2 + |P_a^\gamma(M_{H_a})|^2 \right]$$

$$\begin{aligned} S_{H_1^0}^\gamma &\simeq \mathcal{U}_{11} \left(\frac{-8.3}{\tan\beta} + (-0.025 + i 0.034) \operatorname{Re} \left\{ \frac{\tan\beta}{1 + \kappa_d \tan\beta} \right\} - 0.45 \left(\frac{m_{\tilde{t}_2}^2}{m_{\tilde{t}_1}^2} - 1 \right) \operatorname{Re} \left\{ \frac{\mu m_t \mathcal{R}_{11}^* \mathcal{R}_{21}}{m_{\tilde{t}_2}^2} \right\} \right) + \\ &\mathcal{U}_{12} \left(-6.5 + 0.45 \left(\frac{m_{\tilde{t}_2}^2}{m_{\tilde{t}_1}^2} - 1 \right) \operatorname{Re} \left\{ \frac{A_t^* m_t \mathcal{R}_{11}^* \mathcal{R}_{21}}{m_{\tilde{t}_2}^2} \right\} + 0.45 \left(\frac{m_t^2 |\mathcal{R}_{11}|^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2 |\mathcal{R}_{12}|^2}{m_{\tilde{t}_2}^2} \right) \right) + \\ &\mathcal{U}_{13} \left((-0.025 + i 0.034) \operatorname{Im} \left\{ \frac{\kappa_d \tan^2\beta}{1 + \kappa_d \tan\beta} \right\} + 0.45 \left(\frac{m_{\tilde{t}_2}^2}{m_{\tilde{t}_1}^2} - 1 \right) \operatorname{Im} \left\{ \frac{\mu m_t \mathcal{R}_{11}^* \mathcal{R}_{21}}{m_{\tilde{t}_2}^2} \right\} \right) \end{aligned}$$

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2.- Production cross section

$$\begin{aligned} \sigma(pp \rightarrow H_1) \simeq & \left[0.16 \frac{\tan^2 \beta}{(1 + \kappa_d \tan \beta)^2} \left(|\mathcal{U}_{11}|^2 + |\mathcal{U}_{13}|^2 \right) + \right. \\ & \frac{0.1 \tan^2 \beta}{(1 + \kappa_d \tan \beta)^2} \mathcal{U}_{11}^2 - \frac{1.4 \tan \beta}{1 + \kappa_d \tan \beta} \mathcal{U}_{11} \mathcal{U}_{12} + 13 \mathcal{U}_{12}^2 \\ & \left. + \left(\frac{2}{(1 + \kappa_d \tan \beta)} + \frac{0.1 \tan^2 \beta}{(1 + \kappa_d \tan \beta)^2} + \frac{27}{\tan^2 \beta} \right) \mathcal{U}_{23}^2 \right] \text{ pb} \end{aligned}$$

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3.- Total width

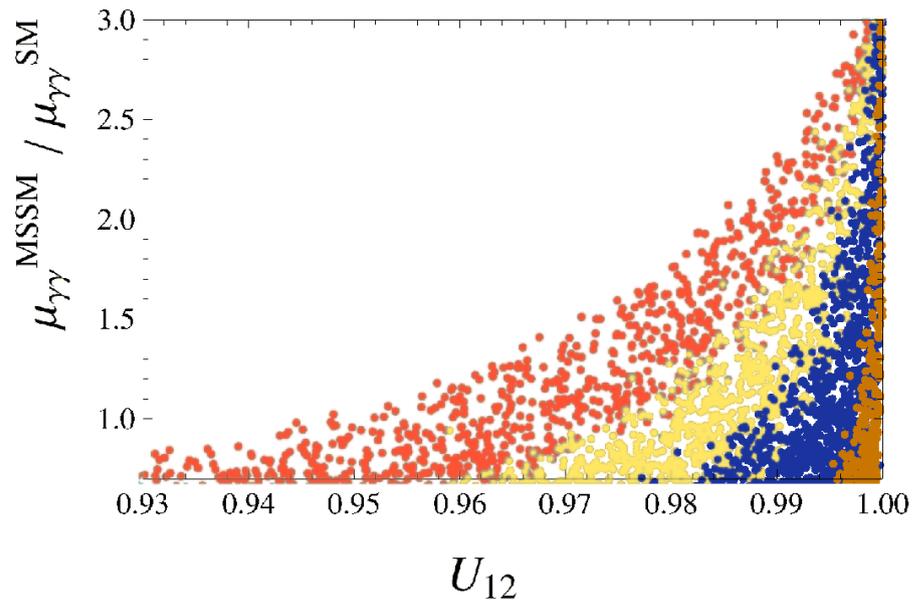
$$\Gamma_{H_1} \simeq \frac{g^2 M_{H_1}}{32\pi M_W^2} \left[\tan^2 \beta (\mathcal{U}_{11}^2 + \mathcal{U}_{13}^2) (3m_b^2 + m_\tau^2) + 6.7 \times 10^{-4} M_{H_1}^2 \left(\mathcal{U}_{12} + \frac{\mathcal{U}_{11}}{\tan \beta} \right)^2 \right]$$

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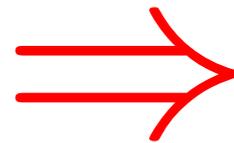
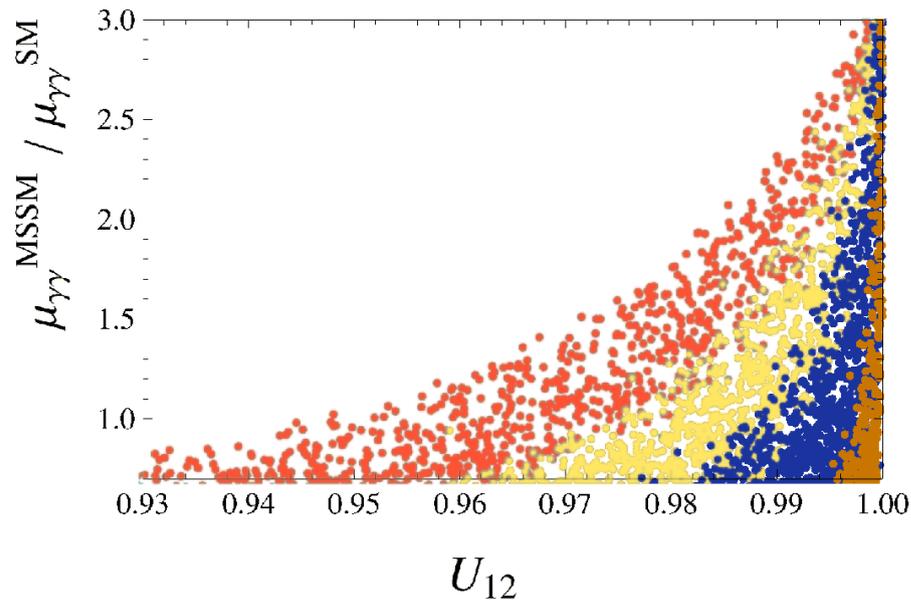


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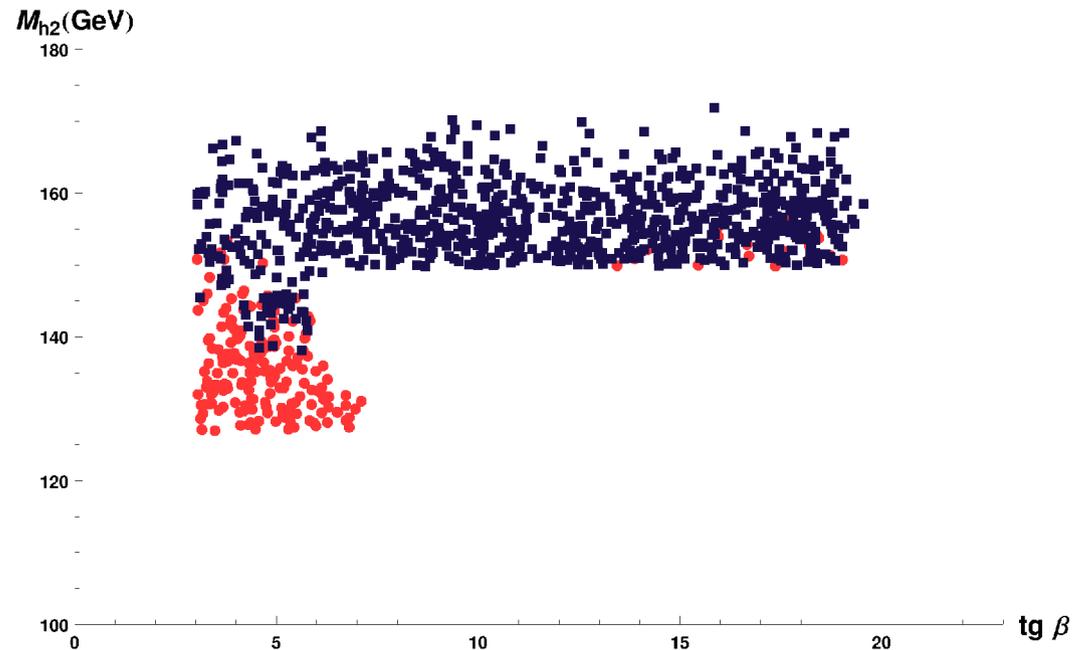
$$U_{12} \sim 1$$

$$U_{11}, U_{13} \sim 1/\text{tg } \beta$$

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"Old" ATLAS/CMS $\tau\tau$ searches

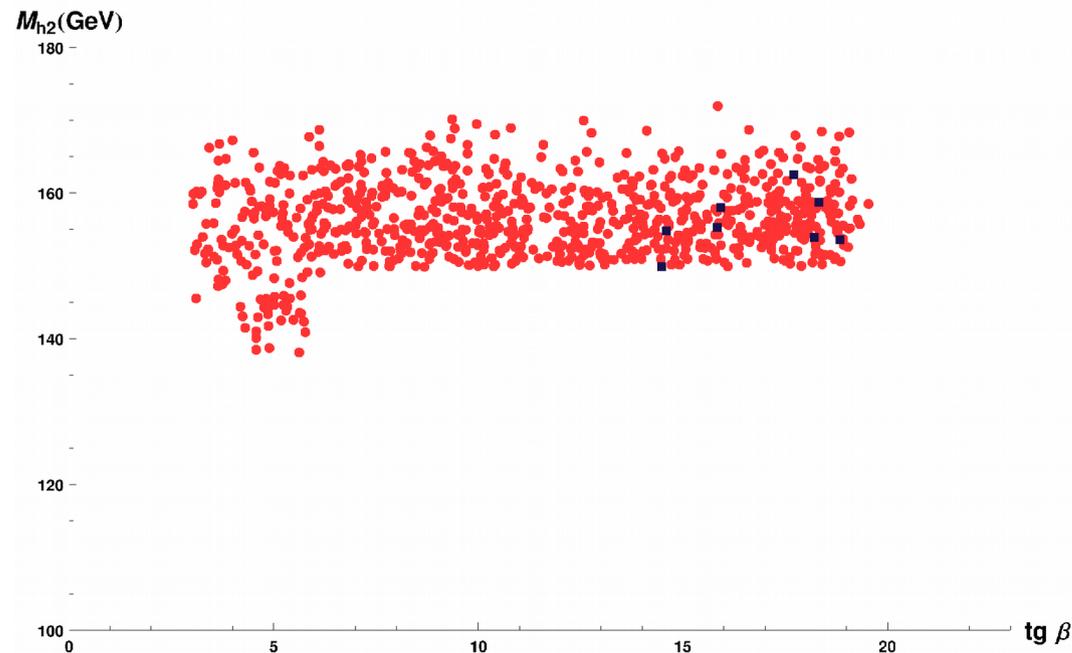


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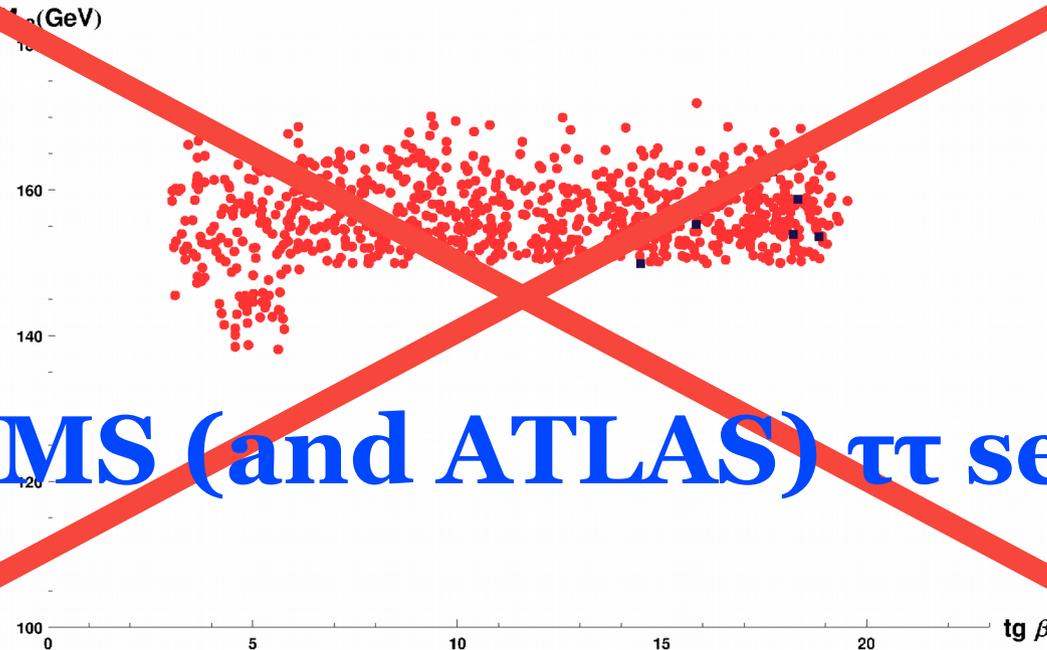


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New CMS (and ATLAS) $\tau\tau$ searches.

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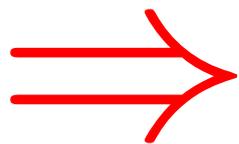
$\text{BR}(B \rightarrow X_s \gamma), B_s \rightarrow \mu^+ \mu^- \dots$

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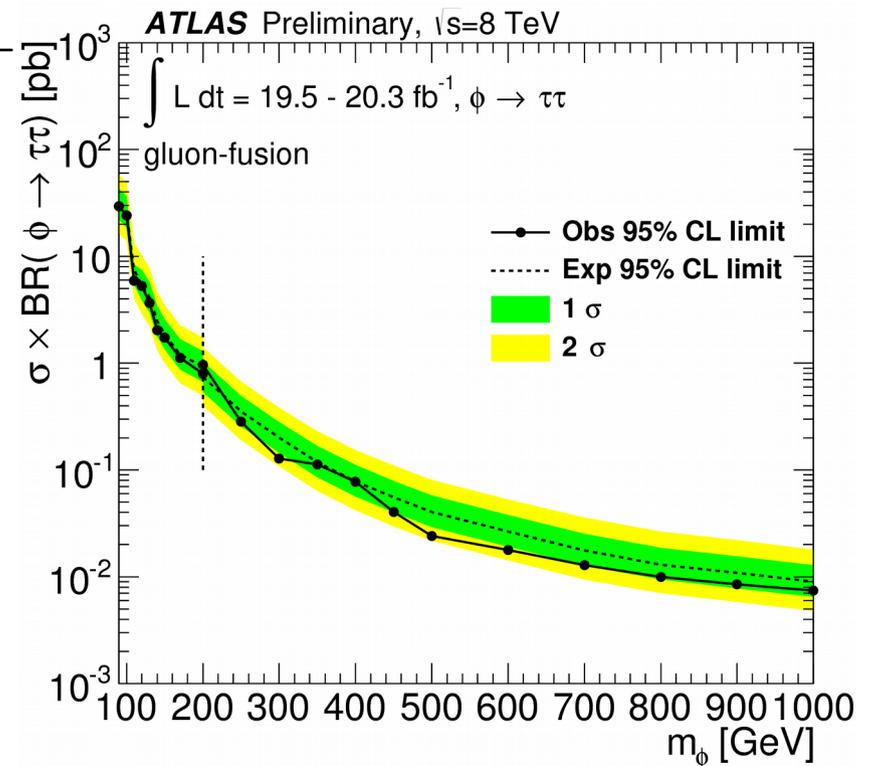
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Importance of b-associated production...

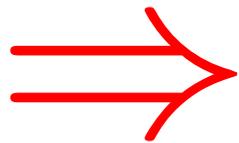


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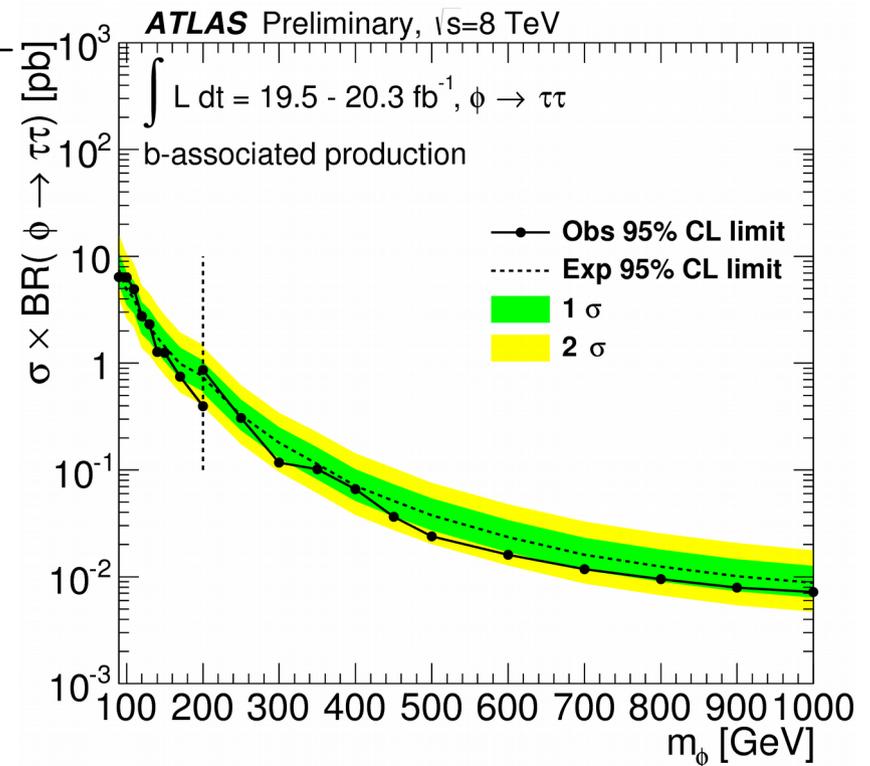
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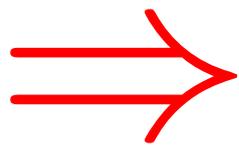


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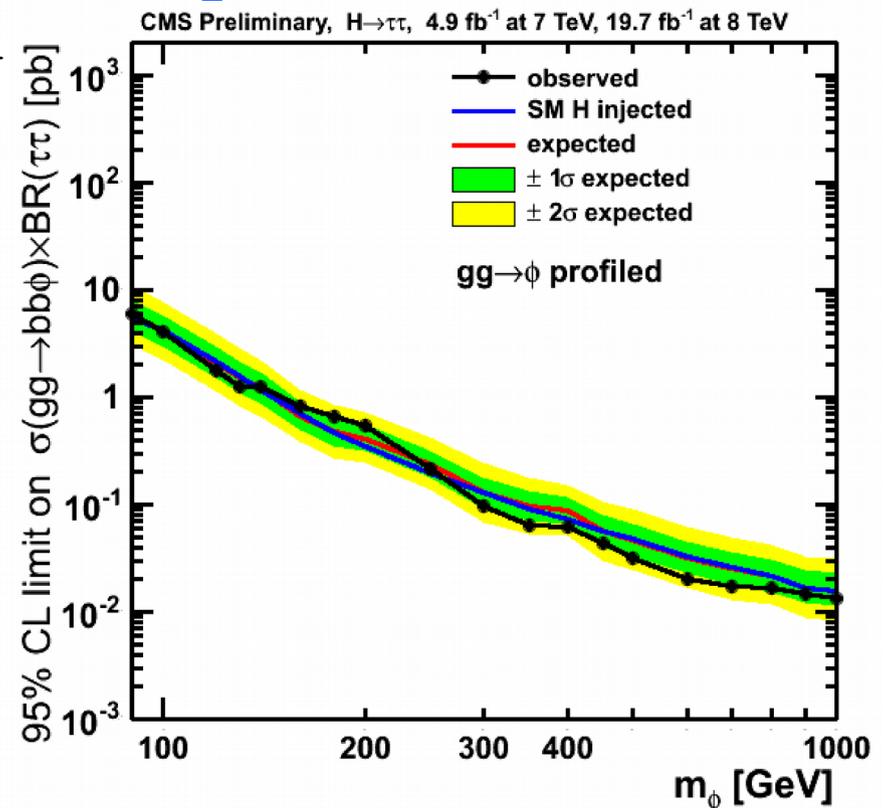
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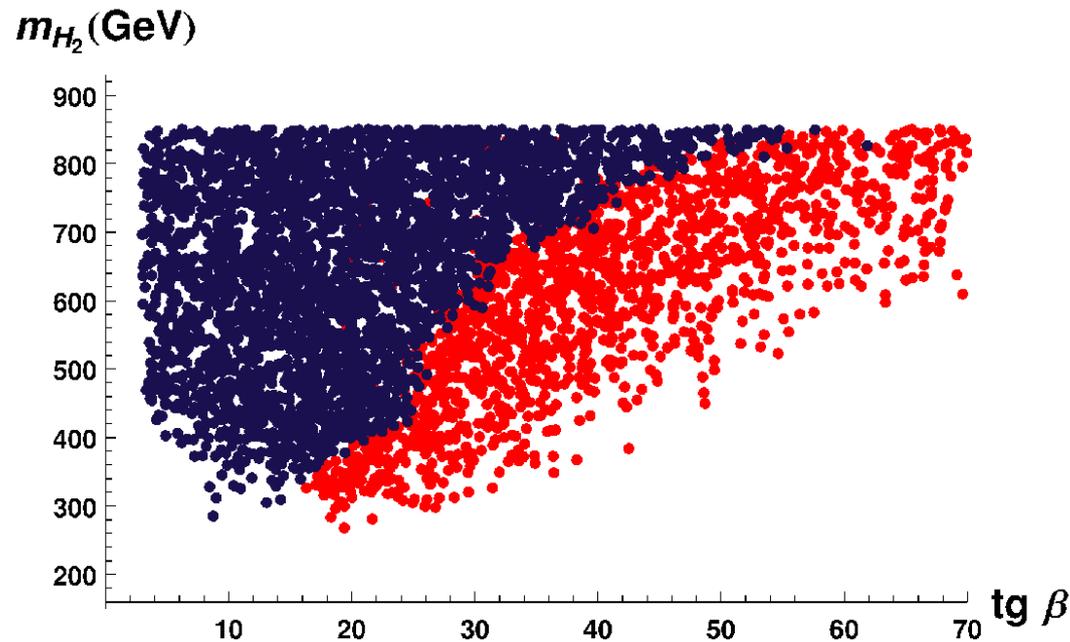


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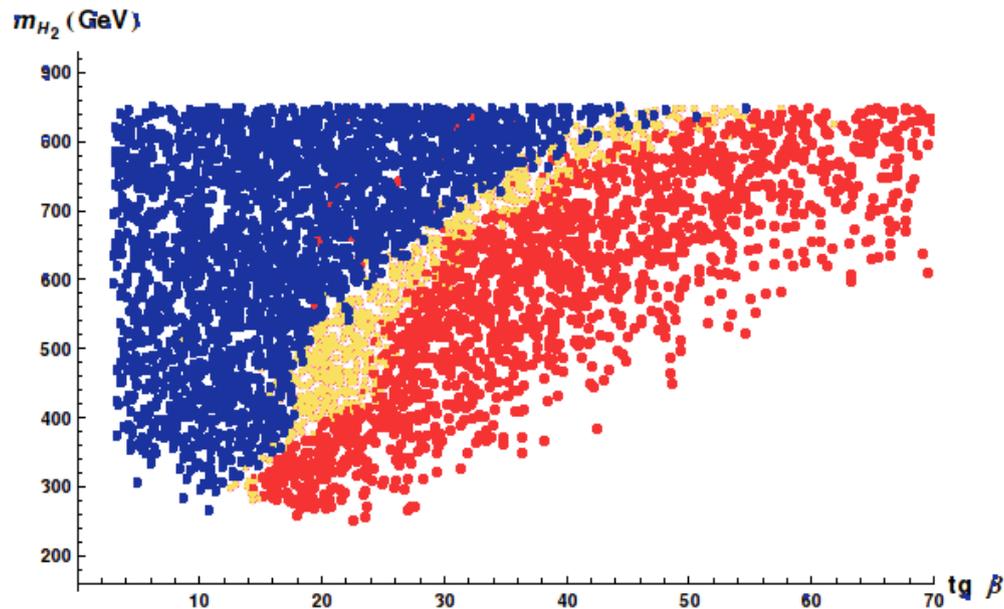
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Conclusions

- Lightest Higgs at 126 GeV boson and the $\gamma\gamma$ channel provide strong constraints on Higgs mixings.
- Combination of direct searches with indirect constraints is an excellent weapon to search for additional Higgs states.
- $\tau\tau$ channel together with indirect constraints forbid extra Higgs states below 250 GeV.
- Strong constraints for large $\tan\beta$ at higher masses.
- Waiting for positive results in the $\tau\tau$ channel !! ...