

Pati-Salam GUT-Flavour Models with Three Higgs Generations

Florian Hartmann

in collaboration with
Wolfgang Kilian and Karsten Schnitter

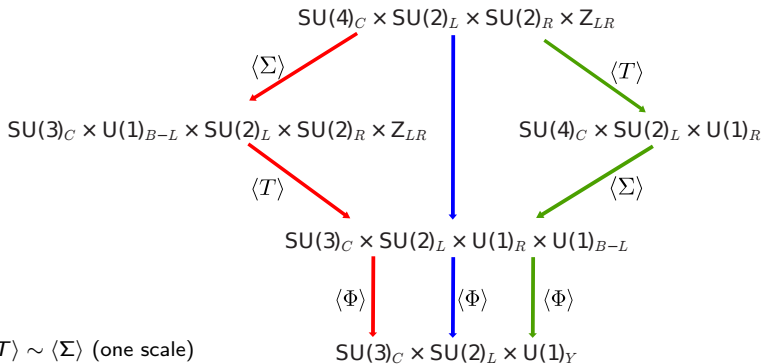
based on:

JHEP **1405** (2014) 064 and arXiv:1405.1901

Universität Siegen
Theoretische Physik I

SUSY14
Manchester, July 24th, 2014

- Consider SUSY Pati-Salam models
 - multi-step breaking, intermediate mass scales
- include 3 generations of MSSM-Higgs doublets
 - consistent with matter-Higgs unification
- embed flavour via flavour symmetry breaking
 - theory of flavour near the Planck scale
 - single $SU(3)_F$ spontaneously broken by flavons
 - consider different flavon representations



Class:

- A: $\langle \Phi \rangle \sim \langle T \rangle \sim \langle \Sigma \rangle$ (one scale)
- B: $\langle \Phi \rangle \ll \langle T \rangle \ll \langle \Sigma \rangle$ (three scales)
- C: $\langle \Phi \rangle \ll \langle \Sigma \rangle \ll \langle T \rangle$ (three scales)
- D: $\langle \Phi \rangle \ll \langle T \rangle \sim \langle \Sigma \rangle$ (two scales)
- E: $\langle \Phi \rangle \ll \langle \Sigma \rangle$ and $\langle T \rangle = 0$ (two scales)
- F: $\langle \Phi \rangle \ll \langle T \rangle$ and $\langle \Sigma \rangle = 0$ (two scales)

PS: $(SU(4)_C, SU(2)_L, SU(2)_R)$ $SO(10)$

Breaking to Intermediate Symmetry

 $\Sigma = (\mathbf{15}, \mathbf{1}, \mathbf{1})$ to break $SU(4)$ $T_R = (\mathbf{1}, \mathbf{1}, \mathbf{3})$ to break $SU(2)_R$ $T_L = (\mathbf{1}, \mathbf{3}, \mathbf{1})$ and $E = (\mathbf{6}, \mathbf{2}, \mathbf{2})$

} 45

Breaking to $SU(3)_C \times SU(2)_L \times U(1)_Y$ $\Phi_R = (\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})$ and $\bar{\Phi}_R = (\mathbf{4}, \mathbf{1}, \mathbf{2})$ to break PS $\Phi_L = (\mathbf{4}, \mathbf{2}, \mathbf{1})$ and $\bar{\Phi}_L = (\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})$ } $16 \oplus \bar{16}$

optional MSSM Higgs

 $h = (\mathbf{1}, \mathbf{2}, \mathbf{2})$: MSSM-higgs $F = (\mathbf{6}, \mathbf{1}, \mathbf{1})$: possibly light triplets

} 10

- construct renormalizable Higgs sector superpotential
- calculate tree-level Higgs masses

Results

- see-saw-like mass terms in Higgs fields
 - new intermediate mass-scale: $M_{\text{IND}} \sim \frac{\langle \Phi \rangle^2}{\langle \Sigma \rangle + \langle T \rangle}$
 - ↪ light colour triplets (down to TeV-scale)
 - ↪ useful for unification

W. Kilian, J.Reuter: arXiv:hep-ph/0606277

- MSSM-Higgs doublet:

$$\langle T \rangle \neq 0: \mu\text{-term} \sim \frac{\langle \Phi \rangle^2}{\langle T \rangle}$$

$\langle T \rangle = 0$: two massless doublets from Φ_L (at MSSM-scale)

↪ makes $h(1, 2, 2)$ optional

↪ $SU(2)_R$ and $SU(2)_L$ breaking from same $SO(10)$ representation

Requirements for unification

- (1) unification condition at the LR-scale

$$\alpha_R^{-1} = \alpha_L^{-1}$$

- (2) unification condition at the PS-scale

$$\alpha_{B-L}^{-1} = \alpha_3^{-1} \quad ; \quad \alpha_3^{-1} = \alpha_4^{-1}$$

- (3) complete unification of couplings near the Planck-scale

- (4) PS hypercharge condition $Y = T_3^R + \frac{B-L}{2}$

$$\alpha_Y^{-1} = \alpha_R^{-1} + \frac{2}{3}\alpha_{B-L}^{-1}$$

Inputs:

- MSSM-scale: 2.5 TeV
- 4 unification-requirements
- but 4 (5) mass-scales

⇒ scales are constrained but not fixed

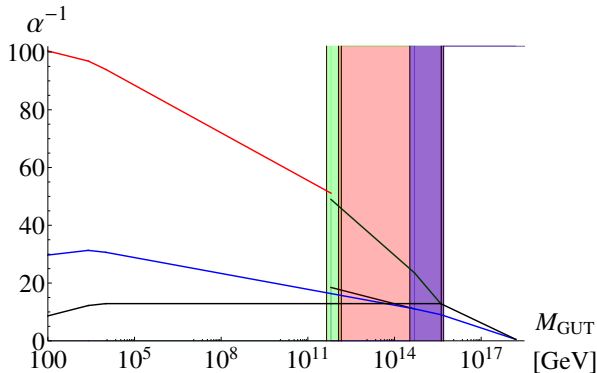
- For each breaking class
 - scan all allowed field configurations
 - each field may appear 0, 1 or 3 times
- For each configuration
 - test whether unification below 10^{19} GeV possible
 - find allowed ranges for all mass scales

- For each breaking class
 - scan all allowed field configurations
 - each field may appear 0, 1 or 3 times
- For each configuration
 - test whether unification below 10^{19} GeV possible
 - find allowed ranges for all mass scales

	class B	class C	class E	class F	Σ
scanned	144	144	324	216	828
GCU	18	57	254	29	358
$M_{\text{GUT}} > 10^{16}$ GeV	18	57	131	29	235
$M_{\text{IND}} < 10$ TeV and $M_{\text{GUT}} > 10^{16}$ GeV	8	34	72	0	114
$M_{\text{LR}} < 100$ TeV	1	11	108	0	120
10^{12} GeV $< M_{N_R} < 10^{14}$ GeV	16	42	123	3	184
$M_{\text{IND}} \in [0.1, 10] \frac{v_\Phi^2}{v_\Sigma + v_T}$	14	20	203	26	263

Field Content

	h	F	$\Phi_{L/R}$	$\bar{\Phi}_{L/R}$	Σ	E	$T_{L/R}$
PS	(1, 2, 2)	(6, 1, 1)	(4, 2, 1) _{L/R}	($\bar{4}$, 2, 1) _{L/R}	(15, 1, 1)	(6, 2, 2)	(1, 3, 1) _{L/R}
#	3	3	1	1	3	0	1



$$M_{\text{IND}} = 10^4 \text{ GeV}$$

$$M_{\text{GUT}} = 10^{18.2} \text{ GeV}$$

green band [GeV]:

$$7 \cdot 10^{11} \leq M_{\text{MSSM}} \leq 1 \cdot 10^{12}$$

red band [GeV]:

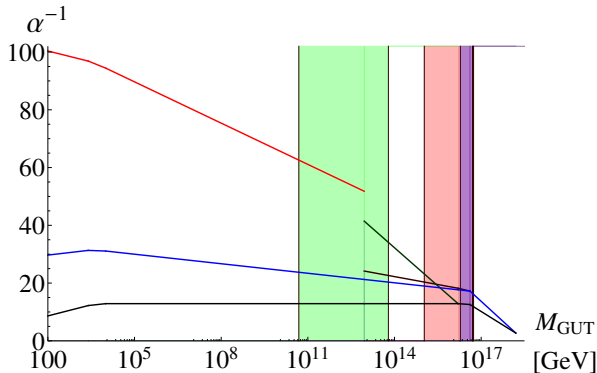
$$2 \cdot 10^{13} \leq M_{\text{LR}} \leq 2 \cdot 10^{14}$$

blue band [GeV]:

$$3 \cdot 10^{15} \leq M_{\text{PS}} \leq 4 \cdot 10^{15}$$

Field Content

	h	F	$\Phi_{L/R}$	$\bar{\Phi}_{L/R}$	Σ	E	$T_{L/R}$
PS	(1, 2, 2)	(6, 1, 1)	(4, 2, 1) _{L/R}	($\bar{4}$, 2, 1) _{L/R}	(15, 1, 1)	(6, 2, 2)	(1, 3, 1) _{L/R}
#	1	3	1	1	1	3	1



$$M_{\text{IND}} = 10^4 \text{ GeV}$$

$$M_{\text{GUT}} = 10^{18.2} \text{ GeV}$$

green band [GeV]:

$$5 \cdot 10^{10} \leq M_{\text{MSSM}} \leq 6 \cdot 10^{13}$$

red band [GeV]:

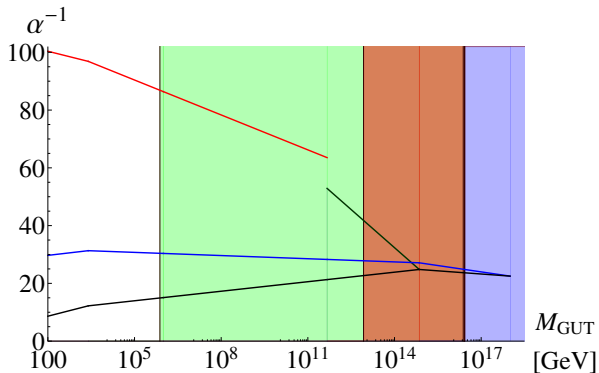
$$1 \cdot 10^{15} \leq M_{\text{QL}} \leq 5 \cdot 10^{16}$$

blue band [GeV]:

$$2 \cdot 10^{16} \leq M_{\text{PS}} \leq 6 \cdot 10^{16}$$

Field Content

	h	F	$\Phi_{L/R}$	$\bar{\Phi}_{L/R}$	Σ	E	$T_{L/R}$
PS	$(1, 2, 2)$	$(6, 1, 1)$	$(4, 2, 1)_{L/R}$	$(\bar{4}, 2, 1)_{L/R}$	$(15, 1, 1)$	$(6, 2, 2)$	$(1, 3, 1)_{L/R}$
#	0	0	1	1	1	0	0



$$3 \cdot 10^3 \leq M_{\text{IND}} \leq 2 \cdot 10^{16}$$

green band [GeV]:

$$8 \cdot 10^8 \leq M_{\text{LR}} \leq 2 \cdot 10^{16}$$

red band [GeV]:

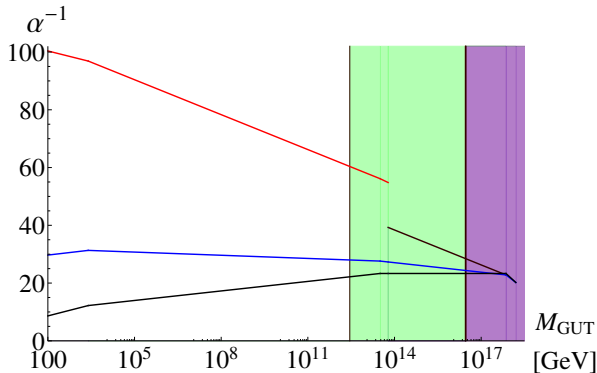
$$8 \cdot 10^{13} \leq M_{\text{PS}} \leq 3 \cdot 10^{16}$$

blue band [GeV]:

$$3 \cdot 10^{16} \leq M_{\text{GUT}} \leq 1 \cdot 10^{19}$$

Field Content

	h	F	$\Phi_{L/R}$	$\bar{\Phi}_{L/R}$	Σ	E	$T_{L/R}$
PS	(1, 2, 2)	(6, 1, 1)	(4, 2, 1) _{L/R}	($\bar{4}$, 2, 1) _{L/R}	(15, 1, 1)	(6, 2, 2)	(1, 3, 1) _{L/R}
#	1	3	3	3	3	1	1



not shown [GeV]:

$$6 \cdot 10^8 \leq M_{\text{IND}} \leq 3 \cdot 10^{16}$$

green band [GeV]:

$$3 \cdot 10^{12} \leq M_{\text{MSSM}} \leq 3 \cdot 10^{16}$$

red band [GeV]:

$$3 \cdot 10^{16} \leq M_{\text{QL}} \leq 9 \cdot 10^{18}$$

blue band [GeV]:

$$3 \cdot 10^{16} \leq M_{\text{PS}} \leq 1 \cdot 10^{19}$$

Flavour symmetry broken by flavons ϕ_k

$$\begin{aligned} \mathcal{L}_{\text{Yuk}} &= -\frac{c}{M^n} \Phi_{ij} h \Psi_i^L \Psi_j^{cR} &> M \sim M_{\text{Planck}} \\ &\rightarrow -\frac{c}{M^n} \langle \Phi_{ij} \rangle h \Psi_i^L \Psi_j^{cR} &< M \end{aligned}$$

with the SM-Yukawa-coupling $y_{ij} = -\frac{c}{M^n} \langle \Phi_{ij} \rangle$ and $\Phi_{ij} = \prod_k^n \phi_k$

Requirements:

(1) following form fits quark-data quite well

$$Y_{u/d} \approx \begin{pmatrix} 0 & O(\epsilon^3) & O(\epsilon^3) \\ O(\epsilon^3) & O(\epsilon^2) & O(\epsilon^2) \\ O(\epsilon^3) & O(\epsilon^2) & O(1) \end{pmatrix}$$

with $\epsilon_u \approx 0.05$ and $\epsilon_d \approx 0.15$

Roberts, Romanino, Ross, Velasco-Sevilla: arXiv:hep-ph/0104088

(2) neutrino-data also described in case of Sequential Right-Handed Neutrino Dominance (SRHND)

King: arXiv:hep-ph/0310204

- consider flavour- $SU(3)_F$ in Pati-Salam GUT
 - consider Higgs as flavour-triplet
- ⇒ consistent with matter-higgs-unification:
↪ $\Psi = (\mathbf{27}, \mathbf{3})$ (E_6 -unification)

Caution: new trivial invariant $\varepsilon^{ijk} \Psi_i^L \Psi_j^R h_k$

↪ large off-diagonal Yukawa matrix entries

⇒ additional discrete symmetry needed

Considered flavon representations

Model 1: $\phi_i \rightarrow \mathbf{3} / \bar{\mathbf{3}}$

Model 2: $\phi_i \rightarrow \bar{\mathbf{10}} = (\bar{\mathbf{3}} \otimes \bar{\mathbf{3}} \otimes \bar{\mathbf{3}})_s$

Ansatz

- consider $SU(3) \otimes PS$ -model with flavour triplet Higgs
(for singlet Higgs similar to existing models) King, Ross: arXiv:hep-ph/0307190
- Flavons are flavour triplets
 $\Rightarrow \langle \phi_3 \rangle = \langle \bar{\phi}_3 \rangle = (0, 0, \sqrt[3]{\epsilon}) M$; $\langle \phi_{23} \rangle = \langle \bar{\phi}_{23} \rangle = (0, \epsilon, \epsilon) M$
- additional symmetry: $U(1) \otimes Z_3$

Results:

- leading contribution from dim 6 term $(\Psi_L \bar{\phi}_3)(\Psi_R \bar{\phi}_3)(h \bar{\phi}_3)$
- dim 6 generates Yukawa 2-3 block ($\phi_3 \rightarrow \phi_{23}$)
- NLO leads to large corrections to Yukawa matrix
- only possible to forbid field-configurations (not invariants)
 \rightarrow fine tuning
- possible, but need to eliminate formally leading terms

Field	$SU(3)_F$	PS	Z_4	Z_2^R
Ψ_L	3	(4, 2, 1)	1	–
Ψ_R	3	($\bar{4}$, 1, 2)	1	–
h	3	(1, 2, 2)	1	+
ϕ_3	$\bar{\mathbf{10}}$	(1, 1, 1)	1	+
ϕ_2	$\bar{\mathbf{10}}$	(1, 1, 1)	1	+
$\bar{\phi}_3$	10	(1, 1, 1)	3	+
$\bar{\phi}_2$	10	(1, 1, 1)	0	+
Φ_R	$\bar{\mathbf{3}}$	(4, 1, 2)	1	+

$$\langle \phi_3 \rangle_{333} = \langle \bar{\phi}_3 \rangle_{333} \sim \epsilon$$

$$\langle \phi_2 \rangle_{ijk} = \langle \bar{\phi}_2 \rangle_{ijk} \sim \epsilon^3 \quad (i, j, k \geq 2)$$

$$\langle h \rangle = (1, 1, 1) \nu_{\text{MSSM}}$$

$$\langle \Phi_R \rangle = (0, 0, 1) \nu_\Phi$$

$$Y_{4\&5}^f \sim \begin{pmatrix} 0 & \epsilon^3 & \epsilon^3 \\ \epsilon^3 & \epsilon^2 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}$$

$$W_{\text{lead}} \sim \frac{1}{M} [\phi_3 \Psi_L \Psi_R h + \phi_2 \Psi_L \Psi_R h]$$

$$W_{\text{dim5}} \sim \frac{1}{M^2} [h \phi_3 \bar{\phi}_2 \Psi_L \Psi_R + h \phi_2 \bar{\phi}_2 \Psi_L \Psi_R]$$

$$W_{\text{Maj}} \sim \frac{1}{M'} (\Psi_R \Phi_R)^2 \left[1 + \frac{1}{M'^2} (\phi_3 \bar{\phi}_3 + \phi_2 \bar{\phi}_3 + \bar{\phi}_2 \bar{\phi}_2) \right. \\ \left. + \frac{1}{M'^3} (\phi_3 \bar{\phi}_2 \bar{\phi}_3 + \phi_2 \bar{\phi}_3 \bar{\phi}_2) \right]$$

$$M_{\text{Maj}} \sim \begin{pmatrix} \epsilon_\nu^6 & \epsilon_\nu^7 & \epsilon_\nu^5 \\ \epsilon_\nu^7 & \epsilon_\nu^4 & \epsilon_\nu^4 \\ \epsilon_\nu^5 & \epsilon_\nu^4 & 1 \end{pmatrix} \frac{\langle \phi_R \rangle^2}{M'}$$

- Eigenvalues of order: $(\epsilon_\nu^6 : \epsilon_\nu^4 : 1)$
- ↪ sequential dominance
- different expansion parameter $\epsilon_\nu < \epsilon_u$
- ↪ PMNS-like lepton mixing

Multi-step breaking of SUSY-Pati-Salam

- unification of gauge-couplings possible
- introduction of intermediate scale
 - ↪ intermediate neutrino scale
- mass scales can vary over quite large mass-ranges

Flavour triplet Pati-Salam Higgs models

- models with $SU(3)_F$ and flavour triplet Higgs are possible
- decuplet flavons reproduce flavour structure
- triplet flavons are challenging but possible

Backup slides

2-D Example

$$M_{RR} = \begin{pmatrix} Y & 0 \\ 0 & X \end{pmatrix} \quad \text{and} \quad M_{LR} = \begin{pmatrix} e & b \\ f & c \end{pmatrix}$$

with: $X \gg Y$

and see-saw $M_{LL} = M_{LR} \cdot M_{RR}^{-1} \cdot M_{LR}$

$$\Rightarrow M_{LL} = \begin{pmatrix} \frac{e^2}{Y} + \frac{b^2}{X} & \frac{ef}{Y} + \frac{bc}{X} \\ \frac{ef}{Y} + \frac{bc}{X} & \frac{f^2}{Y} + \frac{c^2}{X} \end{pmatrix} \approx \begin{pmatrix} \frac{e^2}{Y} & \frac{ef}{Y} \\ \frac{ef}{Y} & \frac{f^2}{Y} \end{pmatrix}$$

$$\Rightarrow m_1 \approx 0 \quad \text{and} \quad m_2 \approx \frac{e^2 + f^2}{Y}$$

$$\Rightarrow \tan(\theta_{23}) \approx \frac{e}{f}$$