

RGEs and group theory calculations with the Susyno program

Renato Fonseca

renato.fonseca@ific.uv.es

AHEP Group, Instituto de Física Corpuscular
CSIC/Universitat de València, Spain



Outline

I

The Susyno program: what is it for?

It was created to calculate the RGEs of SUSY models

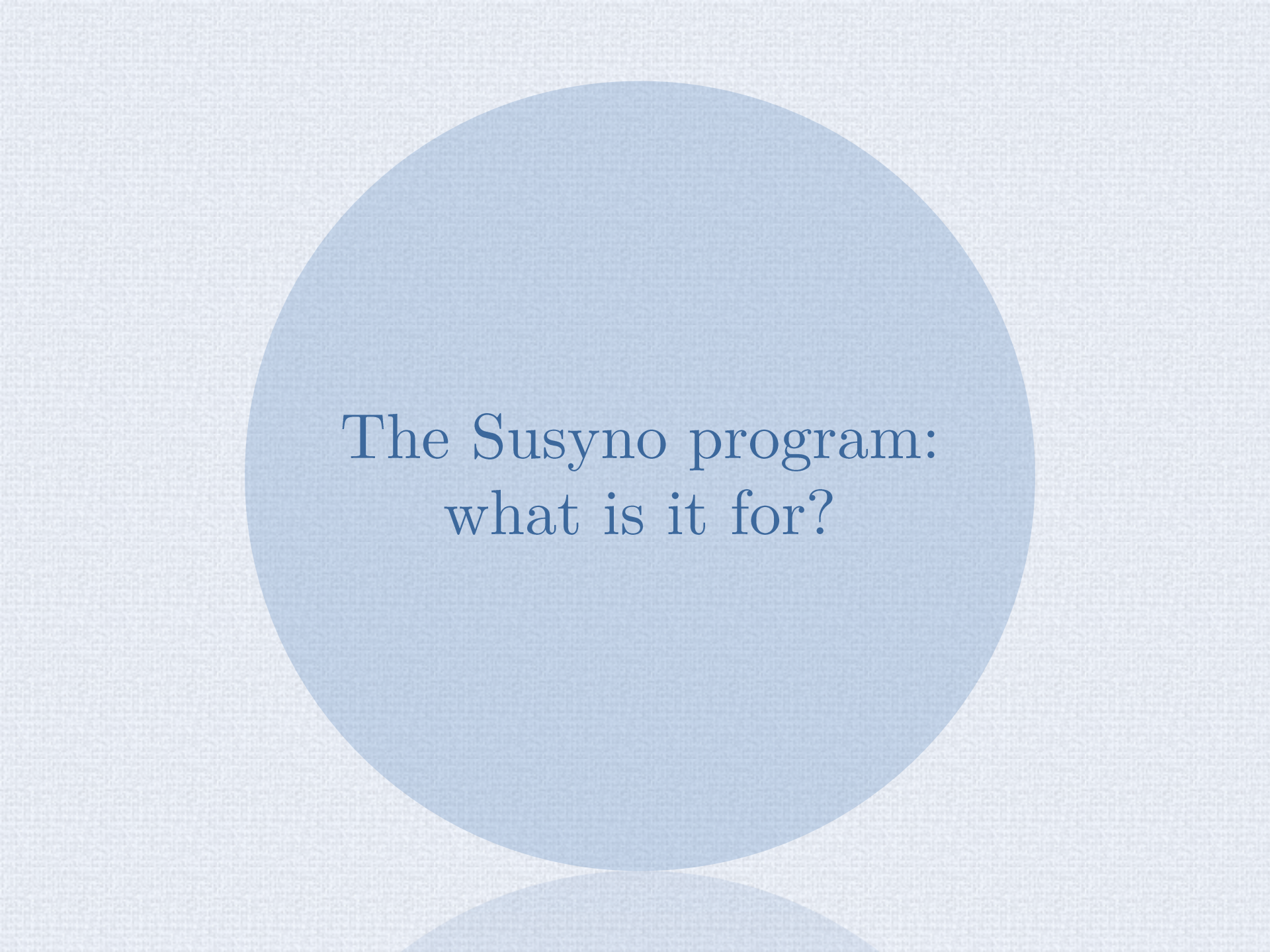
II

Using Susyno to make Group Theory calculations

I will focus mainly on this point in this presentation

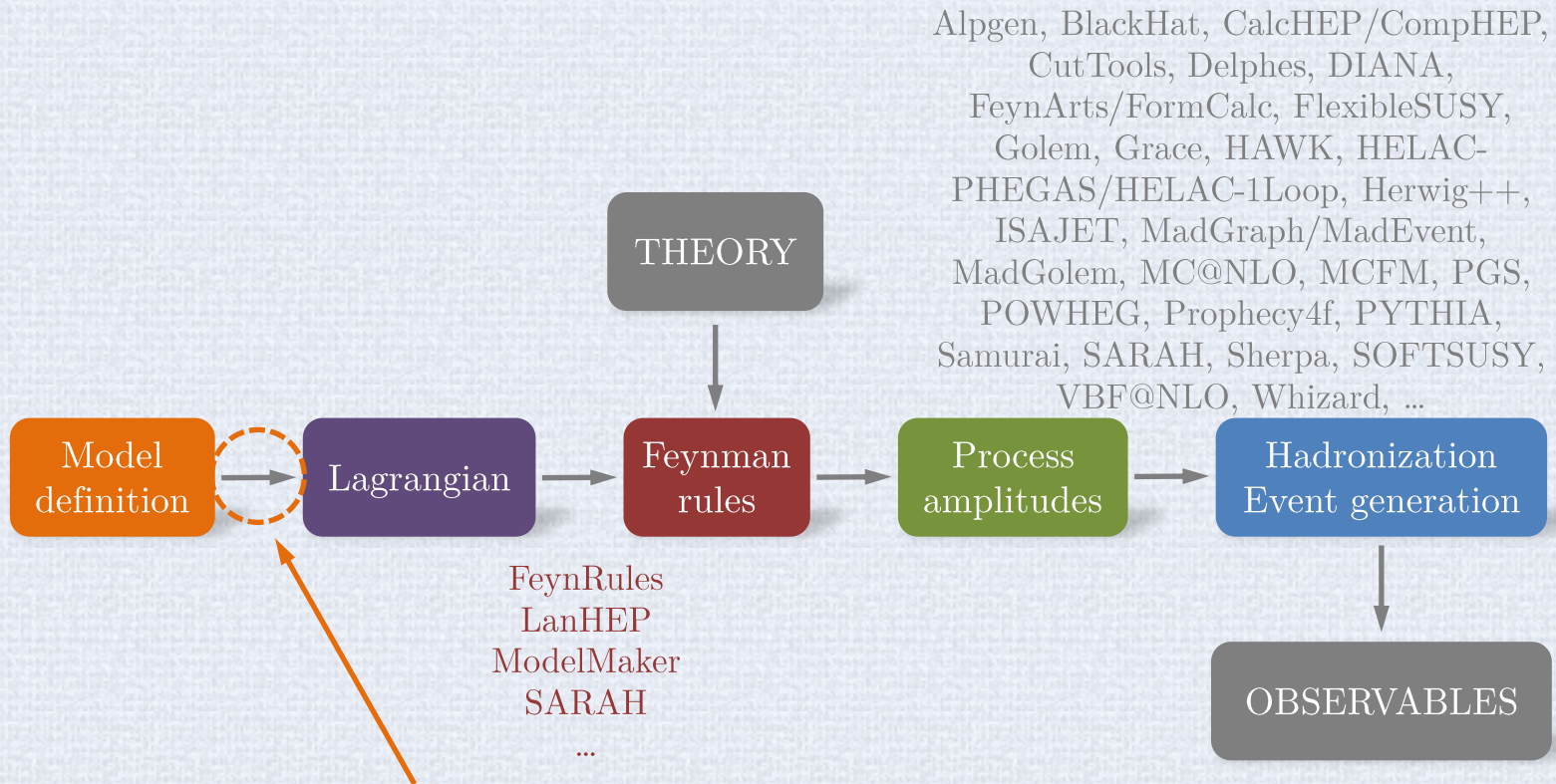
III

Summary



The Susyno program:
what is it for?

Automation in High Energy Physics



Gauge theories are defined by a **gauge group** and fields transforming under some of its **representations**. The rest are usually just conventions/notation (fixing a specific basis for the representations, naming the parameters, ...)

Handling of Group Theory is needed

Susyno

web.ist.utl.pt/renato.fonseca/susyno.html

Susyno is a Mathematica package which was created to compute the **Renormalization Group Equations (RGEs) of SUSY models**.

RF 2012

The two-loop **RGEs of a generic Yang-Mills theory** are known for SUSY (and non-SUSY) models, but in order to apply them to **specific models** it takes some work. One reason for this is that the equations are written **for and as a function of** the following **generic tensors**:

Martin, Vaughn
1994, 2008
Yamada 1994

$$W = \frac{1}{6} Y^{ijk} \Phi_i \Phi_j \Phi_k + \frac{1}{2} \mu^{ij} \Phi_i \Phi_j + L^i \Phi_i$$
$$-\mathcal{L}_{\text{soft}} = \left(\frac{1}{2} M_a \lambda^a \lambda^a + \frac{1}{6} h^{ijk} \phi_i \phi_j \phi_k + \frac{1}{2} b^{ij} \phi_i \phi_j + s^i \phi_i + \text{h.c.} \right) + (m^2)_j^i \phi_i \phi_j^*$$

Note that W and $\mathcal{L}_{\text{soft}}$ are to be expanded in all indices (in principle). This includes the gauge indices.

As such, a big part of Susyno's code is dedicated to Group Theory. It can compute various quantities for **any gauge group and field content**. Such functions are available for other applications.

To my knowledge, **LieART** is the only other Mathematica package with some (not all) of these functions.

Feger, Kephart 2012

SusyNO

Documentation

Published manual

Computer Physics
Communications 183 (2012) 2298
arXiv:1106.5016 [hep-ph]

Built-in
documentation

Up-to-date
More detailed
Easier to use

Susyno

Built-in documentation: easy to use and very handy

The screenshot shows the Wolfram Mathematica interface with the documentation for the `Invariants` function. The browser address bar shows `Susyno/ref/Invariants`. The page title is `Invariants`. Below the title, the function signature is `Invariants[group, {rep1, rep2, ...}]` with the description "Calculates the invariants of rep1 x rep2 x ...". There are sections for `MORE INFORMATION` and `EXAMPLES`. Under `EXAMPLES`, there is a sub-section `Basic Examples (1)`. The first example states: "If a and b are SU(2) doublets [= {1}], they form an invariant:" followed by the input `In[1]:= Invariants[SU2, {{1}, {1}}]` and the output `Out[1]= {a[2] b[1] - a[1] b[2]}`. The second example asks: "How to put together two SU(2) doublets (a and b) and a triplet c?" followed by the input `In[2]:= Invariants[SU2, {{1}, {1}, {2}}]` and the output `Out[2]= { (sqrt(2) a[2] b[2] c[1]) / 3^(1/4) - (a[2] b[1] c[2]) / 3^(1/4) - (a[1] b[2] c[2]) / 3^(1/4) + (sqrt(2) a[1] b[1] c[3]) / 3^(1/4) }`. The third example states: "The overall factor is of course irrelevant, therefore the following is also an invariant:" followed by the input `In[3]:= Expand[3^(1/4) %]`. The interface includes a search bar, navigation arrows, and a zoom level of 100%.

Input example: the MSSM

Output: see extra slides

Pick a name for the model (any)

```
author[MSSM] ^= "Me";  
date[MSSM] ^= "14:50, 25 July 2014";
```

Provide some optional data

```
group[MSSM] ^= {U1, SU2, SU3};
```

Specify the model's gauge group (any)

```
normalization = Sqrt[3/5];  
u = {-2/3 normalization, {0}, {0, 1}};  
d = {1/3 normalization, {0}, {0, 1}};  
Q = {1/6 normalization, {1}, {1, 0}};  
e = {normalization, {0}, {0, 0}};  
L = {-1/2 normalization, {1}, {0, 0}};  
Hu = {1/2 normalization, {1}, {0, 0}};  
Hd = {-1/2 normalization, {1}, {0, 0}};
```

Specify the model's representations (any)

```
reps[MSSM] ^= {u, d, Q, e, L, Hu, Hd};  
fieldNames[MSSM] ^= {"u", "d", "Q", "e", "L", "Hu", "Hd"};
```

```
nFlavs[MSSM] ^= {3, 3, 3, 3, 3, 1, 1};  
discreteSym[MSSM] ^= {-1, -1, -1, -1, -1, 1, 1};
```

Provide the number of flavors and the charges under any abelian symmetry

```
GenerateModel[MSSM, CalculateEverything → True]
```

Tell the program to calculate the model's RGEs (among other things)

Input example: minimal $SO(10)$

Output: see
extra slides

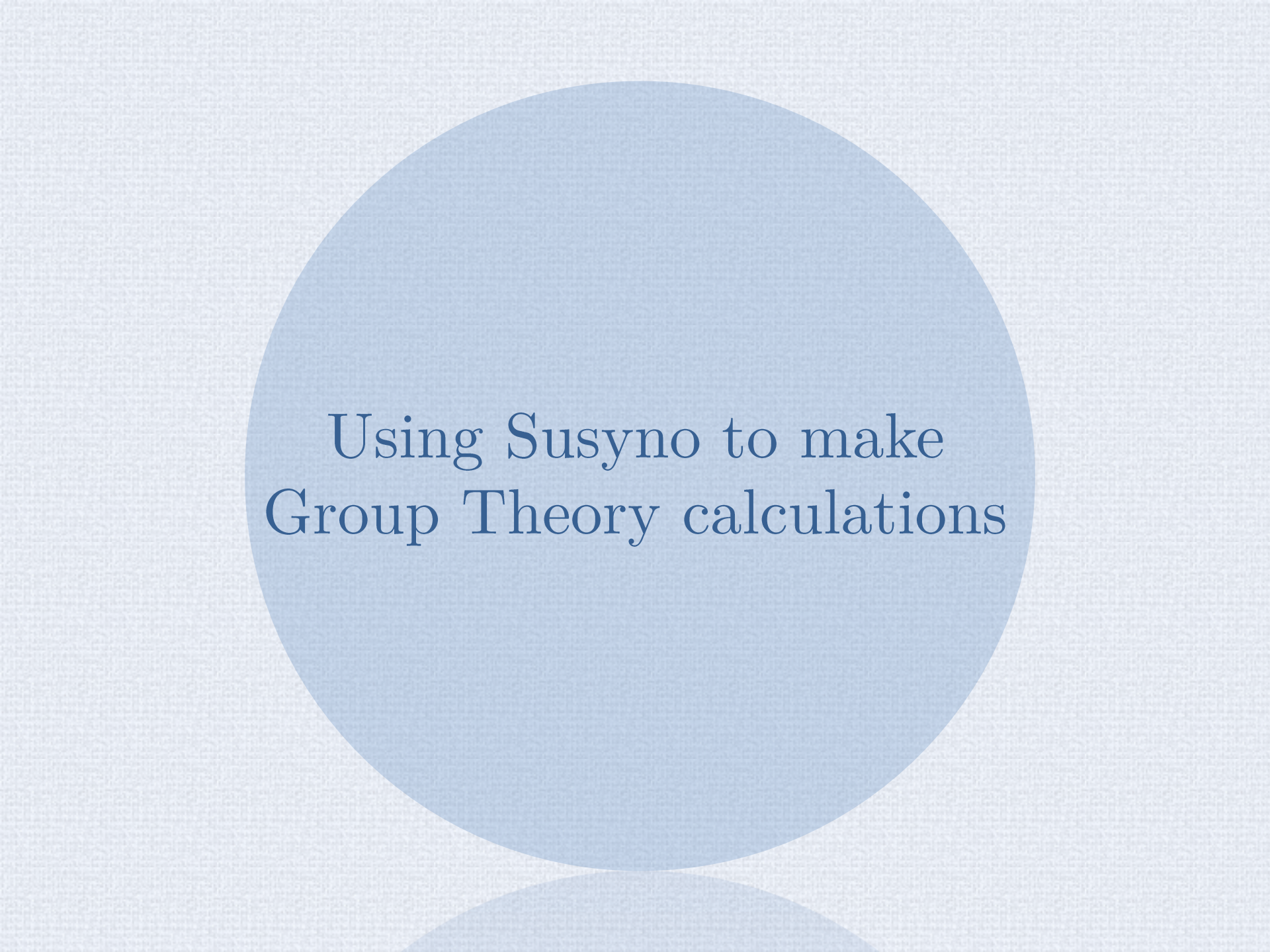
```
group[MinimalSUSYSO10GUT] ^= {SO10};

Ψ = {{0, 0, 0, 0, 1}}; (* 16-dim representation *)
Ξ = {{0, 0, 0, 1, 1}}; (* 210-dim representation *)
Δ = {{0, 0, 0, 0, 2}}; (* 126-dim representation *)
Δb = {{0, 0, 0, 2, 0}}; (* 126-dim representation (conj.) *)
H = {{1, 0, 0, 0, 0}}; (* 10-dim representation *)

reps[MinimalSUSYSO10GUT] ^= {Ψ, Ξ, Δ, Δb, H};
fieldNames[MinimalSUSYSO10GUT] ^= {"Ψ", "Ξ", "Δ", "Δ̄", "H"};

nFlavs[MinimalSUSYSO10GUT] ^= {3, 1, 1, 1, 1};
discreteSym[MinimalSUSYSO10GUT] ^= {1, 1, 1, 1, 1};

GenerateModel[MinimalSUSYSO10GUT, CalculateEverything → True]
```



Using Susyno to make
Group Theory calculations

Overview

Some useful functions

Casimir | ConjugateIrrep | DynkinIndex | DimR |
PermutationSymmetryOfInvariants | ReduceRepProduct |
RepName | RepsUpToDimN | TriangularAnomalyValue ...

Basis-independent
functions

RepMatrices | Invariants

Basis-dependent
functions

HookContentFormula | DecomposeSnProduct |
SnClassCharacter | SnClassOrder |
SnIrrepDim | SnIrrepGenerators

Permutation group
(S_n) functions

DecomposeReps | RegularSubgroupProjectionMatrix |
SubgroupEmbeddingCoefficients

Symmetry breaking
functions

NEW

Specifying a gauge group/representation

Group

Just type the group's name:

U1, SU2, SU3, ..., SO3, SO5, SO6, SO7, ...,
SP4, SP6, SP8, ..., G2, F4, E6, E7, E8

If the group contains more than one factor, use
lists: {U1,SU2,SU3}, {SU5,E8}, ...

Representations

Representations of simple groups have to be
given by their **Dynkin indices**, which are a list
of non-negative integers
(for U(1)'s: provide the hypercharge).

For example, the complete list of representations
of SU(3) is {0,0}, {0,1}, {1,0}, {1,1}, {0,2}, ...

Tables of representations in this notation are
available for example in Slansky 1981

But **Susyno itself can be used to identify
them...**

RepName, DimR, Casimir, DynkinIndex

Identify by name a representation *

Dimension of representation
 $d(R)$

$$\sum_a T_a^2 = C(R)\mathbb{1}$$

$$\text{Tr}(T_a T_b) = T(R)\delta_{ab}$$

Example 1

```
rep = {0, 0, 0, 1};
```

```
RepName[SU5, rep]  
DimR[SU5, rep]  
Casimir[SU5, rep]  
DynkinIndex[SU5, rep]
```

$\bar{5}$

5

$\frac{12}{5}$

$\frac{1}{2}$

Example 2

```
rep = {4, 0, 0, 1};
```

```
RepName[SU5, rep]  
DimR[SU5, rep]  
Casimir[SU5, rep]  
DynkinIndex[SU5, rep]
```

$\overline{315'}$

315

$\frac{88}{5}$

231

* Convention for assigning names to representations: RepName follows the scheme described in Feger, Kephart 2012 (tables in the literature — for example Slansky 1981 — are finite)

RepName, DimR, Casimir, DynkinIndex and RepsUpToDimN

CODE

```
reps = RepsUpToDimN[SO11, 5000];

Grid[
  Prepend[
    {#, RepName[SO11, #], Casimir[SO11, #],
      DynkinIndex[SO11, #]} & /@ reps,
    Style[#, {Darker[Red], Bold}] & /@
    {"Dynkin indices", "Name", "Casimir",
      "DynkinIndex"}]]
```

Dynkin indices	Name	Casimir	DynkinIndex
{0, 0, 0, 0, 0}	1	0	0
{1, 0, 0, 0, 0}	11	5	1
{0, 0, 0, 0, 1}	32	$\frac{55}{8}$	4
{0, 1, 0, 0, 0}	55	9	9
{2, 0, 0, 0, 0}	65	11	13
{0, 0, 1, 0, 0}	165	12	36
{3, 0, 0, 0, 0}	275	18	90
{1, 0, 0, 0, 1}	320	$\frac{99}{8}$	72
{0, 0, 0, 1, 0}	330	14	84
{1, 1, 0, 0, 0}	429	15	117
{0, 0, 0, 0, 2}	462	15	126
{4, 0, 0, 0, 0}	935	26	442
{0, 2, 0, 0, 0}	1144	20	416
{0, 1, 0, 0, 1}	1408	$\frac{135}{8}$	432
{1, 0, 1, 0, 0}	1430	18	468
{2, 0, 0, 0, 1}	1760	$\frac{151}{8}$	604
{2, 1, 0, 0, 0}	2025	22	810
{5, 0, 0, 0, 0}	2717	35	1729
{1, 0, 0, 1, 0}	3003	20	1092
{0, 0, 1, 0, 1}	3520	$\frac{163}{8}$	1304
{0, 0, 0, 0, 3}	4224	$\frac{195}{8}$	1872
{1, 0, 0, 0, 2}	4290	21	1638

OUTPUT

ReduceRepProduct

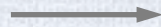
Computes the **decomposition** of a product of group representations **into irreducible parts**

Very useful function in model building

Example

$3 \times 3 \times \bar{6} = ?$ in $SU(3)$

```
RepName[SU3, {1, 0}]  
RepName[SU3, {0, 2}]  
3  
 $\bar{6}$ 
```



```
result = ReduceRepProduct[SU3, {{1, 0}, {1, 0}, {0, 2}}]  
{{{2, 2}, 1}, {{0, 3}, 1}, {{1, 1}, 2}, {{0, 0}, 1}}  
{RepName[SU3, #1], #2} & @@@ result  
{27, 1}, {10, 1}, {8, 2}, {1, 1}}
```

Output is a list of **representations** with **multiplicity**

So $3 \times 3 \times \bar{6} = 27 + \bar{10} + 8 + 8 + 1$

ReduceRepProduct

Code and algorithm
are very fast

$$5 \times 10 \times 15 \times \overline{40} \times 45 =? \text{ in } SU(5)$$

Snow 1990, 1993

```
result = ReduceRepProduct[SU5, {{1, 0, 0, 0}, {0, 1, 0, 0}, {2, 0, 0, 0}, {1, 1, 0, 0}, {0, 1, 0, 1}}];  
{RepName[SU5, #1], #2} &@@@ result
```

```
{  
  { $\overline{15750}$ , 1}, { $\overline{11880}$ , 2}, { $\overline{16170}$ , 8}, { $\overline{2400}$ , 8}, { $\overline{15360}$ , 2}, { $\overline{7425}$ , 6}, { $\overline{4620}$ , 10}, { $\overline{2625}$ , 19},  
  { $\overline{10240}$ , 9}, { $\overline{8750}$ , 16}, { $\overline{4410}$ , 21}, { $\overline{2430}$ , 41}, { $\overline{2475}$ , 2}, { $\overline{2625}$ , 21}, { $\overline{945}$ , 55}, { $\overline{1470}$ , 9},  
  { $\overline{1200}$ , 33}, { $\overline{3780}$ , 10}, { $\overline{2205}$ , 9}, { $\overline{2520}$ , 34}, { $\overline{1120}$ , 54}, { $\overline{420}$ , 10}, { $\overline{175}$ , 15}, { $\overline{450}$ , 14},  
  { $\overline{480}$ , 77}, { $\overline{280}$ , 12}, { $\overline{280}$ , 63}, { $\overline{70}$ , 33}, { $\overline{45}$ , 50}, { $\overline{4410}$ , 1}, { $\overline{1800}$ , 12}, { $\overline{980}$ , 24},  
  { $\overline{720}$ , 72}, { $\overline{105}$ , 54}, { $\overline{50}$ , 36}, { $\overline{2520}$ , 1}, { $\overline{560}$ , 11}, { $\overline{1540}$ , 2}, { $\overline{70}$ , 17}, { $\overline{5}$ , 13}}
```

Products of groups
can also be used

```
d = {1/3, {0}, {0, 1}};  
Q = {1/6, {1}, {1, 0}};  
L = {-1/2, {1}, {0, 0}};
```

```
result = ReduceRepProduct[{U1, SU2, SU3}, {d, Q, L}];  
{RepName[{U1, SU2, SU3}, #1], #2} &@@@ result  
{  
  { $0 \otimes 3 \otimes 8$ , 1}, { $0 \otimes 3 \otimes 1$ , 1}, { $0 \otimes 1 \otimes 8$ , 1}, { $0 \otimes 1 \otimes 1$ , 1}}
```

(contains a gauge singlet)

PermutationSymmetryOfInvariants

Does the **same as ReduceRepProduct** but also provides information on **how the irreducible parts transform under permutations** of the representations being multiplied

Leeuwen, Cohen, Lisser 1992

This is relevant only when there are repeated representations

Example

It is well known that **three triplets of SU(3)** form one invariant which is completely anti-symmetric

In other words, the **invariant is in the {1,1,1} irreducible representation of S_3**

```
PermutationSymmetryOfInvariants[SU3, {{1, 0}, {1, 0}, {1, 0}}]  
{{{1, 2, 3}}, {{{1, 1, 1}}, 1}}
```

Input reps #1, #2, #3 are the same

The given product contains **1 invariant**, in an **{1,1,1} irrep of S_3**

OUTPUT MEANING
(see documentation for details and more complex examples)

RepMatrices

RF 2013

This function builds explicitly the representation matrices of a given gauge group. To do so, a particular basis is chosen by the program

Example 1

Representation matrices of **3** of **SU(3)**

MatrixForm /@ RepMatrices[SU3, {1, 0}]

$$\left\{ \begin{pmatrix} 0 & \frac{1}{2} & 0 \\ \frac{1}{2} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \frac{1}{2} \\ 0 & \frac{1}{2} & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & \frac{i}{2} \\ 0 & 0 & 0 \\ -\frac{i}{2} & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -\frac{i}{2} & 0 \\ \frac{i}{2} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -\frac{i}{2} \\ 0 & \frac{i}{2} & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & -\frac{1}{2} \\ 0 & 0 & 0 \\ -\frac{1}{2} & 0 & 0 \end{pmatrix}, \begin{pmatrix} \frac{1}{\sqrt{3}} & 0 & 0 \\ 0 & -\frac{1}{2\sqrt{3}} & 0 \\ 0 & 0 & -\frac{1}{2\sqrt{3}} \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & \frac{1}{2} & 0 \\ 0 & 0 & -\frac{1}{2} \end{pmatrix} \right\}$$

Example 2

Representation matrices
of **5** of **SU(2)**

MatrixForm /@ RepMatrices[SU2, {4}]

$$\left\{ \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & \sqrt{\frac{3}{2}} & 0 & 0 \\ 0 & \sqrt{\frac{3}{2}} & 0 & \sqrt{\frac{3}{2}} & 0 \\ 0 & 0 & \sqrt{\frac{3}{2}} & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -i & 0 & 0 & 0 \\ i & 0 & -i\sqrt{\frac{3}{2}} & 0 & 0 \\ 0 & i\sqrt{\frac{3}{2}} & 0 & -i\sqrt{\frac{3}{2}} & 0 \\ 0 & 0 & i\sqrt{\frac{3}{2}} & 0 & -i \\ 0 & 0 & 0 & i & 0 \end{pmatrix}, \begin{pmatrix} 2 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & -2 \end{pmatrix} \right\}$$

RepMatrices

```
650 rep of E(6)  
Timing[matrices = RepMatrices[E6, {1, 0, 0, 0, 1, 0}];]  
{77.046875, Null}
```

Code is fast in
most cases

In general, what are the properties of these matrices?
They are **hermitian** and conform to the usual **trace condition** used in Particle Physics

```
And @@ Table[matrix == ConjugateTranspose[matrix], {matrix, matrices}]  
True
```



$$T_a^\dagger = T_a$$

```
DynkinIndex[E6, {1, 0, 0, 0, 1, 0}] IdentityMatrix[DimR[E6, Adjoint[E6]]] ==  
Table[Tr[matrix1.matrix2], {matrix1, matrices}, {matrix2, matrices}]  
True
```



$$\text{Tr}(T_a T_b) = T(R) \delta_{ab}$$

RepMatrices

Real representations: a word of caution

The basis used by **SusyNo** always keeps a maximum number of generators **in diagonal form** (these generators are always the last ones listed)

`MatrixForm /@ RepMatrices[SU2, {2}]`

$$\mathbf{3} \text{ of } \text{SU}(2) \longrightarrow \left\{ \begin{pmatrix} 0 & \frac{1}{\sqrt{2}} & 0 \\ \frac{1}{\sqrt{2}} & 0 & \frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & 0 \end{pmatrix}, \begin{pmatrix} 0 & -\frac{i}{\sqrt{2}} & 0 \\ \frac{i}{\sqrt{2}} & 0 & -\frac{i}{\sqrt{2}} \\ 0 & \frac{i}{\sqrt{2}} & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \right\}$$

This is a **perfectly fine basis**, as long as one is aware of it. In fact, this basis is the **best one to read off the quantum numbers** of the representation components

An important consequence of this is that the matrices of real representations are complex.
Adjust results if needed.

The matrices of a real representation in a real basis must all be anti-symmetric, therefore they cannot be diagonal

EASY TO
CHECK

Invariants

RF 2013

Computes the **Clebsch-Gordon coefficients** of a product of Lie group representations.

In other words, it calculates the **linear combinations of field components which are group invariant**.

CB coefficients are ubiquitous in model building

They are needed to write down a Lagrangian of a gauge theory

Transformation matrices?
As given by RepMatrices

Example

2×2 in $SU(2)$

```
field field
  a    b    ...
Invariants[SU2, {{1}, {1}}]
{a[2] b[1] - a[1] b[2]}
Invariants[SU2, {{1}, {1}}, TensorForm -> True][[1, 1]] // MatrixForm

$$\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

```

Invariants

~16s

27 rep of E(6) 27 rep of E(6) 351 rep of E(6)

Invariants[E6, {{1, 0, 0, 0, 0, 0}, {1, 0, 0, 0, 0, 0}, {0, 0, 0, 1, 0, 0}}]

$$\left\{ -\frac{3^{3/4} a[2] b[1] c[1]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[1] b[2] c[1]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[3] b[1] c[2]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[1] b[3] c[2]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[4] b[1] c[3]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[1] b[4] c[3]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[5] b[1] c[4]}{\sqrt{2} 13^{1/4}} + \dots 1851 \dots + \frac{3^{3/4} a[24] b[27] c[348]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[26] b[25] c[349]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[25] b[26] c[349]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[27] b[25] c[350]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[25] b[27] c[350]}{\sqrt{2} 13^{1/4}} - \frac{3^{3/4} a[27] b[26] c[351]}{\sqrt{2} 13^{1/4}} + \frac{3^{3/4} a[26] b[27] c[351]}{\sqrt{2} 13^{1/4}} \right\}$$

large output show less show more show all set size limit...

IMPORTANT QUESTION
How are the Clebsch-Gordon coefficients normalized?

ANSWER: for $\sum_{i_1, i_2, i_3, \dots, i_n} c_{i_1, i_2, \dots, i_n} \Phi_{i_1}^{(1)} \Phi_{i_2}^{(2)} \dots \Phi_{i_n}^{(n)}$,
 $\sum_{i_1, i_2, i_3, \dots, i_n} |c_{i_1, i_2, \dots, i_n}|^2 = \sqrt{\dim(\Phi^{(1)}) \dim(\Phi^{(2)}) \dots \dim(\Phi^{(n)})}$

ALWAYS

E.g.: if $a[1]b[2]-a[2]b[1]$ is an invariant combination of two SU(2) doublets, so is any multiple of this

Two SU(2) doublets form the invariants $x(a[1]b[2]-a[2]b[1])$ for any x . **Which x does Susyno take?**

$$x^2 + (-x)^2 = \sqrt{2 \times 2} \Rightarrow |x| = 1$$

DecomposeRep

Decomposes some (irreducible) representation of a group G into irreducible components of some subgroup of G . In other words, it **calculates branching rules**

Example 1

What are the $SU(3) \times SU(2) \times U(1)$ representations in the **70** of $SU(5)$?

```
group = {SU5}; subgroup = {SU3, SU2, U1}; representation = {{2, 0, 0, 1}}; (* 70 of {SU(5)} *)
```

```
prjMatrix =  $\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ -2 & -4 & -6 & -3 \end{pmatrix}$ ; (* Use RegularSubgroupProjectionMatrix to find it *)
```

```
DecomposeRep[group, representation, subgroup, prjMatrix, UserName -> True]
```

```
{15 $\otimes$ 1 $\otimes$ -2,  $\bar{6}$  $\otimes$ 2 $\otimes$ -7, 8 $\otimes$ 2 $\otimes$ 3, 3 $\otimes$ 3 $\otimes$ -2, 3 $\otimes$ 1 $\otimes$ -2,  $\bar{3}$  $\otimes$ 3 $\otimes$ 8, 1 $\otimes$ 4 $\otimes$ 3, 1 $\otimes$ 2 $\otimes$ 3}
```


DecomposeRep

Example 2

List the $SO(10) \rightarrow SU(5) \times U(1)$ branching rules for all $SO(10)$ representations up to size 700

```
SO10reps = RepsUpToDimN[SO10, 700];  
  
SOtoSU5U1ProjMatrix =  $\begin{pmatrix} 1 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 1 & 1 & 0 & 0 \\ 2 & 0 & 2 & 1 & -1 \end{pmatrix}$ ; (* Use RegularSubgroupProjectionMatrix *)  
  
data =  
  Table[{Style[RepName[SO10, rep], Darker[Red]],  
    DecomposeRep[{SO10}, {rep}, {SU5, U1}, SOtoSU5U1ProjMatrix, UserName → True]}, {rep, SO10reps}];  
Grid[Prepend[data, Style[#, Bold] & /@ {"SO(10) rep", "SU(5) x U(1) content"}], Frame → All,  
  FrameStyle → LightGray]
```

CODE

DecomposeRep

OUTPUT

SO(10) rep	SU(5) x U(1) content
1	{1 \otimes 0}
10	{5 \otimes 2, $\bar{5}$ \otimes -2}
16	{10 \otimes -1, $\bar{5}$ \otimes 3, 1 \otimes -5}
$\bar{16}$	{5 \otimes -3, $\bar{10}$ \otimes 1, 1 \otimes 5}
45	{24 \otimes 0, 10 \otimes 4, $\bar{10}$ \otimes -4, 1 \otimes 0}
54	{15 \otimes 4, 24 \otimes 0, $\bar{15}$ \otimes -4}
120	{ $\bar{45}$ \otimes -2, 45 \otimes 2, 5 \otimes 2, 10 \otimes -6, $\bar{10}$ \otimes 6, $\bar{5}$ \otimes -2}
$\bar{126}$	{15 \otimes -6, $\bar{45}$ \otimes -2, 50 \otimes 2, 5 \otimes 2, $\bar{10}$ \otimes 6, 1 \otimes 10}
126	{ $\bar{50}$ \otimes -2, 45 \otimes 2, 10 \otimes -6, $\bar{15}$ \otimes 6, $\bar{5}$ \otimes -2, 1 \otimes -10}
144	{15 \otimes -1, $\bar{45}$ \otimes 3, 24 \otimes -5, 5 \otimes 7, 40 \otimes -1, 10 \otimes -1, $\bar{5}$ \otimes 3}
$\bar{144}$	{ $\bar{40}$ \otimes 1, 24 \otimes 5, 45 \otimes -3, 5 \otimes -3, $\bar{15}$ \otimes 1, $\bar{10}$ \otimes 1, $\bar{5}$ \otimes -7}
210	{ $\bar{40}$ \otimes -4, 75 \otimes 0, 24 \otimes 0, 5 \otimes -8, 40 \otimes 4, 10 \otimes 4, $\bar{10}$ \otimes -4, $\bar{5}$ \otimes 8, 1 \otimes 0}
210'	{ $\bar{35}$ \otimes 6, 70 \otimes 2, $\bar{70}$ \otimes -2, 35 \otimes -6}
320	{70 \otimes 2, $\bar{40}$ \otimes 6, $\bar{70}$ \otimes -2, $\bar{45}$ \otimes -2, 45 \otimes 2, 5 \otimes 2, 40 \otimes -6, $\bar{5}$ \otimes -2}
560	{175 \otimes -1, $\bar{50}$ \otimes 3, $\bar{70}$ \otimes 3, $\bar{45}$ \otimes 3, 75 \otimes -5, 24 \otimes -5, 45 \otimes 7, 40 \otimes -1, 10 \otimes -1, 10 \otimes -1, $\bar{10}$ \otimes -9, $\bar{5}$ \otimes 3, 1 \otimes -5}
$\bar{560}$	{70 \otimes -3, $\bar{175}$ \otimes 1, $\bar{40}$ \otimes 1, $\bar{45}$ \otimes -7, 75 \otimes 5, 24 \otimes 5, 50 \otimes -3, 45 \otimes -3, 5 \otimes -3, 10 \otimes 9, $\bar{10}$ \otimes 1, $\bar{10}$ \otimes 1, 1 \otimes 5}
660	{ $\bar{70}$ \otimes 8, 160 \otimes 4, 200 \otimes 0, $\bar{160}$ \otimes -4, 70' \otimes -8}
672	{ $\bar{35}$ \otimes -9, 126 \otimes -5, 210 \otimes -1, 15 \otimes -1, $\bar{175}$ \otimes 3, $\bar{45}$ \otimes 3, 50 \otimes 7, 5 \otimes 7, $\bar{10}$ \otimes 11, 1 \otimes 15}
$\bar{672}$	{175'' \otimes -3, $\bar{210}$ \otimes 1, $\bar{50}$ \otimes -7, $\bar{126}$ \otimes 5, 45 \otimes -3, 35 \otimes 9, 10 \otimes -11, $\bar{15}$ \otimes 1, $\bar{5}$ \otimes -7, 1 \otimes -15}

Takes around 5 minutes to extend this table up to all reps of size 10000 or smaller

Subgroup Embedding Coefficients

Calculates the **relations between the subgroup invariants** (i.e., Clebsch-Gordon coefficients) in a theory which is symmetric under a bigger group

Easier to explain with an example!

Consider the $SO(10)$ invariant combination of the representations $\mathbf{16} \times \mathbf{16} \times \mathbf{10}$. We can see it from a $SU(3) \times SU(2) \times U(1)$ perspective:

$\mathbf{16}$	\times	$\mathbf{16}$	\times	$\mathbf{10}$	$=$	1 $SO(10)$ invariant
$\begin{matrix} 3 \otimes 2 \otimes \frac{1}{6} \\ \bar{3} \otimes 1 \otimes -\frac{1}{2} \\ \bar{3} \otimes 1 \otimes \frac{1}{3} \\ 1 \otimes 2 \otimes -\frac{1}{2} \\ 1 \otimes 1 \otimes 0 \\ 1 \otimes 1 \otimes 1 \end{matrix}$	\times	$\begin{matrix} 3 \otimes 2 \otimes \frac{1}{6} \\ \bar{3} \otimes 1 \otimes -\frac{1}{2} \\ \bar{3} \otimes 1 \otimes \frac{1}{3} \\ 1 \otimes 2 \otimes -\frac{1}{2} \\ 1 \otimes 1 \otimes 0 \\ 1 \otimes 1 \otimes 1 \end{matrix}$	\times	$\begin{matrix} 3 \otimes 1 \otimes -\frac{1}{3} \\ \bar{3} \otimes 1 \otimes \frac{1}{3} \\ 1 \otimes 2 \otimes \frac{1}{2} \\ 1 \otimes 2 \otimes -\frac{1}{2} \end{matrix}$	$=$	Linear combination of the 17 subgroup invariants

SubgroupEmbeddingCoefficients

CODE

```
(* INPUT DATA *)
group = {SO10};
rep16 = {{0, 0, 0, 0, 1}};
rep10 = {{1, 0, 0, 0, 0}};
subgroup = {SU3, SU2, U1};
breakInfo = {{1, {2, 1}}, {1, {4}}, {1, 1/6}};

(* CALCULATE THE EMBEDDING COEFFICIENTS *)
result = SubgroupEmbeddingCoefficients[group, {rep16, rep16, rep10}, subgroup, breakInfo];

(* THE REST OF THE CODE BELOW IS JUST TO FORMAT THE OUTPUT IN A NICE WAY *)
coefficients = result[[2, 5, 1]];

productFields = Map[RepName[subgroup, #] &,
  Flatten[ConstantArray[Extract[result[[2, 3]], #[[1]]], #[[2]]] &/@ result[[2, 4, All, {2, 3}]], 1], {2}];

Print["The product of ",
  {Style[RepName[group, #1], {Darker[Blue], Bold}}, Style[RepName[group, #2], {Darker[Green], Bold}},
  Style[RepName[group, #3], {Darker[Red], Bold}]} &@@ {rep16, rep16, rep10}, " of SO(10) is the same as"];
table =
  Table[Row[{If[coefficients[[i]] > 0, Style["+", Darker[Gray]], ""], Style[coefficients[[i]], Darker[Gray]],
    {Style[#1, {Darker[Blue], Bold}}, Style[#2, {Darker[Green], Bold}], Style[#3, {Darker[Red], Bold}]} &@@
    productFields[[i]]}], {i, Length[coefficients]};
Print[Row[table, " "]];
Print["in SU(3) x SU(2) x U(1)"];

```


Subgroup Embedding Coefficients

OUTPUT

The product of $\{16, 16, 10\}$ of $SO(10)$ is the same as

$$\begin{aligned}
 & -\frac{2^{3/4}}{15^{1/4}} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, 3 \otimes 2 \otimes \frac{1}{6}, 3 \otimes 1 \otimes -\frac{1}{3} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, \bar{3} \otimes 1 \otimes -\frac{2}{3}, 1 \otimes 2 \otimes \frac{1}{2} \right\} \\
 & + \left(\frac{2}{5} \right)^{1/4} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, \bar{3} \otimes 1 \otimes \frac{1}{3}, 1 \otimes 2 \otimes -\frac{1}{2} \right\} - \left(\frac{2}{5} \right)^{1/4} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, 1 \otimes 2 \otimes -\frac{1}{2}, \bar{3} \otimes 1 \otimes \frac{1}{3} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ \bar{3} \otimes 1 \otimes -\frac{2}{3}, 3 \otimes 2 \otimes \frac{1}{6}, 1 \otimes 2 \otimes \frac{1}{2} \right\} \\
 & + \frac{2^{3/4}}{15^{1/4}} \left\{ \bar{3} \otimes 1 \otimes -\frac{2}{3}, \bar{3} \otimes 1 \otimes \frac{1}{3}, \bar{3} \otimes 1 \otimes \frac{1}{3} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ \bar{3} \otimes 1 \otimes -\frac{2}{3}, 1 \otimes 1 \otimes 1, 3 \otimes 1 \otimes -\frac{1}{3} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ \bar{3} \otimes 1 \otimes \frac{1}{3}, 3 \otimes 2 \otimes \frac{1}{6}, 1 \otimes 2 \otimes -\frac{1}{2} \right\} \\
 & + \frac{2^{3/4}}{15^{1/4}} \left\{ \bar{3} \otimes 1 \otimes \frac{1}{3}, \bar{3} \otimes 1 \otimes -\frac{2}{3}, \bar{3} \otimes 1 \otimes \frac{1}{3} \right\} - \left(\frac{2}{5} \right)^{1/4} \left\{ \bar{3} \otimes 1 \otimes \frac{1}{3}, 1 \otimes 1 \otimes 0, 3 \otimes 1 \otimes -\frac{1}{3} \right\} - \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 2 \otimes -\frac{1}{2}, 3 \otimes 2 \otimes \frac{1}{6}, \bar{3} \otimes 1 \otimes \frac{1}{3} \right\} \\
 & + \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 1 \otimes 0, 1 \otimes 2 \otimes \frac{1}{2} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 1 \otimes 1, 1 \otimes 2 \otimes -\frac{1}{2} \right\} - \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 1 \otimes 0, \bar{3} \otimes 1 \otimes \frac{1}{3}, 3 \otimes 1 \otimes -\frac{1}{3} \right\} \\
 & + \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 1 \otimes 0, 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 2 \otimes \frac{1}{2} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 1 \otimes 1, \bar{3} \otimes 1 \otimes -\frac{2}{3}, 3 \otimes 1 \otimes -\frac{1}{3} \right\} + \left(\frac{2}{5} \right)^{1/4} \left\{ 1 \otimes 1 \otimes 1, 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 2 \otimes -\frac{1}{2} \right\}
 \end{aligned}$$

in $SU(3) \times SU(2) \times U(1)$

Why the strange factors?

The answer is simple: **it is due to the program's default normalization of the invariants** (i.e., Clebsch-Gordon coefficients)

Both the $SO(10)$ Clebsch-Gordon coefficients and the subgroup ones

Here, we can opt to quickly **eliminate the $SO(10)$ Clebsch-Gordon's normalization issue** by just looking at **ratios of these embedding coefficients**

On the other hand, most of the $SU(3) \times SU(2) \times U(1)$ invariants shown here are **normalized as usually expected**

The only ones which are not are the ones involving $3 \times 3 \times 3$ of $SU(3)$

Subgroup Embedding Coefficients

$$\begin{array}{c}
 y_u \\
 y_d \\
 + \left(\frac{2}{5}\right)^{1/4} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, \bar{3} \otimes 1 \otimes -\frac{2}{3}, 1 \otimes 2 \otimes \frac{1}{2} \right\} \\
 + \left(\frac{2}{5}\right)^{1/4} \left\{ 3 \otimes 2 \otimes \frac{1}{6}, \bar{3} \otimes 1 \otimes \frac{1}{3}, 1 \otimes 2 \otimes -\frac{1}{2} \right\}
 \end{array}$$

The 4 coefficients are the same

$$\begin{array}{c}
 y_\nu \\
 y_\ell \\
 + \left(\frac{2}{5}\right)^{1/4} \left\{ 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 1 \otimes 0, 1 \otimes 2 \otimes \frac{1}{2} \right\} + \left(\frac{2}{5}\right)^{1/4} \left\{ 1 \otimes 2 \otimes -\frac{1}{2}, 1 \otimes 1 \otimes 1, 1 \otimes 2 \otimes -\frac{1}{2} \right\}
 \end{array}$$

Conclusion

$$SO(10) \Rightarrow y_u = y_d = y_\ell = y_\nu$$

Now, if we just change in the input $\{1,0,0,0,0\}$ (the **10**) to $\{0,0,0,2,0\}$ (the **126**) ...

With minimal code change

$$SO(10) \Rightarrow y_u = y_d = -\frac{1}{3}y_\ell = -\frac{1}{3}y_\nu$$

Georgi-Jarlskog relation
1979



Summary

Summary

The Susyno program can calculate

RGEs of a
SUSY model

Generate list of model
parameters, check anomalies, ...

Group theory
quantities

Useful for model
building in general

1

Download from
<http://web.ist.utl.pt/renato.fonseca/susyno.html>

2

Unpack the folder `Susyno` and drop it inside
(`Mathematica base directory`)/`AddOns/Applications`

3

Type in Mathematica's front the following: `<<Susyno``
The built-in documentation becomes readily available

Quick
start

Thank you



Extra slides

Input example: the MSSM

Pick a name for the model (any)

```
author[MSSM] ^= "Me";  
date[MSSM] ^= "14:50, 25 July 2014";
```

Provide some optional data

```
group[MSSM] ^= {U1, SU2, SU3};
```

Specify the model's gauge group (any)

```
normalization = Sqrt[3/5];  
u = {-2/3 normalization, {0}, {0, 1}};  
d = {1/3 normalization, {0}, {0, 1}};  
Q = {1/6 normalization, {1}, {1, 0}};  
e = {normalization, {0}, {0, 0}};  
L = {-1/2 normalization, {1}, {0, 0}};  
Hu = {1/2 normalization, {1}, {0, 0}};  
Hd = {-1/2 normalization, {1}, {0, 0}};
```

Specify the model's representations
(any)

```
reps[MSSM] ^= {u, d, Q, e, L, Hu, Hd};  
fieldNames[MSSM] ^= {"u", "d", "Q", "e", "L", "Hu", "Hd"};
```

```
nFlavs[MSSM] ^= {3, 3, 3, 3, 3, 1, 1};  
discreteSym[MSSM] ^= {-1, -1, -1, -1, -1, 1, 1};
```

Provide the number of flavors and the
charges under any abelian symmetry

```
GenerateModel[MSSM, CalculateEverything → True]
```

Tell the program to calculate the
model's RGEs (among other things)

Output example: the MSSM

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

^ Model name

MSSM

^ Author

Me

^ Date

14:50, 25 July 2014

Output example: the MSSM

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

```
U1 x SU2 x SU3
```

```
GOOD NEWS: The model is gauge anomaly free.
```

```
>>> Extra information
```

```
This data is contained in the group[MSSM] variable.
```

Output example: the MSSM

Model Information	Gauge group		Representations				Parameters in model	Lagrangian	BetaFunctions
	u	d	Q	e	L	Hu	Hd		
U1	$-\frac{2}{\sqrt{15}}$	$\frac{1}{\sqrt{15}}$	$\frac{1}{2\sqrt{15}}$	$\sqrt{\frac{3}{5}}$	$-\sqrt{\frac{3}{5}}$	$\frac{\sqrt{3}}{2}$	$-\frac{\sqrt{3}}{2}$		
SU2	{0}	{0}	{1}	{0}	{1}	{1}	{1}		
SU3	{0, 1}	{0, 1}	{1, 0}	{0, 0}	{0, 0}	{0, 0}	{0, 0}		
#Flavors	3	3	3	3	3	1	1		
R-Charges	-1	-1	-1	-1	-1	1	1		

>>> Extra information
 This data is contained in the `reps[MSSM]` variable.

Output example: the MSSM

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

^ Gauge coupling constants

$g[1]$

$g[2]$

$g[3]$

^ Gaugino masses

$M[1]$

$M[2]$

$M[3]$

^ Superpotential trilinear parameters

$y[\{u, Q, H_u\}, \{f[1], f[2]\}]$

$y[\{d, Q, H_d\}, \{f[1], f[2]\}]$

$y[\{e, L, H_d\}, \{f[1], f[2]\}]$

^ Superpotential bilinear parameters

$\mu[\{H_u, H_d\}]$

^ Superpotential linear parameters

^ Soft trilinear parameters

$h[\{u, Q, H_u\}, \{f[1], f[2]\}]$

$h[\{d, Q, H_d\}, \{f[1], f[2]\}]$

$h[\{e, L, H_d\}, \{f[1], f[2]\}]$

^ Soft bilinear parameters

$b[\{H_u, H_d\}]$

^ Soft linear parameters

^ Soft masses

$m2[\{u, u\}, \{f[1], f[2]\}]$

$m2[\{d, d\}, \{f[1], f[2]\}]$

$m2[\{Q, Q\}, \{f[1], f[2]\}]$

$m2[\{e, e\}, \{f[1], f[2]\}]$

$m2[\{L, L\}, \{f[1], f[2]\}]$

$m2[\{H_u, H_u\}]$

$m2[\{H_d, H_d\}]$

Output example: the MSSM

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

^ Superpotential (trilinear terms)

```
y[{{e, L, Hd}}, {f[1], f[2]}] (e[f[1]] Hd[2] L[f[2]] [1] - e[f[1]] Hd[1] L[f[2]] [2]) +
y[{{d, Q, Hd}}, {f[1], f[2]}] (Hd[2] d[f[1]] [1] Q[f[2]] [1, 1] + Hd[2] d[f[1]] [2] Q[f[2]] [1, 2] +
Hd[2] d[f[1]] [3] Q[f[2]] [1, 3] - Hd[1] d[f[1]] [1] Q[f[2]] [2, 1] - Hd[1] d[f[1]] [2] Q[f[2]] [2, 2] - Hd[1] d[f[1]] [3] Q[f[2]] [2, 3]) +
y[{{u, Q, Hu}}, {f[1], f[2]}] (Hu[2] Q[f[2]] [1, 1] u[f[1]] [1] - Hu[1] Q[f[2]] [2, 1] u[f[1]] [1] + Hu[2] Q[f[2]] [1, 2] u[f[1]] [2] -
Hu[1] Q[f[2]] [2, 2] u[f[1]] [2] + Hu[2] Q[f[2]] [1, 3] u[f[1]] [3] - Hu[1] Q[f[2]] [2, 3] u[f[1]] [3])
```

^ Superpotential (bilinear terms)

```
mu[{{Hu, Hd}}] (Hd[2] Hu[1] - Hd[1] Hu[2])
```

^ Superpotential (linear terms)

0

^ Soft SUSY breaking Lagrangian (trilinear terms)

```
h[{{e, L, Hd}}, {f[1], f[2]}] (e[f[1]] Hd[2] L[f[2]] [1] - e[f[1]] Hd[1] L[f[2]] [2]) +
h[{{d, Q, Hd}}, {f[1], f[2]}] (Hd[2] d[f[1]] [1] Q[f[2]] [1, 1] + Hd[2] d[f[1]] [2] Q[f[2]] [1, 2] +
Hd[2] d[f[1]] [3] Q[f[2]] [1, 3] - Hd[1] d[f[1]] [1] Q[f[2]] [2, 1] - Hd[1] d[f[1]] [2] Q[f[2]] [2, 2] - Hd[1] d[f[1]] [3] Q[f[2]] [2, 3]) +
h[{{u, Q, Hu}}, {f[1], f[2]}] (Hu[2] Q[f[2]] [1, 1] u[f[1]] [1] - Hu[1] Q[f[2]] [2, 1] u[f[1]] [1] + Hu[2] Q[f[2]] [1, 2] u[f[1]] [2] -
Hu[1] Q[f[2]] [2, 2] u[f[1]] [2] + Hu[2] Q[f[2]] [1, 3] u[f[1]] [3] - Hu[1] Q[f[2]] [2, 3] u[f[1]] [3])
```

^ Soft SUSY breaking Lagrangian (bilinear terms)

```
b[{{Hu, Hd}}] (Hd[2] Hu[1] - Hd[1] Hu[2])
```

^ Soft SUSY breaking Lagrangian (linear terms)

0

^ Soft SUSY breaking Lagrangian (mass terms)

```
m2[{{e, e}}, {f[1], f[2]}] Conjugate[e][f[2]] e[f[1]] +
m2[{{Hd, Hd}}] (Conjugate[Hd][1] Hd[1] + Conjugate[Hd][2] Hd[2]) + m2[{{Hu, Hu}}] (Conjugate[Hu][1] Hu[1] + Conjugate[Hu][2] Hu[2]) +
m2[{{d, d}}, {f[1], f[2]}] (Conjugate[d][f[2]] [1] d[f[1]] [1] + Conjugate[d][f[2]] [2] d[f[1]] [2] + Conjugate[d][f[2]] [3] d[f[1]] [3]) +
m2[{{L, L}}, {f[1], f[2]}] (Conjugate[L][f[2]] [1] L[f[1]] [1] + Conjugate[L][f[2]] [2] L[f[1]] [2]) +
m2[{{Q, Q}}, {f[1], f[2]}] (Conjugate[Q][f[2]] [1, 1] Q[f[1]] [1, 1] - Conjugate[Q][f[2]] [1, 2] Q[f[1]] [1, 2] + Conjugate[Q][f[2]] [1, 3] Q[f[1]] [1, 3] +
Conjugate[Q][f[2]] [2, 1] Q[f[1]] [2, 1] - Conjugate[Q][f[2]] [2, 2] Q[f[1]] [2, 2] + Conjugate[Q][f[2]] [2, 3] Q[f[1]] [2, 3]) +
m2[{{u, u}}, {f[1], f[2]}] (Conjugate[u][f[2]] [1] u[f[1]] [1] + Conjugate[u][f[2]] [2] u[f[1]] [2] + Conjugate[u][f[2]] [3] u[f[1]] [3])
```

>>> Parameters are shown in dark orange; field heads are shown in blue.

Output example: the MSSM

Model Information	Gauge group	Representations	Parameters in model	Lagrangian	BetaFunctions
<ul style="list-style-type: none"> ▼ g[1] ▼ g[2] ▲ g[3] <ul style="list-style-type: none"> ▲ ~~~~~ β⁽¹⁾ ~~~~~ -3 g[3]³ ▲ ~~~~~ β⁽²⁾ ~~~~~ $\frac{11}{5} g[1]^2 g[3]^3 + 9 g[2]^2 g[3]^3 + 14 g[3]^5 - 4 \text{Conjugate}[y[(d, Q, Hd), \{f[1], f[3]\}]] g[3]^3 y[(d, Q, Hd), \{f[1], f[3]\}] - 4 \text{Conjugate}[y[(u, Q, Hu), \{f[1], f[3]\}]] g[3]^3 y[(u, Q, Hu), \{f[1], f[3]\}]$ ▼ M[1] ▼ M[2] ▼ M[3] ▼ y[{u, Q, Hu}, {f[1], f[2]}] ▼ y[{d, Q, Hd}, {f[1], f[2]}] ▼ y[{e, L, Hd}, {f[1], f[2]}] ▼ μ[{Hu, Hd}] ▼ h[{u, Q, Hu}, {f[1], f[2]}] ▼ h[{d, Q, Hd}, {f[1], f[2]}] ▼ h[{e, L, Hd}, {f[1], f[2]}] ▼ b[{Hu, Hd}] ▼ m2[{u, u}, {f[1], f[2]}] ▼ m2[{d, d}, {f[1], f[2]}] ▼ m2[{Q, Q}, {f[1], f[2]}] ▼ m2[{e, e}, {f[1], f[2]}] ▼ m2[{L, L}, {f[1], f[2]}] ▼ m2[{Hu, Hu}] ▼ m2[{Hd, Hd}] 					
>>> Extra information					

Input example: minimal $SO(10)$

```
group[MinimalSUSYSO10GUT] ^= {SO10};

Ψ = {{0, 0, 0, 0, 1}};    (* 16-dim representation *)
Ξ = {{0, 0, 0, 1, 1}};    (* 210-dim representation *)
Δ = {{0, 0, 0, 0, 2}};    (* 126-dim representation *)
Δb = {{0, 0, 0, 2, 0}};   (* 126-dim representation (conj.) *)
H = {{1, 0, 0, 0, 0}};    (* 10-dim representation *)

reps[MinimalSUSYSO10GUT] ^= {Ψ, Ξ, Δ, Δb, H};
fieldNames[MinimalSUSYSO10GUT] ^= {"Ψ", "Ξ", "Δ", "Δ̄", "H"};

nFlavs[MinimalSUSYSO10GUT] ^= {3, 1, 1, 1, 1};
discreteSym[MinimalSUSYSO10GUT] ^= {1, 1, 1, 1, 1};

GenerateModel[MinimalSUSYSO10GUT, CalculateEverything → True]
```

Output example: minimal SO(10)

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

^ Model name

MinimalSUSYSO10GUT

Output example: minimal SO(10)

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

SO10

GOOD NEWS: The model is gauge anomaly free.

>>> Extra information

This data is contained in the `group[MinimalSUSYSO10GUT]` variable.

Output example: minimal SO(10)

Model Information	Gauge group		Representations		Parameters in model	Lagrangian	BetaFunctions
	Ψ	$\bar{\Psi}$	Δ	$\bar{\Delta}$	H		
S010	{0, 0, 0, 0, 1}	{0, 0, 0, 1, 1}	{0, 0, 0, 0, 2}	{0, 0, 0, 2, 0}	{1, 0, 0, 0, 0}		
#Flavours	3	1	1	1	1		
R-Charges	1	1	1	1	1		

>>> Extra information
 This data is contained in the `reps[MinimalSUSYSO10GUT]` variable.

Output example: minimal SO(10)

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

^ Gauge coupling constants

$g[1]$

^ Gaugino masses

$M[1]$

^ Superpotential trilinear parameters

$y[\{\Psi, \Psi, \bar{\Delta}\}, \{f[1], f[2]\}]$ (symmetric under a permutation of the flavor indices $\{f[1], f[2]\}$)

$y[\{\Psi, \Psi, H\}, \{f[1], f[2]\}]$ (symmetric under a permutation of the flavor indices $\{f[1], f[2]\}$)

$y[\{\Phi, \Phi, \Phi\}]$

$y[\{\Phi, \Delta, \bar{\Delta}\}]$

$y[\{\Phi, \Delta, H\}]$

$y[\{\Phi, \bar{\Delta}, H\}]$

^ Superpotential bilinear parameters

$\mu[\{\Phi, \Phi\}]$

$\mu[\{\Delta, \bar{\Delta}\}]$

$\mu[\{H, H\}]$

^ Superpotential linear parameters

^ Soft trilinear parameters

$h[\{\Psi, \Psi, \bar{\Delta}\}, \{f[1], f[2]\}]$ (symmetric under a permutation of the flavor indices $\{f[1], f[2]\}$)

$h[\{\Psi, \Psi, H\}, \{f[1], f[2]\}]$ (symmetric under a permutation of the flavor indices $\{f[1], f[2]\}$)

$h[\{\Phi, \Phi, \Phi\}]$

$h[\{\Phi, \Delta, \bar{\Delta}\}]$

$h[\{\Phi, \Delta, H\}]$

$h[\{\Phi, \bar{\Delta}, H\}]$

^ Soft bilinear parameters

$b[\{\Phi, \Phi\}]$

$b[\{\Delta, \bar{\Delta}\}]$

$b[\{H, H\}]$

^ Soft linear parameters

^ Soft masses

...

Output example: minimal SO(10)

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

Superpotential (trilinear terms)

$y[\{\bar{\Phi}, \bar{\Delta}, H\}]$

$$\begin{aligned}
 & \left(\frac{H[10] \bar{\Delta}[91] \Phi[1]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[116] \Phi[1]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[123] \Phi[1]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[125] \Phi[1]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[6] \bar{\Delta}[126] \Phi[1] - \frac{H[10] \bar{\Delta}[90] \Phi[2]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[115] \Phi[2]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[122] \Phi[2]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[7] \bar{\Delta}[124] \Phi[2] - \frac{H[6] \bar{\Delta}[125] \Phi[2]}{e^{1/4}} - \right. \\
 & \frac{H[10] \bar{\Delta}[89] \Phi[3]}{e^{1/4}} + \frac{H[9] \bar{\Delta}[114] \Phi[3]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[121] \Phi[3]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[125] \Phi[3]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[4] \bar{\Delta}[126] \Phi[3] + \frac{H[10] \bar{\Delta}[88] \Phi[4]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[113] \Phi[4]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[120] \Phi[4]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[5] \bar{\Delta}[124] \Phi[4] - \frac{H[4] \bar{\Delta}[125] \Phi[4]}{e^{1/4}} + \\
 & \frac{H[10] \bar{\Delta}[87] \Phi[5]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[112] \Phi[5]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[8] \bar{\Delta}[119] \Phi[5] - \frac{H[7] \bar{\Delta}[122] \Phi[5]}{e^{1/4}} + \frac{H[6] \bar{\Delta}[123] \Phi[5]}{e^{1/4}} + \frac{H[10] \bar{\Delta}[86] \Phi[6]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[111] \Phi[6]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[121] \Phi[6]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[123] \Phi[6]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[3] \bar{\Delta}[126] \Phi[6] - \\
 & \frac{H[10] \bar{\Delta}[83] \Phi[7]}{2 \cdot e^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[84] \Phi[7] + \frac{H[9] \bar{\Delta}[108] \Phi[7]}{2 \cdot e^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[9] \bar{\Delta}[109] \Phi[7] - \frac{H[8] \bar{\Delta}[118] \Phi[7]}{2^{3/4} 3^{1/4}} - \frac{H[7] \bar{\Delta}[120] \Phi[7]}{2^{3/4} 3^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[5] \bar{\Delta}[122] \Phi[7] - \frac{H[4] \bar{\Delta}[123] \Phi[7]}{2^{3/4} 3^{1/4}} - \\
 & \frac{H[3] \bar{\Delta}[125] \Phi[7]}{2^{3/4} 3^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[83] \Phi[8] - \frac{H[10] \bar{\Delta}[84] \Phi[8]}{2 \cdot e^{1/4}} + \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[9] \bar{\Delta}[108] \Phi[8] - \frac{H[9] \bar{\Delta}[109] \Phi[8]}{2 \cdot e^{1/4}} - \frac{3^{1/4} H[8] \bar{\Delta}[118] \Phi[8]}{2^{3/4}} + \frac{3^{1/4} H[7] \bar{\Delta}[120] \Phi[8]}{2^{3/4}} - \frac{2^{1/4} H[6] \bar{\Delta}[121] \Phi[8]}{3^{3/4}} - \\
 & \frac{H[4] \bar{\Delta}[123] \Phi[8]}{e^{3/4}} + \frac{H[3] \bar{\Delta}[125] \Phi[8]}{e^{3/4}} - \frac{H[10] \bar{\Delta}[85] \Phi[9]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[110] \Phi[9]}{e^{1/4}} - \left(\frac{2}{3}\right)^{3/4} H[6] \bar{\Delta}[121] \Phi[9] + \left(\frac{2}{3}\right)^{3/4} H[4] \bar{\Delta}[123] \Phi[9] - \left(\frac{2}{3}\right)^{3/4} H[3] \bar{\Delta}[125] \Phi[9] + \frac{H[10] \bar{\Delta}[82] \Phi[10]}{e^{1/4}} - \\
 & \frac{H[9] \bar{\Delta}[107] \Phi[10]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[120] \Phi[10]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[122] \Phi[10]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[3] \bar{\Delta}[124] \Phi[10] - \frac{H[10] \bar{\Delta}[81] \Phi[11]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[106] \Phi[11]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[8] \bar{\Delta}[117] \Phi[11] - \frac{H[5] \bar{\Delta}[120] \Phi[11]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[121] \Phi[11]}{e^{1/4}} - \\
 & \frac{H[10] \bar{\Delta}[80] \Phi[12]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[105] \Phi[12]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[118] \Phi[12]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[5] \bar{\Delta}[119] \Phi[12] + \frac{H[3] \bar{\Delta}[125] \Phi[12]}{e^{1/4}} - \frac{H[2] \bar{\Delta}[126] \Phi[12]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[79] \Phi[13]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[104] \Phi[13]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[115] \Phi[13]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[117] \Phi[13]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[121] \Phi[13]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[122] \Phi[13]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[123] \Phi[13]}{e^{1/4}} - \\
 & \frac{H[3] \bar{\Delta}[122] \Phi[13]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[78] \Phi[14]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[103] \Phi[14]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[7] \bar{\Delta}[117] \Phi[14] + \frac{H[6] \bar{\Delta}[121] \Phi[14]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[122] \Phi[14]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[123] \Phi[14]}{e^{1/4}} - \frac{H[3] \bar{\Delta}[124] \Phi[14]}{e^{1/4}} - \frac{H[2] \bar{\Delta}[125] \Phi[14]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[77] \Phi[15]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[102] \Phi[15]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[114] \Phi[15]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[116] \Phi[15]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[120] \Phi[15]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[121] \Phi[15]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[122] \Phi[15]}{e^{1/4}} - \frac{H[3] \bar{\Delta}[123] \Phi[15]}{e^{1/4}} - \\
 & \frac{H[4] \bar{\Delta}[118] \Phi[15]}{e^{1/4}} - \frac{H[3] \bar{\Delta}[120] \Phi[15]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[76] \Phi[16]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[9] \bar{\Delta}[101] \Phi[16] - \frac{H[8] \bar{\Delta}[116] \Phi[16]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[117] \Phi[16]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[121] \Phi[16]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[122] \Phi[16]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[123] \Phi[16]}{e^{1/4}} - \frac{H[3] \bar{\Delta}[124] \Phi[16]}{e^{1/4}} - \frac{H[2] \bar{\Delta}[125] \Phi[16]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[75] \Phi[17]}{e^{1/4}} - \frac{H[9] \bar{\Delta}[100] \Phi[17]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[115] \Phi[17]}{e^{1/4}} - \frac{H[7] \bar{\Delta}[116] \Phi[17]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[120] \Phi[17]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[121] \Phi[17]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[122] \Phi[17]}{e^{1/4}} - \frac{H[3] \bar{\Delta}[123] \Phi[17]}{e^{1/4}} - \\
 & \frac{H[5] \bar{\Delta}[116] \Phi[17]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[2] \bar{\Delta}[126] \Phi[17] + \frac{H[10] \bar{\Delta}[72] \Phi[18]}{2 \cdot e^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[73] \Phi[18] - \frac{H[10] \bar{\Delta}[73] \Phi[18]}{2 \cdot e^{1/4}} - \frac{3^{1/4} H[9] \bar{\Delta}[100] \Phi[18]}{2^{3/4}} + \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[8] \bar{\Delta}[108] \Phi[18] + \frac{H[8] \bar{\Delta}[109] \Phi[18]}{2 \cdot e^{1/4}} - \\
 & \frac{3^{1/4} H[7] \bar{\Delta}[113] \Phi[18]}{2^{3/4}} + \frac{2^{1/4} H[6] \bar{\Delta}[114] \Phi[18]}{3^{3/4}} + \frac{H[4] \bar{\Delta}[116] \Phi[18]}{e^{3/4}} - \frac{H[2] \bar{\Delta}[125] \Phi[18]}{e^{3/4}} - \frac{H[10] \bar{\Delta}[74] \Phi[19]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[110] \Phi[19]}{e^{1/4}} + \left(\frac{2}{3}\right)^{3/4} H[6] \bar{\Delta}[114] \Phi[19] - \left(\frac{2}{3}\right)^{3/4} H[4] \bar{\Delta}[116] \Phi[19] + \\
 & \left(\frac{2}{3}\right)^{3/4} H[2] \bar{\Delta}[125] \Phi[19] - \frac{H[10] \bar{\Delta}[71] \Phi[20]}{e^{1/4}} - \frac{H[8] \bar{\Delta}[107] \Phi[20]}{e^{1/4}} - \frac{H[6] \bar{\Delta}[113] \Phi[20]}{e^{1/4}} + \frac{H[4] \bar{\Delta}[115] \Phi[20]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[2] \bar{\Delta}[124] \Phi[20] - \frac{H[10] \bar{\Delta}[70] \Phi[21]}{e^{1/4}} + \left(\frac{2}{3}\right)^{1/4} H[9] \bar{\Delta}[99] \Phi[21] - \\
 & \frac{H[8] \bar{\Delta}[106] \Phi[21]}{e^{1/4}} - \frac{H[5] \bar{\Delta}[113] \Phi[21]}{e^{1/4}} - \frac{H[4] \bar{\Delta}[114] \Phi[21]}{e^{1/4}} - \frac{H[10] \bar{\Delta}[67] \Phi[21]}{2 \cdot e^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[68] \Phi[21] + \frac{H[9] \bar{\Delta}[98] \Phi[21]}{2^{3/4} 3^{1/4}} + \frac{H[8] \bar{\Delta}[105] \Phi[21]}{2^{3/4} 3^{1/4}} - \frac{H[7] \bar{\Delta}[108] \Phi[21]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[5] \bar{\Delta}[112] \Phi[21] - \\
 & \frac{H[3] \bar{\Delta}[116] \Phi[21]}{2^{3/4} 3^{1/4}} - \frac{H[2] \bar{\Delta}[123] \Phi[21]}{2^{3/4} 3^{1/4}} - \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[67] \Phi[22] + \frac{H[10] \bar{\Delta}[68] \Phi[22]}{2 \cdot e^{1/4}} - \frac{3^{1/4} H[9] \bar{\Delta}[98] \Phi[22]}{2^{3/4}} - \frac{3^{1/4} H[8] \bar{\Delta}[105] \Phi[22]}{2^{3/4}} - \frac{H[7] \bar{\Delta}[109] \Phi[22]}{e^{1/4}} - \frac{2^{1/4} H[6] \bar{\Delta}[111] \Phi[22]}{3^{3/4}} - \\
 & \frac{H[3] \bar{\Delta}[116] \Phi[22]}{e^{3/4}} - \frac{H[2] \bar{\Delta}[123] \Phi[22]}{e^{3/4}} - \frac{H[10] \bar{\Delta}[69] \Phi[22]}{e^{1/4}} + \frac{H[7] \bar{\Delta}[110] \Phi[22]}{e^{1/4}} - \left(\frac{2}{3}\right)^{3/4} H[6] \bar{\Delta}[111] \Phi[22] - \left(\frac{2}{3}\right)^{3/4} H[3] \bar{\Delta}[116] \Phi[22] - \left(\frac{2}{3}\right)^{3/4} H[2] \bar{\Delta}[123] \Phi[22] + \frac{H[10] \bar{\Delta}[64] \Phi[23]}{2 \cdot e^{1/4}} - \\
 & \frac{1}{2} \left(\frac{3}{2}\right)^{1/4} H[10] \bar{\Delta}[65] \Phi[23] - \frac{H[9] \bar{\Delta}[97] \Phi[23]}{2^{3/4} 3^{1/4}} - \frac{H[8] \bar{\Delta}[104] \Phi[23]}{2^{3/4} 3^{1/4}} + \frac{H[6] \bar{\Delta}[108] \Phi[23]}{e^{1/4}} - \left(\frac{2}{3}\right)^{1/4} H[4] \bar{\Delta}[112] \Phi[23] + \frac{H[3] \bar{\Delta}[115] \Phi[23]}{2^{3/4} 3^{1/4}} - \frac{H[2] \bar{\Delta}[122] \Phi[23]}{2^{3/4} 3^{1/4}} - \frac{H[10] \bar{\Delta}[64] \Phi[23]}{2 \cdot e^{1/4}} - \frac{H[10] \bar{\Delta}[65] \Phi[23]}{e \cdot e^{1/4}} +
 \end{aligned}$$

Peculiar numerical factors in the Lagrangian: this is due to the Clebsch-Gordon normalization convention of the program, which is used consistently for all groups, all products of representations. (see slide 22)

Output example: minimal SO(10)

Model Information

Gauge group

Representations

Parameters in model

Lagrangian

BetaFunctions

```

v g[1]
^ M[1]
  ^ ~~~~ beta^(1) ~~~~
    218 g[1]^2 M[1]
  ^ ~~~~ beta^(2) ~~~~
    1624 sqrt(2/3) Conjugate[y[{Phi, Delta, H}]] g[1]^2 h[{Phi, Delta, H}] +
    1624 sqrt(2/3) Conjugate[y[{Phi, Delta, H}]] g[1]^2 h[{Phi, Delta, H}] -
    2072 sqrt(42/5) Conjugate[y[{Phi, Delta, Delta}]] g[1]^2 h[{Phi, Delta, Delta}] + 112 sqrt(210) Conjugate[y[{Phi, Phi, Phi}]] g[1]^2 h[{Phi, Phi, Phi}] +
    56 sqrt(2/5) Conjugate[y[{Psi, Psi, H}, {f[1], f[3]}]] g[1]^2 h[{Psi, Psi, H}, {f[1], f[3]}] +
    152/sqrt(3) sqrt(14) Conjugate[y[{Psi, Psi, Delta}, {f[1], f[3]}]] g[1]^2 h[{Psi, Psi, Delta}, {f[1], f[3]}] +
    1624 sqrt(2/3) Conjugate[y[{Phi, Delta, H}]] g[1]^2 M[1] y[{Phi, Delta, H}] - 1624 sqrt(2/3) Con
    2072 sqrt(42/5) Conjugate[y[{Phi, Delta, Delta}]] g[1]^2 M[1] y[{Phi, Delta, Delta}] - 112 sqrt(210) Co
    56 sqrt(2/5) Conjugate[y[{Psi, Psi, H}, {f[1], f[3]}]] g[1]^2 M[1] y[{Psi, Psi, H}, {f[1], f[3]}] -
    152/sqrt(3) sqrt(14) Conjugate[y[{Psi, Psi, Delta}, {f[1], f[3]}]] g[1]^2 M[1] y[{Psi, Psi, Delta}, {f[1], f[3]}]
v y[{Psi, Psi, Delta}, {f[1], f[2]}]
v y[{Psi, Psi, H}, {f[1], f[2]}]
v y[{Phi, Phi, Phi}]
v y[{Phi, Delta, Delta}]
v y[{Phi, Delta, H}]
v y[{Phi, Delta, H}]
v mu[{Phi, Phi}]

```

Peculiar numerical factors in the RGEs are a consequence of the normalization of the Clebsch-Gordan factors used by the program. [Conversion to other normalizations is easy](#) (see documentation)