

Getting the correct Higgs mass in R-symmetric models

Thomas Grégoire



based on work done with Enrico Bertuzzo, Claudia Frugiuele and Eduardo Ponton

MSSM

Superpotential

$$W = \lambda_u Q U^c H_u + \lambda_d Q D^c H_d + \lambda_e L E^c H_d + \mu H_u H_d$$

R-parity, subgroup of U(1)-R with charge assignments

$$R(Q, U^c, D^c, L, E) = 1 \quad R(H_u, H_d) = 0 \quad R(W_\alpha^a) = 0$$

$$W_{\cancel{R}} = \lambda'' D^c D^c U^c + \lambda' Q D^c L + \lambda L L E^c + \mu_l H_u L$$

Forbidden by R-parity

Dirac gauginos

In the MSSM gauginos are Majorana

$$M\lambda\lambda$$

$$F_X \theta^2 \leftarrow \int d^2\theta \textcircled{X} W_\alpha W^\alpha$$

Can be Dirac if new superfields are added

$$W_\alpha^1, W_\alpha^2, W_\alpha^3 \quad S, T, G$$

$$M_D \lambda \Psi$$

N=2 supersymmetry
extra-dimension

Supersoft SUSY breaking

Fox, Nelson, Weiner '02

$D'\theta_\alpha$ ← D-term breaking

$$\int d^2\theta W'_\alpha W_i^\alpha \Phi_i$$

Dirac gauginos do not feed into scalar masses through **renormalization**

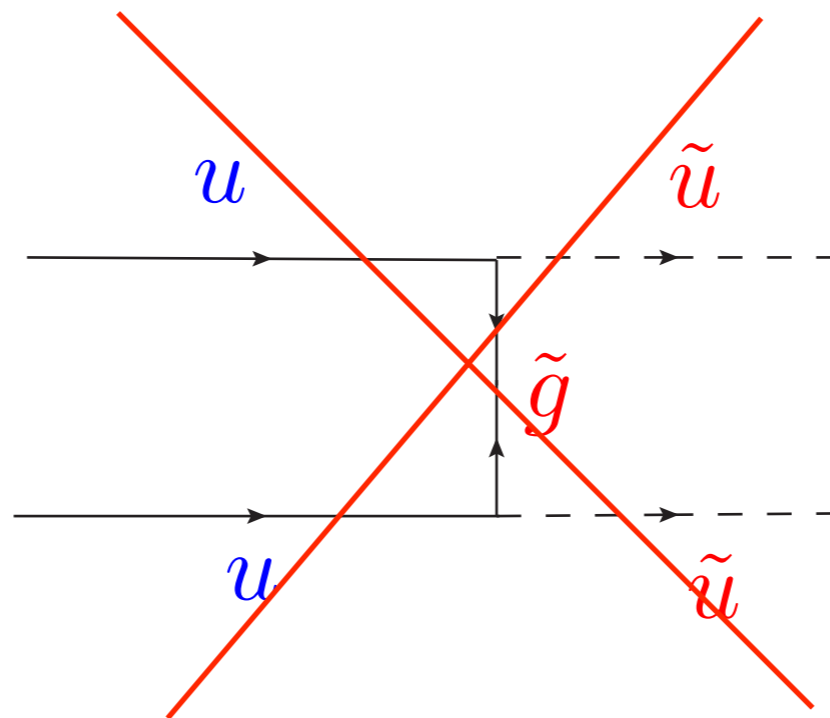
$$m^2 = \frac{C_i(r) \alpha_i m_i^2}{\pi} \log \left(\frac{\delta^2}{m_i^2} \right)$$

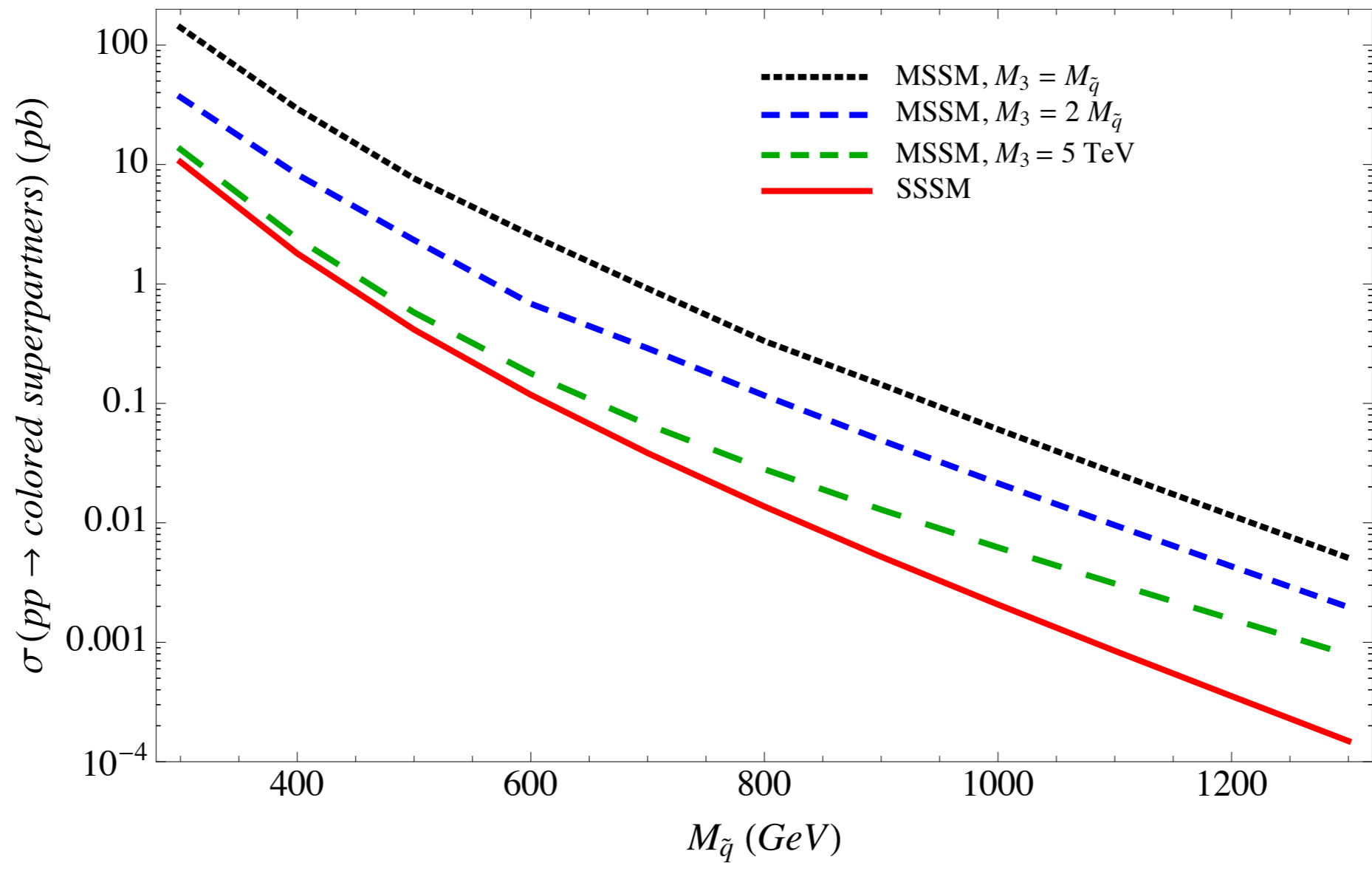
They can be naturally **heavier than scalars**

LHC will have a harder time seeing the
gluino...

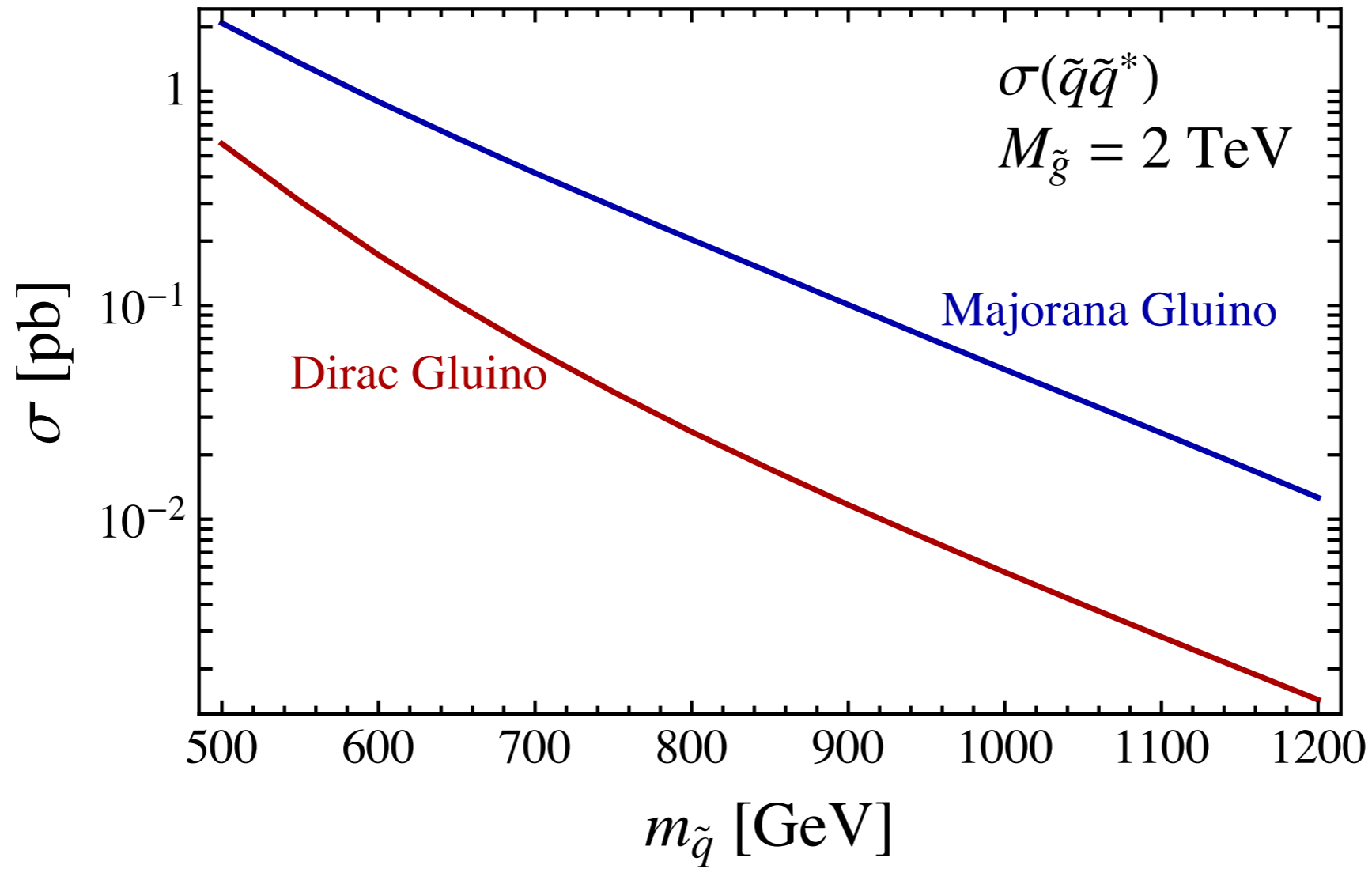
M. Heikinheimo, M. Kellerstein, V. Sanz '12
Kribs, Martin '12

...and squarks





Squark production



Frugieuele, T.G., Kumar, Ponton

R-symmetry

With Dirac gaugino: possible to impose an $U(1)$
R-symmetry

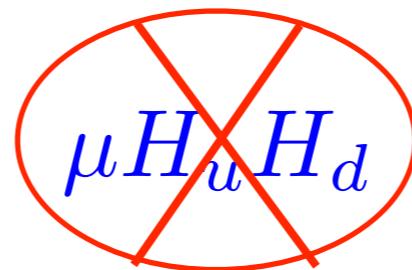
$$M_D \lambda \Psi$$

Kribs, Poppitz, Weiner '02

- Bounds from FCNC are weaker: off diagonal m_{ij}

$$R[Q, U^c, D^c, L, E^c] = 1$$

$$R[H_u, H_d] = 0$$



$\mu H_u H_d$

MRSSM:

add two additional Higgs doublets: $R_u R_d$

$$W = \sqrt{2}\lambda_T^u H_u (T) R_d + 2\lambda_T^d R_u (T) H_d + \lambda_S^u H_u (S) R_d + \lambda_S^d R_u (S) H_d + \mu_u H_u (R_d) + \mu_d H_d (R_u)$$

The diagram shows the following connections:

- Arrows from the circled T in $\sqrt{2}\lambda_T^u H_u (T) R_d$ and T in $2\lambda_T^d R_u (T) H_d$ point to "Adjoint partners".
- Arrows from the circled S in $\lambda_S^u H_u (S) R_d$ and S in $\lambda_S^d R_u (S) H_d$ point to "Adjoint partners".
- Arrows from the circled R_d in $\mu_u H_u (R_d)$ and R_u in $\mu_d H_d (R_u)$ point to "R-charge 2, do not get vets".

R-charge 2, do not get vets

Other possibility:

$U(1)_R$ can be identified with a lepton number

Superpartners have different lepton numbers

e.g.: quarks have lepton number 0

squarks have lepton number 1

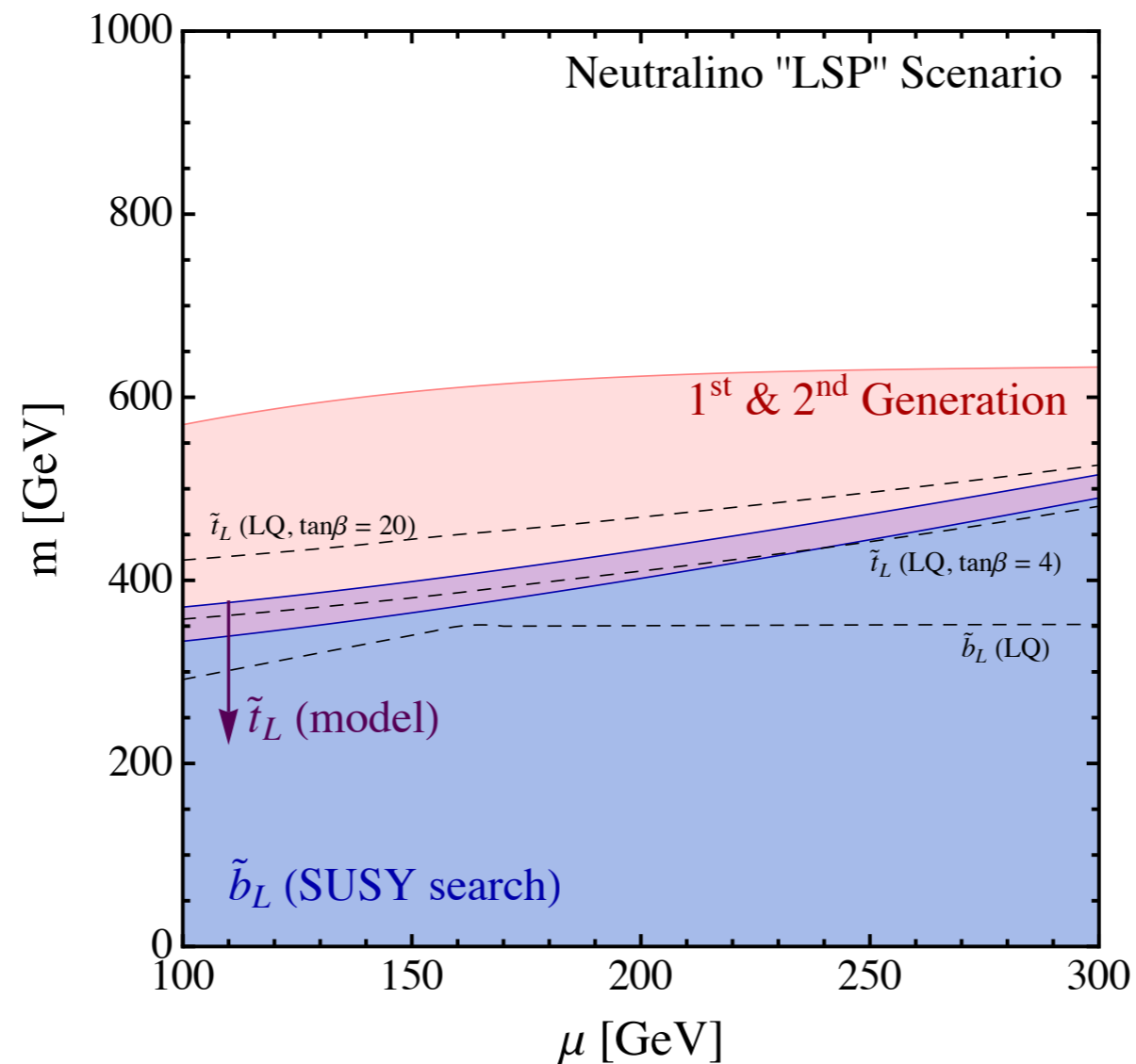
$$R(Q, U^c, D^c) = 1 \quad R(H_u, L) = 0 \quad R(R_d, E) = 2$$

$$W = \lambda_u Q H_u U + \lambda_d Q L D + \lambda_e L L E + \\ \mu H_u R_d + \lambda_T H_u T R_d + \lambda_S H_u S R_d$$

Sneutrino plays the role of the down quark
squarks are lepto-quark

Unusual phenomenology

Frugiuele, T.G, Kumar, Ponton '12



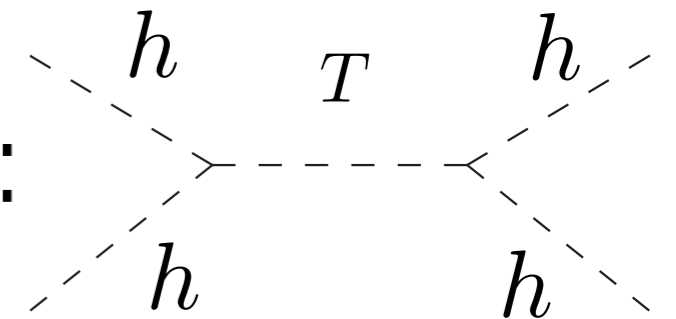
Higgs mass

Tree-level:

Reduced quartic, usual of Dirac gauginos

$$\int d^2\theta W'_\alpha W_i^\alpha \Phi_i \quad \longrightarrow \quad D_2 = M_2 T^a + H_u^\dagger \sigma^a H_u + \dots$$

When the scalar T is integrated out:



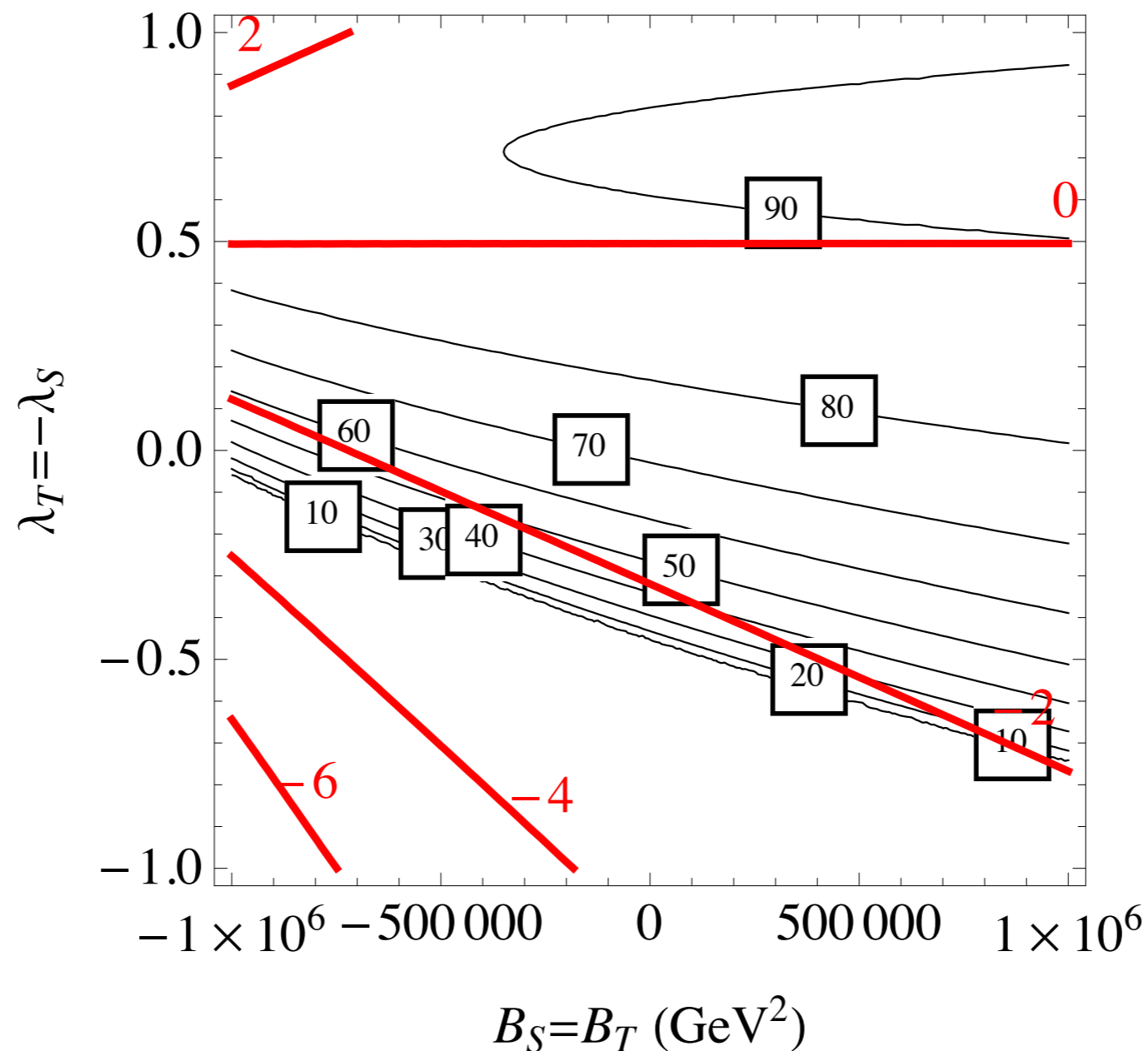
$$\lambda \rightarrow 0$$

Higgs quartic

If the mass of T is set by M_2 and $\lambda_T = 0$

With soft SUSY breaking mass terms for T (m_T)
and $\lambda_T \neq 0$

$$\delta\lambda = -\frac{(-\sqrt{2}gM_2 + 2\lambda_T\mu)^2}{m_{T_R}^2}$$



$M_2 = 600 \text{ GeV}$
 $m_T = 1500 \text{ GeV}$

No help (at tree-level) from

$$\lambda_T H_u T(R_d) + \lambda_S H_u S(R_d)$$

don't get a vev (In the limit of exact R-symmetry)

But do help in models without an R-symmetry

Benakli, Goodsell, Staub 1211.0552

Conclusions

Dirac gauginos can lead to reduced tree-level Higgs mass

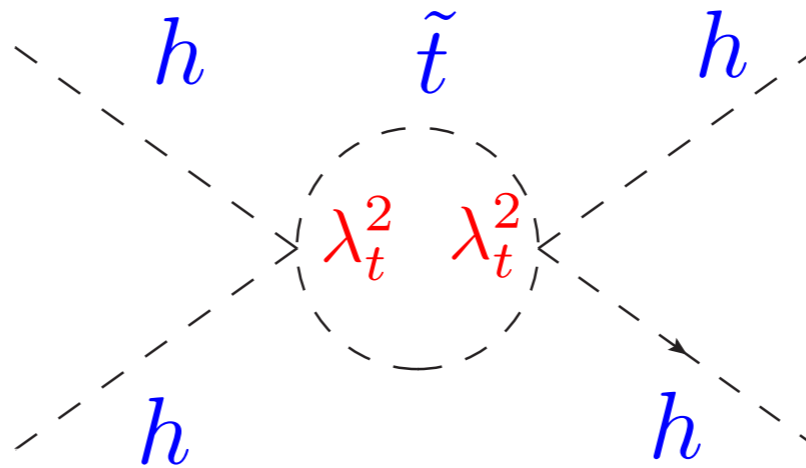
Presence of the new adjoint superfields can be used to overcome this problem

With R-symmetry, one has to rely on loop effects that are very sensitive to the coupling

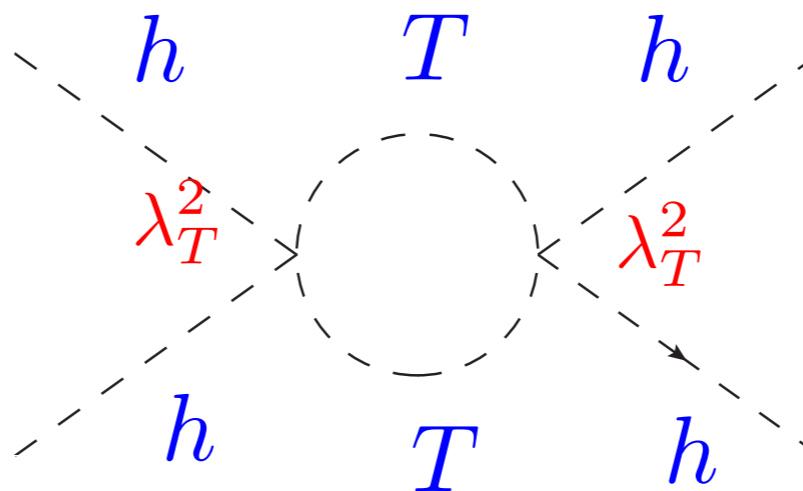
Possible tension with EWPM, but possible to raise the Higgs mass with reasonable masses.

Loop-level

Usual stop correction (but A-terms are 0)



Similar loop from the triplet



$$V_{\text{CW}} \sim \frac{1}{16\pi^2} \left(5\lambda_T^4 \log \frac{m_T^2}{M_2^2} + 3\lambda_t^4 \log \frac{m_{\tilde{t}}^2}{m_t^2} \right)$$



Very sensitive to λ_T

....but so are electroweak precision measurements

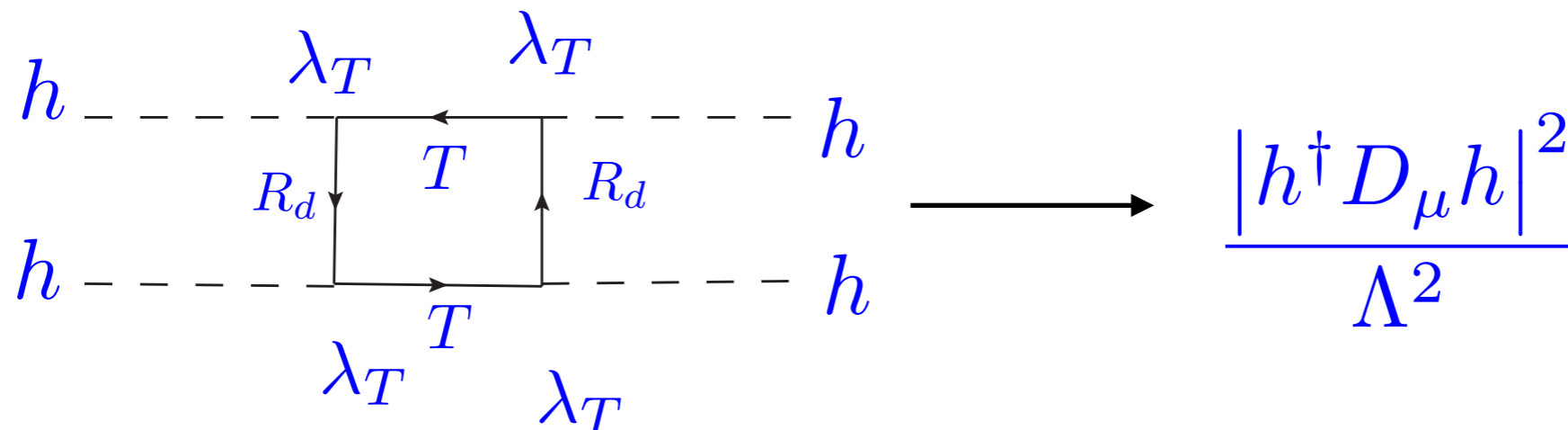
Tree-level

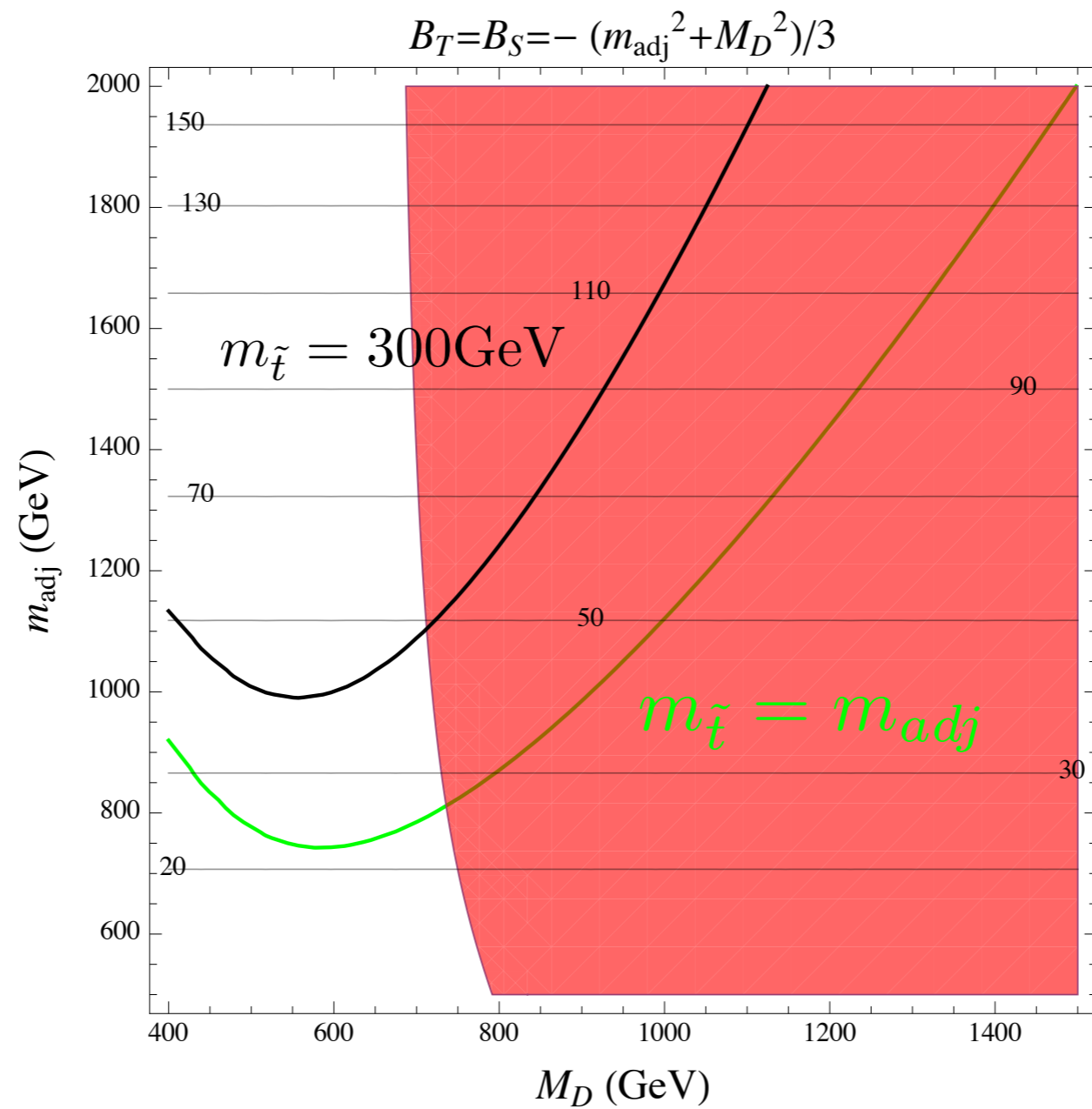
vev of the triplet

$$\hat{T} = 4 \frac{v_T^2}{v^2}$$

$$v_T = \frac{\sqrt{2}gM_2 - 2\lambda_T\mu}{2m_{T_R}^2} v^2$$

loop effect also important

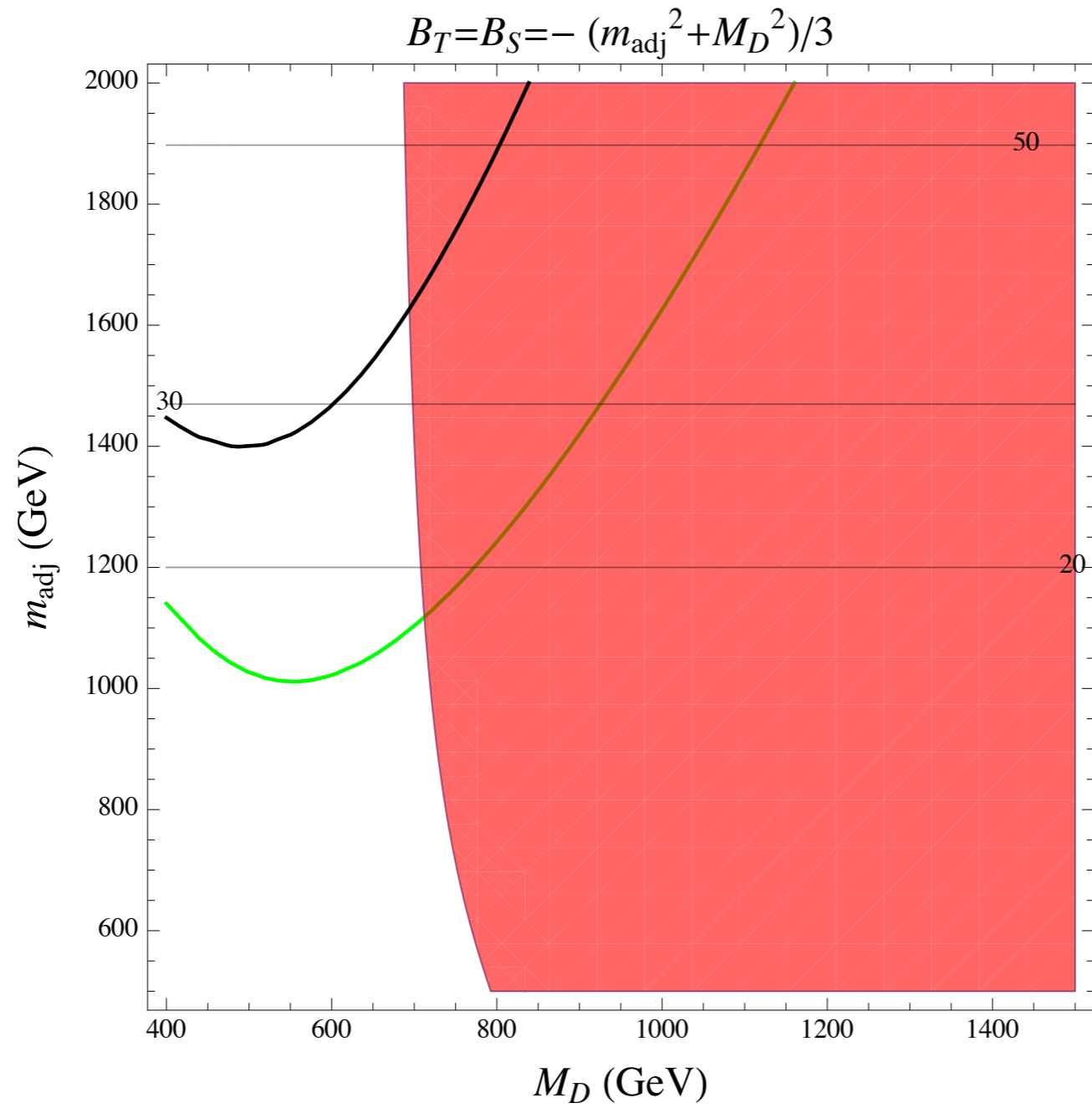




$$\lambda_T = -\lambda_S = 1$$

$$m_{R_d} = m_{adj}$$

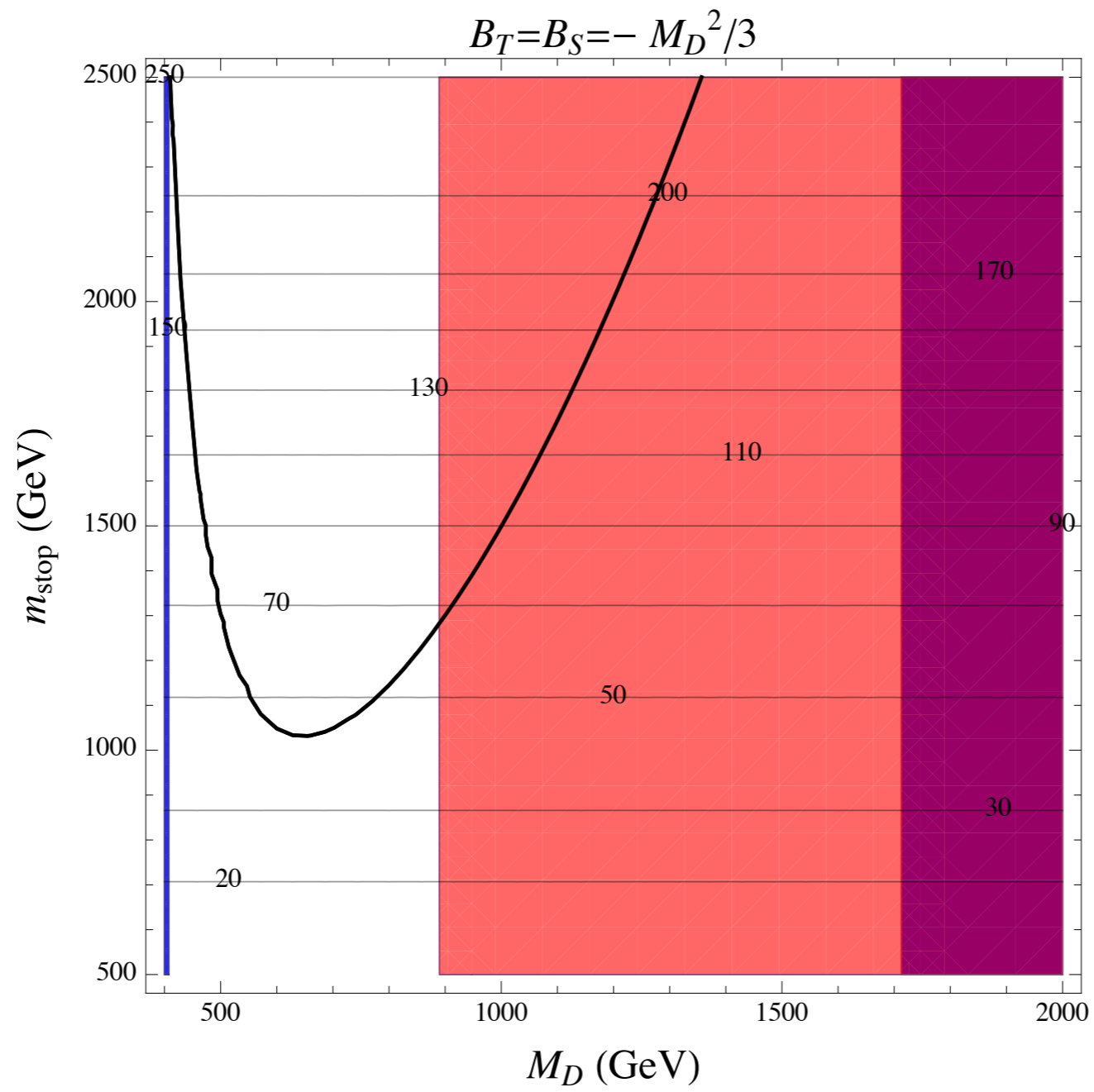
Tuning dominated by R_d



$$\lambda_T = -\lambda_S = 1$$

$$m_{R_d} = m_{adj}/2$$

$$m_T = 0$$



$$m_{R_d} = m_{\tilde{t}}$$