## Hierarchies in the Hidden sector and Gravitational rescue of gauge mediation

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## Gauge mediated Supersymmetry Breaking(GMSB) (see lectures by Dedes in Presusy)

2 Gravitational rescue of Gauge mediation

- Two spurion SUSY scale breaking and model building possibilities
- **4** An application to GMSB

- GMSB models assume that SUSY is broken in a secluded sector
- Messenger fields communicate susy breaking from hidden sector to MSSM sector



• Messenger fields transform vectorially under SM gauge group

 In mGMSB, soft terms are generated through 1-loop and 2-loop diagrams involving messenger fields coupled to the MSSM fields through gauge interactions. The messengers are connected through the hidden sector which is parameterised by a spurion field X<sub>1</sub> as

$$W_{\rm mess} = \lambda X_1 \Phi \bar{\Phi}$$

• The soft terms due to gauge mediation at the messenger scale *M*<sub>1</sub> are given by:

$$M'_{a} \approx \frac{\alpha_{a}}{4\pi} \left(\frac{F_{1}}{M_{1}}\right)$$
$$m'^{2} \approx 2\sum_{a=1}^{3} \left(\frac{\alpha_{a}}{4\pi}\right)^{2} C_{a} \left(\frac{F_{1}^{2}}{M_{1}^{2}}\right)$$
$$A' \approx 0$$

- Typically  $\Lambda \equiv F_1/M_1$  is chosen to be around 100 TeV to get low energy soft terms around at TeV
- $M_1$  can have a range between 200 TeV- GUT scale (10<sup>16</sup>GeV)
- *F*<sub>1</sub> is scaled accordingly while keeping Λ fixed.
- The low energy spectrum is computed by running the renormalisation group equations

$$egin{aligned} m_{Q_3}^2(M_{weak}) &pprox & 9.4 imes 10^{-5}\Lambda^2 \ m_{U_3}^2(M_{weak}) &pprox & 8.7 imes 10^{-5}\Lambda^2 \ A_t(M_{weak}) &pprox & -0.002\Lambda \end{aligned}$$

• The expressions are evaluated for  $M_1 = 10^6$  GeV. Keeping  $\Lambda$  fixed, if  $M_1$  is increased from  $10^6$  to  $10^{14}$ , then the magnitude of  $A_t$  only increases from 0.002 $\Lambda$  to 0.005 $\Lambda$ 

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## Character of gravitational contribtion in mGMSB

- Can gravity mediation act as additional source?
- Gravity mediation will always contribute to the soft terms at the high scale

$$\begin{split} \tilde{m}_{ij}^{2} &= k_{ij} \int d^{2}\theta d^{2}\bar{\theta} \left( \frac{F_{1}^{\dagger}F_{1}}{M_{Pl}^{2}} \right) \Phi^{\dagger}\Phi \\ \tilde{M}_{ab} &= f_{ab} \int d^{2}\theta \left( \frac{F_{1}}{M_{Pl}} \right) \mathcal{W}^{a}\mathcal{W}^{b} \\ \tilde{A}_{ijk} &= \eta_{ijk} \int d^{2}\theta \left( \frac{F_{1}}{M_{Pl}} \right) \Phi_{i}\Phi_{j}\Phi_{k} \end{split}$$

• For a fixed messenger scale,  $M_1$  (typical value of  $10^5$  GeV),  $F_1 \sim 10^{11}$ , so gravitational contributions are negligible

- Owing to the fact that A-terms are zero at the boundary mGMSB is not a viable SUSY breaking model to get 125 GeV Higgs mass
- In single spurion case, Gravitational mediation contributions are negligible

With two hierarchical spurions in the hidden sector: Is it possible to build acceptable model?

- Owing to the fact that A-terms are zero at the boundary mGMSB is not a viable SUSY breaking model to get 125 GeV Higgs mass
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With two hierarchical spurions in the hidden sector: Is it possible to build acceptable model?

- If answer is Yes, what are sources of SUSY breaking?
- Is there a UV complete model which can give rise to SUSY breaking with Hierarchical spurions?

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- Two hidden sector chiral multiplets  $X_1$  and  $X_2$
- M, A,  $\tilde{m}$  receive gravity contribution at scale  $X_1$
- Out put at scale X<sub>2</sub> receives extra contributio from gauge mediation (A terms are no longer zero at the boundary)
- Finally, weak scale out put is confronted with experiment.

## Semi-Analytical expressions for gauginos and A terms

$$egin{aligned} M_{a}(\textit{weak}) &= M_{a}(\textit{weak}) + M'_{a}(\textit{weak}) \ A_{t}(\textit{weak}) &= A_{t}(\textit{weak}) + A'_{t}(\textit{weak}) \end{aligned}$$



$$m^2_{ ilde{f}}(\textit{weak}) = m^2_{ ilde{f}}(\textit{weak}) + m^{2'}_{ ilde{f}}(\textit{weak}) + \textit{Interference terms}$$



- Due to the closed form expressions of two spurion SUSY breaking, finding dependence of low energy spectrum on SUSY breaking, in this case, is more tractable.
- One can embed various SUSY breaking models in this frame work and study the low energy spectrum
- One can also build Hidden sector models which can break SUSY dynamically with two spurions.
- We present an application to GMSB

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- We consider a model of single warped extradimension called Randall-Sundrum models
- The warp factor in this case is chosen to be

$$\epsilon = \frac{M_{IR}}{M_{PI}} \sim 0.01$$

• We assume the two Higgs doublet  $H_u$  and  $H_d$  to be localized on the low energy brane (IR brane). Supersymmetry breaking spurion  $X = \theta^2 F$  is introduced on the IR brane. SUSY breaking terms are generated by the brane local interactions of the spurion with the bulk fields

The soft breaking terms are of the form:

$$\begin{split} m_{H_u,H_d}^2 &= \hat{\beta}_{u,d} \ m_{3/2}^2 \\ (m_{\tilde{f}}^2)_{ij} &= m_{3/2}^2 \ \hat{\beta}_{ij} \ e^{(1-c_i-c_j)kR\pi}\xi(c_i)\xi(c_j) \\ A_{ij}^{u,d} &= m_{3/2}A_{ij}'e^{(1-c_i-c_j')kR\pi}\xi(c_i)\xi(c_j') \\ m_{1/2} &= k_h m_{3/2} \end{split}$$

 $\hat{eta}_{ij}, {\cal A}', k_h$  are dimensionless  ${\cal O}(1)$  parameters.  $\xi(c_i)$  is defined as

$$\xi(c_i) = \sqrt{\frac{(0.5-c_i)}{e^{(1-2c_i)\pi kR}-1}}$$

Table : Soft spectrum for RS set-up  $m_{susy}=1.66$  TeV,  $m_{\tilde{g}}=2.24$  TeV,  $tan\beta=15$ 

| Parameter             | Mass(TeV) | Paramater         | Mass(TeV) | Paramater       | Mass(TeV) |
|-----------------------|-----------|-------------------|-----------|-----------------|-----------|
| $\tilde{t}_1$         | 1.26      | $\widetilde{b}_1$ | 2.07      | $	ilde{	au}_1$  | 3.13      |
| $\tilde{t}_2$         | 2.12      | $	ilde{b}_2$      | 2.20      | $	ilde{	au}_2$  | 3.412     |
| <i>Ĉ</i> <sub>R</sub> | 2.17      | $\widetilde{s}_R$ | 2.20      | $\tilde{\mu}_R$ | 0.490     |
| $\tilde{c}_L$         | 2.31      | ŝL                | 2.31      | $	ilde{\mu}_L$  | 0.796     |
| ũ <sub>R</sub>        | 2.17      | $\tilde{d}_R$     | 2.20      | ẽ <sub>R</sub>  | 0.485     |
| ũ <sub>L</sub>        | 2.31      | $\tilde{d}_L$     | 2.31      | $\tilde{e}_L$   | 0.725     |
| $m_{A^0}$             | 3.01      | $m_H^{\pm}$       | 3.01      | m <sub>h</sub>  | 0.1243    |

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