

Covariant techniques in projective and harmonic superspace

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- Harmonic superspace was largely developed in 1984.
Galperin, Ivanov, Kalitsyn, Ogievetsky, Sokatchev
(see monograph [GIOS])
 - Projective superspace was developed also in 1984 by efforts of
Karlhede, Lindström, Roček
- Work in later years included...
- Gates, Gonzalez-Rey, Hitchin, Kuzenko, Wiles, von Unge
- Supergravity developments based on 5D work of
[Kuzenko, Tartaglino-Mazzucchelli '08]

Motivation: Why study these superspaces?

Matter actions with $\mathcal{N} = 2$ SUSY involve

- vector multiplets describing a special Kähler manifold
- hypermultiplets parametrizing a hyperkähler manifold (rigid SUSY) or quaternion-Kähler manifold (local SUSY)

Key distinction:

- Vector multiplets are off-shell but hypermultiplets are on-shell.

Some important ramifications of hypers being on-shell:

- Hyperkähler / QK are harder to construct than special Kähler.
- Higher-derivative actions easier for vector multiplets.
- Localization easily applied to vector multiplets (even in SUGRA) but trickier for hypermultiplets.

Harmonic and projective superspace allow **off-shell** hypermultiplets.

- 1 Review: Hypermultiplets in harmonic and projective superspace
- 2 Connecting projective to harmonic superspace
- 3 Applications: Sigma models and supergravity

Hypermultiplet superfields and off-shell representations

The free $N = 2$ (Fayet-Sohnius) hypermultiplet consists of f^i , ψ_α , $\bar{\chi}_{\dot{\alpha}}$. SUSY closes **on-shell**. Its superfield is given by

$$q^i(\theta, \bar{\theta}) = f^i + \theta^i \psi + \bar{\theta}_{\dot{i}} \bar{\chi} + (x\text{-derivative terms})$$
$$D_\alpha^{(i} q^{j)} = \bar{D}_{\dot{\alpha}}^{(i} q^{j)} = 0 \quad \implies \quad \square q^i = 0$$

Idea of **harmonic** and **projective superspace**:

Separate the constraint into a **kinematic** and a **dynamical** piece by introducing auxiliary manifold with coordinate v^{i+} .

- 1 $D_\alpha^+ Q^+ = \bar{D}_{\dot{\alpha}}^+ Q^+ = 0$ for $D_\alpha^+ = v_i^+ D_\alpha^i$ **analyticity condition**
- 2 $Q^+ = q^i v_i^+$ **equation of motion**

idea goes back to Rosly
(related to idea of Witten)

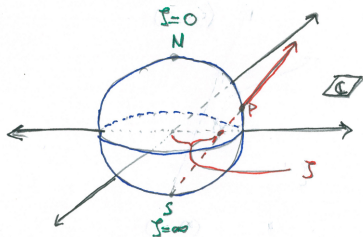
Harmonic coordinates on $S^2 \cong \mathbb{C}P^1$

Both **harmonic** and **projective superspace** use $S^2 \cong \mathbb{C}P^1$.

- Introduce **harmonics** v^{i+} and v_i^- with $v_i^- = (v^{i+})^*$ and $v^{i+} v_i^- = 1$.
- Identify $v^{i\pm} \sim e^{\pm i\psi} v^{i\pm}$.

A useful choice: $v^{i+} \sim \frac{(1, \zeta)}{\sqrt{1 + \zeta \bar{\zeta}}}$.

ζ describes $\mathbb{C}P^1 \cong S^2$.



North pole

$\zeta = 0$ or $v^{i+} \sim (1, 0)$

South pole

$\zeta = \infty$ or $v^{i+} \sim (0, 1)$

Differences between harmonic and projective superspace

The differences between harmonic and projective superspace lie in the dependence on the S^2 .

Harmonic superspace

Functions are **globally-defined**

$$Q^+ = q_i v^{i+} + q_{ijk} v^{i+} v^{j+} v^{k-} + \dots$$

Free EOM: $D^{++} Q^+ = 0$ where
 $D^{++} v^{i+} = 0$, $D^{++} v_i^- = v_i^+$,

$$\implies Q^+ = q_i v^{i+}$$

Projective superspace

Functions are holomorphic on S^2 ,
locally defined near N or S.

e.g. Q^+ is **holomorphic** near N.
It is **arctic**.

$$Q^+ = v^{1+} \sum_{n=0}^{\infty} q_n \zeta^n$$

Free EOM: Q^+ is also holomorphic
near south pole (**antarctic**).

$$\implies Q^+ = q_1 v^{1+} + q_2 v^{2+}$$

Harmonic vs projective: Comparisons

- **Harmonic superspace** nicely accommodates **gauge** and **supergravity prepotentials**. (Good for quantum calculations.)
- **Projective superspace** has useful **covariant formulations** where prepotentials are hidden within covariant derivatives. (Good for derivation of covariant component actions.)

Compare with $4D$ $N = 1$ superspace: chiral multiplet with charge q

ϕ is conventionally chiral

Φ is covariantly chiral

$$\mathcal{L} = \int d^4\theta \bar{\phi} e^{qV} \phi, \quad \bar{D}_{\dot{\alpha}}\phi = 0$$

$$\mathcal{L} = \int d^4\theta \Phi^\dagger \Phi, \quad \bar{D}_{\dot{\alpha}}\Phi = 0$$

Impose Wess-Zumino gauge on V .

No need for any gauge choice.

$$\mathcal{L} = -\partial_m \bar{\phi} \partial^m \phi - iq A^m (\bar{\phi} \overleftrightarrow{\partial}_m \phi) - q^2 A^2 \bar{\phi} \phi + \dots$$

$$\mathcal{L} = -\mathcal{D}_m \bar{\phi} \mathcal{D}^m \phi + \dots$$

Supergravity case is similar but even more tricky with prepotentials!

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Connecting projective to harmonic superspace

Main issue: they involve the **same space** ($S^2 \cong \mathbb{C}P^1$) but **different fields** (globally defined vs. holomorphic)

Two alternative ways of addressing this:

① Deform the fields. [Kuzenko '98]

Embed **projective multiplets** into globally defined **harmonic multiplets**
– holomorphic everywhere except at the poles.

② Deform the space. [Jain, Siegel '09; DB '12]

Complexify the S^2 of **harmonic superspace** to $\mathbb{C}P^1 \times \mathbb{C}P^1$.

Identify **projective superspace** S^2 with the first $\mathbb{C}P^1$.

Second $\mathbb{C}P^1$ is additional auxiliary structure.

We will take this approach.

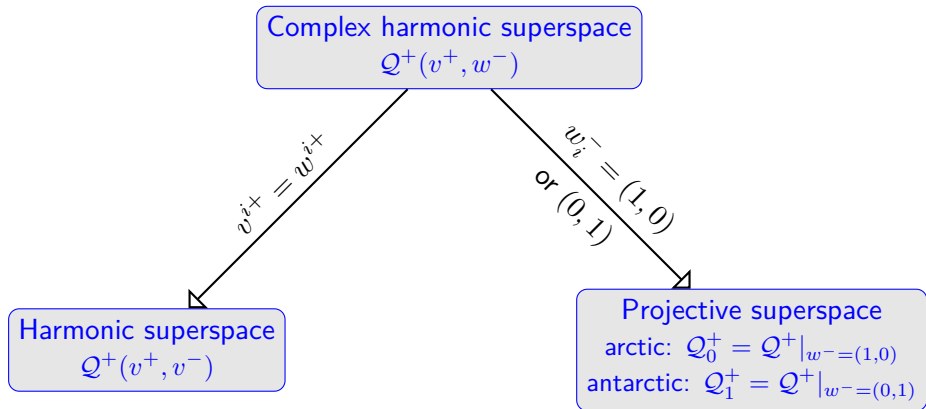
The second method is similar to an approach in twistor theory.

[Newman '86].

How do we complexify S^2 to $\mathbb{C}P^1 \times \mathbb{C}P^1$?

- Take harmonics u^{i+} and u_i^- , keep $u^{i+}u_i^- = 1$, but now $u_i^- \neq (u^{i+})^*$:

$$u^{i+} = v^{i+}, \quad u_i^- = \frac{w_i^-}{(v^+w^-)}, \quad (v^{i\pm}, w^{i\pm}) \in \mathbb{C}P^1 \times \mathbb{C}P^1$$



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Harmonic solution [GIOS]

- Action

$$\mathcal{L}^{+4} = -\frac{1}{2} \mathcal{P}^+ \overleftrightarrow{D}^{++} Q^+ + \mathcal{H}^{+4}$$

- Equations of motion

$$D^{++} Q^+ = \frac{\partial \mathcal{H}^{+4}}{\partial \mathcal{P}^+}$$

$$D^{++} \mathcal{P}^+ = -\frac{\partial \mathcal{H}^{+4}}{\partial Q^+}$$

Projective solution

[Gates, Kuzenko; Lindström, Roček]

- Action

$$\mathcal{L}^{++} = \mathcal{F}^{++}(Q_0, P_1)$$

- Equations of motion

$$P_0^+ = \frac{\partial \mathcal{F}^{++}}{\partial Q_0^+}$$

dual arctic

$$Q_1^+ = \frac{\partial \mathcal{F}^{++}}{\partial P_1^+}$$

dual antarctic

Compare to classical mechanics:

- Action principle: $F = \frac{1}{2}(q_0 p_0 + q_1 p_1) + \int_{t_0}^{t_1} dt \left(-\frac{1}{2} p \frac{\overleftrightarrow{d}}{dt} q + H \right)$
- Using Hamilton's equations, $F(q_0, p_1)$ is a canonical transformation:

$$p_0 = \frac{\partial F}{\partial q_0}, \quad q_1 = \frac{\partial F}{\partial p_1}$$

- Projective actions / solutions can be derived from harmonic ones.

Unifies two generating schemes for hyperkähler manifolds.

[DB '12]

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Supergravity: Projective superspace (review)

- Key idea: [Kuzenko '07]

$SU(2)_R$ of superconformal group $\equiv SU(2)$ isometry group of S^2
*this tells us that we must **geometrize** the R -symmetry group!*

Curved projective superspace: $\mathcal{M}^{4|8} \times S^2$

[Kuzenko, Lindström, Roček, Tartaglino-Mazzucchelli '08] and [DB '14]

$$E_{\mathcal{M}}^{\mathcal{A}} = \begin{pmatrix} E_M^A & \mathcal{V}_M^a \\ 0 & \underline{\mathcal{V}}_{\underline{m}}^a \end{pmatrix}$$

Analogous to the placement of the gravitino in the supervielbein $E_M^{\mathcal{A}}$

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Diagram illustrating the vielbein $E_{\mathcal{M}}^{\mathcal{A}}$ on curved projective superspace $\mathcal{M}^{4|8} \times S^2$. The matrix is block-structured. The top-left block is E_M^A , labeled as the vielbein on $\mathcal{M}^{4|8}$. The top-right block is \mathcal{V}_M^a , labeled as the $SU(2)_R$ connection on $\mathcal{M}^{4|8}$. The bottom-right block is $\underline{\mathcal{V}}_{\underline{m}}^a$, labeled as the vielbein, spin connection on S^2 . The bottom-left block is 0.

Analogous to the placement of the gravitino in the supervielbein $E_M^{\mathcal{A}}$

- Follow the flat case and embed curved projective superspace ($\mathcal{M}^{4|8} \times S^2$) into curved harmonic ($\mathcal{M}^{4|8} \times \mathbb{C}P^1 \times \mathbb{C}P^1$).
- Complication: “extra” $\mathbb{C}P^1$. Solution: attach “flat” $SU(2)$ group to sugra.

$$E_{\mathcal{M}}^{\mathcal{A}} = \begin{pmatrix} E_M^A & \mathcal{V}_M^a & 0 \\ 0 & \mathcal{V}_{\underline{m}}^a & 0 \\ 0 & 0 & \mathcal{W}_{\underline{m}}^{\dot{a}} \end{pmatrix}$$

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Take the framework for a test drive

- Projective and harmonic descriptions of sigma model coupled to conformal supergravity

$$S = -\frac{1}{2\pi} \oint_{\mathcal{C}} v_i^+ dv^{i+} \int d^4x d^4\theta^+ \mathcal{E} \mathcal{F}^{++} ,$$

$$S = \frac{i}{2\pi} \oint_{\mathcal{S}} v_i^+ dv^{i+} \wedge w_j^- dw^{j-} \int d^4x d^4\theta^+ \mathcal{E} \left(-\frac{1}{2} \mathcal{P}^+ \overleftrightarrow{D}^{(0,2)} \mathcal{Q}^+ + \mathcal{H}^{(2,2)} \right)$$

- Component reduction gives sigma model (a hyperkähler cone) coupled to conformal supergravity.
- Results agree with each other and with known component results of [\[de Wit, Kleijn, Vandoren '99\]](#).

Conclusion and open questions

- **Harmonic** and **projective superspaces** are not intrinsically different formulations of off-shell $\mathcal{N} = 2$ superspace but are rather **complementary**.
- Understanding projective superspace tells us how to introduce **covariant formulation** of harmonic superspace coupled to conformal supergravity.
- **Covariant formulation** readily admits **higher-derivative actions**. Can we construct these with hypermultiplets using either projective superspace or harmonic superspace? *see e.g. [DB, Kuzenko '10]*
- Can we learn (more) about **prepotentials** in projective superspace using harmonic? Advances in understanding **gauge prepotential** already due to [Jain,Siegel]. What about **supergravity**?

stay tuned...