

Conformal Extensions of the Standard Model with Veltman Conditions

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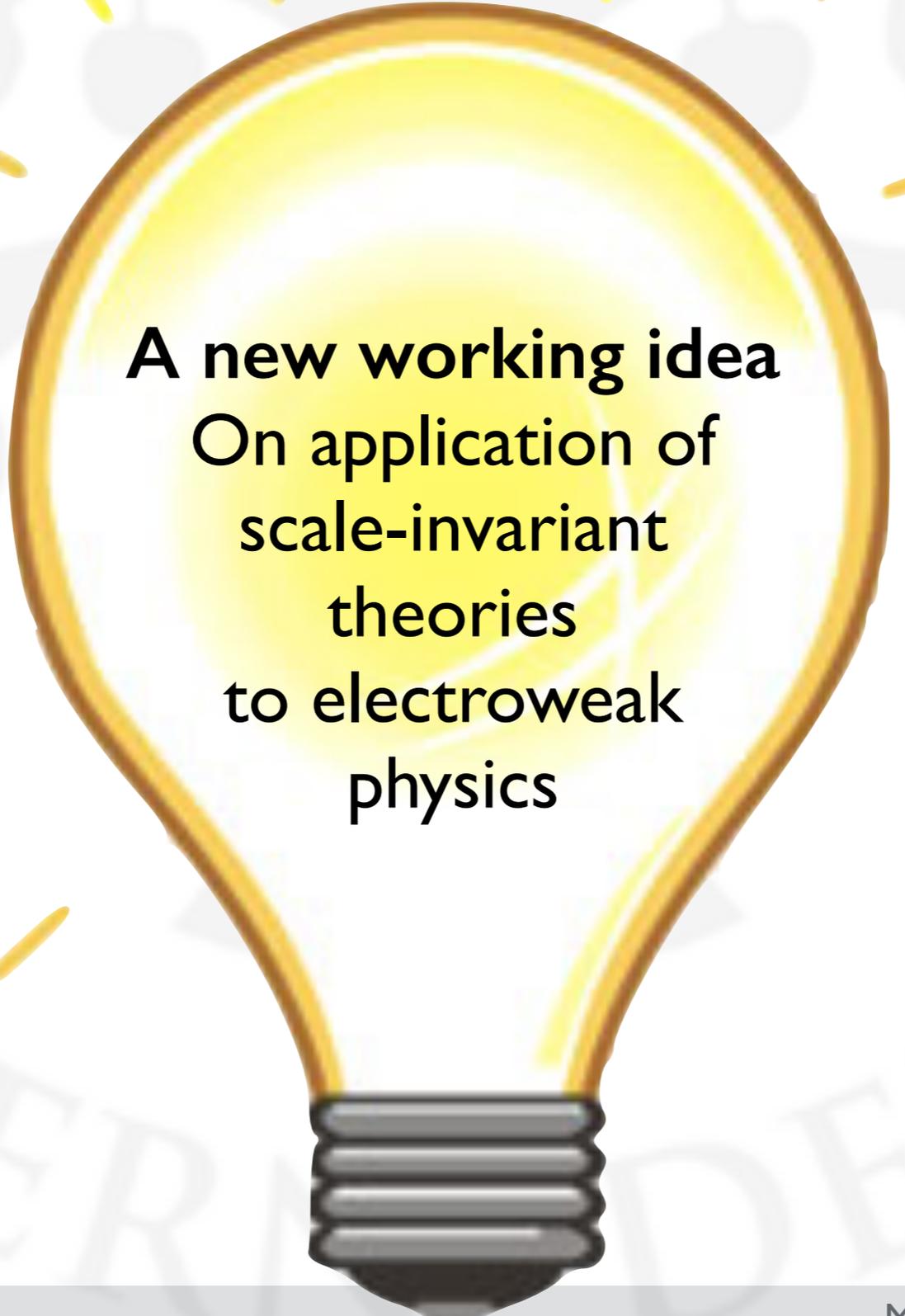
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CP³ Origins

Cosmology & Particle Physics

SUSY14, Manchester, July 21, 2014

Presenting...



**A new working idea
On application of
scale-invariant
theories
to electroweak
physics**

Coleman-Weinberg (CW) in Electroweak (EW) Theory

Has gained popularity with the recent Higgs-discovery; Why?

It naturally produces a *light* scalar compared to the new physics scale

Requires *conformality* at the classical level;

At some scale all dimensionfull couplings vanish

Spont. sym. breaking through quantum corrections

Does *not* solve the hierarchy-problem in *reality*

$$S = \int d^4x \mathcal{L}_{\text{Universe}} = \int^{\Lambda} d^4x \mathcal{L}_{\text{SM}} + \int_{\Lambda} d^4x \mathcal{L}_{\text{NP}}$$

Cutoff is physical, since SM is neither asympt. free nor safe

Fine-Tuning the Hierarchy Problem

Cutoff regularized scalar sector

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi_r)^2 - \frac{1}{2}m^2 \phi_r^2 - \frac{\lambda}{4!} \phi_r^4 + \frac{\delta_Z}{2}(\partial_\mu \phi_r)^2 - \frac{\delta_m}{2} \phi_r^2 - \frac{\delta_\lambda}{4!} \phi_r^4$$

$$m^2 = m_0^2 \left(1 + f_1(\lambda, g_i) \log \frac{\Lambda^2}{m_0^2}\right) - f_2(\lambda, g_i) \Lambda^2$$

Coleman-Weinberg Potential

$$m^2(\mu_0) \approx 0, \quad \lambda(\mu_0) \approx 0 \quad V(\phi_c) \approx V_1(\phi_c) = A_1 \phi_c^4 + B_1 \phi_c^4 \ln \frac{\phi_c^2}{\mu_0^2},$$

$$\ln \frac{\langle \phi_c^2 \rangle}{\mu_0^2} = -\frac{1}{2} - \frac{A_1}{B_1}, \quad m_{CW}^2 = \left. \frac{d^2 V}{d\phi^2} \right|_{\langle \phi_c \rangle} = 8B_1 \langle \phi_c^2 \rangle$$

Without fine-tuning:

$$m_0^2 = 0, \quad m_{CW}^2 = 8B_1 \langle \phi_c^2 \rangle - f_2(\lambda, g_i) \Lambda^2$$

Perturbative Natural Conformality

Coleman-Weinberg without fine-tuning;

Perturbative Natural Conformality (PNC)

$$m_0^2 = 0,$$

$$m_{CW}^2 = 8B_1 \langle \phi_c^2 \rangle - f_2(\lambda, g_i) \Lambda^2$$

Veltman meets Coleman-Weinberg;

Set $f_2 = 0$ at the same perturbative order as the CW-analysis

1. Classical Conformality and tree-level flatness

$$V_0(\phi_i) = \frac{\lambda_{ijkl}}{24} \phi_i \phi_j \phi_k \phi_l \quad \lambda_{ijkl}(\mu_0) n_i n_j n_k n_l \approx 0$$

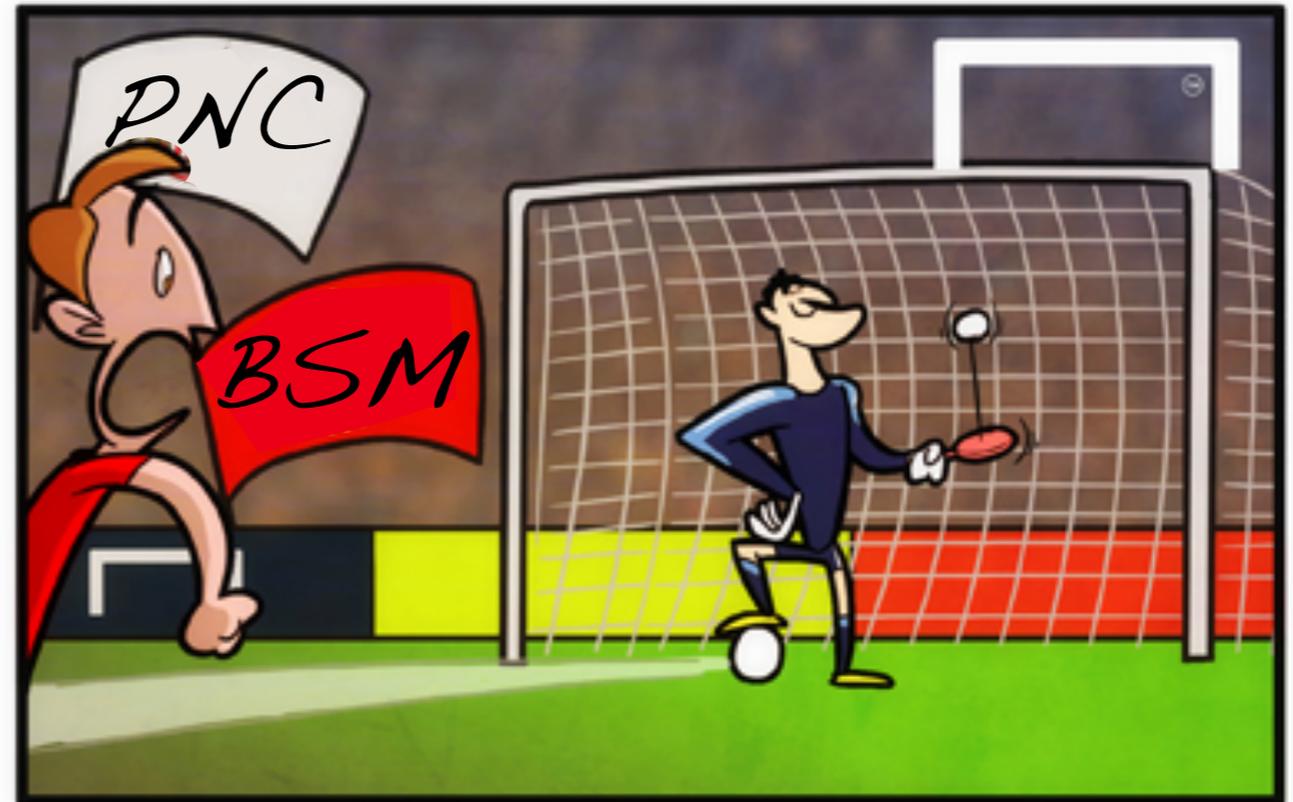
2. Quantum corrections retaining conformality

Delayed Naturalness

$$V_1(\phi_c) = \frac{1}{64\pi^2} \text{Str} \left[\underbrace{\Lambda^4 \left(\ln \Lambda^2 - \frac{1}{2} \right)}_{\text{Cosm. const.}} + \underbrace{2M^2(\phi_c) \Lambda^2}_{\text{Quad. div.}} + \underbrace{M^4(\phi_c) \left(\ln \frac{M^2(\phi_c)}{\Lambda^2} - \frac{1}{2} \right)}_{\text{(CW) phi-4 correction}} \right] + c.t.$$

$$\frac{1}{2} \frac{\partial^2 \text{Str}[M^2(\phi_i)]}{\partial \phi_i^2} \Big|_{\mu_0} = 0 \quad \text{(Veltman Condition)}$$

The PNC Shootout



PNC vs Standard Model (SM)

SM without mass term is classically conformal

$$V_0^{SM} = \lambda (H^\dagger H)^2 - \frac{1}{2} \left(g^2 W_\mu^+ W^{-\mu} + \frac{g^2 + g'^2}{2} Z_\mu Z^\mu \right) H^\dagger H + y_t (\bar{t}_L, 0) (i\sigma^2 H^*) t_R + \text{h.c.} + \text{c.t.}$$

$$H = \frac{1}{\sqrt{2}} (\pi_2 + i\pi_1, v + h - i\pi_3)$$

Breaking EW spont. *and* the delaying naturalness scale yields

CW:

$$\lambda(\mu_0) \approx 0$$

Veltman: $\frac{1}{2} \frac{\partial^2 \text{Str}[M^2(h)]}{\partial h^2} \Big|_{\mu_0} = 6\lambda(\mu_0) + \frac{9}{4}g^2(\mu_0) + \frac{3}{4}g'^2(\mu_0) - 6y_t^2(\mu_0) = 0$

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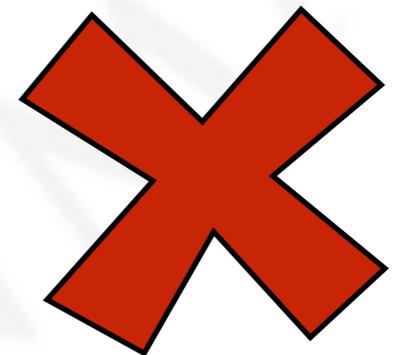
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$$\implies m_t \approx 73 \text{ GeV}$$

$$\implies m_h \approx 5 \text{ GeV}$$



Doesn't work... as we already knew

PNC vs SM + Singlet Scalar

Simplest extension of the SM potential

$$V_0 = V_0^{SM} + \lambda_{HS} H^\dagger H S^2 + \frac{\lambda_S}{4} S^4 + c.t.$$

Stability: $\lambda \geq 0$, $\lambda_S \geq 0$, and if $\lambda_{HS} < 0$: $\lambda \lambda_S \geq \lambda_{HS}^2$

2 x Veltman: $\frac{1}{2} \frac{\partial^2 \text{Str}[M^2(S)]}{\partial S^2} \Big|_{\mu_0} = 3\lambda_S(\mu_0) + 4\lambda_{HS}(\mu_0) = 0$

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Considering the moduli, $\langle H \rangle \neq 0$, $\langle S \rangle \neq 0$

CW: $\sqrt{\lambda(\mu_0)\lambda_S(\mu_0)} + \lambda_{HS}(\mu_0) = 0$

$$\phi = \sqrt{\frac{4}{7}}h + \sqrt{\frac{3}{7}}s, \quad \Phi = \sqrt{\frac{4}{7}}s - \sqrt{\frac{3}{7}}h$$

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$$\phi = \sqrt{\frac{4}{7}}h + \sqrt{\frac{3}{7}}s, \quad \Phi = \sqrt{\frac{4}{7}}s - \sqrt{\frac{3}{7}}h$$

$$\implies m_\phi \approx 95 \text{ GeV}, \quad m_\Phi \approx 541 \text{ GeV}$$



PNC vs SM + Singlet Scalar + Singlet Fermion

Singlet fermions (sterile neutrinos and/or dark matter)

$$V_0 = V_0^{SM} + \lambda_{HS} H^\dagger H S^2 + \frac{\lambda_S}{4} S^4 + y_\chi S(\chi\chi + \bar{\chi}\bar{\chi}) + c.t.$$

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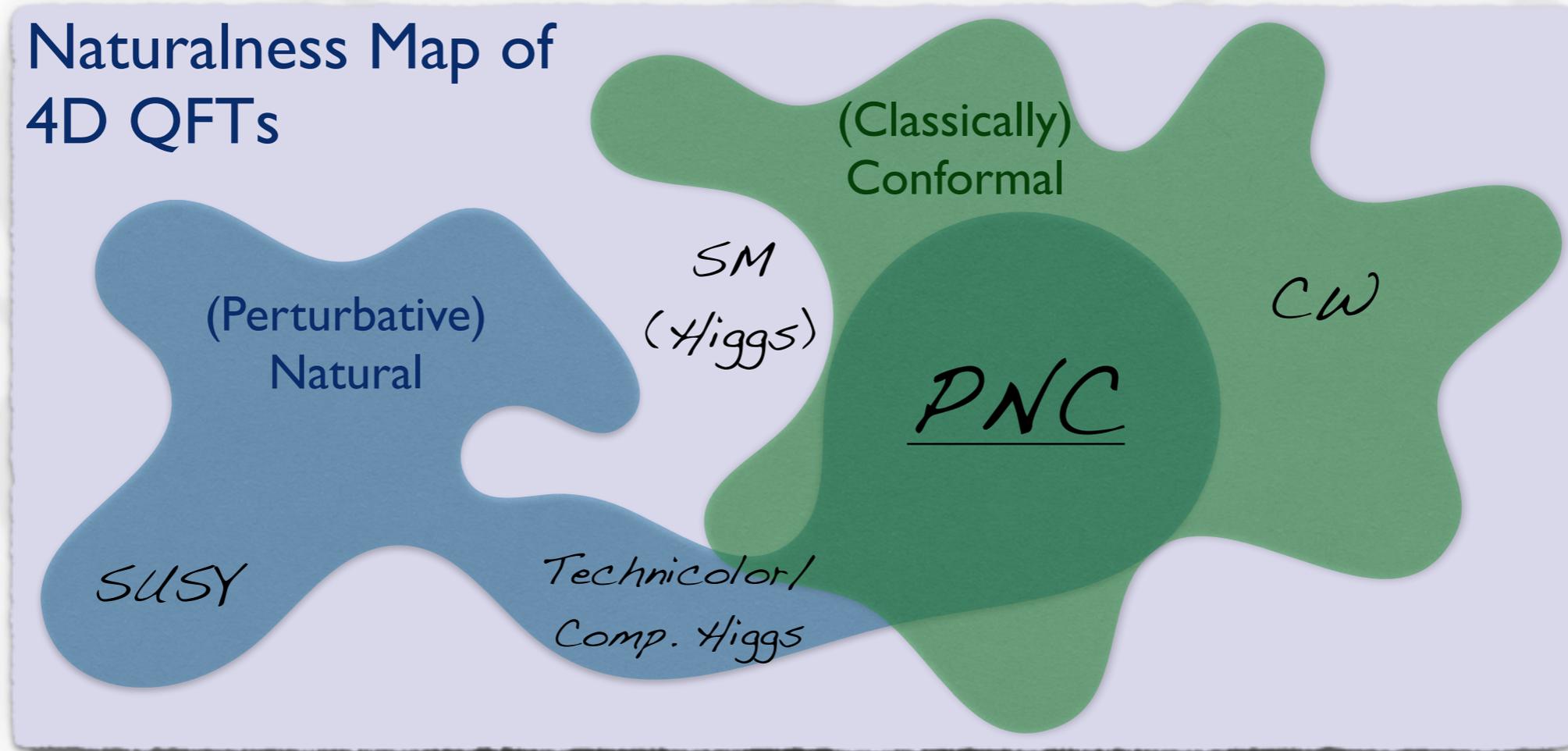
$$m_h^2 = \frac{3}{8\pi^2} \left[\frac{1}{16} (3g^4 + 2g^2g'^2 + g'^4) - y_t^4 + \frac{\lambda_{HS}^2}{3} \right] v^2 \quad \Rightarrow \quad m_h \approx 126 \text{ GeV},$$

$$m_S^2 = \lambda_{HS} v^2 \quad \Rightarrow \quad m_S \approx 541 \text{ GeV}.$$

Testable at the LHC - A natural CW scenario



Conclusions

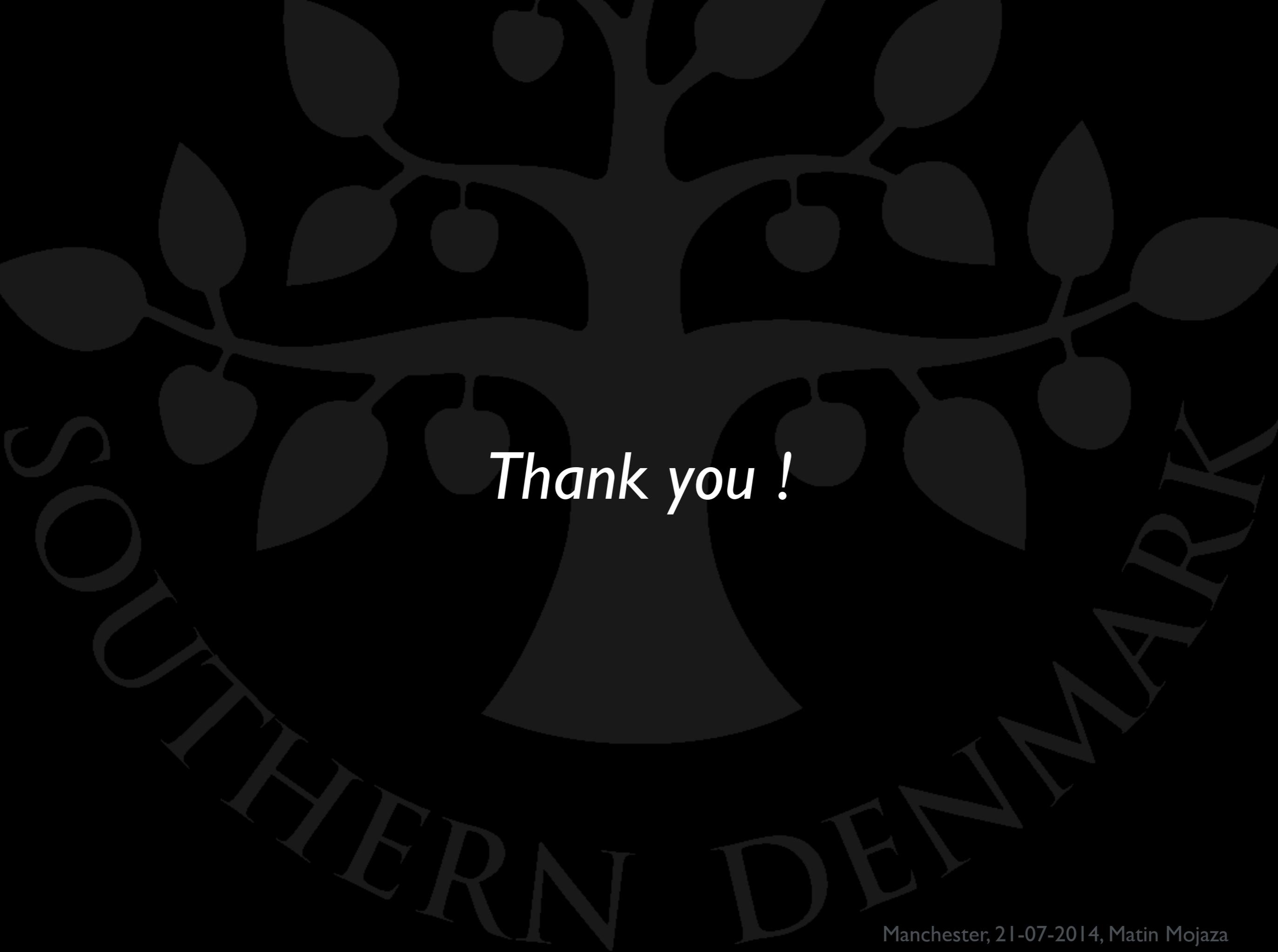


Naturalness scale could *accidentally* be orders of mag. higher

Requires new physics at LHC

Simplest realization predicts a dark sector

Highly constrained!



Thank you !