

*Based on PSB Dev, PM, A Pilaftsis & D Teresi
(cf. previous talk), arXiv: 1404.1003 [hep-ph];
accepted for publication in Nucl. Phys. B*



Flavour Covariance in Semi-Classical and Field-Theoretic Transport Phenomena

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+ Outline



- Introduction
- Flavour-Covariant Quantization
- Generalized Optical Theorem
- Implications for Resonant Leptogenesis
- Conclusions



Introduction

Why flavour?

- Observed mixing in
 - SM quark sector (**C**abibbo-**K**obayashi-**M**askawa)
 - SM neutrino sector (**P**ontecorvo-**M**aki-**N**akagawa-**S**akata)

⇒ global $U(3)_Q \times U(3)_L$ structure
- Predicted mixing
 - BSM right-handed (Majorana) neutrino sector
 - Squark and slepton sector

⇒ **complex couplings** permit **CP violating phases**



Introduction

Why transport phenomena?

- Development of the observed structure of the Universe requires processes occurring **out-of-equilibrium**

$$\Gamma(X \rightarrow Y) < H = \left(\frac{4\pi^3}{45} \right)^{1/2} g_*^{1/2} \frac{T^2}{M_{\text{Pl}}} \quad t = \frac{M_{\text{Pl}}}{34T^2}$$

- Core collapse supernovae
[Zhang and Burrows, arXiv:1310.2164 [hep-ph].]
- QCD phase diagram; jet modification in Quark Gluon Plasma
[Blaizot, Dominguez, Iancu and Mehtar-Tani, arXiv:1311.5823 [hep-ph].]
- CPT-violation from neutrino propagating in gravitational backgrounds
[Mavromatos and Sarkar, Eur. Phys. J. C 73, 2359 (2013).]

+ Goal

- Quest for a “first-principles” description of **quantum transport phenomena** that includes
 - non-Markovianity (memory effects)
 - spatial inhomogeneity
 - finite particle widths
 - spin coherence
 - ‘interspecial’ coherence (not discussed here)
 - **flavour effects: oscillations, mixing and coherence** [see *D. Teresi’s talk, Flavour Covariant Transport Equations for Resonant Leptogenesis*]
- **Non-equilibrium QFT/Kadanoff-Baym equations** vs. **semi-classical Boltzmann equations** [see e.g. extensive list of references in Dev, Millington, Pilaftsis, Teresi arXiv: 1404.1003 [hep-ph].]



Flavour-Covariant Quantization

Creation and Annihilation Operators

- \mathcal{N} -flavour model of **Dirac fermions**

$$\psi_k(x) = \begin{pmatrix} \xi_{\mathbf{a}, k}(x) \\ \bar{\eta}_{\dot{\mathbf{a}}}^k(x) \end{pmatrix} \quad (\text{Weyl basis})$$



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Creation and Annihilation Operators



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$$\psi_k(x) = \begin{pmatrix} \xi_{\mathbf{a}, k}(x) \\ \bar{\eta}_k^{\dot{\mathbf{a}}}(x) \end{pmatrix} \quad \text{(Weyl basis)}$$

$$\psi_k(x) = \sum_s \int_{\mathbf{p}} [(2E(\mathbf{p}))^{-1}]_k^l \left([e^{-ip \cdot x}]_l^m [u(\mathbf{p}, s)]_m^n b_n(\mathbf{p}, s, 0) \right. \\ \left. + [e^{ip \cdot x}]_l^m [v(\mathbf{p}, s)]_m^n d_n^\dagger(\mathbf{p}, s, 0) \right) \quad \text{(Dirac basis)}$$

$$\{b_k(\mathbf{p}, s, \tilde{t}), b^l(\mathbf{p}', s', \tilde{t})\} = (2\pi)^3 [2E(\mathbf{p})]_k^l \delta^{(3)}(\mathbf{p} - \mathbf{p}') \delta_{ss'}$$



Flavour-Covariant Quantization

Spinorial Structure (Helicity Amplitude Formalism)

■ Flavour rank-2 spinors

[cf. Bouchiat and Michel, Nucl. Phys. 5, 416 (1958); Michel, Suppl. Nuovo Cim. 14, 95 (1959); see also Pokorski, Gauge field theories, 2° ed., Cambridge University Press (2000) and Haber, in *Proceedings of the 21st SLAC Summer Institute on Particle Physics*, p. 231 (1993).]

$$[u(\mathbf{p}, s)]_k^l = \frac{1}{\sqrt{2}} \begin{pmatrix} [(E(\mathbf{p}) + m)^{1/2}]_k^l u_s(\mathbf{s}) \\ s [(E(\mathbf{p}) - m)^{1/2}]_k^l u_s(\mathbf{s}) \end{pmatrix}$$



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$$s^\mu = (0, \mathbf{s})$$

$$[s'^\mu]_k^l = \Lambda^\mu{}_\nu (\beta_k^l) s^\nu = \left(|\mathbf{p}| [m^{-1}]_k^l, [E(\mathbf{p})]_k^m [m^{-1}]_m^l \frac{\mathbf{p}}{|\mathbf{p}|} \right)$$



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■ And we cannot resist writing:

$$[E(\mathbf{p})]_k^l = [\gamma(\mathbf{p})]_k^m m_m^l \quad [\gamma^2(\mathbf{p})]_k^l = \delta_k^l + \mathbf{p}^2 [m^{-2}]_k^l$$



Flavour-Covariant Quantization

Non-Homogeneous Flavour-Covariant CTP Propagators

- **Schwinger-Keldysh CTP formalism** in a nut shell: replace **vacuum expectation values (VEVs)** with **ensemble expectation values (EEVs)**:

[see e.g. Schwinger, J. Math. Phys. **2**, 407 (1961); Keldysh, Zh. Eksp. Teor. Fiz. **47**, 1515 (1964); Calzetta and Hu, Phys. Rev. D **35**, 495 (1987); Phys. Rev. D. **37**, 2878 (1988); Berges, AIP Conf. Proc. **739** (2005) 3-62; Millington and Pilaftsis, Phys. Rev. D **88**, 085009 (2013).]

$$\left\langle \psi_k(x; \tilde{t}_i) \bar{\psi}^l(y; \tilde{t}_i) \right\rangle_t = \text{Tr} \left[\rho(\tilde{t}; \tilde{t}_i) \psi_k(x; \tilde{t}_i) \bar{\psi}^l(y; \tilde{t}_i) \right]$$

$$t = \tilde{t} - \tilde{t}_i$$

$$\left\langle b^l(\mathbf{p}', s', \tilde{t}) b_k(\mathbf{p}, s, \tilde{t}) \right\rangle_t = [(2E)^{1/2}]_k^m [f_{ss'}(\mathbf{p}, \mathbf{p}', t)]_m^n [(2E')^{1/2}]_n^l$$

+

Flavour-Covariant Quantization

Non-Homogeneous Flavour-Covariant CTP Propagators

$$\begin{aligned}
 [iS_{\geq}(p, p', \tilde{t})]_k^l &= \sum_{s, s'} 2\pi |2p_0|^{1/2} [\delta(p^2 - m^2)]_k^i \\
 &\times 2\pi |2p'_0|^{1/2} [\delta(p'^2 - m^2)]_m^l \\
 &\times e^{i(p_0 - p'_0)\tilde{t}} \\
 &\times [\mathcal{P}(p, s; p', s')]_i^{jm}{}_n \\
 &\times \left(\theta(\pm p_0)\theta(\pm p'_0)(2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{p}') \delta_{ss'} \delta_j^n \right. \\
 &\quad \left. - [\tilde{f}_{ss'}(p, p', t)]_j^n \right)
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Spectral –
 time-translational
 invariance violated

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 &\quad \left. - [\tilde{f}_{ss'}(p, p', t)]_j^n \right) \quad \text{Statistical – spin and flavour coherences; three-momentum non-conservation}
 \end{aligned}$$

Spectral –
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Flavour-Covariant Quantization

Resummation and Quasi-Particle Approximation



- **Violation of time-translation invariance** controlled by the phase:

$$(p_0 - p'_0)t \approx \frac{m_k^2 - m_l^2}{2|\mathbf{p}|}$$

- **Infeasibility of exact inversion of the Schwinger-Dyson equation:**
 - violation of time-translational invariance means infinite nesting of convolution integrals over zeroth-component momenta.
 - energy non-conservation at interaction vertices due to finite time effects.
- How are we to **extract number densities** from **resummed propagators**?



Generalized Optical Theorem

Schematics



- **Optical theorem:**

$$S^\dagger S = S S^\dagger = \mathbb{I} \quad S = \mathbb{I} + iT \quad 2 \operatorname{Im} \mathcal{T} = \mathcal{T}^\dagger \mathcal{T}$$

- Completeness of the **Fock space** is a **flavour, spin** and **isospin singlet** and a **Lorentz scalar**.

⇒ From where does all the non-trivial structure arise?



Generalized Optical Theorem

Schematics



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$$\langle 2\text{Im } \mathcal{T} \rangle = \langle \mathcal{T}^\dagger \mathcal{T} \rangle = \text{Tr } \mathcal{T} \rho \mathcal{T}^\dagger$$

$$\langle 2\text{Im } \mathcal{T} \rangle = \langle \mathcal{T}^\dagger \mathcal{T} \rangle = \sum_{A,B,C} \langle A | \mathcal{T} | B \rangle \langle B | \rho | C \rangle \langle C | \mathcal{T}^\dagger | A \rangle$$

+

Generalized Optical Theorem

Example: Heavy-Neutrino Production Rate

$$-\mathcal{L}_N = h_l^\alpha \bar{L}^l \tilde{\Phi} N_{R,\alpha} + \text{H.c.}$$

- Choose from the inserted completeness relations the term containing **an ingoing charged-lepton, Higgs state** and **an outgoing heavy Neutrino state**:



Generalized Optical Theorem

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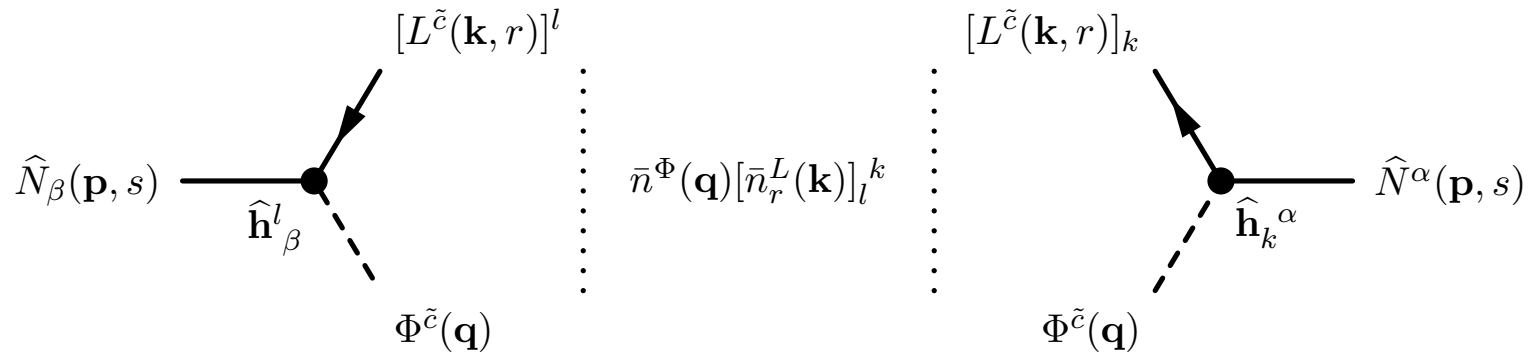
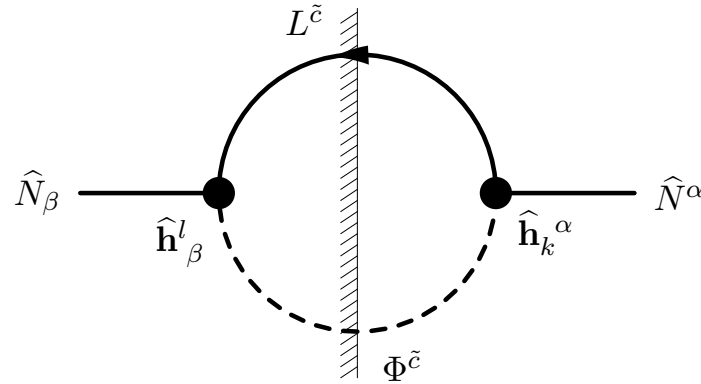
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$$\begin{aligned} [\mathcal{T}_{rs'}^0(L\Phi \rightarrow N; \mathbf{k}, \mathbf{p}', \mathbf{q}')]_\alpha^l &\equiv \langle \mathbf{k}, r, N_\alpha | \mathcal{T}^0 | \mathbf{p}', s', L^l; \mathbf{q}', \Phi^\dagger \rangle \\ &= [\mathcal{M}^0(L\Phi \rightarrow N)]^{i'}_{\gamma'} [(2E_L(\mathbf{p}'))^{-1/2}]_{i'}^{j'} \\ &\times [(2E_N(\mathbf{k}))^{-1/2}]^{\gamma'}_{\delta'} (2E_\Phi(\mathbf{q}'))^{-1/2} \\ &\times (2\pi)^4 [\delta^{(4)}(k - p' - q')]_{j'}^{m' \delta'} \epsilon'_{\epsilon'} [\bar{u}(\mathbf{k}, r)]^{\epsilon'}_\alpha P_L [u(\mathbf{p}', s')]_{m'}^l \end{aligned}$$

+

Generalized Optical Theorem

Diagrammatically

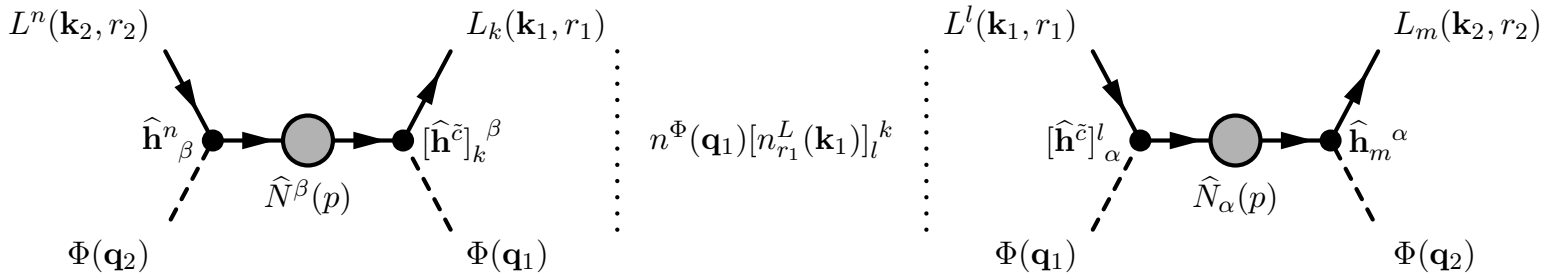
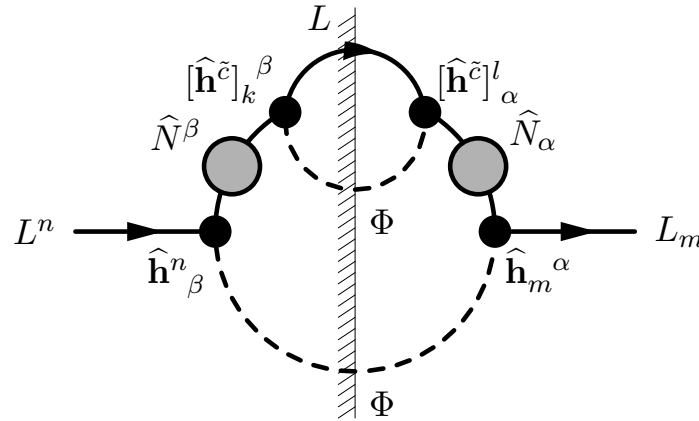


⇒ Cuts of partial self-energies leads to rank-4 tensor rates $[\gamma(L\Phi \rightarrow N)]_{k\alpha}^{l\beta}$

+

Generalized Optical Theorem

Diagrammatically

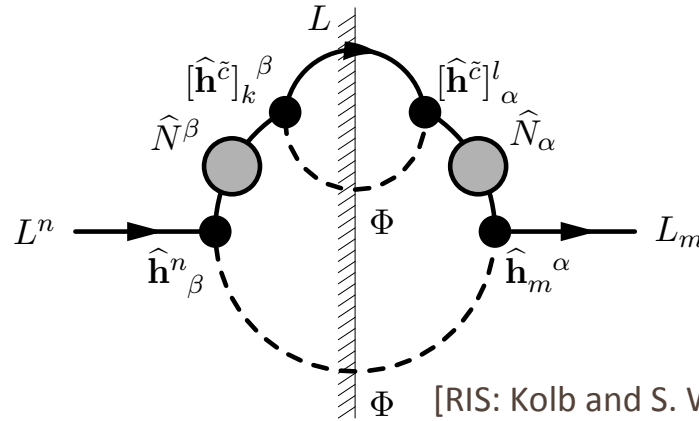


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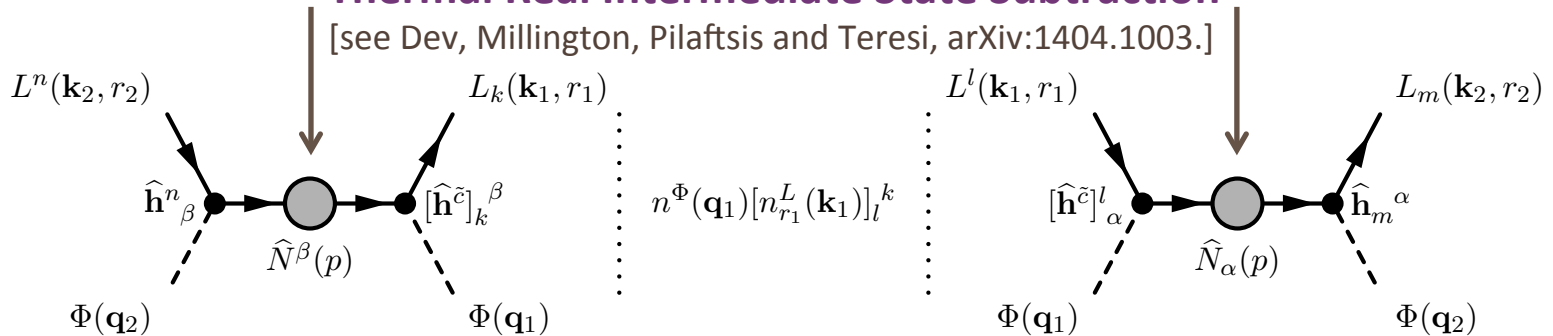
Diagrammatically



[RIS: Kolb and S. Wolfram, Nucl. Phys. B172 (1980) 224.]

Thermal Real Intermediate State Subtraction

[see Dev, Millington, Pilaftsis and Teresi, arXiv:1404.1003.]



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Implications for Resonant Leptogenesis

[see e.g. Pliافتsis and Underwood, Nucl. Phys. B692 (2004) 303-345]

- Inclusion of **higher-rank tensor rates in flavour space** allows us to account for processes that depend upon
 1. **flavour oscillations**
 2. **flavour mixing**
 3. **flavour coherences**
- Consistently accounting for these effects can lead to **enhancement** of CP asymmetry in RL by **up to an order of magnitude** [*see D. Teresi's talk, Flavour Covariant Transport Equations for Resonant Leptogenesis*]

+ Conclusions



- **Interdependence of flavour, spinorial and momentum structure** in the case of non-equilibrium statistical backgrounds.
- Potential importance for “first-principle” treatments of quantum transport phenomena, including **quasi-particle approximations**.
- Emergence of **higher-rank tensor rates** from a generalization of the optical theorem.
- Potential applications: in early Universe particle cosmology, astro-particle physics and high-temperature QCD.
- **Why still use semi-classical approaches?**