General NMSSM with a mixed Higgs sector and implication of light higgsinos

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Outline

- The motivation of our work
- How the higgsino mass is constrained
- Constraint from the Higgs signal strength
- Numerical results
- Summary



Where is SUSY?



That's all we know...

Where is SUSY?



Where is SUSY?







(L) $200 \text{GeV} < m_{\tilde{t}} < 600 \text{GeV}$



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(H) $600 \text{GeV} < m_{\tilde{t}} < 1 \text{TeV}$



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We **need additional** contributions to the Higgs mass.

Add a SM gauge **singlet** supermultiplet!

Z3-invariant NMSSM, nMSSM, PQ-NMSSM, U(1)-extended MSSM, ...

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F-term potential of S

But it can be sizable only when 2 < aneta < 4

(perturbative bounds)

$$\tan\beta = \left\langle H_{\rm u}^{\,0} \right\rangle / \left\langle H_{\rm d}^{\,0} \right\rangle$$

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With a light singlet boson



Mixing with a singlet

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$$m_s < m_h = 126 \text{GeV}$$

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- We don't need to consider the decoupling regime to get "126".
- The singlet-like boson is discoverable because of the mixing.
- The Higgs signal strength can deviate from the SM value.
- But, it is still alive.

(perturbative bounds)

With a light singlet boson



Mixing with a singlet

In order to push up the Higgs mass, we need

$$m_S < m_h = 126 \text{GeV}$$

 $h = \hat{h}\cos\theta_1\cos\theta_2 - \hat{H}\sin\theta_1 - \hat{s}\cos\theta_1\sin\theta_2$

observed

 $h = \hat{h}\cos\theta_1\cos\theta_2 - \hat{H}\sin\theta_1 - \hat{s}\cos\theta_1\sin\theta_2$

observed

SM-like

MSSM Higgs

singlet



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The LEP Higgs search



 $h = \hat{h}\cos\theta_1\cos\theta_2 - \hat{H}\sin\theta_1 - \hat{s}\cos\theta_1\sin\theta_2$ $H = \cdots$

 $s = \cdots$



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 $\left(m_0^2 + (\lambda^2 v^2 - m_Z^2)\sin^2 2\beta - (\lambda^2 v^2 - m_Z^2)\sin 2\beta\cos 2\beta - \lambda v(2\mu - \Lambda\sin 2\beta)\right)$ $- (\lambda^2 v^2 - m_Z^2) \sin^2 2\beta + \frac{2b}{\sin 2\beta} \qquad \lambda v \Lambda \cos 2\beta$ $m_{\hat{s}}^2$

 X_t

 \hat{h}

 $b, \Lambda, m_{\hat{s}}^2$: parameters depending on the model m_0^2 :Z boson mass + Q.C. (top, stop, ...) *m*¹[TeV] m² 130 $W = \lambda SH_{\mu}H_{d} + \cdots$ -22

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The higgsino mass μ is constrained by the Higgs results!

 $W = \lambda SH_{\mu}H_{d} + \cdots$

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Anything is OK.

No CPV is assumed.

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 $m_s < m_h = 126 \text{GeV}$ 350GeV $< m_H < 1 \text{TeV}$

 $X_t = (A_t - \mu \cot \beta) / m_{\tilde{t}}$

 $W = (MSSM - Yukawas) + \lambda SH_uH_d + f(S)$

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 $m_S < m_h = 126 \text{GeV}$ 350GeV $< m_H < 1 \text{TeV}$

Let us take $m_H = 800 \text{GeV}$

 $W = (MSSM - Yukawas) + \lambda SH_{u}H_{d} + f(S)$

 $b \rightarrow s \gamma$

 $|X_{t}| < 1$ $0.01 < \lambda < 1$ $m_{s} < m_{h} = 126 \text{GeV}$ $M_{H} < 1 \text{TeV} \Rightarrow \begin{array}{l} \text{Anything is OK.} \\ \text{No CPV is assumed.} \\ \text{Let us take} \\ m_{H} = 800 \text{GeV} \end{array}$

L) $200 \text{GeV} < m_{\tilde{t}} < 600 \text{GeV}$ (H) $600 \text{GeV} < m_{\tilde{t}} < 1 \text{TeV}$ $100 \text{GeV} < m_0 < 115 \text{GeV}$ $105 \text{GeV} < m_0 < 120 \text{GeV}$

When m_H is rather heavy,

$$|\lambda\mu| \cong \frac{1}{2\nu} \left[\frac{2|\theta_1 \theta_2|}{\theta_1^2 + \theta_2^2} \frac{1}{\tan\beta} \left(m_H^2 - m_S^2 \right) + |\theta_2| \left(m_h^2 - m_S^2 \right) \right]$$

$$\lambda^{2} \cong \frac{m_{Z}^{2}}{v^{2}} + \frac{\tan^{2}\beta}{4v^{2}} \left[\left(m_{h}^{2} - m_{0}^{2} \right) - \left(\theta_{1}^{2} + \theta_{2}^{2} \right) \left(m_{h}^{2} - m_{S}^{2} \right) \right]$$

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Lower bound

Bounded above

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When m_H is rather heavy,

Fixed (not so large)

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Bounded below Upper bound

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Fixed (not so large)

$$|\lambda\mu| \cong \frac{1}{2\nu} \left[\frac{2|\theta_1 \theta_2|}{\theta_1^2 + \theta_2^2} \frac{1}{\tan \beta} \left(m_H^2 - m_S^2 \right) + |\theta_2| \left(m_h^2 - m_S^2 \right) \right]$$

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Bounded below
Upper bound
Thus µ is bounded above!

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- No systematic errors
- Independent Gaussian distributions
- Additional contributions to the effective couplings
 - are dominated by those of the higgsinos and the stops
 - are taken freely, but consistently with our parameter space

Then we obtain the $\mathbf{1}\sigma$ and $\mathbf{2}\sigma$ preferred regions

- using the chi-squared distribution with 5 or 6 DOF
- taking the parameters of higgsinos and stops that minimize the chi-squared.

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	WW/ggF	ZZ/ggF	bb/VH-VBF	ττ/VH-VBF	γγ/Χ	γγ/Υ
ATLAS	0.99 ± 0.30	1.43 ± 0.38	1.09 <u>+</u>	0.34	1.66 ± 0.37	1.72 ± 0.79
CMS	0.68 ± 0.20	0.92±0.28	1.15 ± 0.62	1.10 ± 0.41	0.65 ± 0.34	1.80 ± 0.89

$$\mu^{X} - 1 = (\mu^{ggF} - 1)\cos\varphi + (\mu^{VH/VBF} - 1)\sin\varphi$$

$$\mu^{Y} - 1 = -(\mu^{ggF} - 1)\sin\varphi + (\mu^{VH/VBF} - 1)\cos\varphi$$

$$\cos\varphi = 0.98(\text{ATLAS}), 0.97(\text{CMS})$$

hep-ex/1307.1427, ATLAS-CONF-2014-009, CMS-PAS-HIG-13-005



Numerical results (tan β =5)



Numerical results (tan β =10)



Numerical results (tan β =15)



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Then, in much of the parameter region, the higgsinos tend to be light!

Find the stop, find the deviation, then the higgsinos are just beside you!

Backups

gamma gamma channels





How small is μ? - Simple example -

(L) If $\sqrt{\theta_1^2 + \theta_2^2} < 0.58$ (H) If $\sqrt{\theta_1^2 + \theta_2^2} < 0.44$

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How small is μ? - Simple example -

(L) If $\sqrt{\theta_1^2 + \theta_2^2} < 0.58$ (H) If $\sqrt{\theta_1^2 + \theta_2^2} < 0.44$ $\lambda^2 \cong \frac{m_z^2}{v^2} + \frac{\tan^2 \beta}{4v^2} [(m_h^2 - m_0^2) - (\theta_1^2 + \theta_2^2)(m_h^2 - m_s^2)] > 0$ $\lambda > \frac{m_z}{v} \cong 0.52$ $|\lambda\mu| \cong \frac{1}{2v} [\frac{2|\theta_1\theta_2|}{\theta_1^2 + \theta_2^2} \frac{1}{\tan \beta} (m_H^2 - m_s^2) + |\theta_2| (m_h^2 - m_s^2)]$ $|\mu| < 350 \text{GeV}$ (tan $\beta = 10$)