

Identification of τ Leptons at The DØ Experiment

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**THE 11TH INTERNATIONAL WORKSHOP ON
TAU LEPTON PHYSICS**

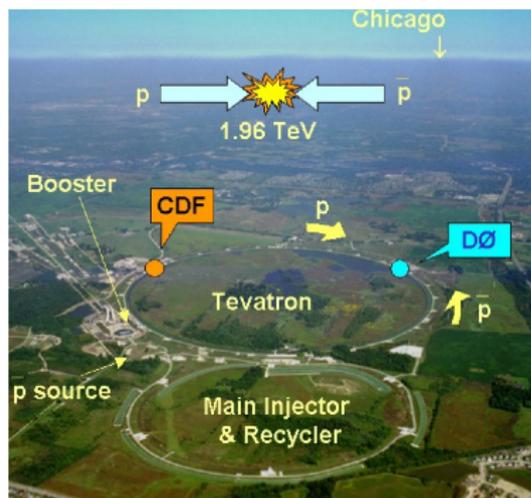
Manchester, UK – 16th of September 2010



The Tevatron

Fermilab collider :

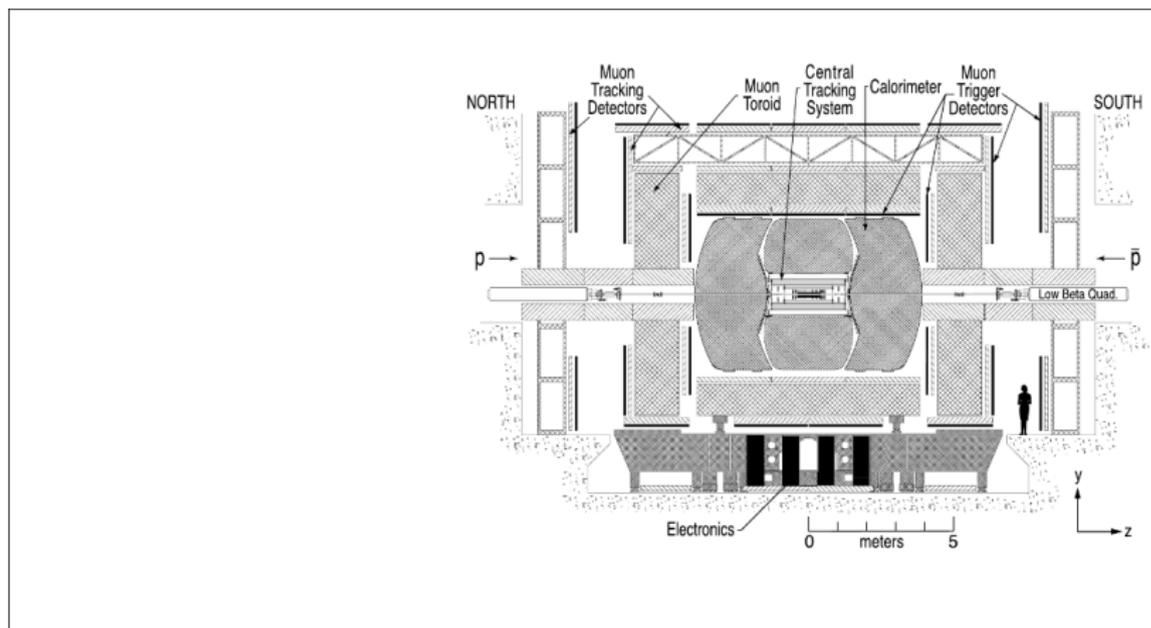
- $p\bar{p}$ collisions
- $\sqrt{s} = 1.96 \text{ TeV}$
- $\mathcal{L}_{\text{max}} \sim 400 \times 10^{30} \text{ cm}^{-2} \cdot \text{s}^{-1}$
- 8 fb^{-1} of delivered collisions
 ~ 6 millions Z bosons into leptons



Two interaction points with detectors
 CDF & DØ

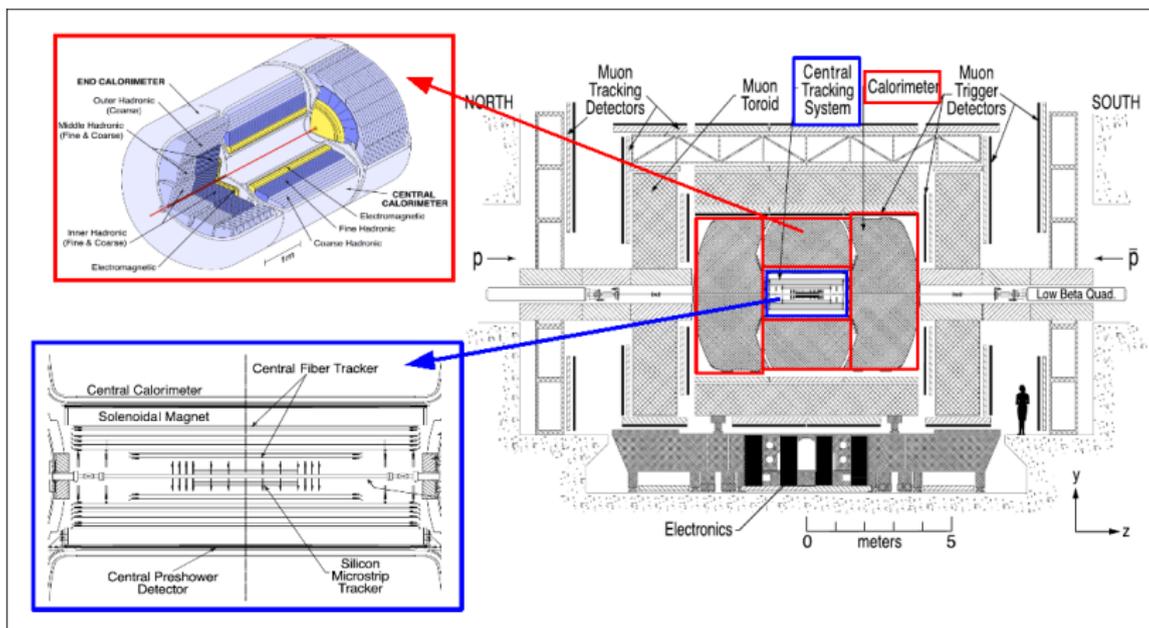
The DØ detector

Multi purpose detector :
electrons, muons, taus, photons ID, (b-)jets, mET



The DØ detector

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Why τ at hadron colliders ?

Potential acceptance gain for leptonic final states :

$(e, \mu) \Rightarrow (e, \mu, \tau)$: single lepton $\times 1.5$, dilepton $\times 2.0$, trilepton $\times 3.0$

- **Electroweak physics** : Test of lepton universality with $Z \rightarrow \tau\tau$ and $W \rightarrow \tau\nu_\tau$ cross section measurement.
- **Top quark physics** : top quark property measurements in τ final state are sensitive to new physics and test the Standard Model (SM) consistency.
- **Higgs searches** : Many decay chains initiated by Higgs boson (Electroweak Symmetry breaking origin) involve τ leptons and allow to increase the sensitivity.
- **New physics** : Supersymmetric extensions of SM predict new particles that can decay in τ leptons. τ final state acts as a probe of new physics.

... But experimentally challenging !

Impact of neutrino(s) involved in τ decay :

- ① Invisible energy : ν escapes the detector without interaction.
- ② Visible decay products are soft : more sensitive to backgrounds from soft QCD processes.

Impact of various τ decay modes :

- ① leptonic decays ($\sim 35\%$) are indistinguishables from e/μ leptons produced in W/Z direct decays which are much more abundant and suffer from poor stat. ($\mathcal{BR}(\tau\tau \rightarrow e\mu) = 6\%$).
- ② hadronic decays ($\sim 65\%$) :
 - ① different signatures depending on the hadronic final state.
 - ② Large bkg from direct QCD interactions in hadronic collisions.
- ③ Need to combine several channels.

Hadronically decaying τ leptons require sophisticated algorithms to deal with all these difficulties.

Overview

1 τ reconstruction at DØ

- Tracks and calorimeter of τ object
- τ candidate definition
- Reconstruction efficiencies

2 τ /jet discrimination

- Problematics and strategy
- Algorithm performances
- Further optimizations

3 Energy measurement

- Problematics
- Strategy : track propagation
 - Absolute correction
 - Relative correction

4 Conclusions and outlooks

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Tracks and calorimeter objects of τ_{cand}

Calorimeter cluster :

found by Simple Cone Algorithm
in a $\Delta R \leq 0.5$ cone.

CAL clu

Electromagnetic subcluster :

found by Nearest Neighbour
Algorithm with seed in the 3rd
EM layer (finer segmentation).
 $E_{\text{EMsubclu}} \geq 800$ MeV.

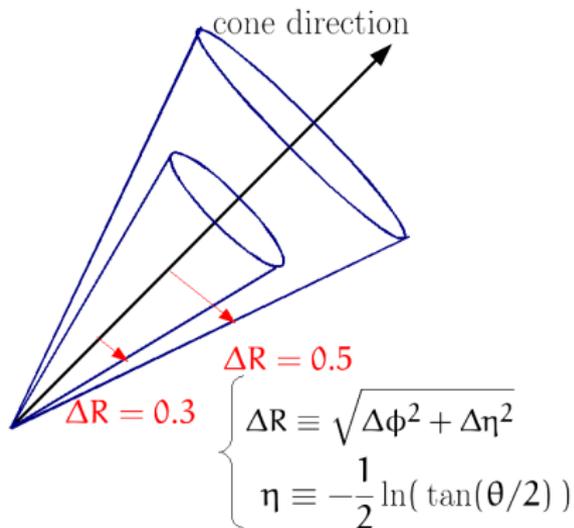
EM sub clu

Tracks :

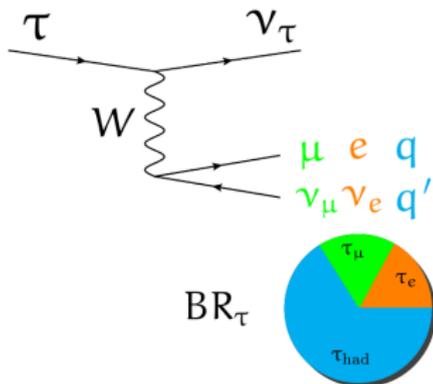
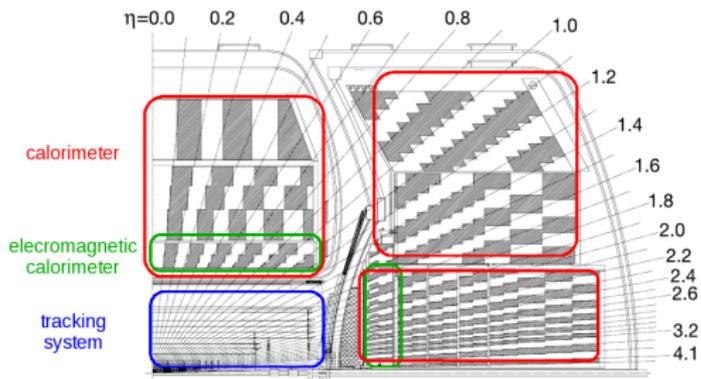
All tracks in a $\Delta R \leq 0.3$ cone around
the cal cluster compatible with τ decay
(inv. mass cut).

Highest track $p_T \geq 1.5$ GeV.

trk(s)



Type of τ candidate

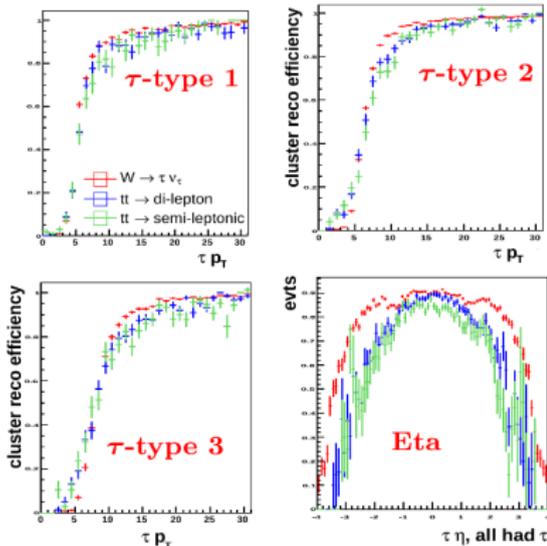
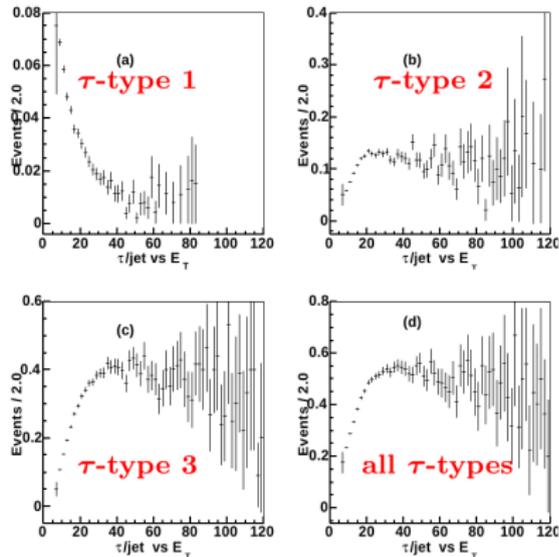


We will focus on **hadronic decay of τ** : τ_{had}

Reconstruction and DØ τ type definition for hadronic decay :

- DØ type 1 \equiv 1 **trk** , **CAL clu** $\sim \tau^\pm \rightarrow \pi^\pm \nu_\tau$
- DØ type 2 \equiv 1 **trk** , **CAL clu**, **EM sub clu** $\sim \tau^\pm \rightarrow \rho^\pm (\rightarrow \pi^0 \pi^\pm) \nu_\tau$
- DØ type 3 $\equiv \geq 2$ **trks**, **CAL clu** $\sim \tau^\pm \rightarrow a_1^\pm (\rightarrow 3\pi^\pm) \nu_\tau$

Reconstruction efficiencies

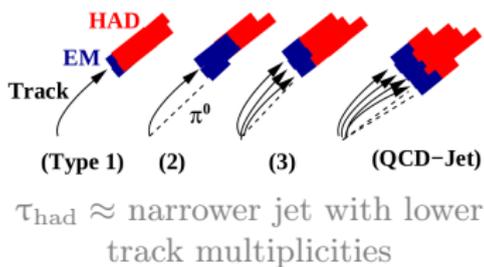
True τ (MC) τ_{cand} from QCD evts (data)

A large fraction of QCD objects (jets) can pass τ_{cand} reconstruction.

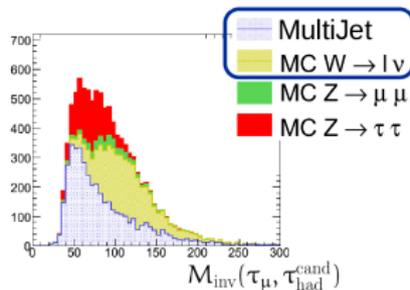
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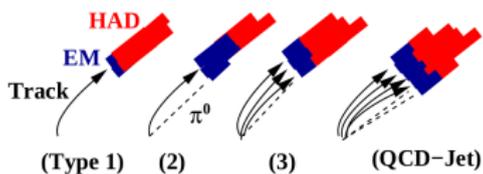
Identification of true τ



Jets could have the same experimental signature as hadronic τ and have to be removed.

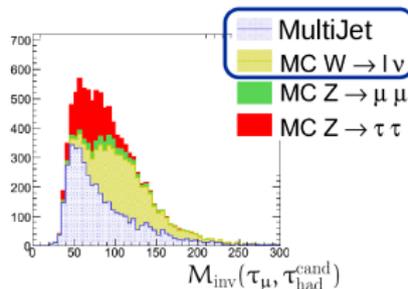


Identification of true τ



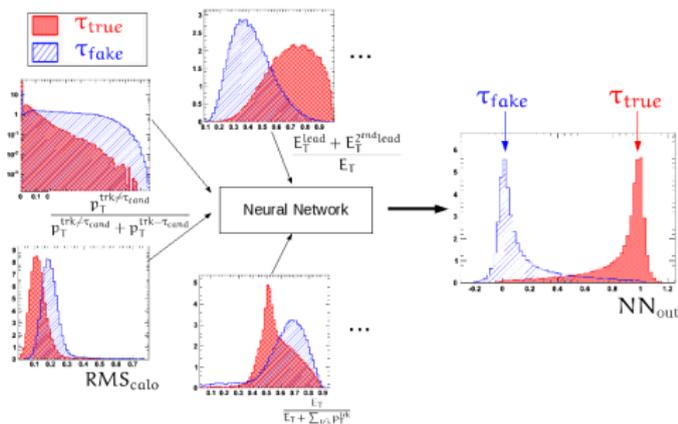
$\tau_{\text{had}} \approx$ narrower jet with lower track multiplicities

Jets could have the same experimental signature as hadronic τ and have to be removed.



τ /jet separation

Several observables having different shape for true τ and jets are combined in a Neural Network (NN).

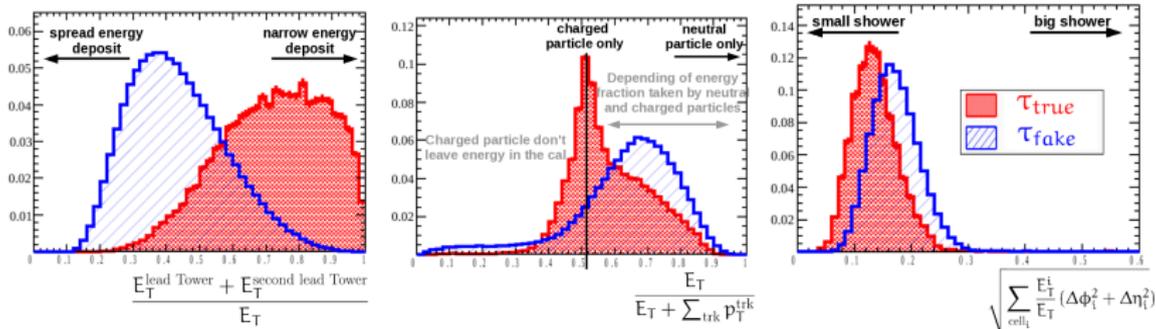


Discriminating observables

Which observables ?

- Isolation in the tracking system
- Isolation in the calorimeter
- Shower shape variables
- Correlations between tracks and calorimeter objects

Example of input variables and their physical meaning :

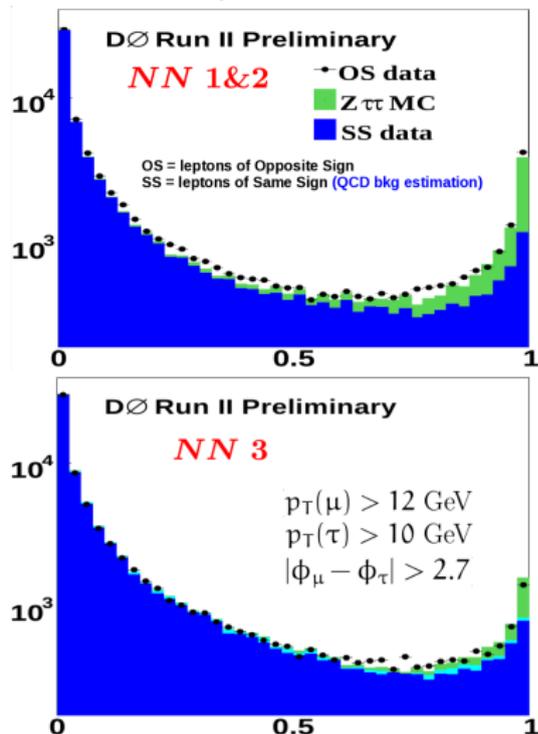


$Z \rightarrow \tau\tau$ events

Efficiencies (%)

 $20 < E_T^{\tau} < 40$ GeV, $|\eta^{\tau}| < 2.5$

τ -type	1	2	3	all
jets	2	12	38	52
τ	11	60	24	95
$NN > 0.9$				
jets	0.06	0.24	0.80	1.1
τ	7	44	16	67

Data in $\mu\tau$ final state :

$W \rightarrow \tau\nu$ events

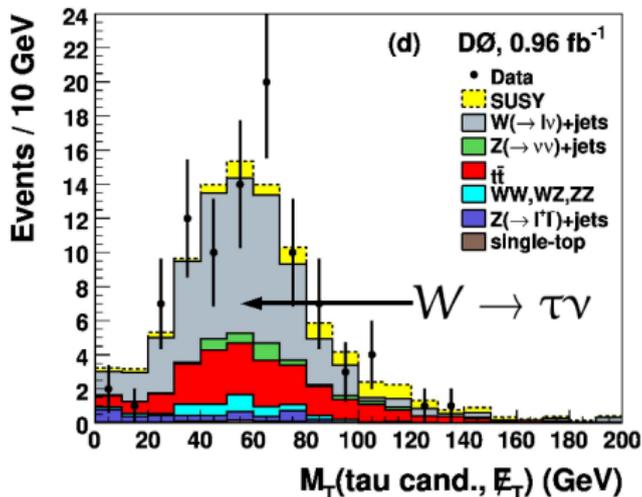
Context of search for new physics :

e.g. search for squark production in $\tau + 2\text{jets} + \text{mET}$ events.

$W \rightarrow \tau\nu \equiv \text{SM background}$

Main event selection

- $\text{mET} \geq 75 \text{ GeV}$
- $p_T(j_i) \geq 35 \text{ GeV}$
- $p_T(\tau) \geq 35 \text{ GeV}$
- $\Delta R(j_i, \tau) \geq 0.5$



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Optimization strategy

General point of view : Neural Network output $\eta^{\text{NN}}(\vec{X})$ converges to

$$\eta^{\text{true}}(\vec{X}) \equiv \frac{\mathcal{S}(\vec{X})}{\mathcal{S}(\vec{X}) + \mathcal{B}(\vec{X})}$$

where $\vec{X} \equiv (x_1, x_2, \dots, x_n)$ describes the discriminating variables space.

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In the τ identification context :

A lot of ideas were tested to optimize the identification of τ leptons :

- Include preshower detector measurement ✗
 - Exploit the long τ life time (like for b-jets) ✓
 - Tune NN parameters (epoch, nodes, statistics) ✓
 - Dedicated training for τ of high p_{τ} ✓
 - Dedicated training for high luminosity events ✗
- }

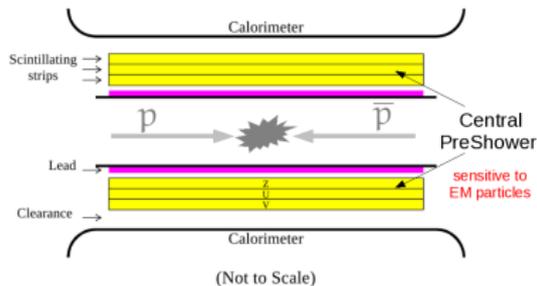
improve $\eta^{\text{true}}(\vec{X})$
- }

minimize
 $|\eta^{\text{NN}} - \eta^{\text{true}}|$

Central PreShower (CPS) for type 2

Physical idea. Exploit specific resonance of τ **type 2** decay : $\tau^\pm \rightarrow \rho^\pm \nu \rightarrow \pi^\pm \pi^0 \nu$.
Use Central PreShower detector with fine segmentation : $\Delta\phi_{\text{CPS}} \simeq 0.1 \times \Delta\phi_{\text{calo}}$

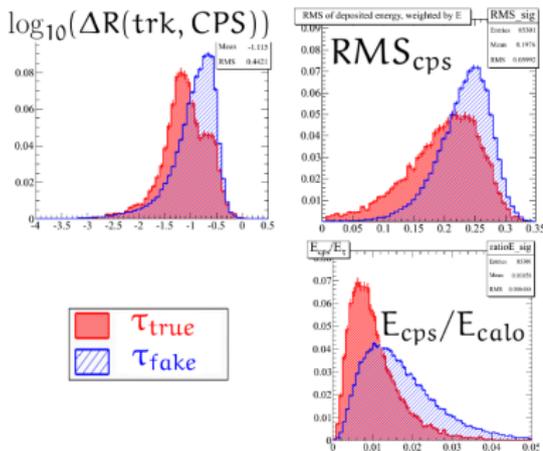
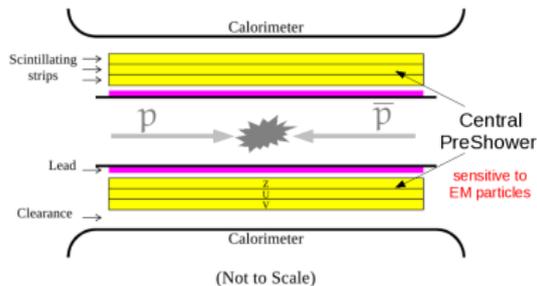
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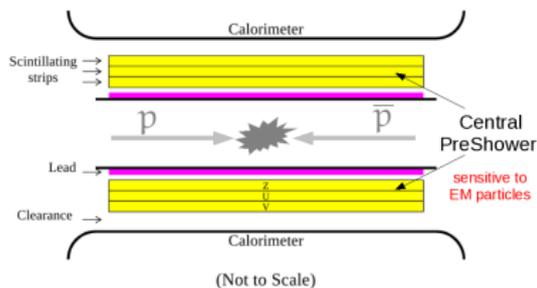
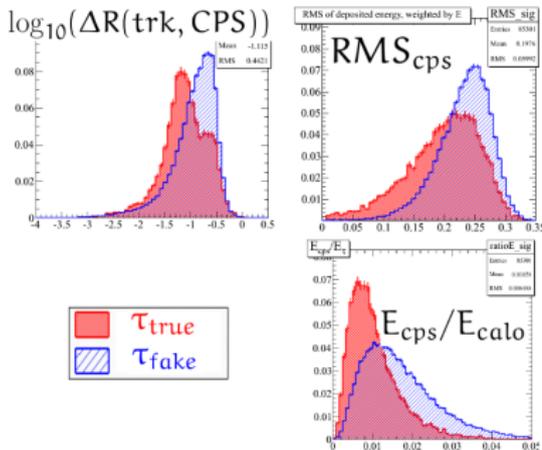
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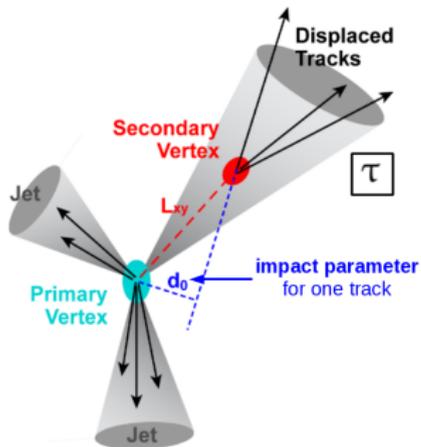
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After adding these variables in the NN **No significant improvement** was observed.

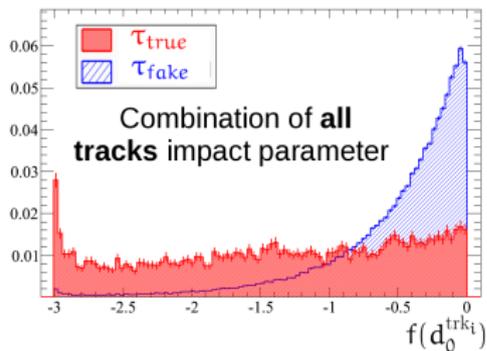
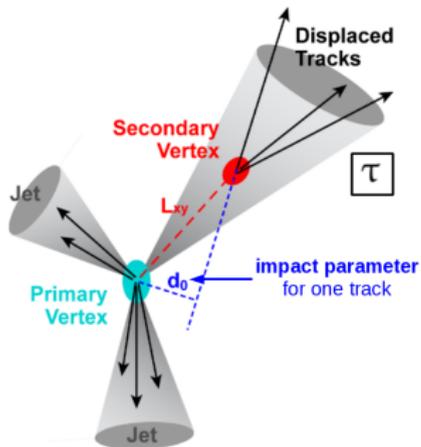
Reason : these informations were already included via calorimeter measurement.

τ is a long lived particle



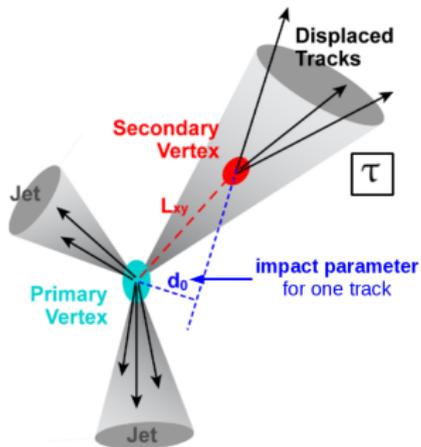
Use impact parameter to remove jets faking τ more efficiently.
 (large $c\tau_{\text{life}} \Rightarrow$ large d_0)

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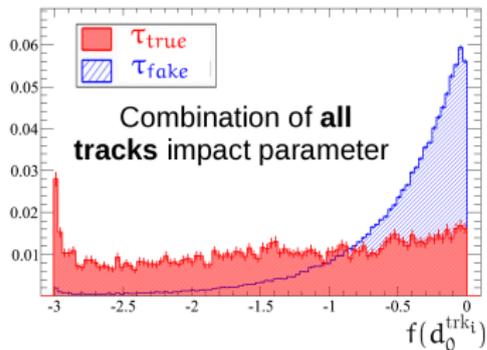


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Use impact parameter to remove jets faking τ more efficiently.
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After adding these variables in the NN
clear improvement
was observed :
 $\sim 10\%$ more signal for the same bkg

Impact of optimizations

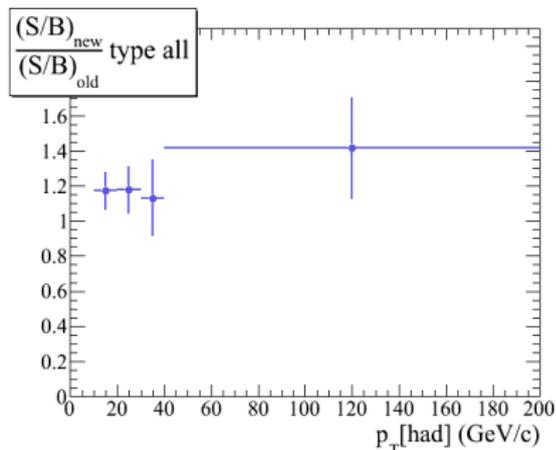
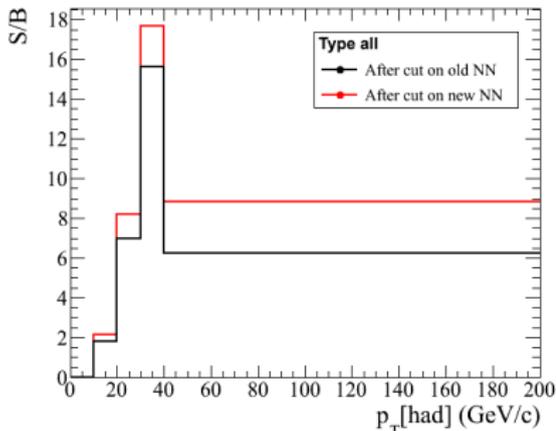
Consequences of optimizations : comparison of $S/B(p_T^{\tau\text{cand}})$ after a cut

- 1 on NN[without opt.] (old NN)
- 2 on NN[with opt.] (**new NN**)
- 3 ratio of **new/old**

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Consequences of optimizations : comparison of $S/B(p_T^{\tau_{\text{cand}}})$ after a cut

- ❶ on NN[whitout opt.] (old NN)
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- ❸ ratio of **new/old**



Optimizations bring $\sim 15\%$ improvement on $N(\tau_{\text{true}})/N(\tau_{\text{fake}})$ ratio

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Why the E_τ measurement is tricky ?

Challenges of τ energy calibration from $Z \rightarrow \tau\tau$:

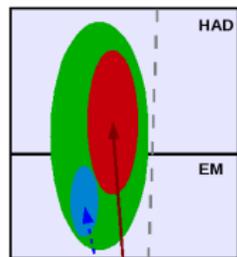
- Z peak suffers from escaping ν 's energy : broad and shifted,
- low statistics because of $\text{BR}_{\tau \rightarrow X} : o(10^3)$ vs $o(10^5)$ for $Z \rightarrow ee$.

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Alternative approach :



$$E^{\text{true}} = R_e E_{\pi^0}^{\text{meas}} + R_\pi E_{\pi^\pm}^{\text{meas}}$$

But :

1. calo not compensated : $R_e \neq R_\pi$
2. Can't separate the two showers

$$E^{\text{meas}} \equiv E_{\pi^0}^{\text{meas}} + E_{\pi^\pm}^{\text{meas}}$$

is the **only** measured observable

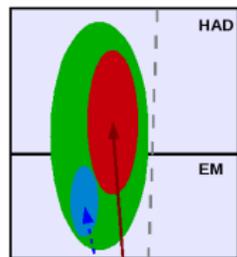
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Energy scale

- $R_e \Rightarrow Z \rightarrow ee$ calib
- $R_\pi \Rightarrow ?$
- how to use R_e, R_π ?

Energy resolution

- tracker ?
- calorimeter ?
- both ?

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Absolute correction

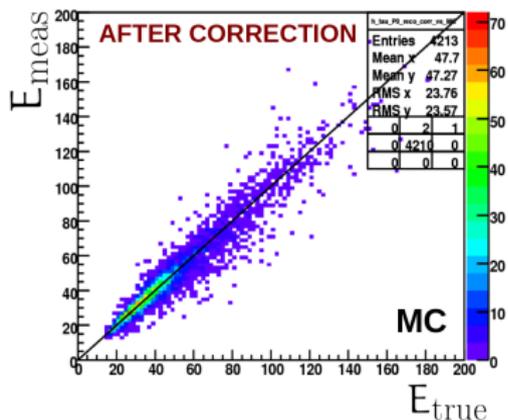
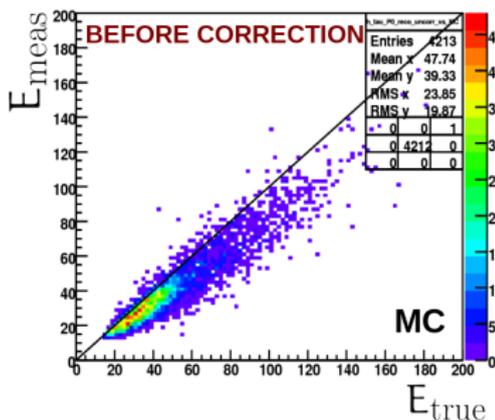
Strategy : measure the π^\pm with the tracker and π^0 with the calorimeter.
To avoid double counting, the **average** calorimeter π^\pm energy $\langle E_{\text{cal}}^{\pi^\pm} \rangle$ is subtracted (for type 2 and 3) :

$$E_{\text{corr}} \equiv E_{\text{trk}} + E_{\text{cal}} - \underbrace{\langle R_{\pi}^{\text{data}/\text{mc}}(E_{\text{trk}}, \eta) \rangle}_{\langle E_{\text{cal}}^{\pi^\pm} \rangle} \cdot E_{\text{trk}}$$

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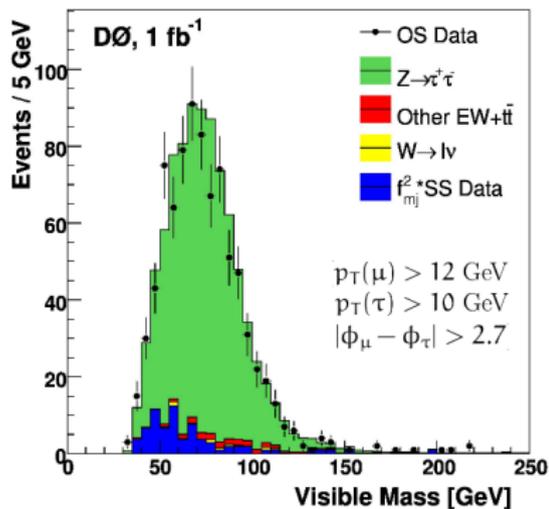
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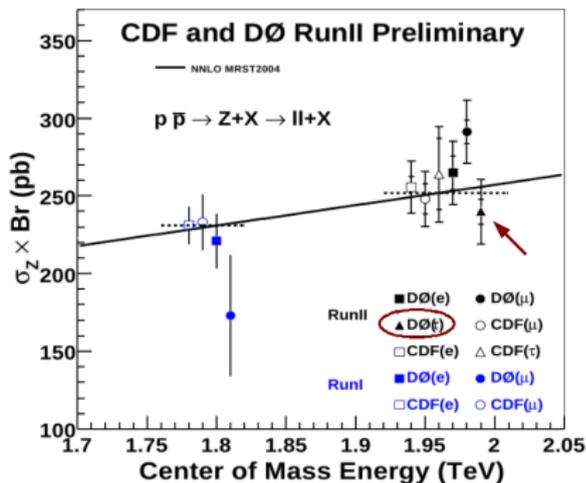
Comment : for type 1 (one track and one Cal Clu) \Leftrightarrow particle flow

$\sigma(p\bar{p} \rightarrow Z \rightarrow \tau\tau)$ measurement

$\mu\tau_{\text{had}}$ final state using this energy correction method :



$$M_{\text{vis}}^2 \equiv (p^\tau + p^\mu + mET)^2$$



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Relative correction

Motivation : the *in situ* measurement of R_π **for data in hadronic environment** is difficult (requiring isolated pions). Track energy propagation can be done differently.

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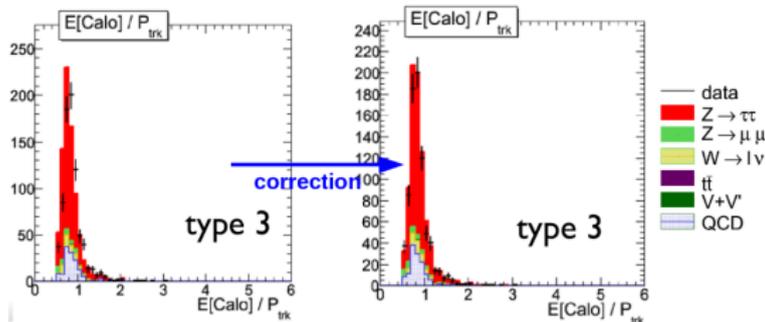
Correction method :

Use the track energy as reference to correct simulation event by event :

$$\left(\frac{E}{p}\right)_{\text{MCcorr}} = \left(\frac{E}{p}\right)_{\text{MC}} \times \frac{\langle E/p \rangle_{\text{data}[Z \rightarrow \tau\tau]}}{\langle E/p \rangle_{\text{MC}}}$$

with

- $E/p \equiv E^{\text{calo}}/p^{\text{trk}}$
- $\langle E/p \rangle \equiv$ average value.

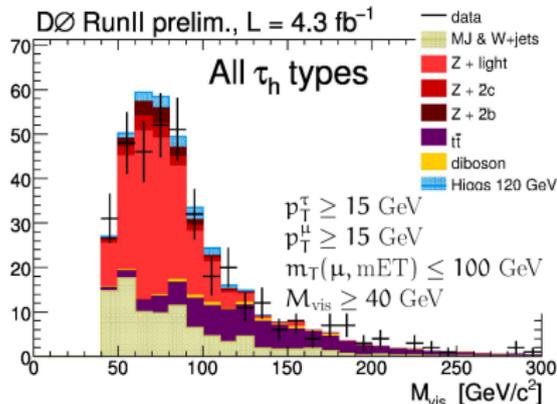


Higgs searches

Higgs searches using this energy correction method :

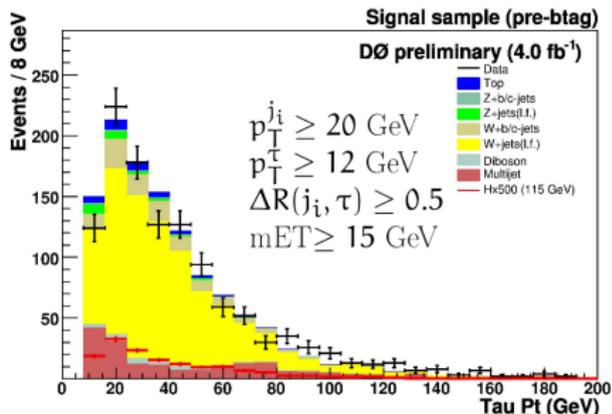
SUSY neutral Higgs ϕ

$p\bar{p} \rightarrow b\phi \rightarrow \tau\tau b$



SM Higgs

$p\bar{p} \rightarrow W(\rightarrow \tau\nu)H(\rightarrow \bar{b}b)$



DØ preliminary

See next talk by Tammy on Higgs boson searches in τ final state at DØ

Conclusions and outlooks

Tau identification in hadronic environment requires sophisticated algorithms.

Promising improvements at DØ using additional τ properties and kinematic dependences.

In spite of experimental challenges, τ are well understood and allow to

- test SM consistency at high energy
 - search for new phenomena
 - search for the origin of electroweak symmetry breaking
- next talk by Tammy*

BACKUP SLIDES

Minimal Supersymmetric SM and τ 's

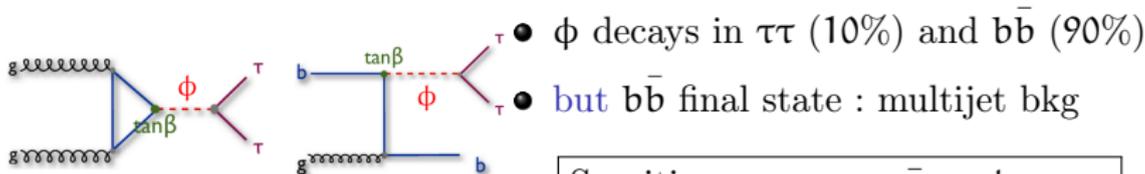
MSSM extension :

- $\tilde{q}, \tilde{g},$
 - weak gauginos, ...
- }
- cascade decays can end with τ 's

Higgs sector of MSSM After $SU(2)_L \times U(1)_Y$ symmetry breaking :

- ❶ 3 neutral Higgs fields $\phi \equiv (H^0, h, A),$
- ❷ 2 charged Higgs fields $H^+, H^-.$

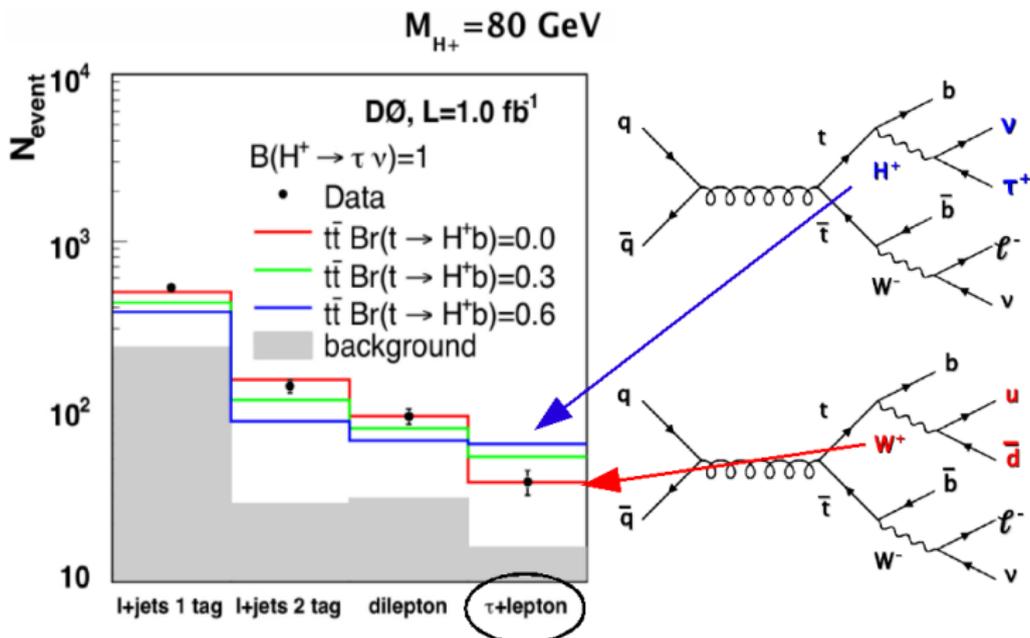
For the neutral Higgs search :



Sensitive process : $p\bar{p} \rightarrow \phi \rightarrow \tau\tau$

MSSM charged Higgs

Charged higgs bosons via $t\bar{t}$ events



R_π measurement in data

Strategy : fit an improved MC in a restricted region where R_π is measured in data. Trust MC for extrapolation.

