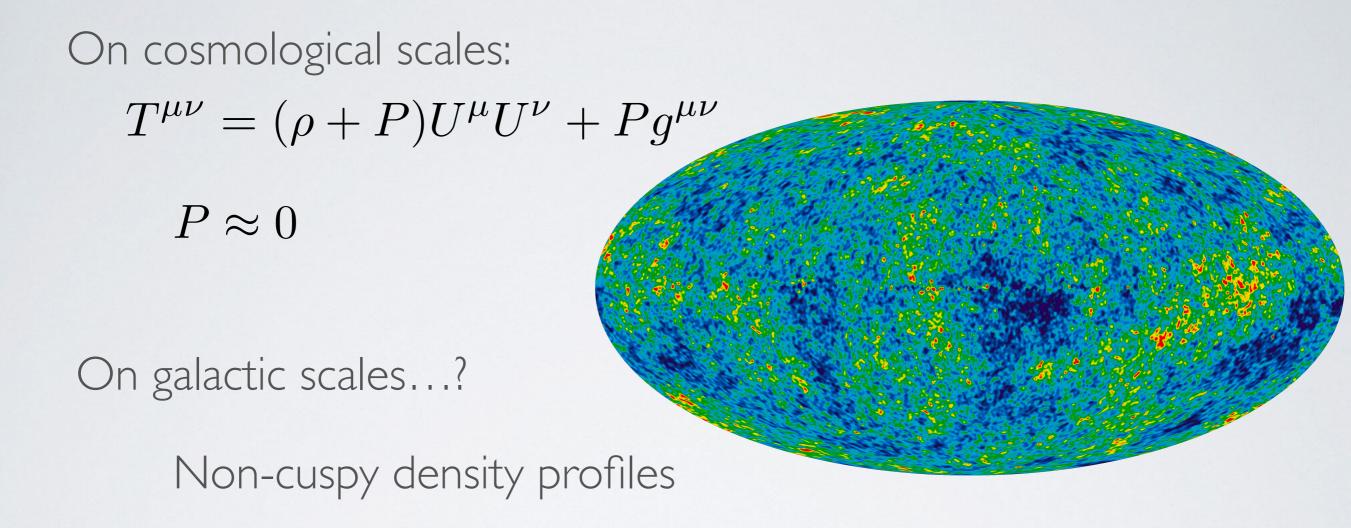
DYNAMICAL FRICTION IN A SUPERFLUID

Benjamin Elder University of Nottingham

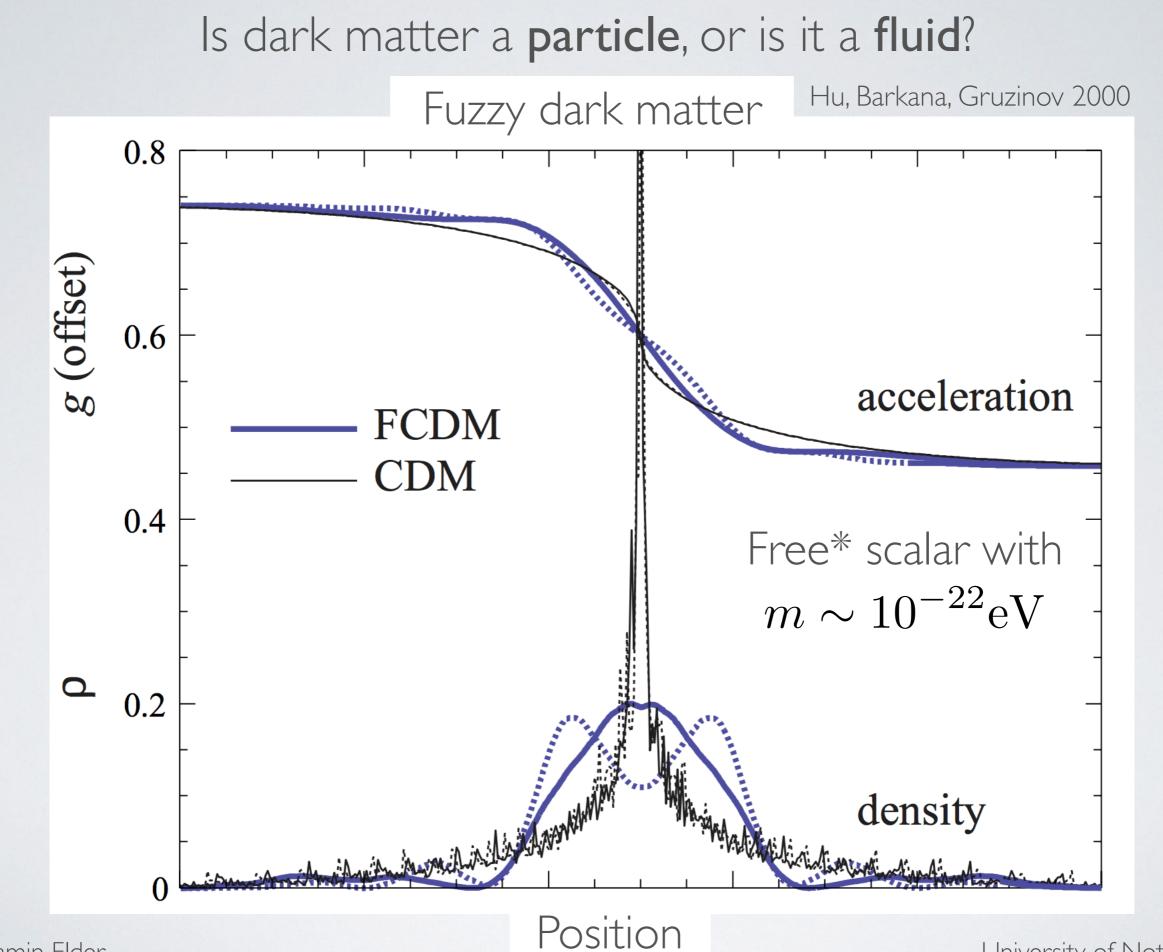
with Lasha Berezhiani & Justin Khoury arXiv:1905.09297

PASCOS 2019

Is dark matter a **particle**, or is it a **fluid**?



Baryonic mass correlated to circular velocity

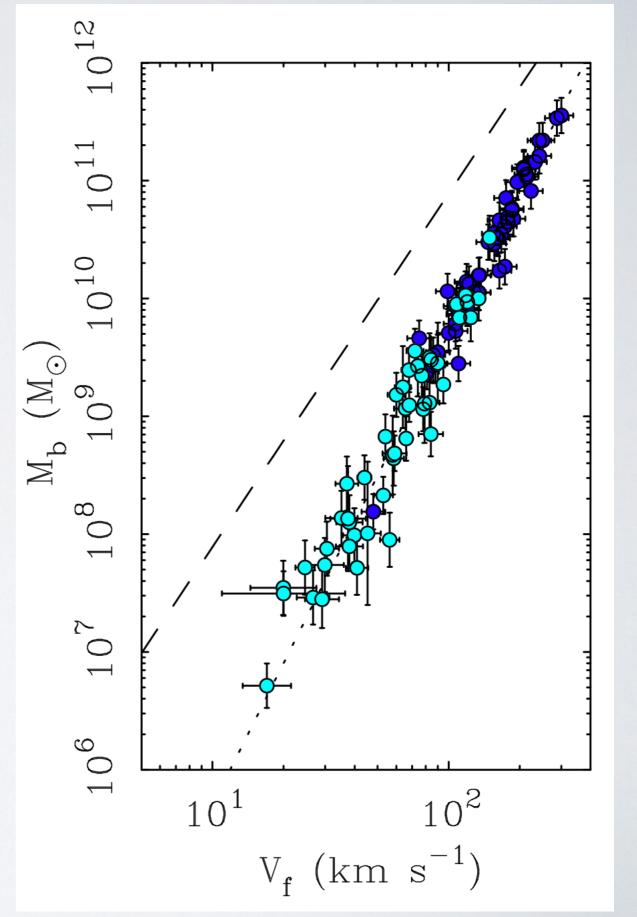


Is dark matter a **particle**, or is it a **fluid**?

Superfluid dark matter Berezhiani & Khoury 2015

DM forms a superfluid condensate around galaxies

Phonons mediate a fifth force that reproduces the Baryonic Tully-Fisher Relation



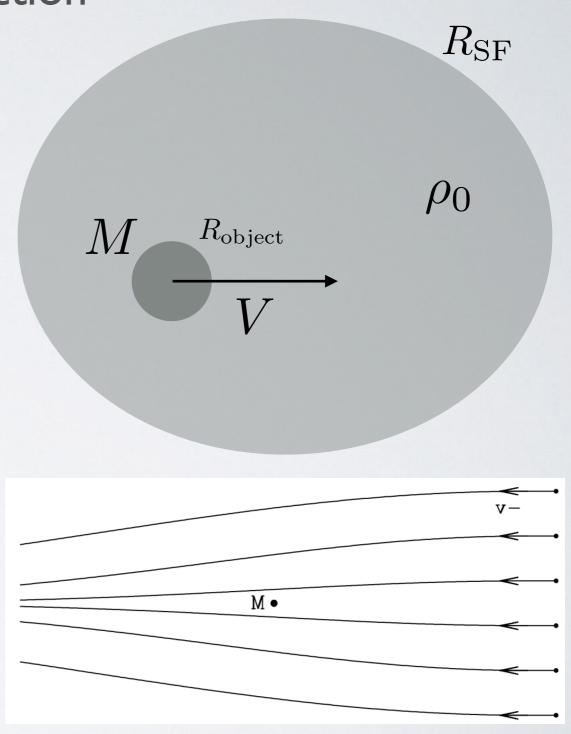
How to distinguish particle dark matter from dark matter as a fluid?

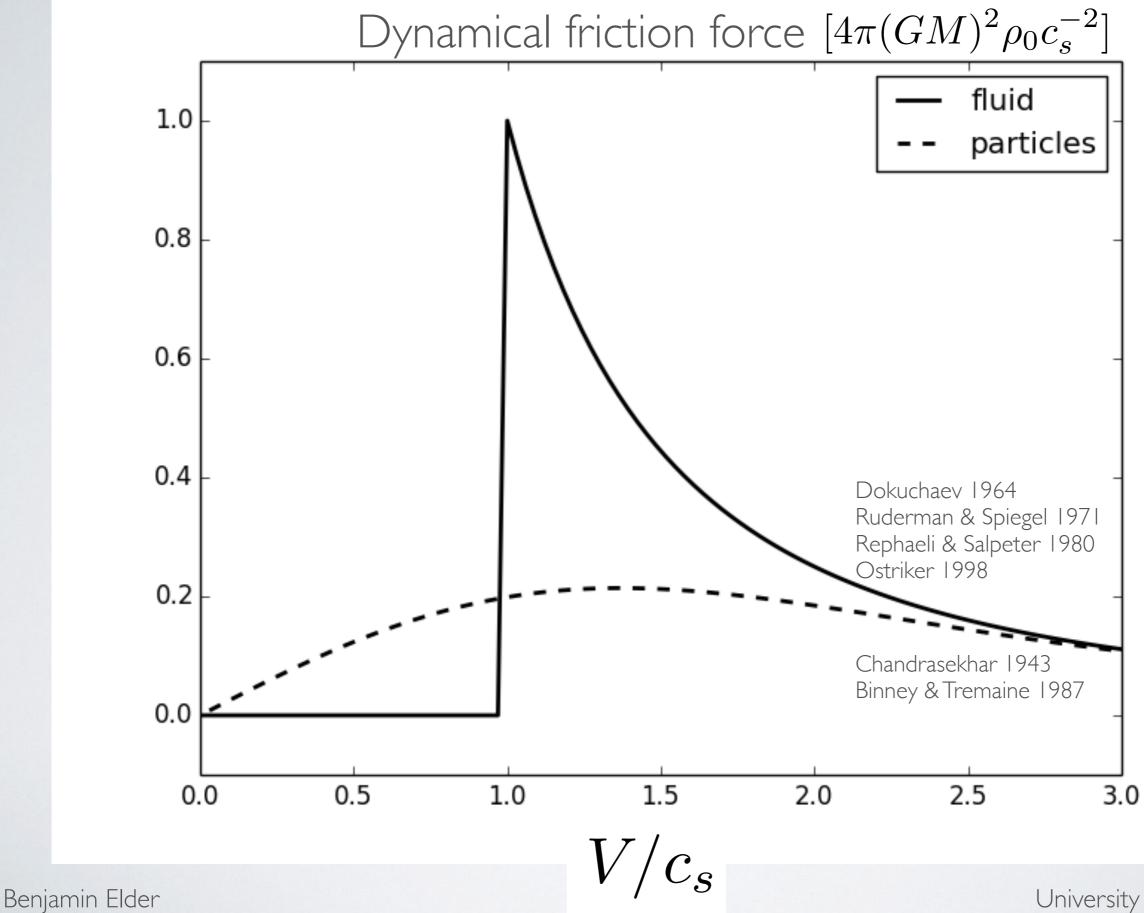
Dynamical friction

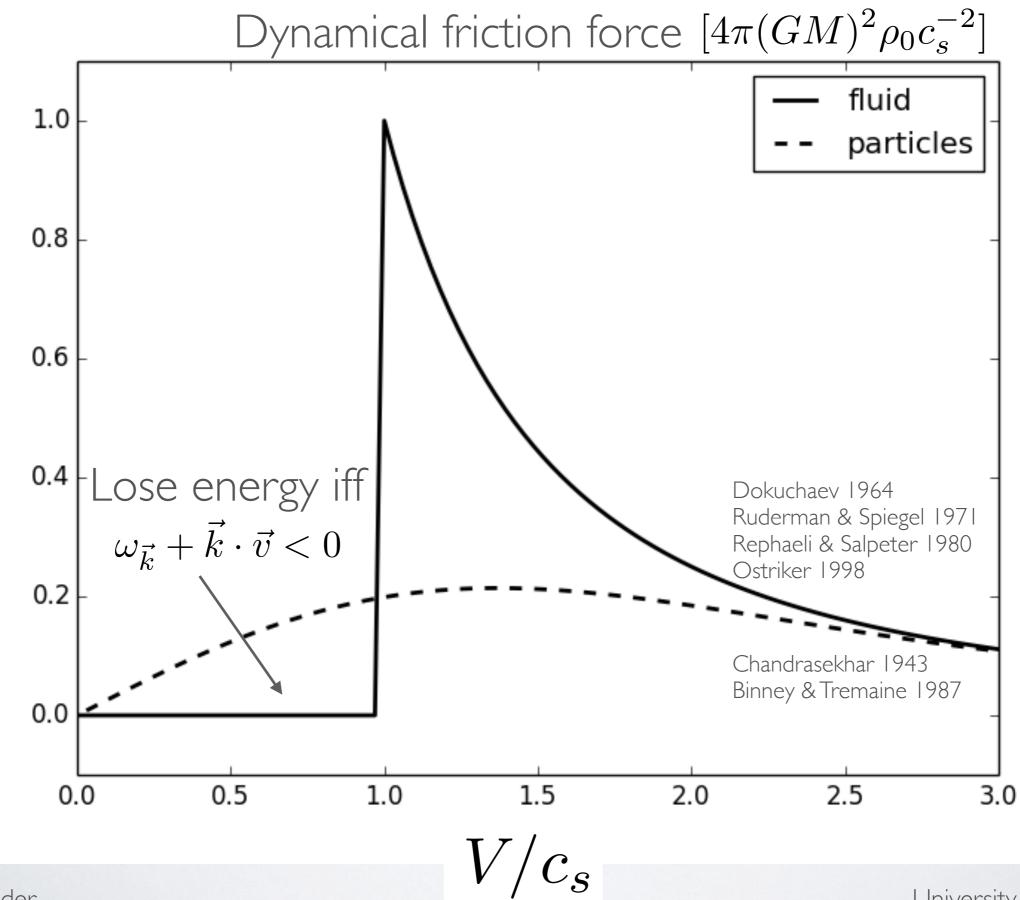
(i) A perturber moves through a uniform cloud of particles

(ii) An overdensity builds up behind the perturber

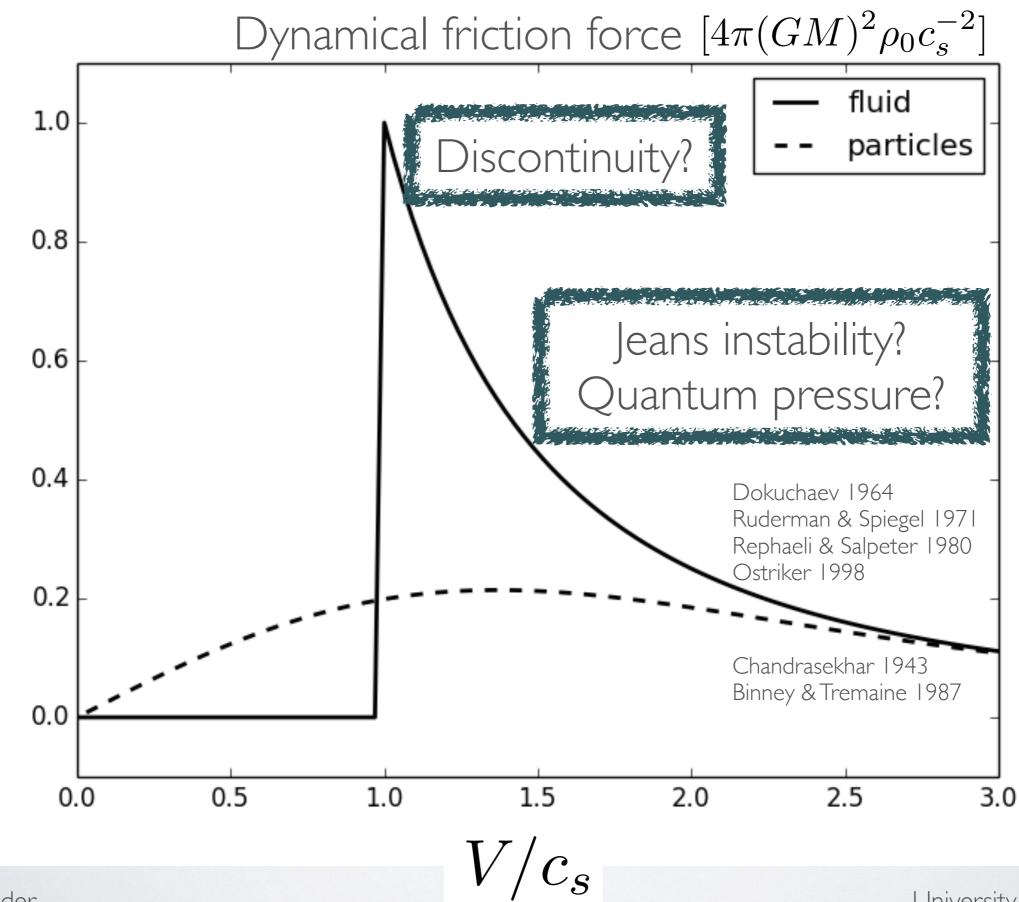
(iii) The gravitational attraction of the overdensity slows the perturber







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Hydrodynamical description of superfluid dynamical friction

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{16\pi G} - |\partial \Psi|^2 - m^2 |\Psi|^2 - \frac{\lambda}{2} |\Psi|^4 \right)$$

Non-relativistic limit:

$$\Psi = \frac{\psi}{\sqrt{2m}} \,\mathrm{e}^{-\mathrm{i}mt}$$

Cast GP into hydro form:

 $\mathrm{i}\partial_t\psi=-\frac{\Delta}{2m}\psi+\frac{\lambda}{4m^2}|\psi|^2\psi+m\Phi\psi$

$$\psi = \sqrt{\frac{\rho}{m}} \mathrm{e}^{\mathrm{i}\theta}$$
 $P = \frac{\lambda}{8m^4} \rho^2$ $\vec{v} = \frac{\vec{\nabla}\theta}{m}$

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) &= 0; \\ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} &= -\frac{1}{\rho} \vec{\nabla} P - \vec{\nabla} \Phi + \frac{1}{2m^2} \vec{\nabla} \left(\frac{\Delta \sqrt{\rho}}{\sqrt{\rho}} \right) \end{aligned}$$

Hydrodynamical description of superfluid dynamical friction

$$\begin{aligned} \frac{\partial \rho}{\partial t} &+ \vec{\nabla} \cdot (\rho \vec{v}) &= 0; \\ \frac{\partial \vec{v}}{\partial t} &+ (\vec{v} \cdot \vec{\nabla}) \vec{v} &= -\frac{1}{\rho} \vec{\nabla} P - \vec{\nabla} \Phi + \frac{1}{2m^2} \vec{\nabla} \left(\frac{\Delta \sqrt{\rho}}{\sqrt{\rho}} \right) \end{aligned}$$

Combine:

$$\ddot{\alpha} - c_s^2 \Delta \alpha - m_g^2 \alpha + \frac{1}{4m^2} \Delta^2 \alpha = 4\pi G \rho_{\text{ext}}$$
Tachyonic mass
$$m_g^2 \equiv 4\pi G \rho_0$$
Quantum pressure
$$Perturber sourcing overdensity$$

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Hydrodynamical description of superfluid dynamical friction

$$\ddot{\alpha} - c_s^2 \Delta \alpha - m_g^2 \alpha + \frac{1}{4m^2} \Delta^2 \alpha = 4\pi G \rho_{\text{ext}} \qquad \rho_{\text{ext}}(x) = M\delta(x)\delta(y)\delta(z - Vt)$$

Dynamical friction is the work done by gravity: $F = \frac{M}{V} \dot{\phi}_{\alpha}$

$$\phi_{\alpha}(k_0, \vec{k}) = -\frac{4\pi G \rho_0}{\vec{k}^2} \alpha(k_0, \vec{k}) \qquad \alpha(k_0, \vec{k}) = -\frac{4\pi G \rho_{\text{ext}}(k)}{k_0^2 - \omega_k^2}$$

$$F = \frac{M}{V} \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \,\mathrm{e}^{\mathrm{i}k_0 t - \mathrm{i}\vec{k}\cdot\vec{x}} \,\mathrm{i}k_0 \,\frac{4\pi G\rho_0}{\vec{k}^2} \,\frac{4\pi G\rho_{\mathrm{ext}}(k)}{k_0^2 - \omega_k^2}$$

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 $R_{\rm SF}$

 ρ_0

Hydrodynamical description of superfluid dynamical friction Low-momentum cutoffs

Cutoff I: size of superfluid cloud

 $k_{\rm min} = 2\pi R_{\rm SF}^{-1}$

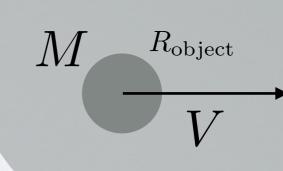
Cutoff II: Jeans scale

$$k_{\min} = k_{\rm J}$$

$$\omega_{k_{\rm J}}^2 \equiv 0 = -m_g^2 + c_s^2 k_J^2 + \frac{k_J^4}{4m^2}$$

Use whichever cutoff whose scale is higher:

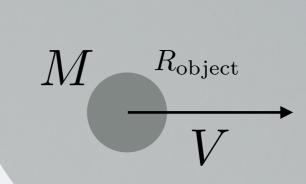
$$k_{\min} = \max\left(2\pi R_{\rm SF}^{-1}, k_{\rm J}\right)$$



Hydrodynamical description of superfluid dynamical friction High-momentum cutoffs $R_{\rm SF}$

(or linearization breaks down)

 $k_{\rm max} = 2\pi R_{\rm object}^{-1}$



"Cutoff" II: integrand is non-zero

$$F = -\frac{4\pi G^2 M^2 \rho_0}{V^2} \int \frac{\mathrm{d}k}{k} \,\mathrm{d}\cos\theta \,\mathrm{e}^{\mathrm{i}\omega_k t - \mathrm{i}kVt\cos\theta} \delta\left(\frac{\omega_k}{kV} - \cos\theta\right)$$

Integrand is only non-zero for modes $k < k_{\star}$

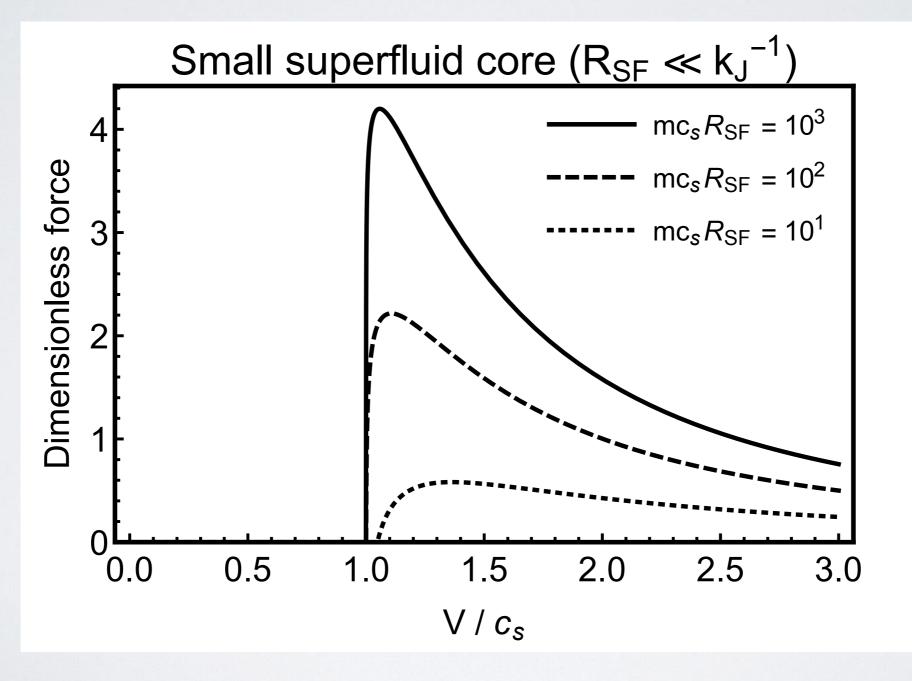
$$k_{\max} = \min\left(2\pi R_{\text{object}}^{-1}, k_{\star}\right)$$

 ρ_0

 $\frac{\omega_{k_{\star}}}{k_{\perp}V} \equiv 1$

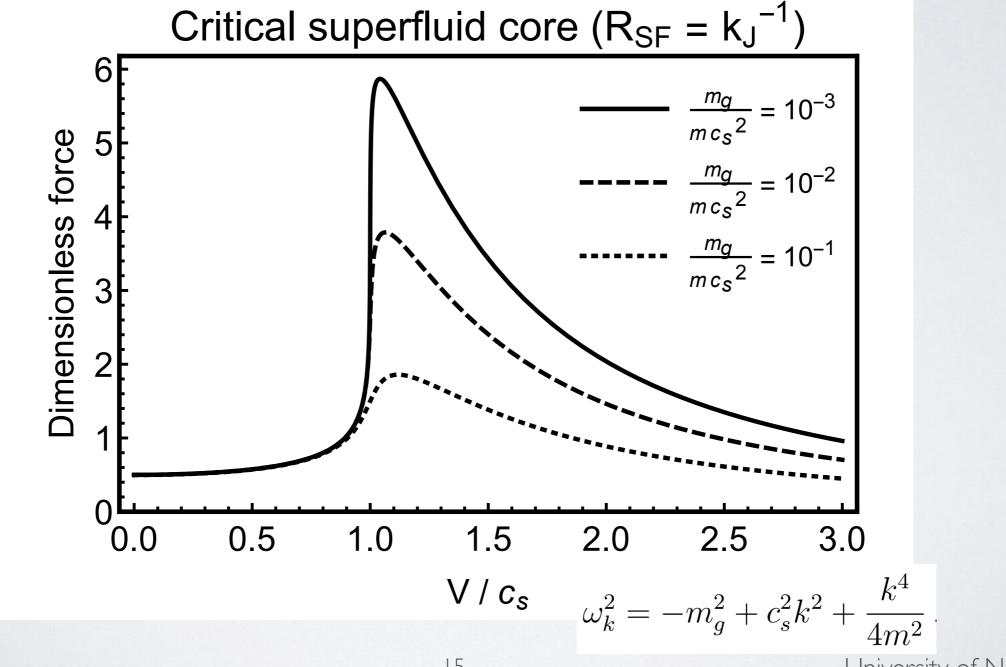
Results I: small superfluid core

$$|F| = \frac{4\pi G^2 M^2 \rho_0}{V^2} \ln\left(\frac{\min\left(2\pi R_{\rm object}^{-1}, k_\star\right)}{\max\left(2\pi R_{\rm SF}^{-1}, k_{\rm J}\right)}\right)$$



Results II: large superfluid core

$$|F| = \frac{4\pi G^2 M^2 \rho_0}{V^2} \ln\left(\frac{\min\left(2\pi R_{\text{object}}^{-1}, k_\star\right)}{\max\left(2\pi R_{\text{SF}}^{-1}, k_{\text{J}}\right)}\right)$$



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Quasiparticle description of superfluid dynamical friction

$$S = \int d^{4}x \sqrt{-g} \left(\frac{1}{16\pi G} R - |\partial \Psi^{2}| - m^{2}|\Psi|^{2} - \frac{\lambda}{2}|\Psi|^{4} - \frac{1}{2}(\partial \chi)^{2} - \frac{1}{2}M^{2}\chi^{2} \right)$$
Superfluid
$$Perturber$$

$$\Psi$$

$$R_{SF}$$

$$\chi$$

$$\chi$$

$$\chi$$

Quasiparticle description of superfluid dynamical friction

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{16\pi G} R - |\partial \Psi^2| - m^2 |\Psi|^2 - \frac{\lambda}{2} |\Psi|^4 - \frac{1}{2} (\partial \chi)^2 - \frac{1}{2} M^2 \chi^2 \right)$$

Superfluid Perturber
Perturb the condensate:

$$\Psi = (v + h) e^{i\sqrt{m^2 + \lambda v^2}t + i\pi}$$

Effective theory of phonons:

$$\mathcal{L} = \frac{1}{2}\dot{\pi}^2 + \frac{1}{2}\pi \left(m_g^2 + c_s^2 \Delta - \frac{\Delta^2}{4m^2} \right) \pi - 4\pi G M^2 \sqrt{\rho_0} \chi^2 \frac{1}{\sqrt{\Delta} (m_g + c_s^2 \Delta - \frac{\Delta^2}{4m^2})} \dot{\pi}}$$

Rate of energy loss: $|F| = \frac{\dot{E}}{V} = \int \omega_k \, d\Gamma \longrightarrow |F| = \frac{4\pi G^2 M^2 \rho_0}{V^2} \ln \left(\frac{\min(2\pi R_{object}^{-1}, k_j)}{\max(2\pi R_{obj}^{-1}, k_j)} \right)$

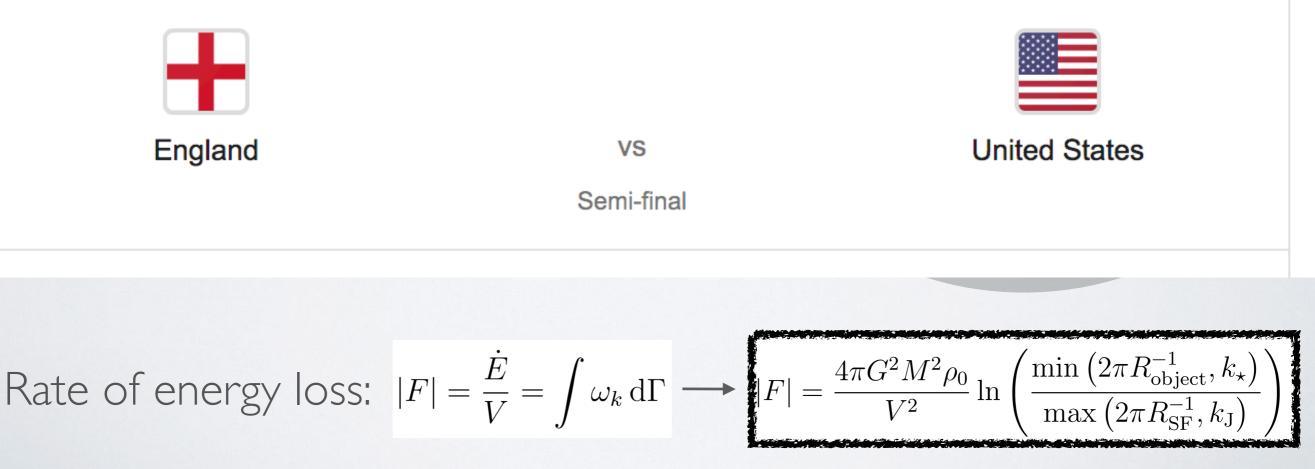
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Quasiparticle description of superfluid dynamical friction

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{16\pi G} R - |\partial \Psi^2| - m^2 |\Psi|^2 - \frac{\lambda}{2} |\Psi|^4 - \frac{1}{2} (\partial \chi)^2 - \frac{1}{2} M^2 \chi^2 \right)$$

England vs United States

Women's World Cup · Today, 20:00



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