

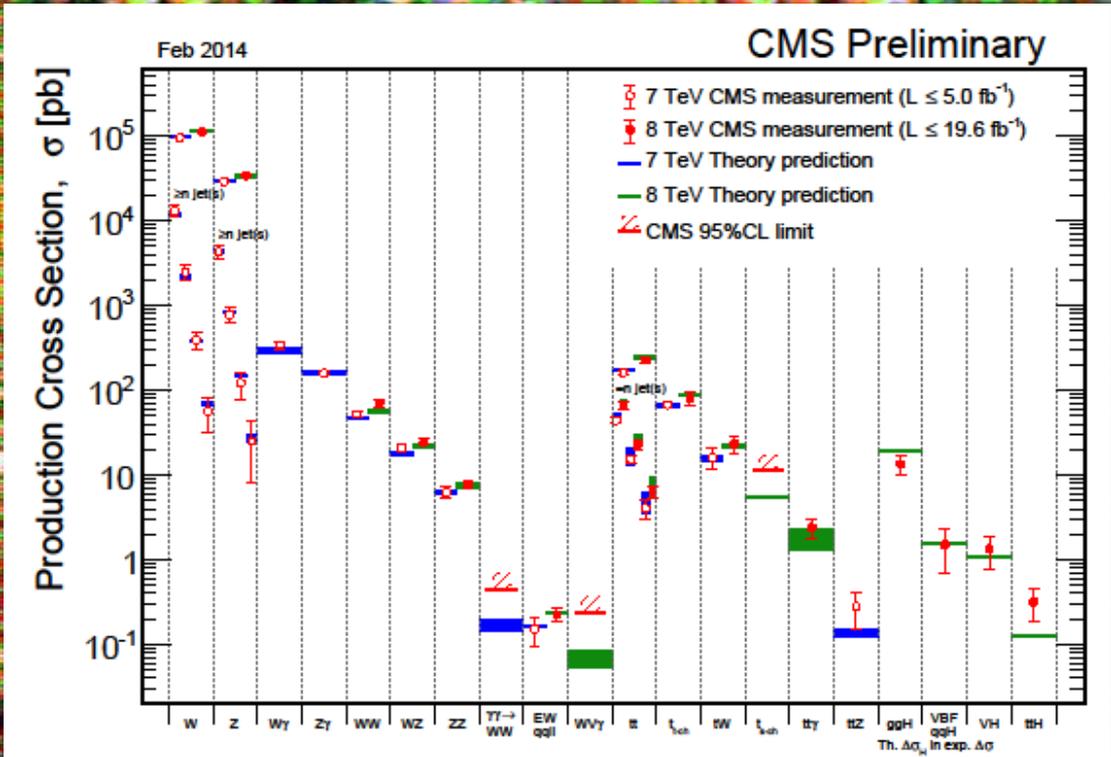
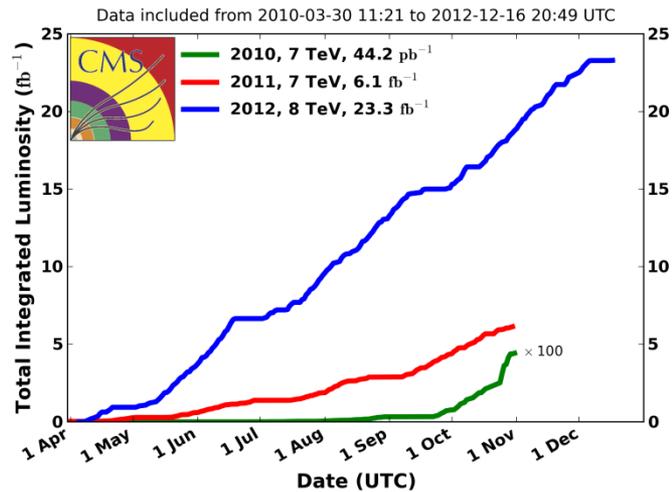


Higgs results from CMS

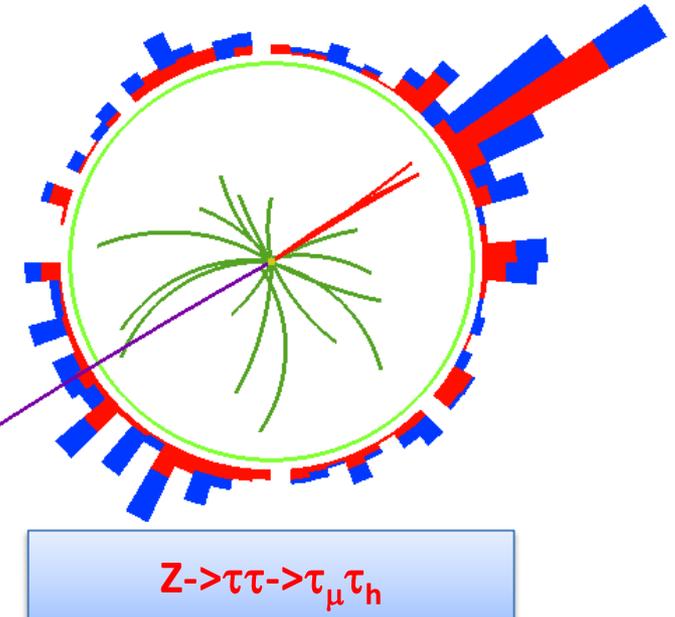
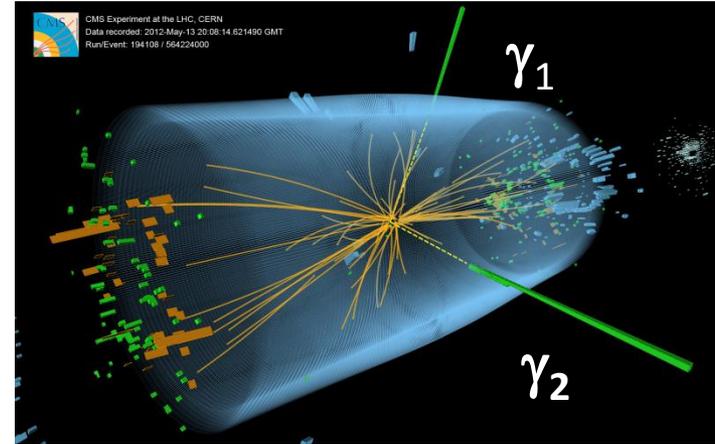
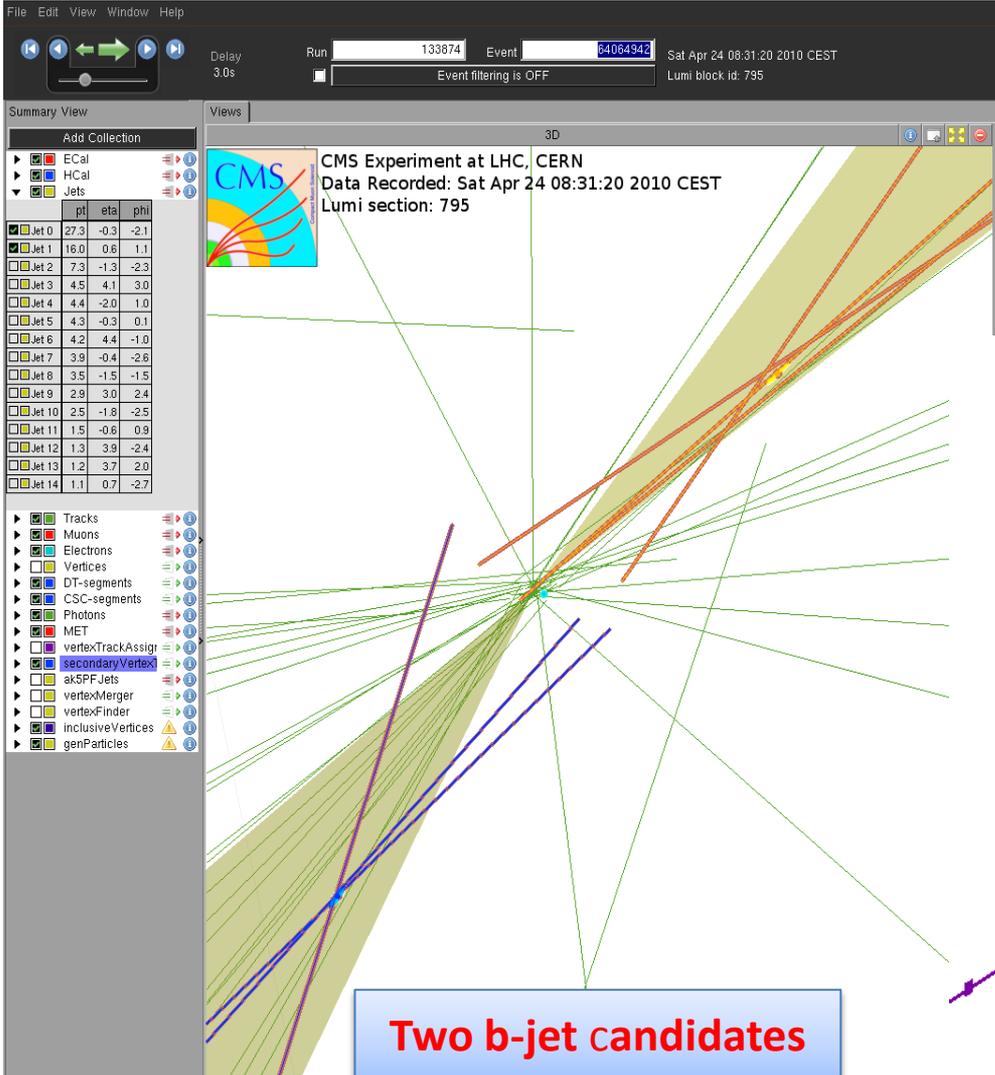


**A. Nikitenko, Imperial College, London
SUSY 2014, 22 July.**

CMS Integrated Luminosity, pp

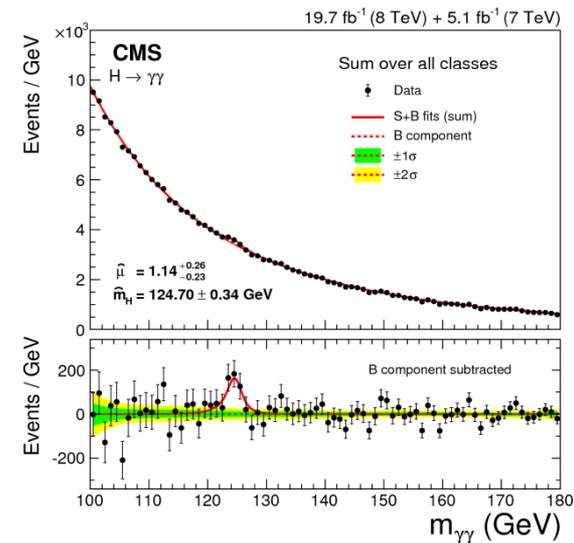
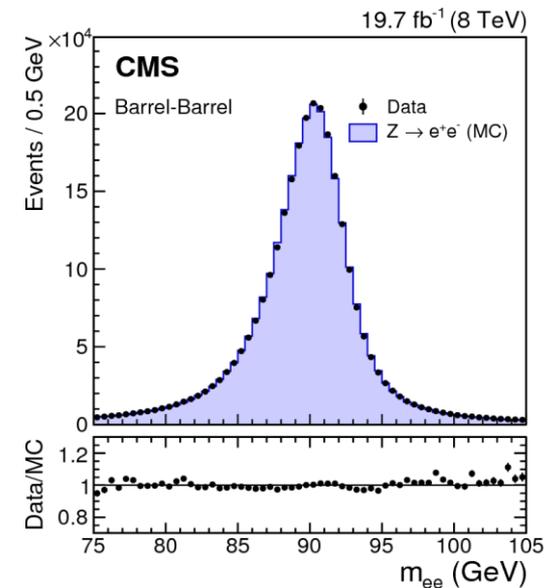


Few words about CMS reconstruction



Reconstruction: *photons*

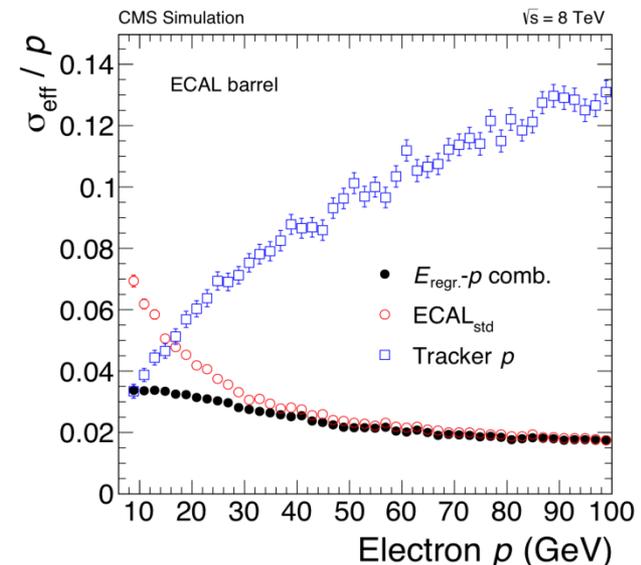
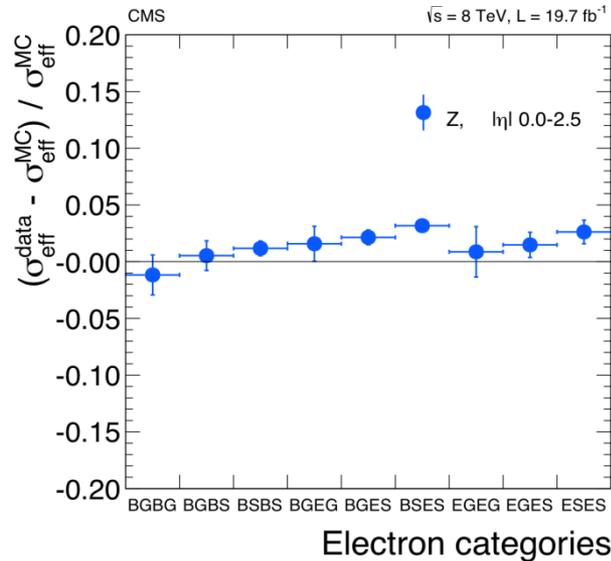
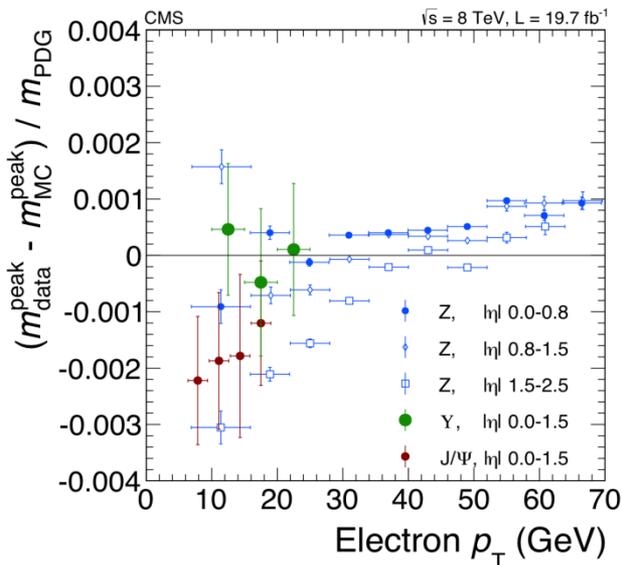
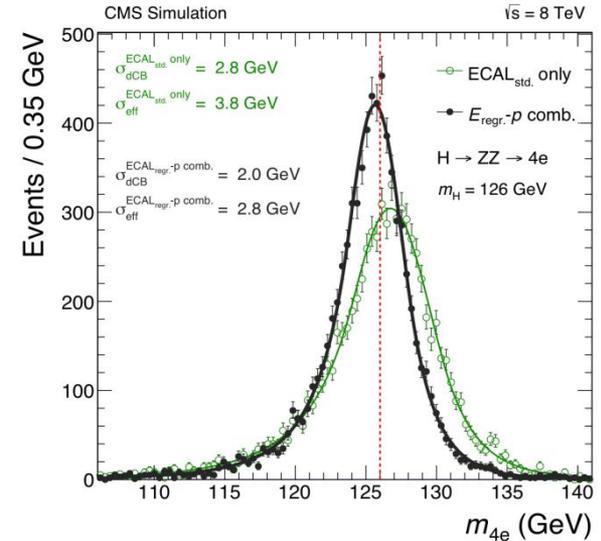
- variation of crystal transparency is monitored every 40 min with laser system
- energy scale is obtained with MC photons using multivariate regression technique
 - input: cluster shape variables, E/H , (η, ϕ) , N_{vtx} , PU density
- Finally, reconstructing e 's as γ 's in $Z \rightarrow ee$ events obtain:
 - corrections of energy in data to reproduce MC m_{ee} peak position
 - smearing of MC energy to reproduce m_{ee} resolution
 - take into account imperfect modelling of e/γ difference in MC:
 - Imperfect description of tracker material
 - Non uniformity of light collection
 - G4 model



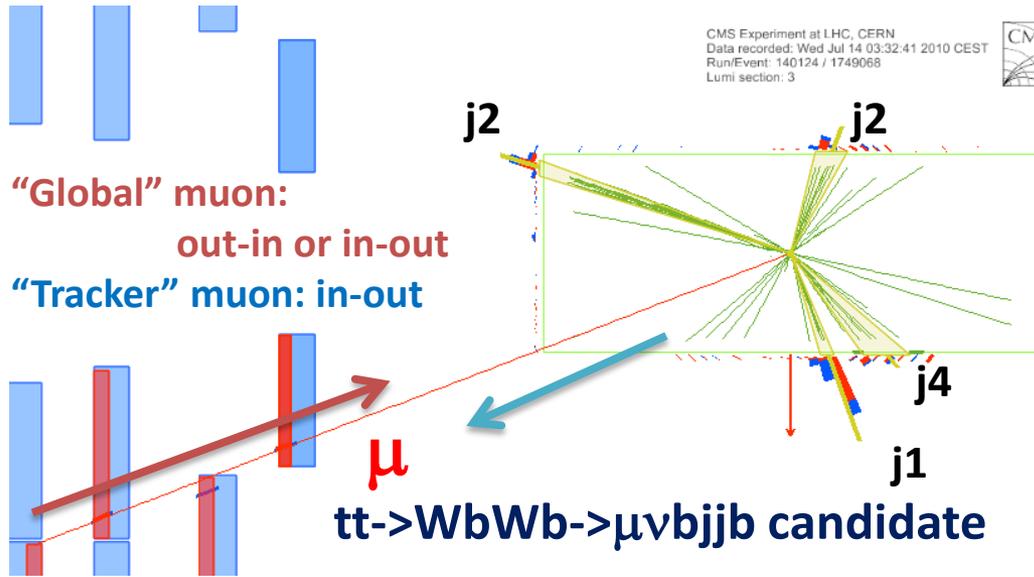
$$\sigma_{\text{eff}}(m_{\gamma\gamma}) = [1.0-2.6] \text{ GeV}$$

Reconstruction: *electrons*

- Momentum scale is obtained with MC electrons:
 - combined ECAL + tracker information using MVA regression technique
- Correct discrepancy between data and simulation on the momentum scale and resolution using $Z \rightarrow ee$ events
 - dominant sources of discrepancy
 - residual tracker misalignment
 - imperfect corrections of the crystal transparency loss
- Validate with $J/\psi, \Upsilon(1S) \rightarrow ee$ for e 's $p_T < 20$ GeV



Reconstruction: muons

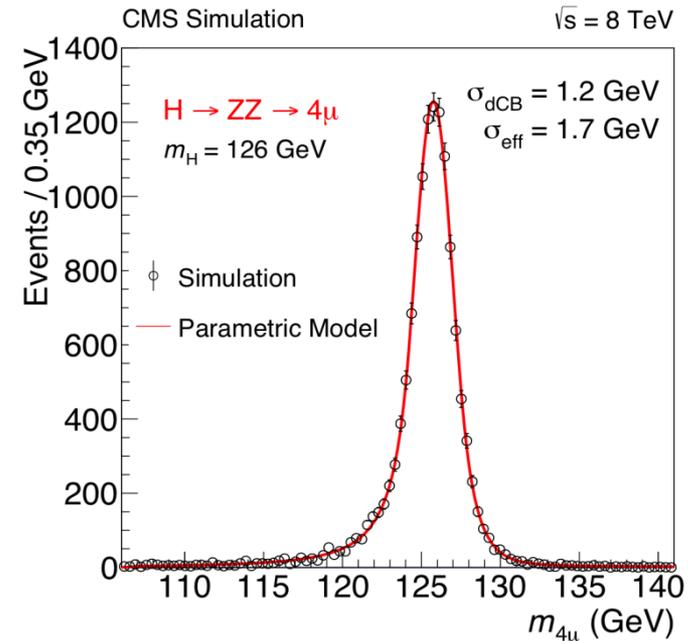
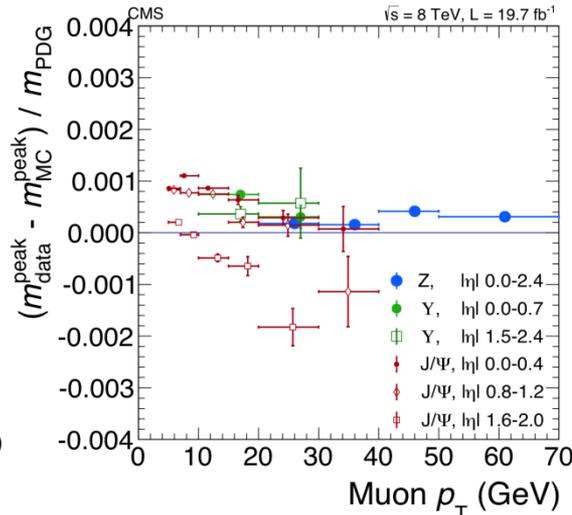
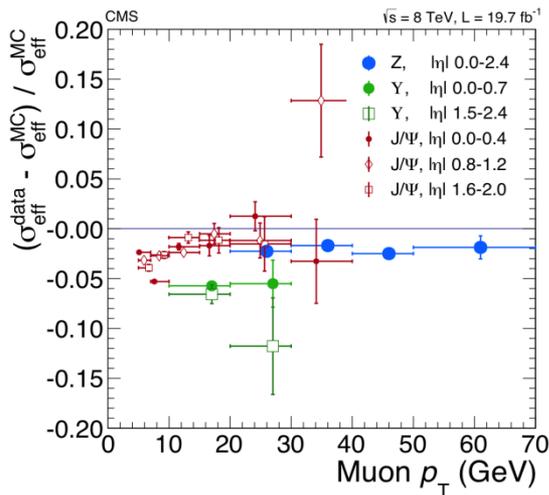


- Absolute scale and resolution is obtained with $Z \rightarrow \mu\mu$ and low mass di-muon resonances

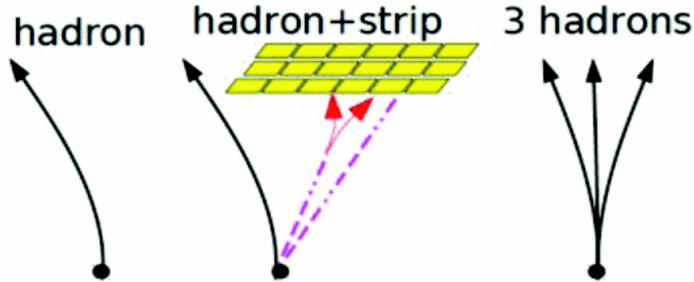
- main source of data-MC disagreement is residual tracker misalignment

p_T resolution in the momentum range relevant for $H \rightarrow ZZ \rightarrow 4\ell$ analysis:

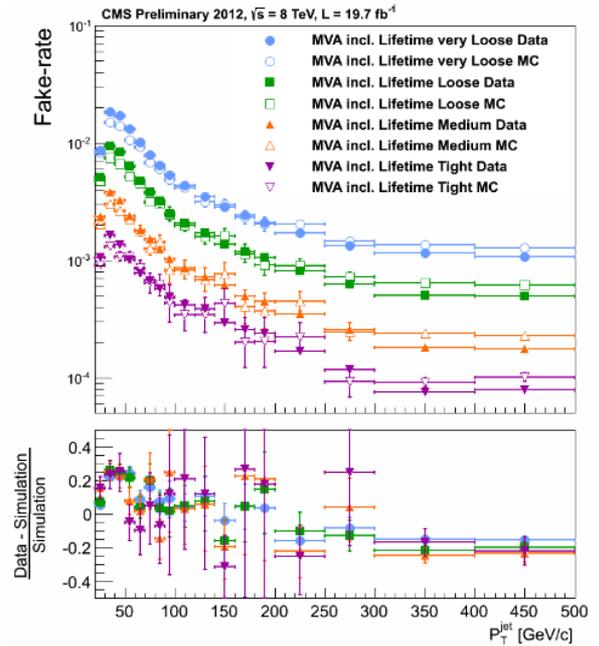
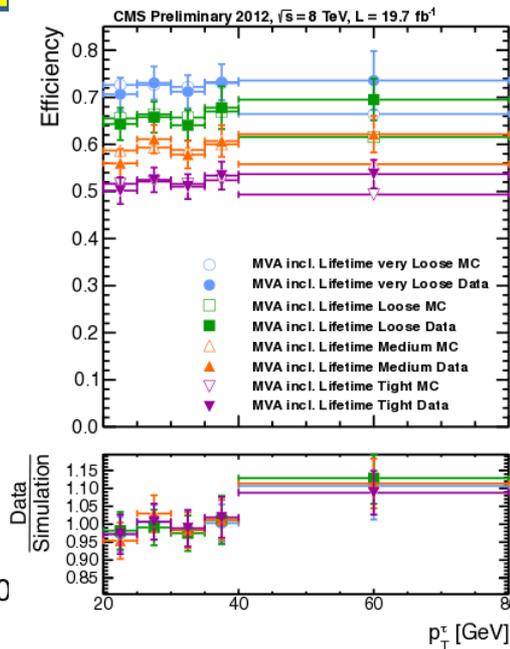
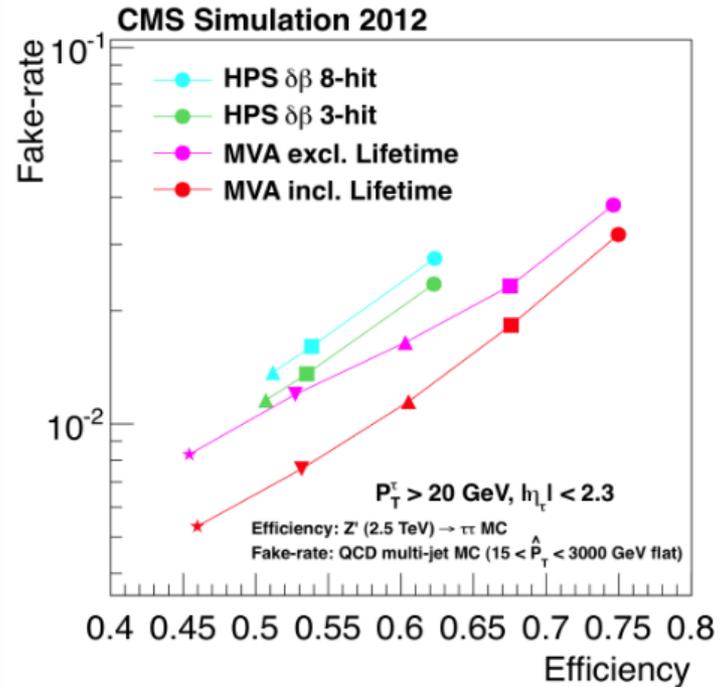
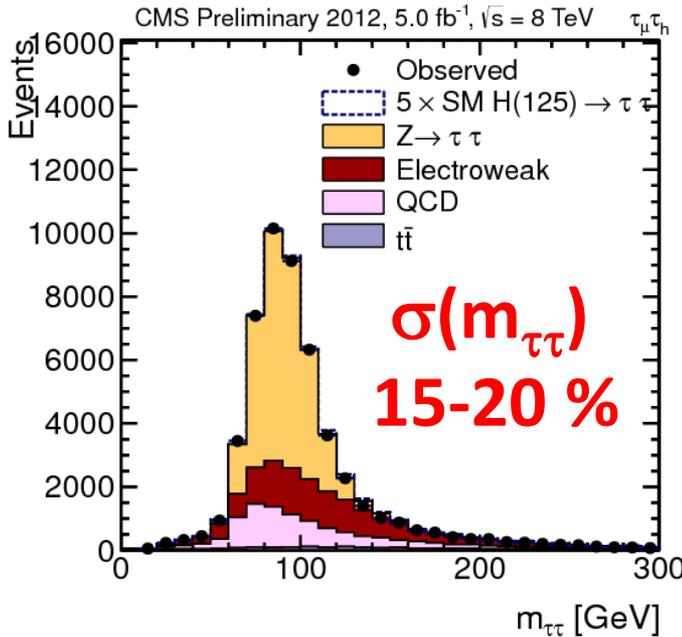
- 1.3-2.0 % barrel
- 6 % endcap



Reconstructions: $\tau \rightarrow \text{hadrons} (\tau_h)$



τ ID uncertainty – 7 %
 τ JES uncertainty < 3 %

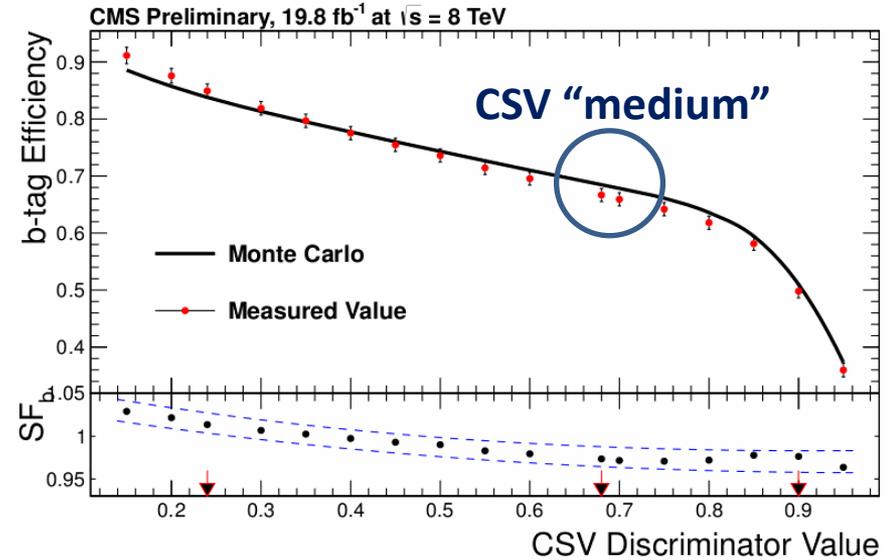


• Reconstruction: *b*-jets

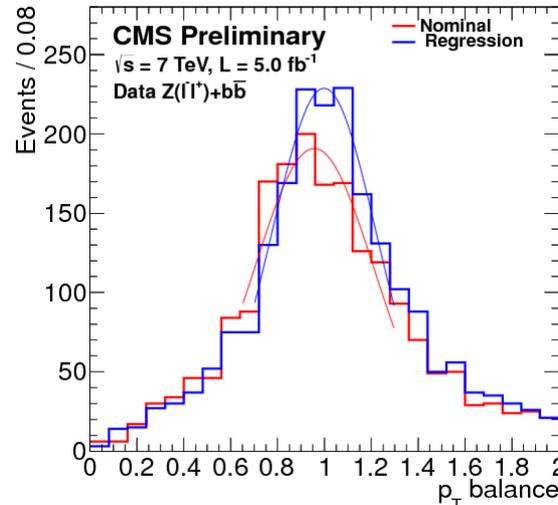
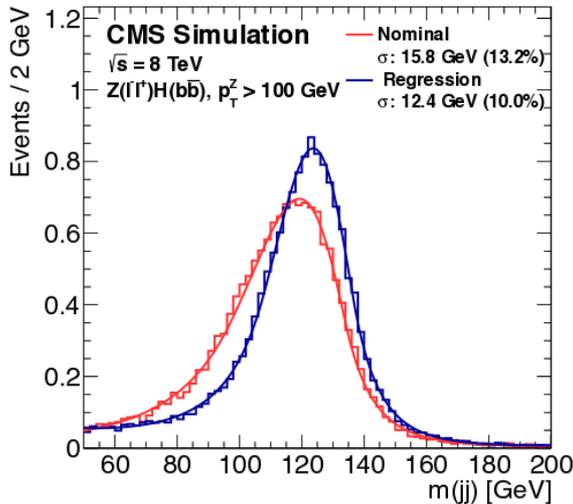
Fake rate: for $80 < p_T^j < 120$ GeV, $|\eta| < 2.4$

b tagger	misidentification probability	SF_{light}
JPL	0.0944 ± 0.0004	$1.03 \pm 0.01 \pm 0.07$
CSVL	0.0990 ± 0.0004	$1.10 \pm 0.01 \pm 0.05$
JPM	0.0105 ± 0.0002	$1.10 \pm 0.02 \pm 0.20$
CSVM	0.0142 ± 0.0002	$1.17 \pm 0.02 \pm 0.15$
TCHPI	0.0026 ± 0.0001	$1.27 \pm 0.06 \pm 0.27$
JPT	0.0013 ± 0.0001	$1.11 \pm 0.07 \pm 0.31$
CSVT	0.0016 ± 0.0001	$1.26 \pm 0.07 \pm 0.28$

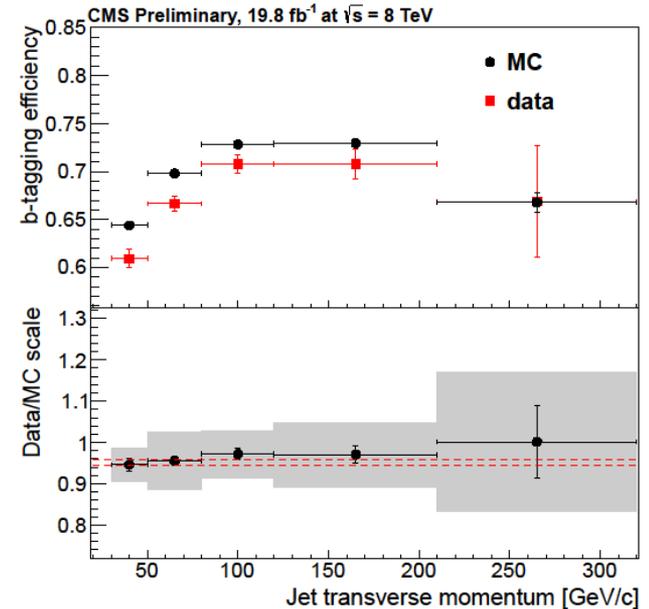
three operational points for *b*-tagging with Combined Secondary Vertex algorithm



m_{bb} resolution $\sim 10\%$

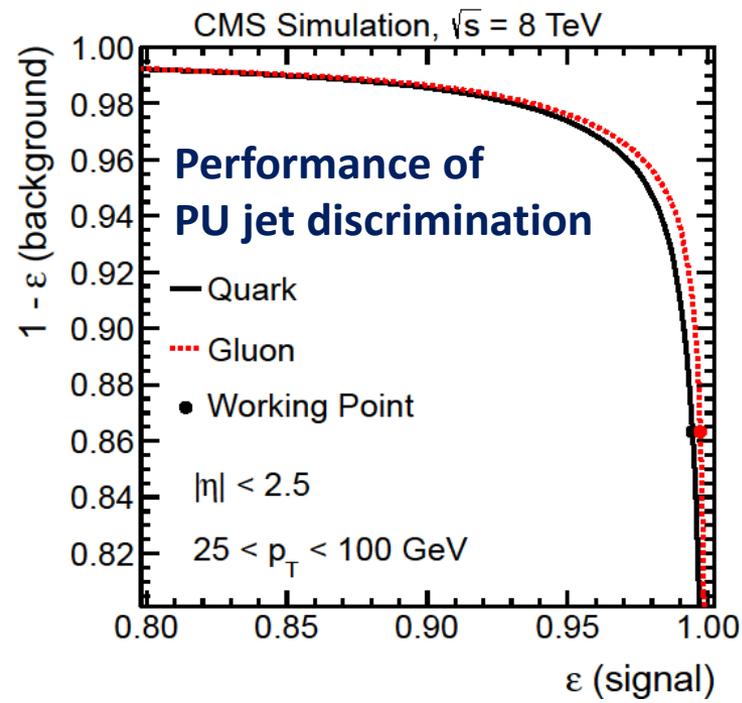
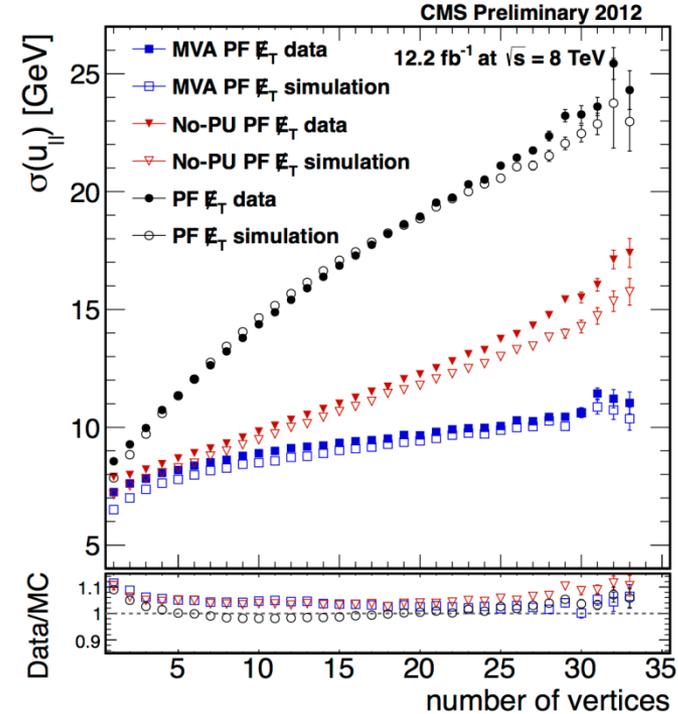
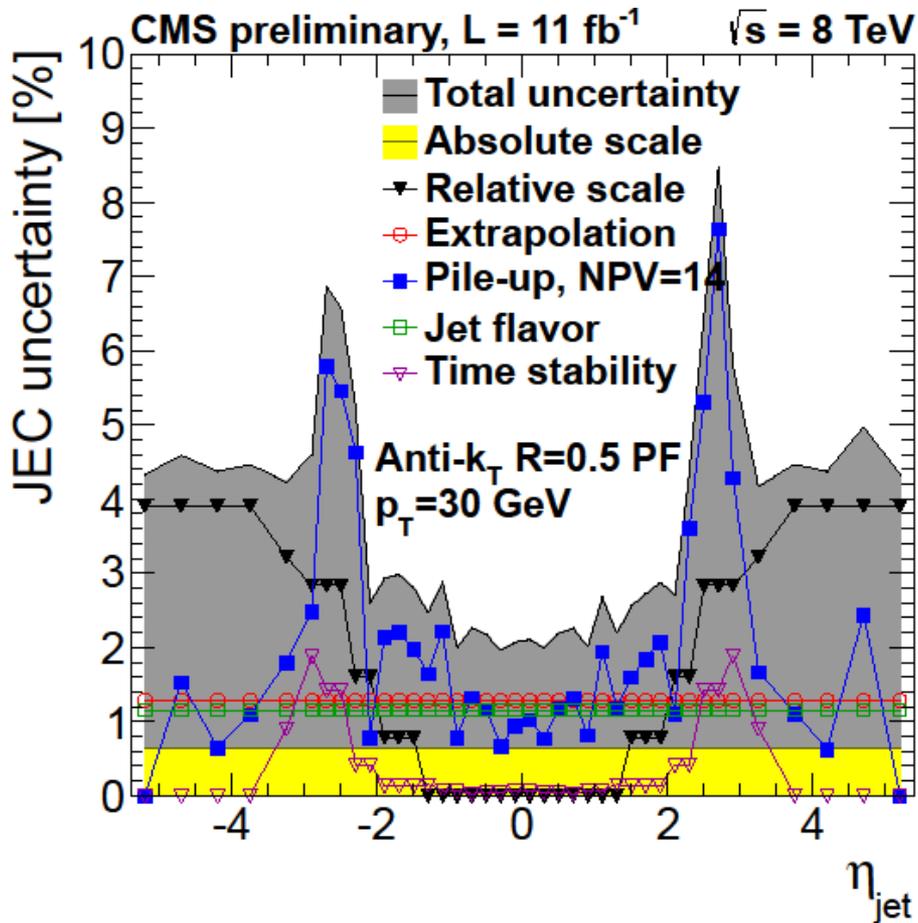


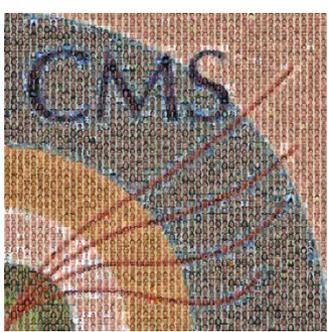
CSVM efficiency vs p_T of *b*-jets



Reconstruction: jets, E_T^{miss}

- PF (+JPT, Trk, Calo) jets - anti- k_T R=0.5
- PU jet reductions
 - most useful within the tracker acceptance
- PU mitigation algos for E_T^{miss}

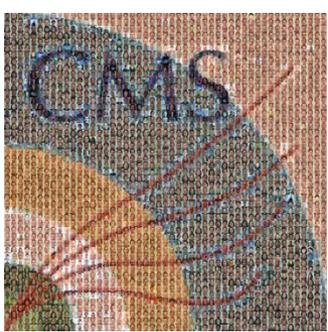




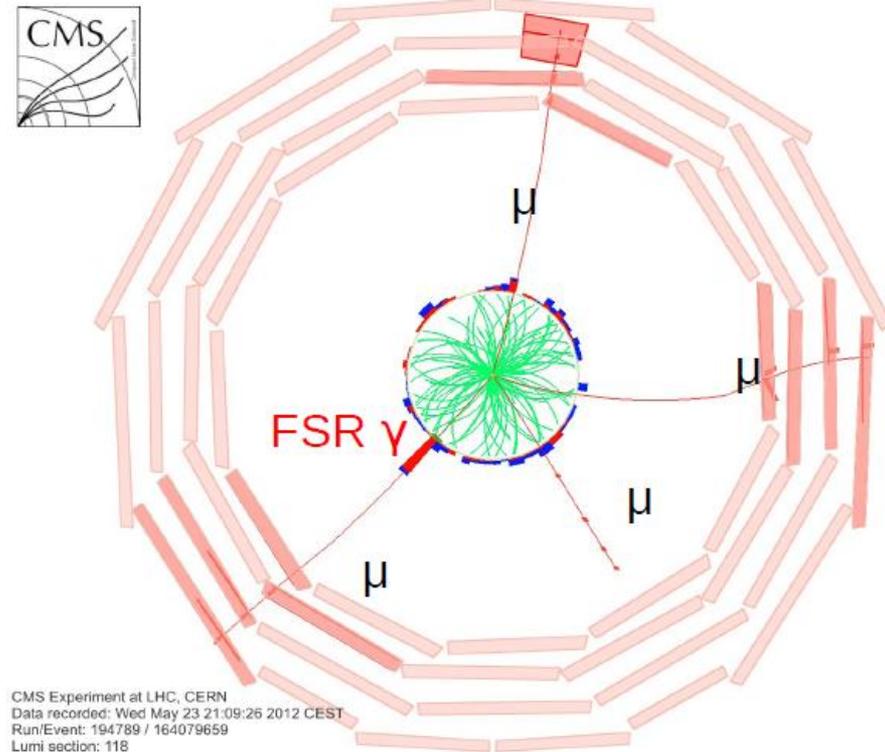
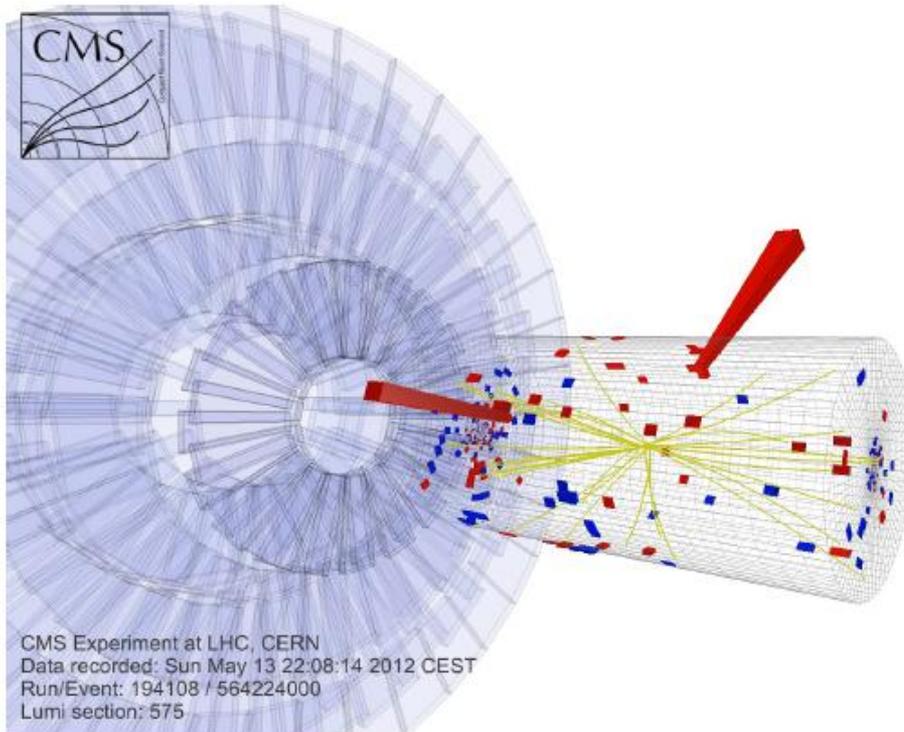
Layout for Higgs result presentation

- **measurements of h properties**
 - mass, width, spin/CP
 - consistency with SM couplings to bosons and fermions
- **Searches for BSM Higgs boson(s)**
 - non-SM h decays
 - other Higgs bosons
- **Rare SM h decays**

References to CMS publications are in backup slides



Measurement of h properties



How mass, spin/CP are measured in $h \rightarrow ZZ \rightarrow 4\ell$ and $h \rightarrow \gamma\gamma$ analyses

- $h \rightarrow ZZ \rightarrow 4\ell$

- mass: fit with 3D likelihood

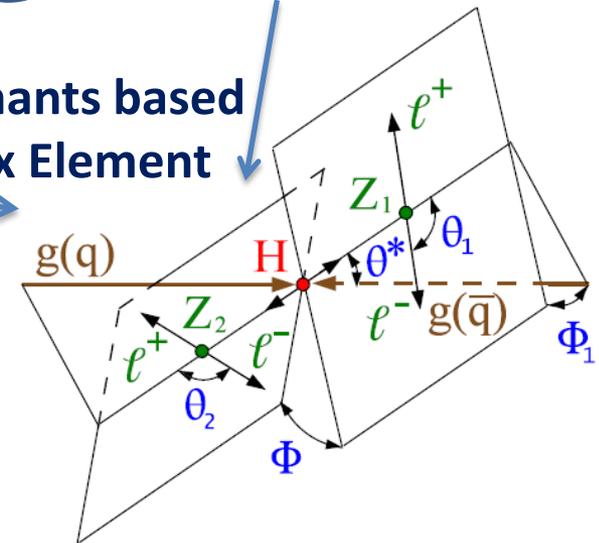
$$\mathcal{L}_{3D}^{m,\Gamma} \equiv \mathcal{L}_{3D}^{m,\Gamma}(m_{4\ell}, \mathcal{D}_m, \mathcal{D}_{\text{bkg}}^{\text{kin}}) = \mathcal{P}(m_{4\ell} | m_H, \Gamma, \mathcal{D}_m) \mathcal{P}(\mathcal{D}_m | m_{4\ell}) \times \mathcal{P}(\mathcal{D}_{\text{bkg}}^{\text{kin}} | m_{4\ell}).$$

- spin/CP; CL_s using 2D likelihood ratio

$$\mathcal{L}_{2D}^{J^P} \equiv \mathcal{L}_{2D}^{J^P}(\mathcal{D}_{\text{bkg}}, \mathcal{D}_{J^P})$$

Discriminants based on Matrix Element

per-event mass uncertainty



- $h \rightarrow \gamma\gamma$

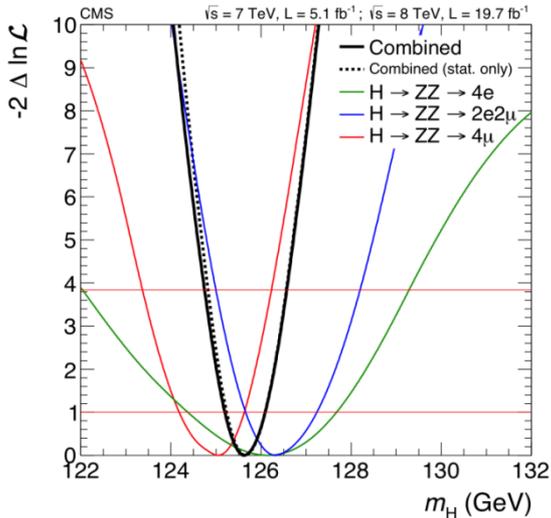
- mass: fit $m_{\gamma\gamma}$ floating $\mu_{\gamma\gamma}$
- spin/CP: CL_s with ratio of likelihoods using $\cos(\theta_{\text{CS}})$ to discriminate between hypotheses

$$\cos(\theta_{\text{CS}}^*) = 2 \times \frac{E_2 p_{z1} - E_1 p_{z2}}{m_{\gamma\gamma} \sqrt{m_{\gamma\gamma}^2 + (p_{\text{T}}^{\gamma\gamma})^2}}$$

For spin/CP HJU generator is used in both $\gamma\gamma$ and $ZZ \rightarrow 4\ell$ modes

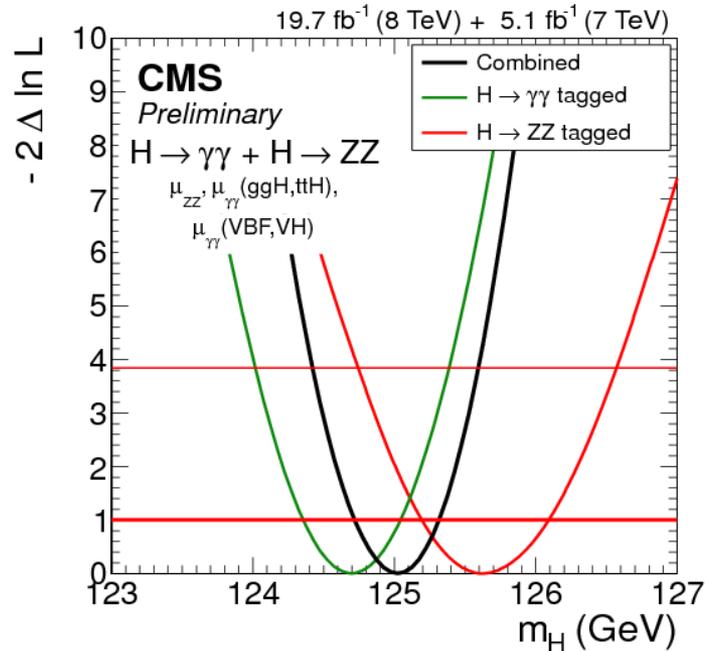
h boson mass

- h → ZZ → 4e



$$m_h^{ZZ} = 125.6 \pm 0.4(\text{stat.}) \pm 0.2(\text{syst.})$$

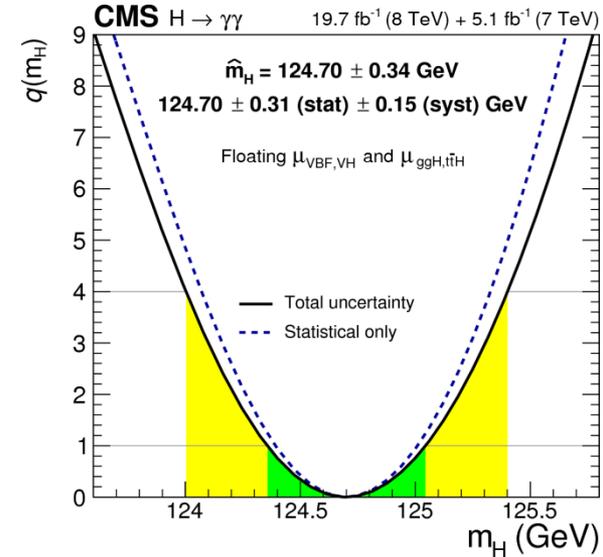
$$\mu = 0.94 \pm_{0.23}^{0.26}(\text{stat.}) \pm_{0.09}^{0.13}(\text{syst.})$$



$$m_h^{\gamma\gamma} = 124.72 \pm 0.31(\text{stat.}) \pm 0.15(\text{syst.})$$

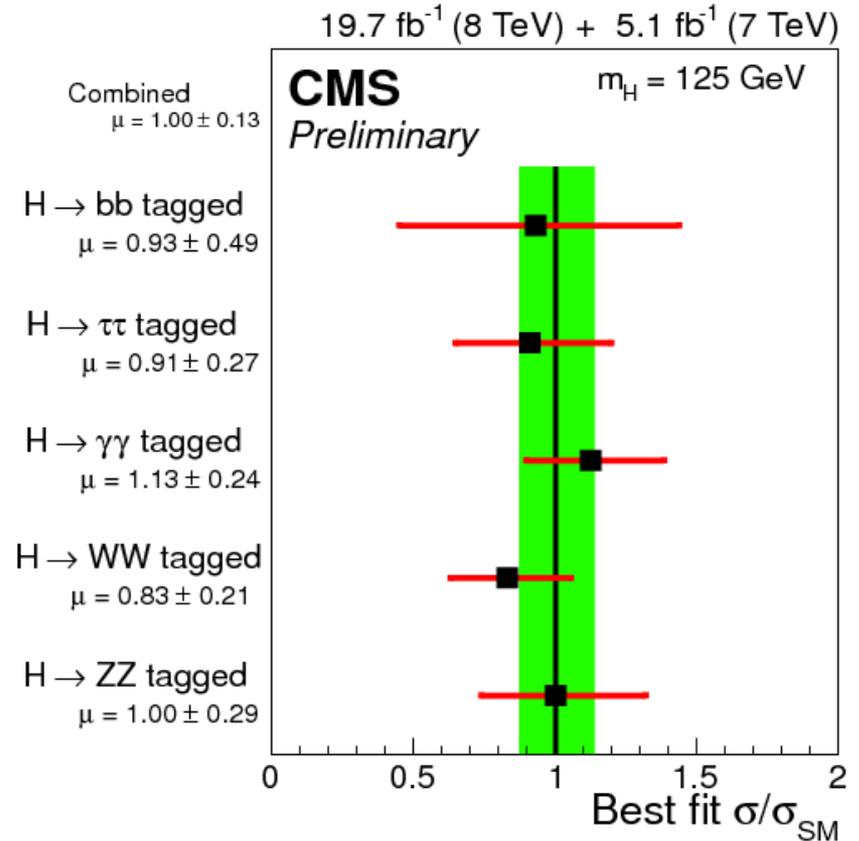
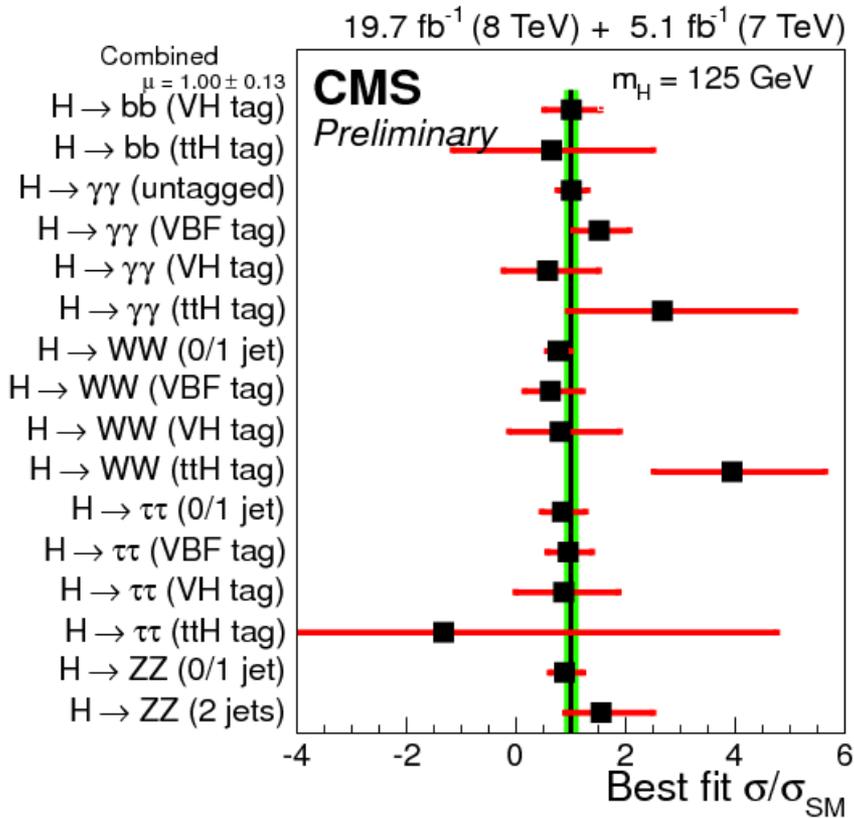
$$\mu = 1.14 \pm_{0.23}^{0.26}$$

- h → γγ



Combined $m_h = 125.0 \pm 0.30$ GeV

h signal strength at best fit mass



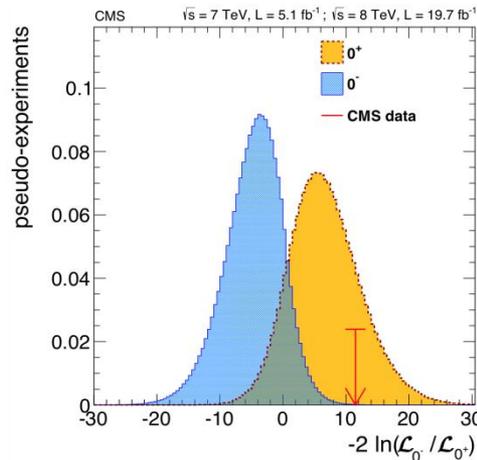
- $\mu = 1.00 \pm 0.13$

h spin/CP

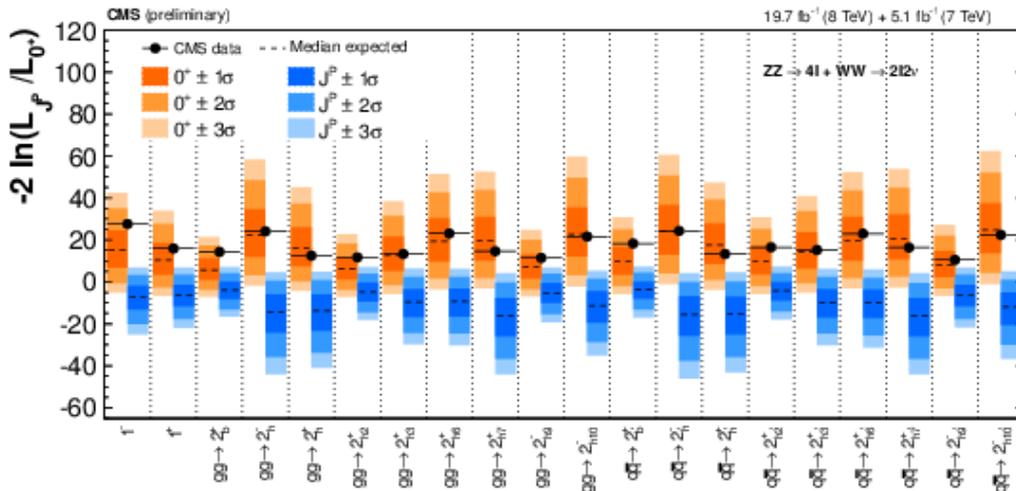
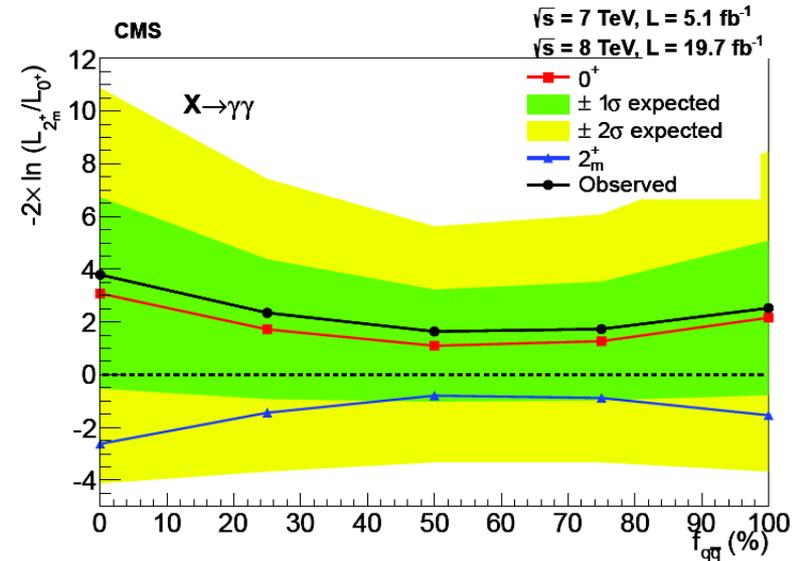
CP even vs CP odd; spin 0 vs spin 1 or spin 2

- $h \rightarrow ZZ \rightarrow 4\ell + H \rightarrow WW \rightarrow \ell\nu\ell\nu$

Pseudoscalar boson, spin-one and spin-two hypotheses are excluded at 99.9% C.L. or higher



- $h \rightarrow \gamma\gamma: 0^+ \text{ vs } 2^+$



$f_{q\bar{q}}$	$1 - CL_s$	
	expected	observed
0	0.92	0.94
0.25	0.78	0.83
0.50	0.64	0.71
0.75	0.69	0.75
1	0.83	0.85

Measuring CP-odd fraction with $h \rightarrow ZZ \rightarrow 4\ell$ (I)

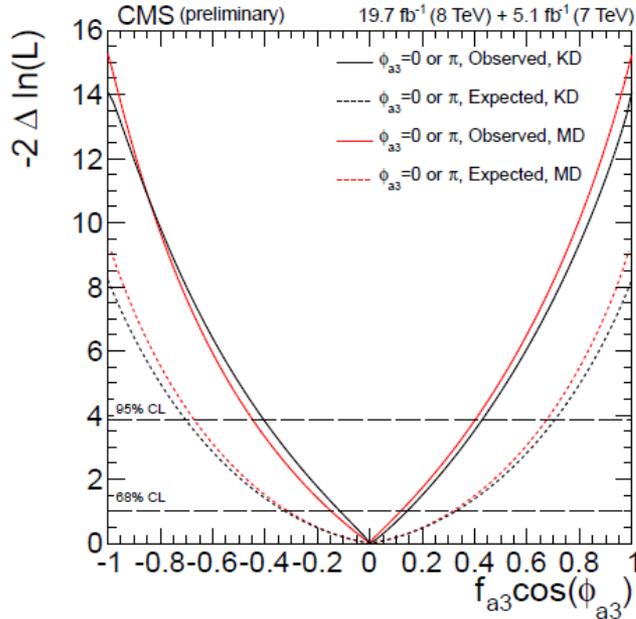
$$A(X_{J=0} \rightarrow V_1 V_2) \sim v^{-1} \left(\left[a_1 - e^{i\phi_{\Lambda_1}} \frac{q_{Z_1}^2 + q_{Z_2}^2}{(\Lambda_1)^2} \right] m_Z^2 \epsilon_{Z_1}^* \epsilon_{Z_2}^* \right.$$

a_1, a_2 - CP-even couplings

a_3 - CP-odd couplings

Λ_1 - scale of new physics

$$\begin{aligned} &+ a_2 f_{\mu\nu}^{*(Z_1)} f^{*(Z_2),\mu\nu} + a_3 f_{\mu\nu}^{*(Z_1)} \tilde{f}^{*(Z_2),\mu\nu} \\ &+ a_2^{Z\gamma} f_{\mu\nu}^{*(Z)} f^{*(\gamma),\mu\nu} + a_3^{Z\gamma} f_{\mu\nu}^{*(Z)} \tilde{f}^{*(\gamma),\mu\nu} \\ &+ a_2^{\gamma\gamma} f_{\mu\nu}^{*(\gamma_1)} f^{*(\gamma_2),\mu\nu} + a_3^{\gamma\gamma} f_{\mu\nu}^{*(\gamma_1)} \tilde{f}^{*(\gamma_2),\mu\nu} \end{aligned}$$



fraction of CP-odd contribution:

$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda_1} / (\Lambda_1)^4}$$

$$\frac{|a_i|}{|a_1|} = \sqrt{\frac{f_{ai}}{f_{a1}}} \times \sqrt{\frac{\sigma_1}{\sigma_i}} \quad \sigma_1/\sigma_3=6.36$$

- assuming a_3 to be constant
- $a_2=0$ (no CP-even high dim.)

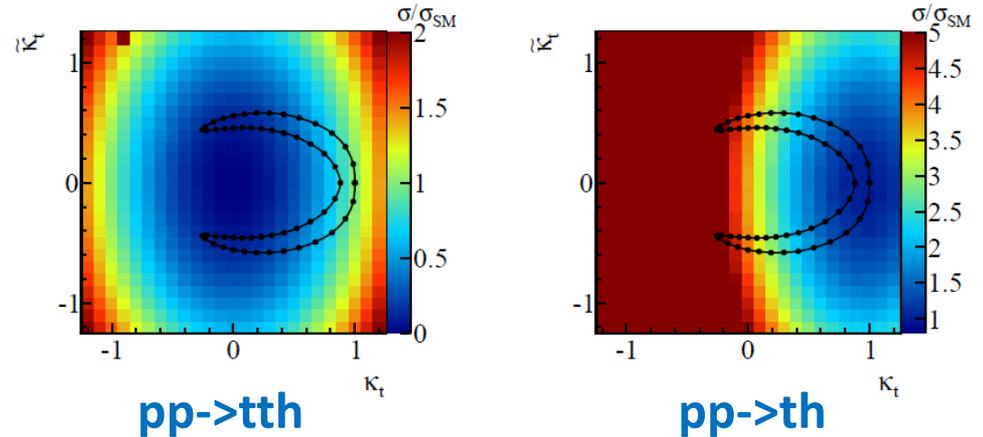
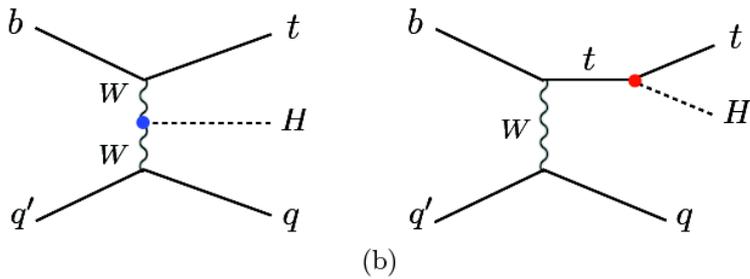
$a_3/a_1 = [-2.05, 2.19]$ observed

Measuring CP-odd fraction with $h \rightarrow ZZ \rightarrow 4\ell$ (II)

- Small f_{a3} does not necessary mean small CP-odd admixture
 - observables involving the Higgs-gauge couplings project on the CP-even component
 - better to look into observables with Higgs-fermion couplings

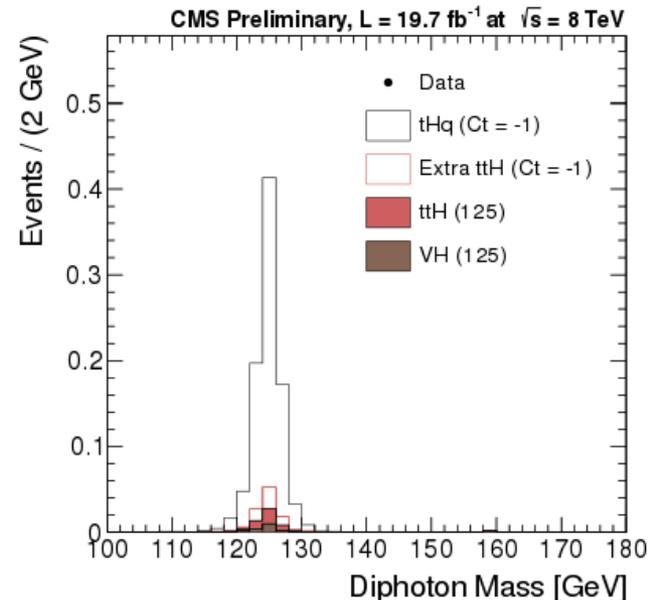
Single top+h, $h \rightarrow \gamma\gamma$ analysis

J. Ellis, D.S.Hwang, M. Takeuchi arXiv:1312.5736



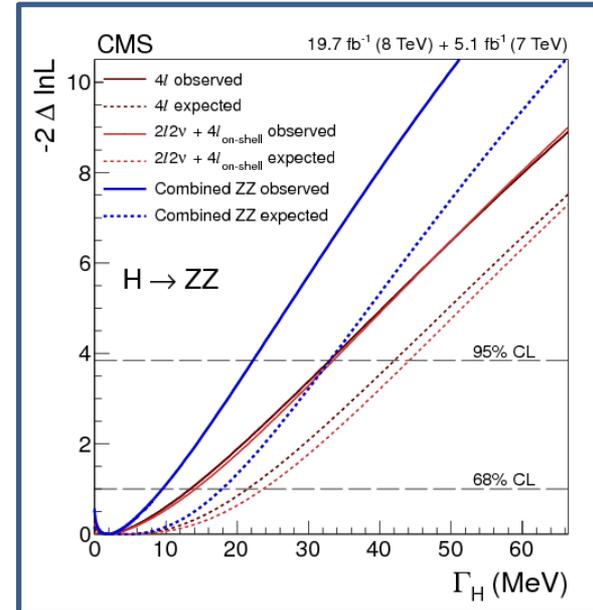
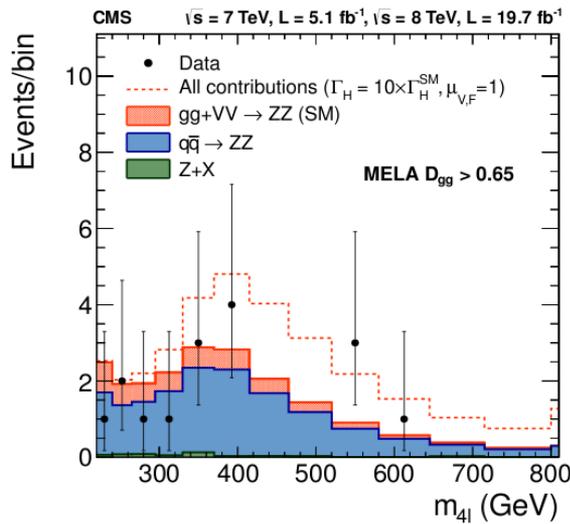
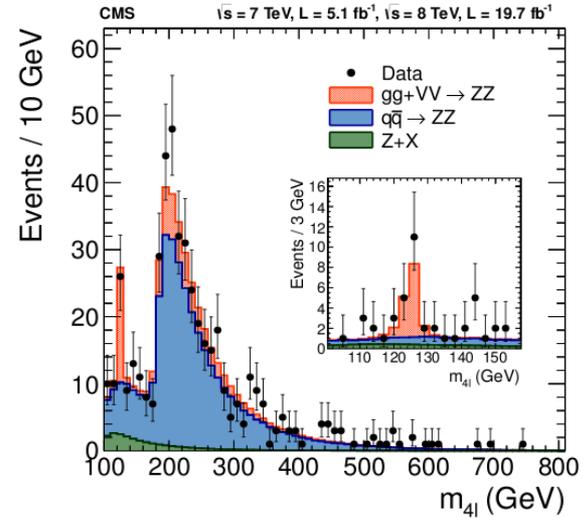
- Due to interference $pp \rightarrow tqh$ process is sensitive to both magnitude and sign of γ_{tth}
- tth and th can be used to measure scalar and pseudo-scalar h -top couplings

Upper limit = 4.1 x expected cross-section with $k_t = -1$.



Width of h(125) boson using off-shell H production and ZZ decay

- $\sigma_{i \rightarrow h(\text{on-shell}) \rightarrow f} \sim g_i^2 g_f^2 / \Gamma_h$
- $\sigma_{i \rightarrow h(\text{off-shell}) \rightarrow f} \sim g_i^2 g_f^2$
- $\sigma_{h \rightarrow gg \rightarrow ZZ} \text{ interference} \sim g_i g_f$
- Fit on- and off-shell mass region
 - varying $g_i^2 g_f^2$ and Γ_h
 - assuming SM off/on-shell yield ratio



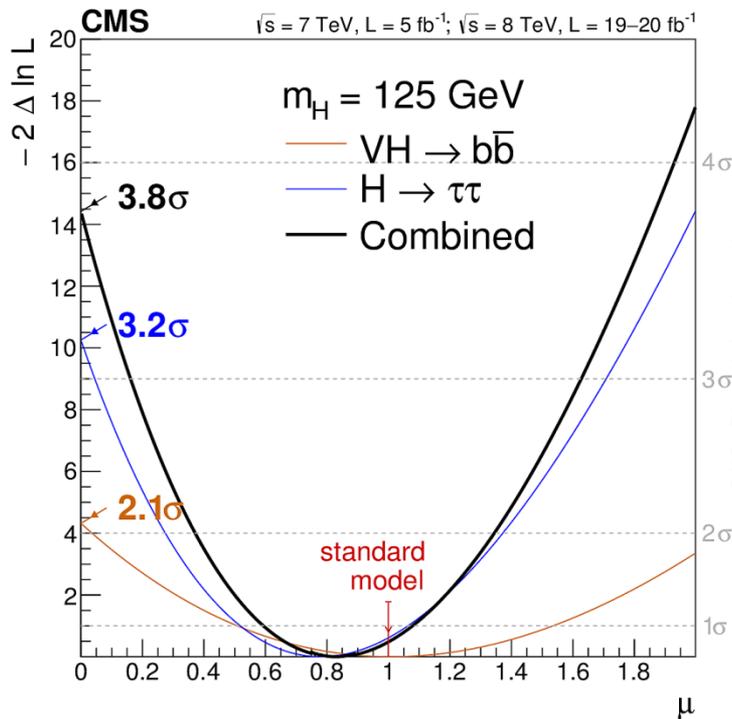
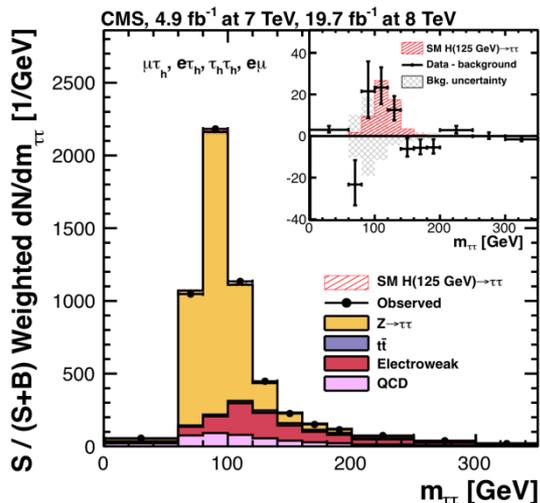
Higgs boson couplings at LHC

- determination of total width impossible without further assumptions
- not all final states are accessible
- what is done so far:
 - perform fits to coupling strength modifiers to check consistency with SM
 - *LHC HXSWG Recommendations, benchmark models*

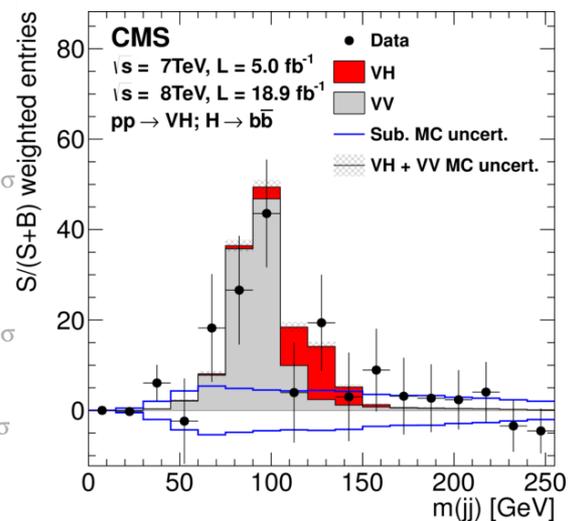
October 2013 – January 2014 :

h(125) boson couples to τ and b

$h \rightarrow \tau\tau$



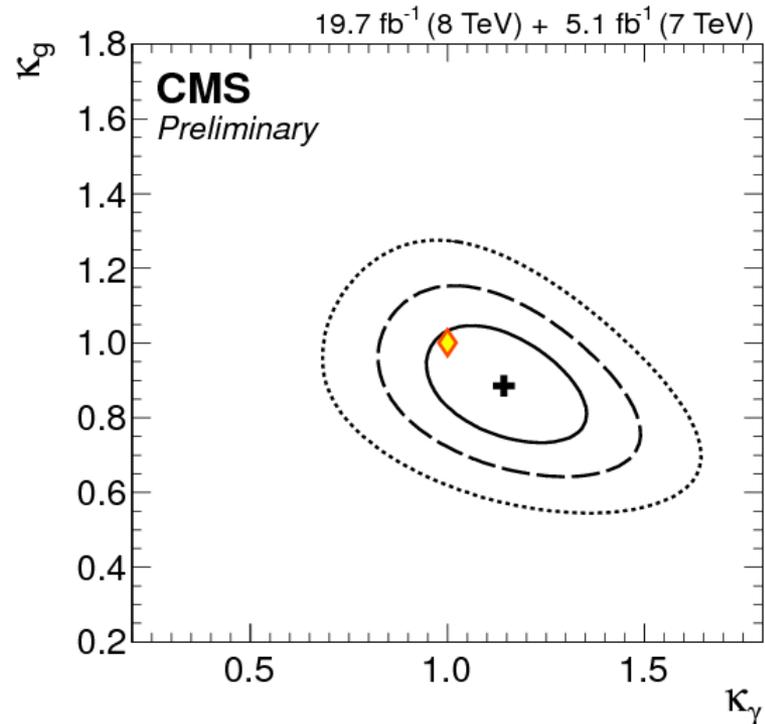
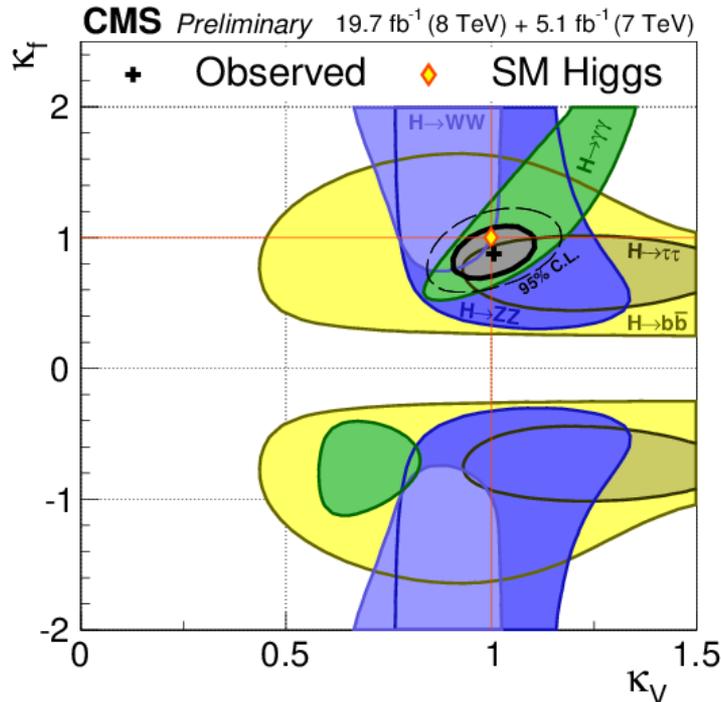
Vh, $h \rightarrow b\bar{b}$



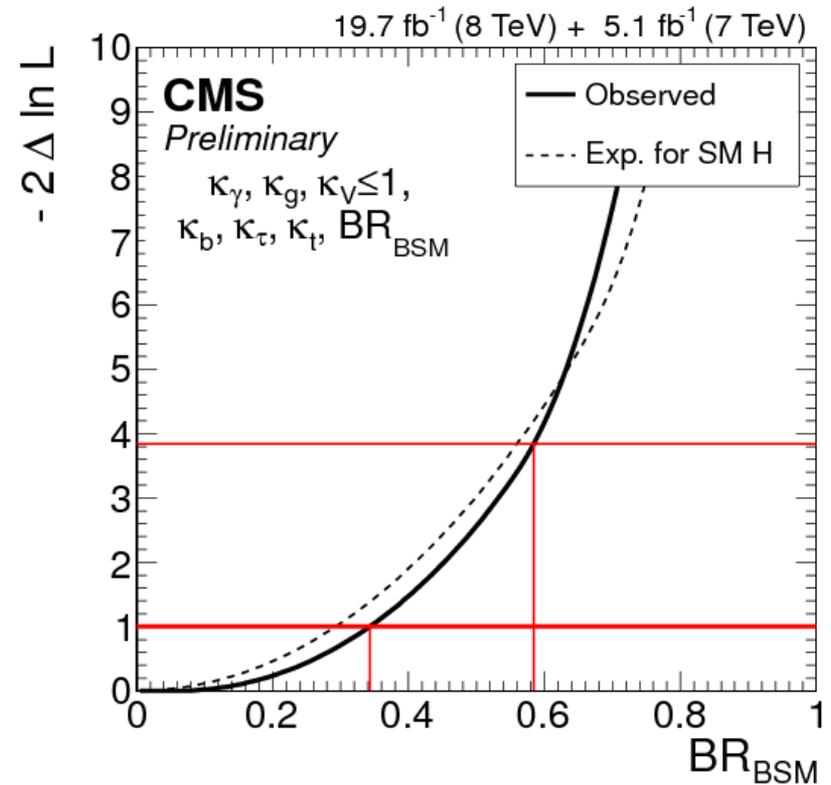
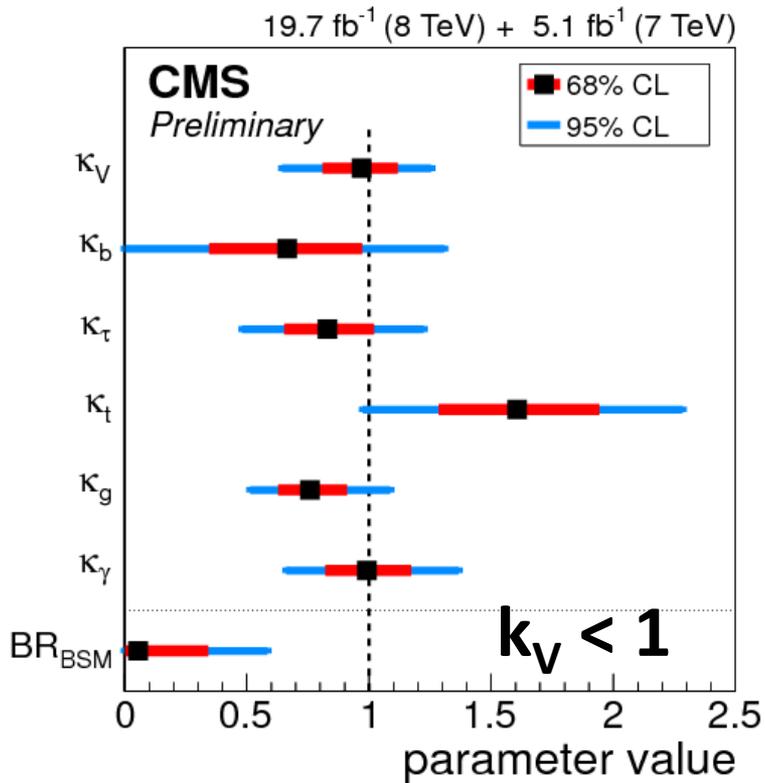
Channel	Significance (σ)		Best-fit
($m_H = 125 \text{ GeV}$)	Expected	Observed	μ
$VH \rightarrow b\bar{b}$	2.3	2.1	1.0 ± 0.5
$H \rightarrow \tau\tau$	3.7	3.2	0.78 ± 0.27
Combined	4.4	3.8	0.83 ± 0.24

Some benchmark parameterizations of Higgs boson couplings

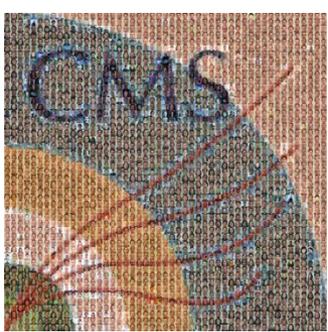
- k_V – common scale factor for vector boson couplings
- k_f – common scale factor for fermion couplings
- assuming no invisible and undetectable widths
- k_γ, k_g – probing loop structure assuming no invisible or undetectable width



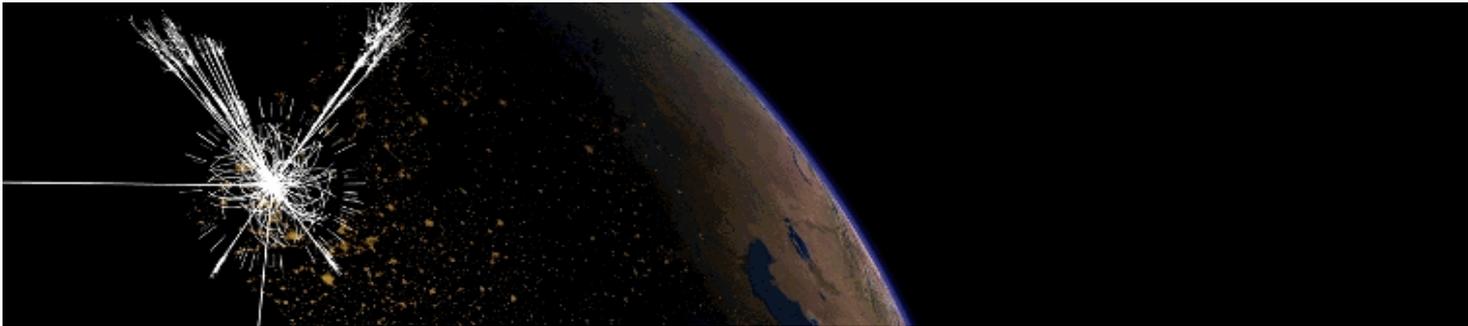
Constraining $h \rightarrow \text{BSM}$ with “visible” modes



$\text{BR}_{\text{BSM}} < 0.58$ at 95% CL



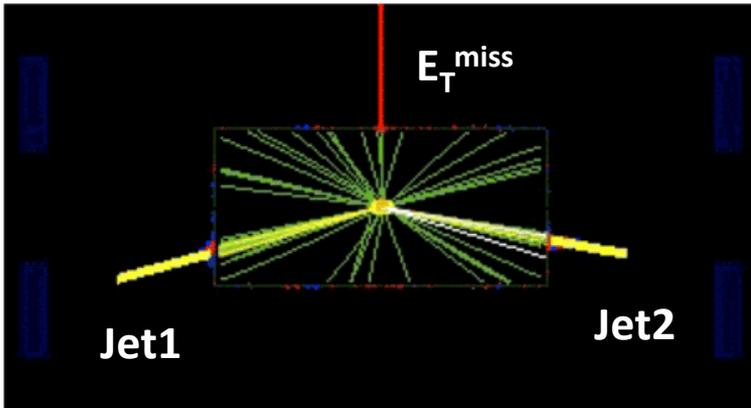
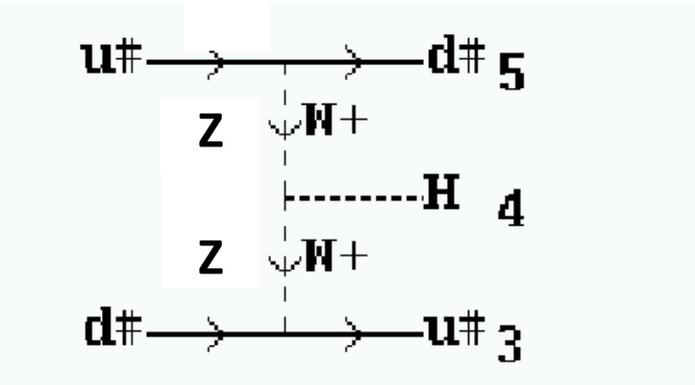
Non-SM h decays:
*searches for $h \rightarrow \text{invisible}$
with VBF h , Zh ($Z \rightarrow \ell\ell, bb$)*



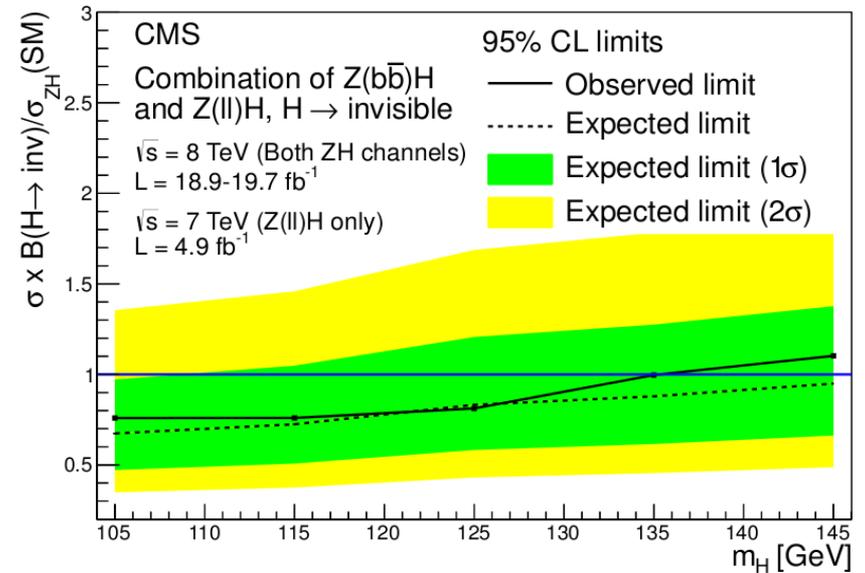
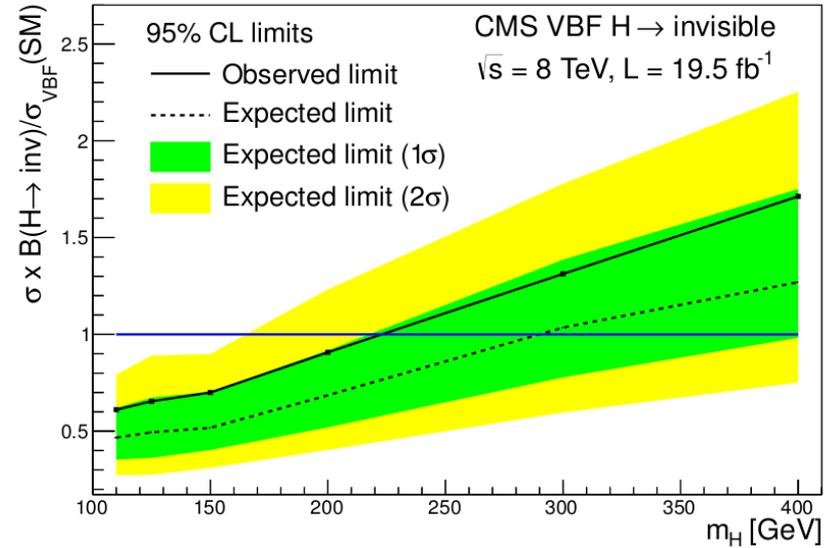
Detection of Dark Matter

most sensitive mode

$qq' \rightarrow qq'h$ (VBF h)

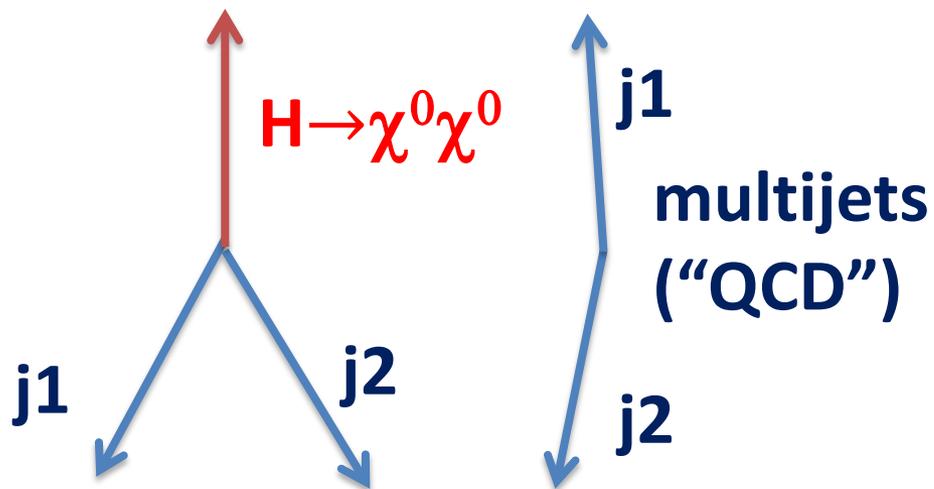


Event 191202:51:82701983

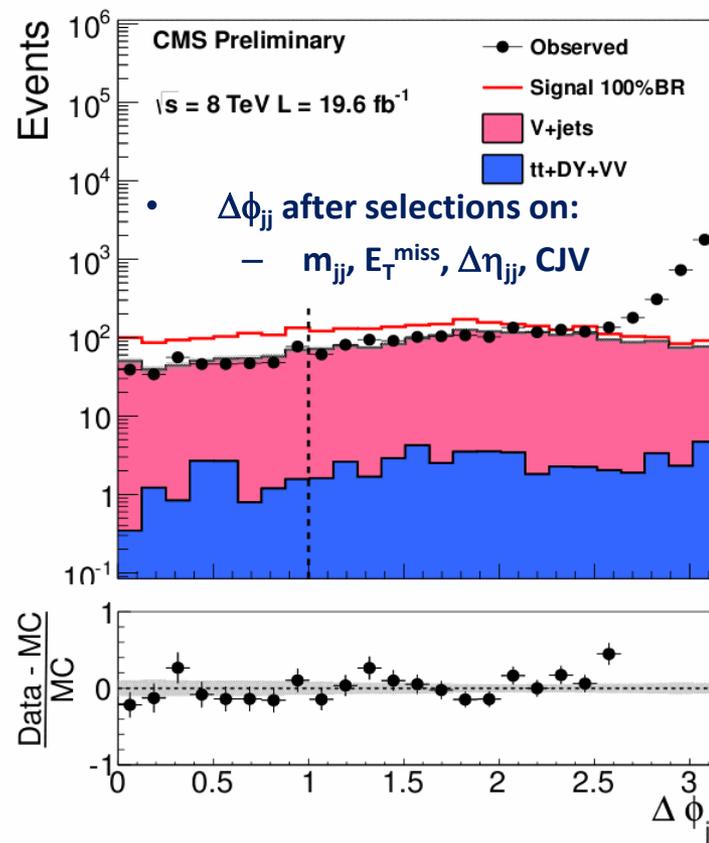


VBF $h \rightarrow$ invisible: offline signal selections and topology

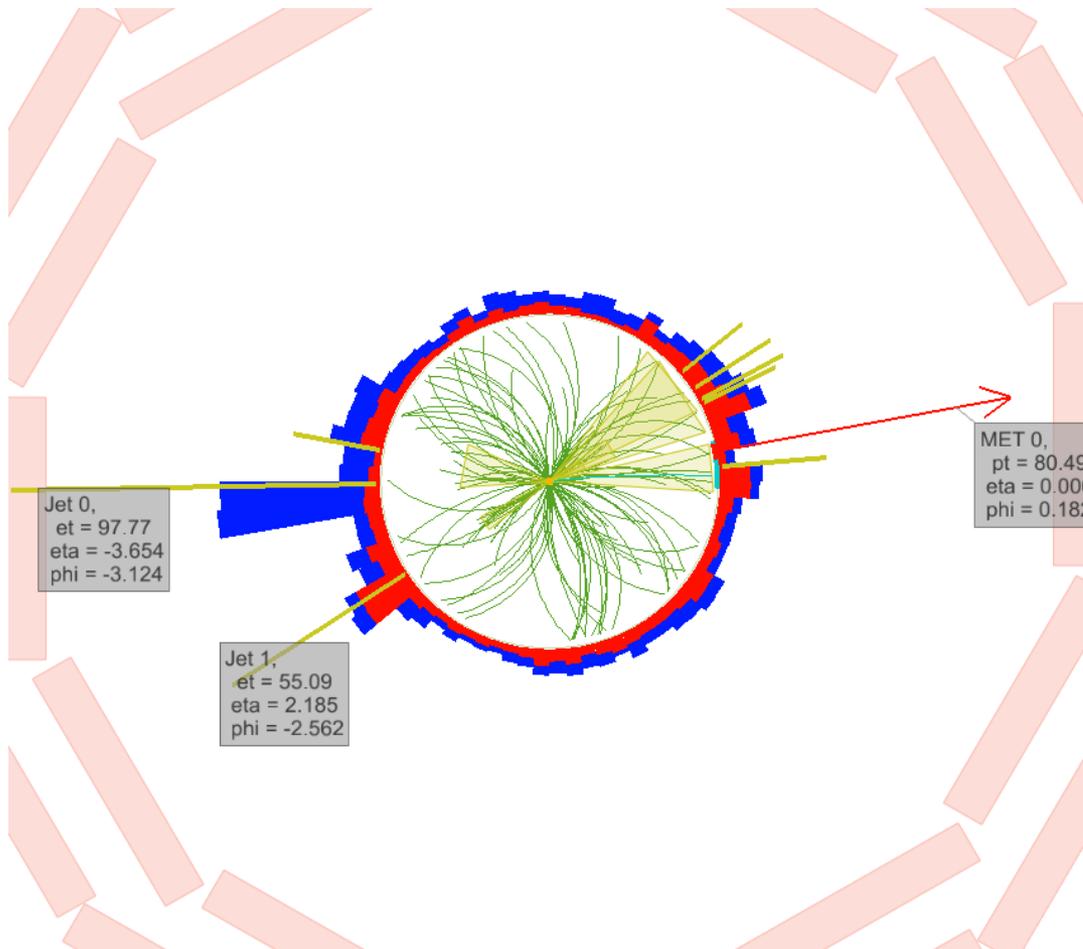
- two jets $p_T > 50$ GeV, $|\eta| < 4.7$
- $m_{jj} > 1100$ GeV
- $\Delta\eta_{jj} > 4.2$
- $E_T^{\text{miss}} > 130$ GeV
- $\Delta\phi_{jj} < 1.0$
- Central Jet Veto (CJV)



Signal: small $\Delta\phi_{jj}$
QCD: large $\Delta\phi_{jj}$

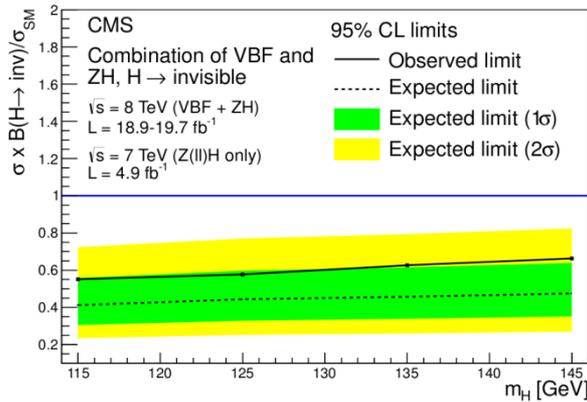


Signal region, with CJV (x,y view)

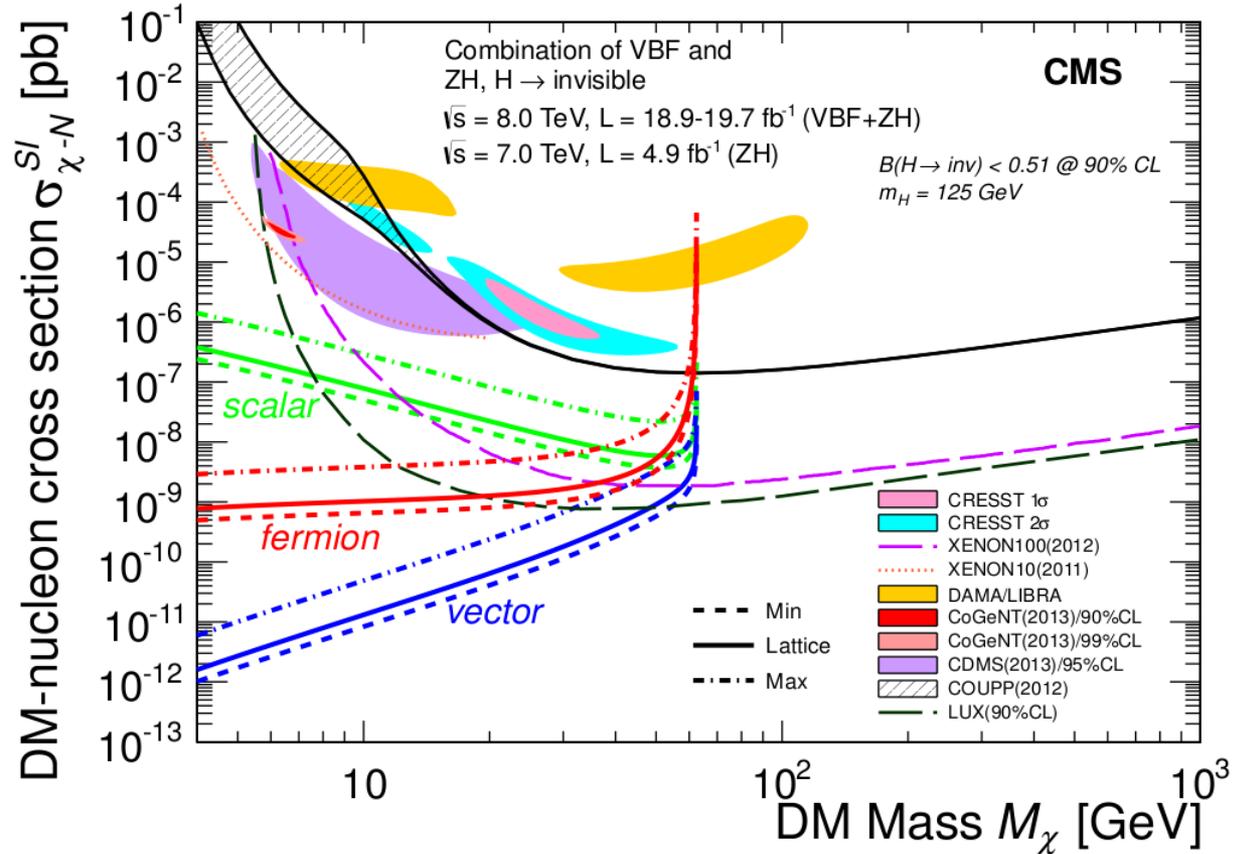


	∇pT	η	ϕ
0	97.8	-3.654	-3.124
1	55.1	2.185	-2.562
2	27.0	-0.200	0.084
3	24.8	2.286	0.502
4	23.4	-2.044	2.958
5	23.0	0.359	0.575
6	21.6	-2.901	0.472
7	20.4	1.349	0.691

Interpretation of H->invisible direct search in Higgs-portal Dark Matter models



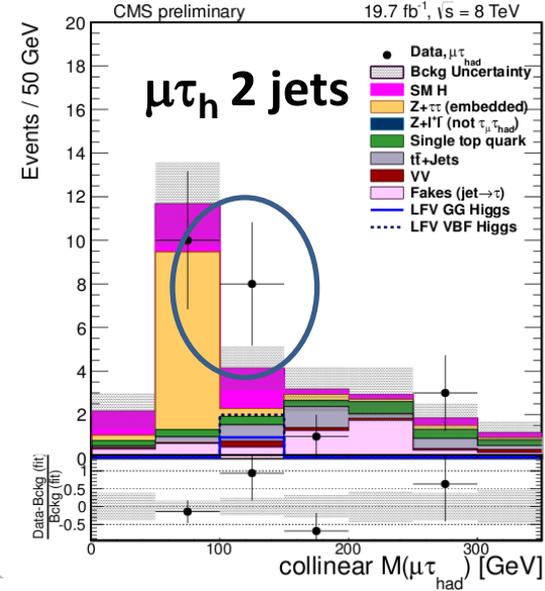
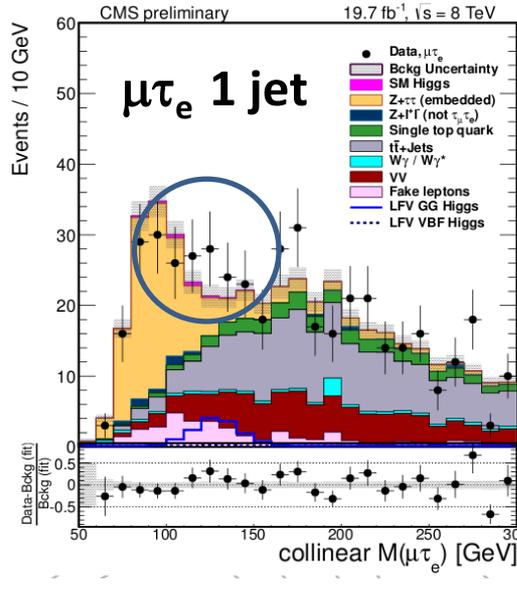
**$BR(h_{125 \text{ GeV}} \rightarrow \text{inv})$
 < 0.58 at 95 % CL
 (expected < 0.44)
 combining
 VBF and ZH modes**



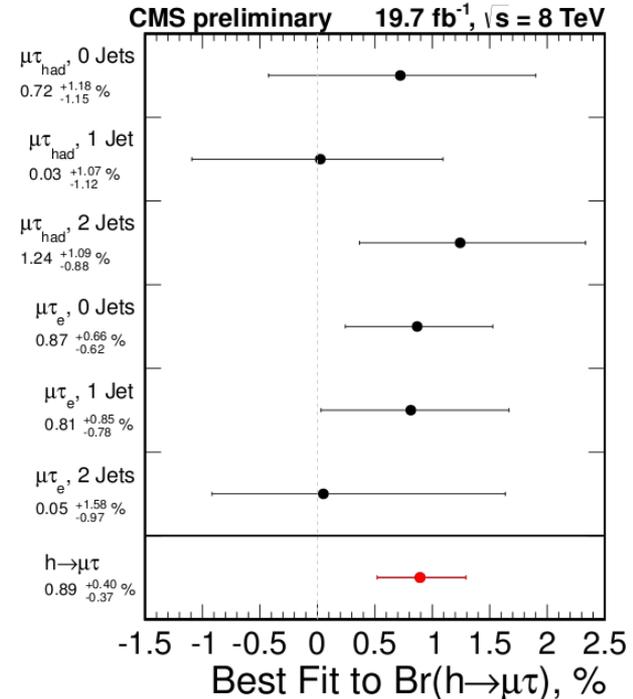
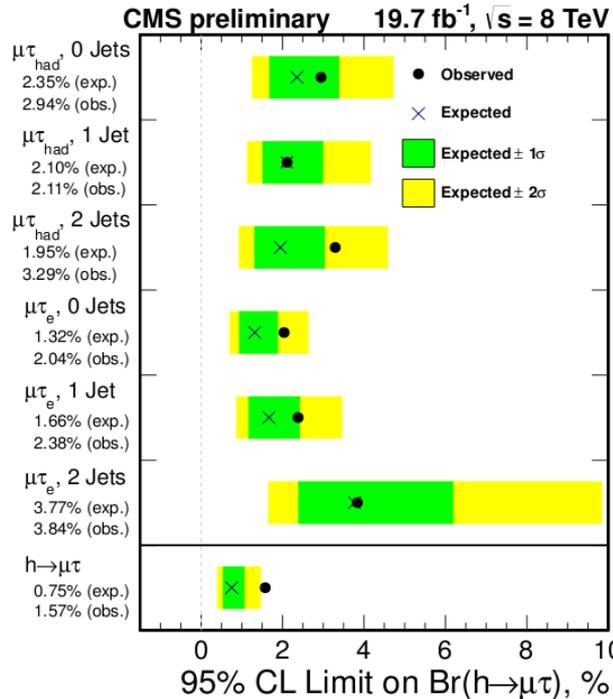
Non-SM h decays:

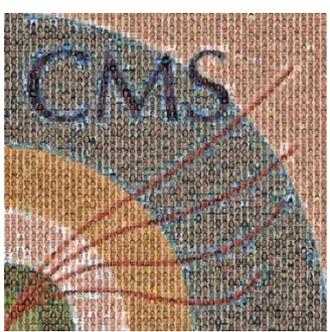
LFV with $h \rightarrow \mu\tau$

- Events are subdivided on 0, 1, 2 jet categories
- $\mu\tau_e$, $\mu\tau_h$ topologies
- m_h is reconstructed in collinear approximation



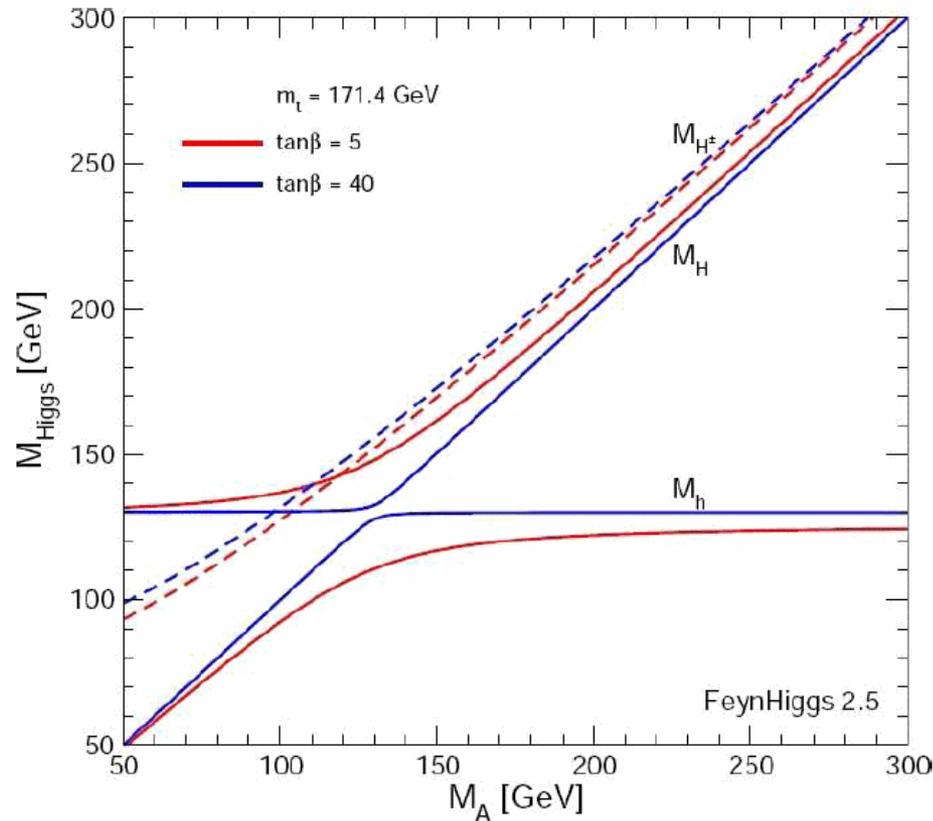
- **Upper limit:**
– BR < 1.57 %
- **~ 2.5 σ access**





Searches for other Higgs bosons

MSSM Higgs bosons

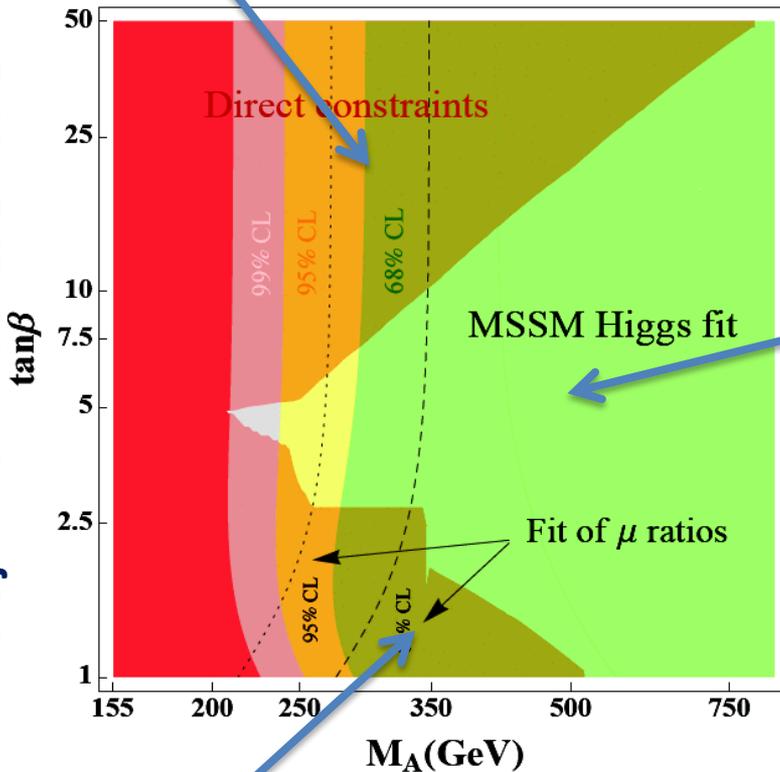


- **five Higgs bosons in MSSM:**
 - two CP-even h, H ; one CP-odd A , two charged $H^{+/-}$
 - *what is discovered as state of 125 GeV, h or H ?*

Landscape of BSM Higgs channels in MSSM

- *High $\tan\beta$* : $\phi \rightarrow \tau\tau$, $\phi \rightarrow \mu\mu$; $H^+ \rightarrow \tau\nu$, tb

A. Djouadi et.al. arXiv:1307.5205



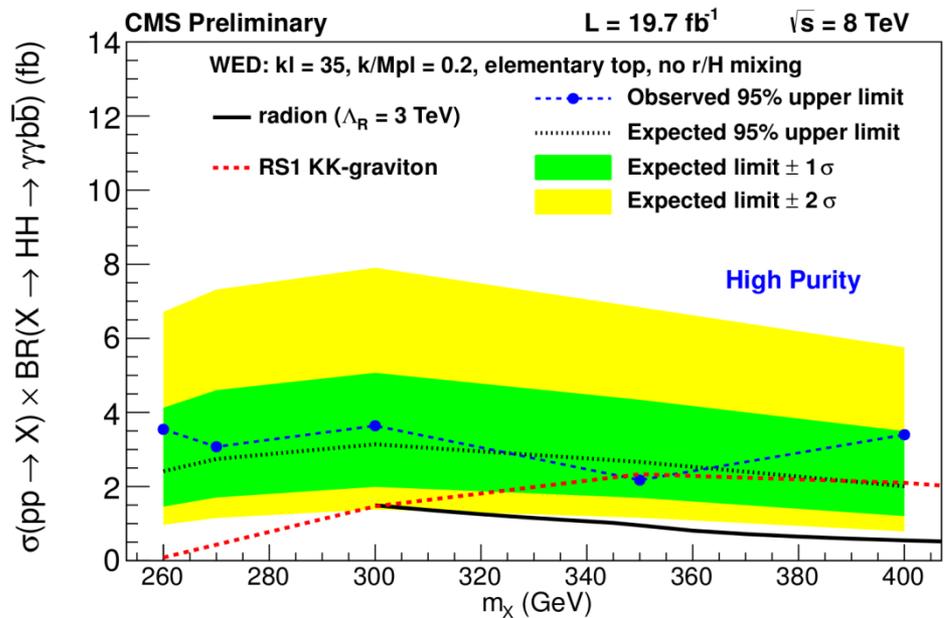
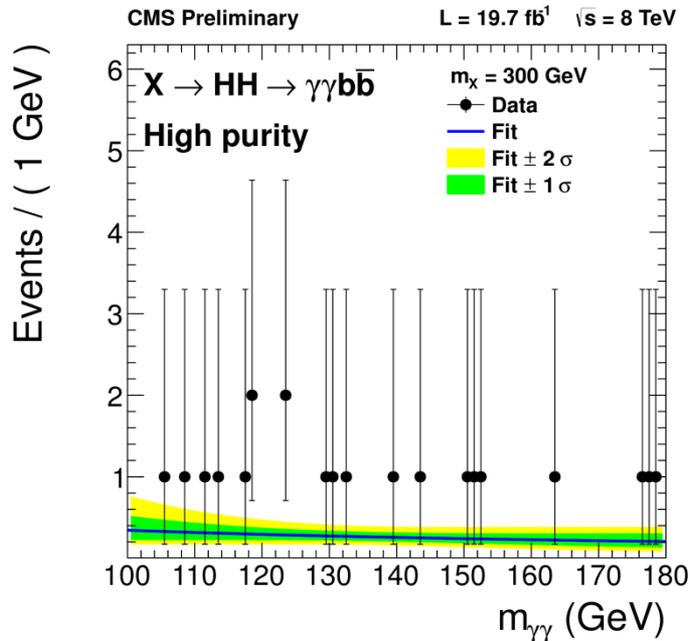
- *Intermediate $\tan\beta$* :

- $H/A \rightarrow \chi_i^0 \chi_j^0, \chi_i^+ \chi_j^-$
- $H^+ \rightarrow \chi_i^+ \chi_j^0$

- *Low $\tan\beta$* : $A \rightarrow Zh$; $H \rightarrow hh, tt$; $H^+ \rightarrow cs, cb, \tau\nu, tb, Wh$

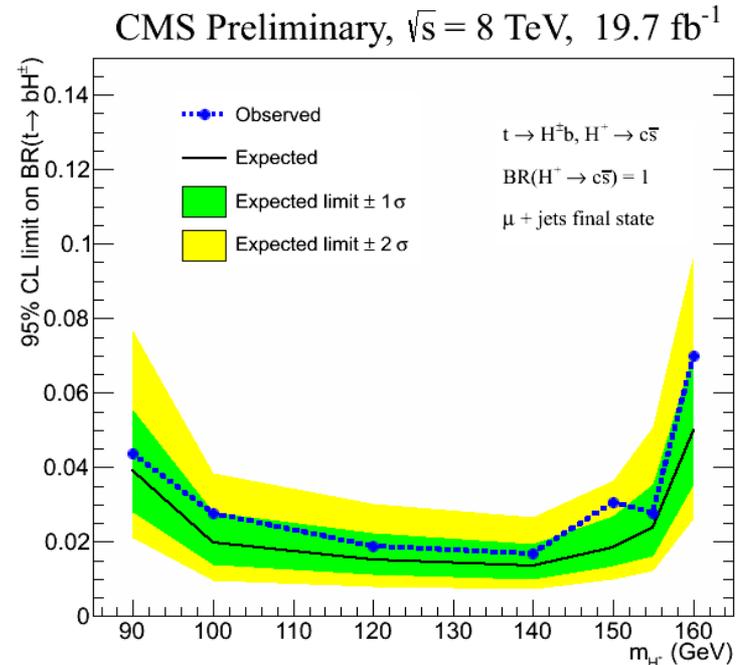
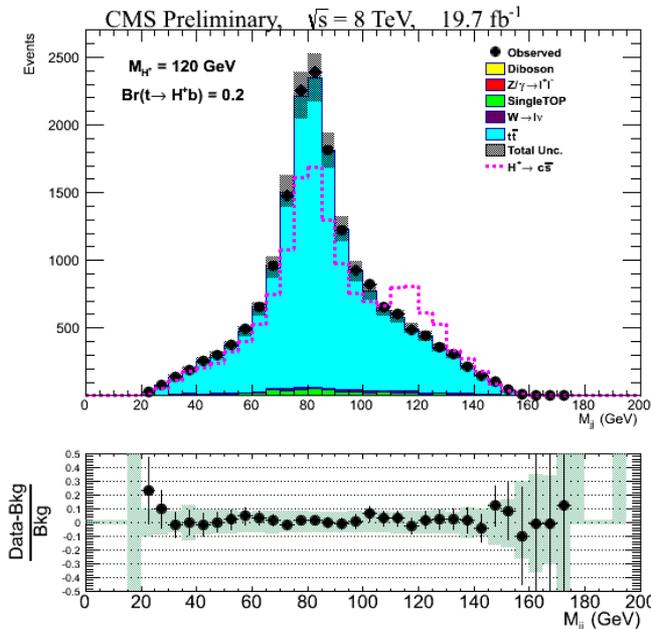
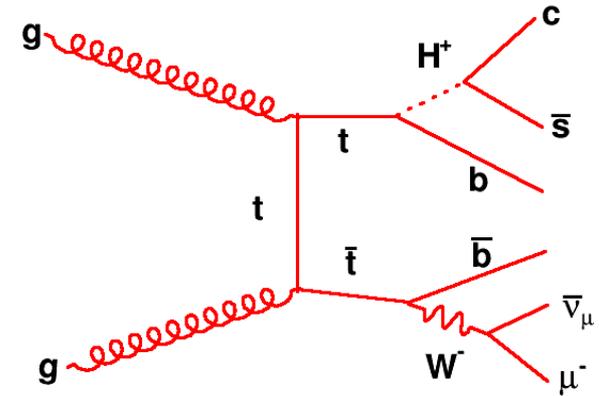
H → hh → γγbb

- Search strategy:
 - looking for signal in $m_{\gamma\gamma}$ distribution for $\gamma\gamma bb$ events selected within m_{bb} and $m_{\gamma\gamma bb}$ mass windows
- In hMSSM (A. Djouadi et.al. arXiv:1307.5205):
 - $\sigma(gg \rightarrow H) \times BR(H \rightarrow hh \rightarrow \gamma\gamma bb) = 2.9 \text{ fb}$ for $m_A = 300 \text{ GeV}$, $\tan\beta = 2$
 - close to observed limit !

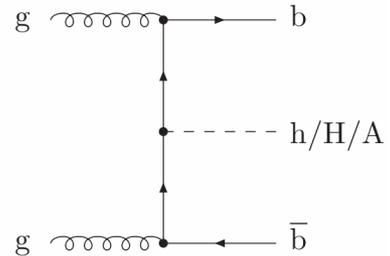
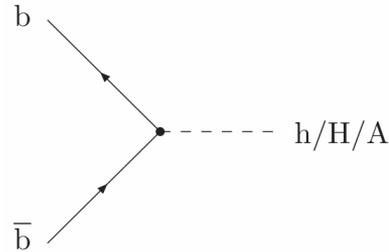
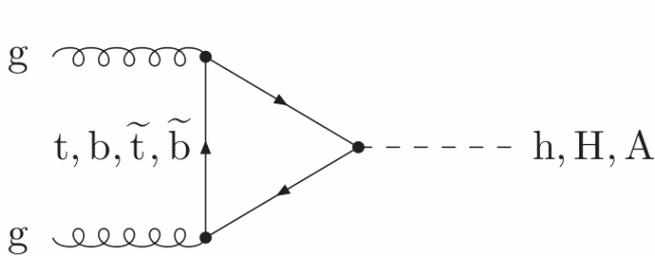


$tt \rightarrow bWbH^+, H^+ \rightarrow cs, m_{H^+} < m_t$

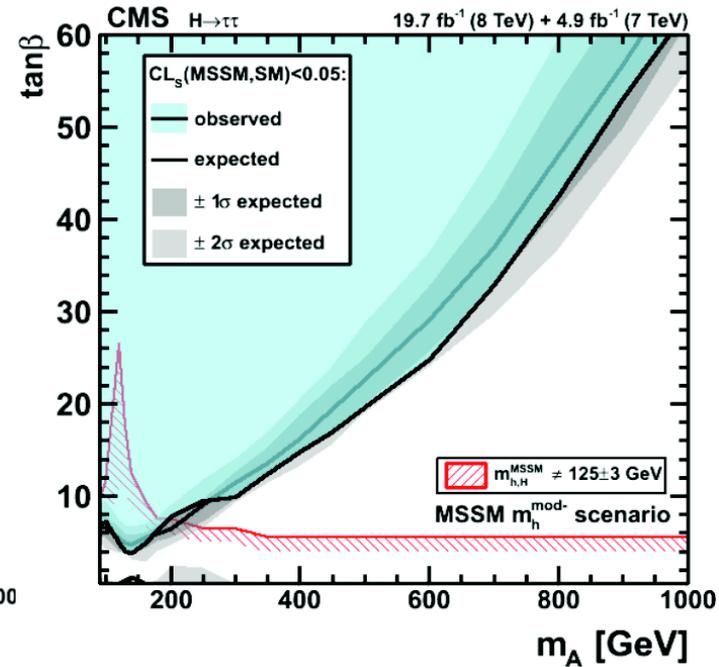
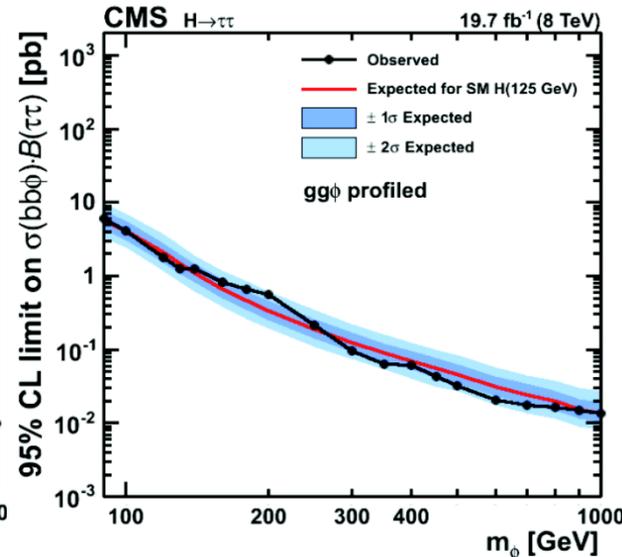
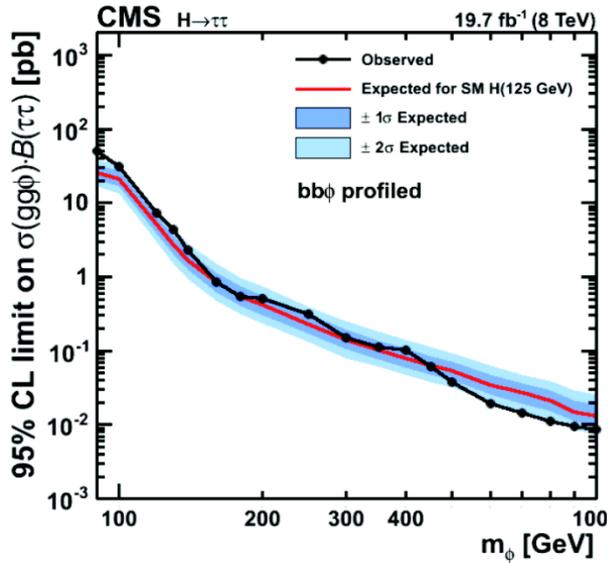
- Search strategy:
 - trigger with muon from $t \rightarrow Wb \rightarrow \mu\nu b$ decay
 - search for bump in di-jet mass distribution for jets from 2nd top decay
 - kinematic fit with top mass constraint



$h, H, A \rightarrow \tau\tau$



- split events into b-tag and no-b-tag categories
- consider $\tau_\mu\tau_h, \tau_e\tau_h, \tau_h\tau_h, \tau_e\tau_\mu, \tau_\mu\tau_\mu$ final states



- model independent limits
- m_A - $\tan\beta$ exclusions in new benchmark scenarios (M.Carena et al. arXiv:1302.7033)

Searches for heavy H decays into ZZ and WW

Interpretation in the Additional EW singlet model (h, H):

$$\sigma_h = c \times \sigma_{h,SM}$$

$$\Gamma_h = c \times \Gamma_{h,SM}$$

$$BR_{h,i} = BR_{h,SM,i}$$

$$\sigma_H = c' \times \sigma_{H,SM}$$

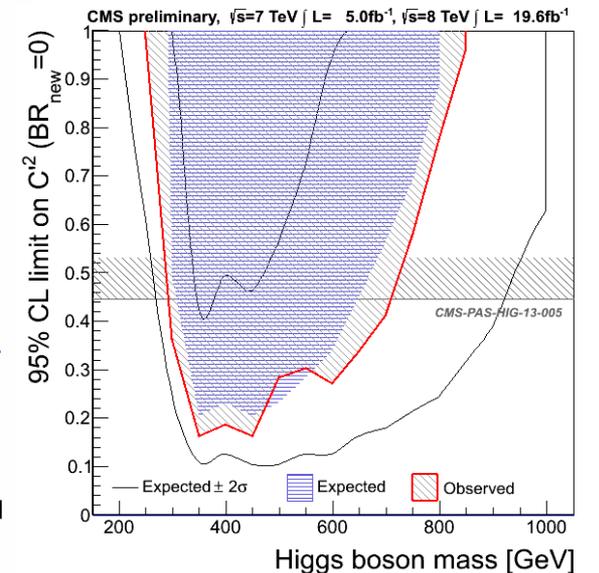
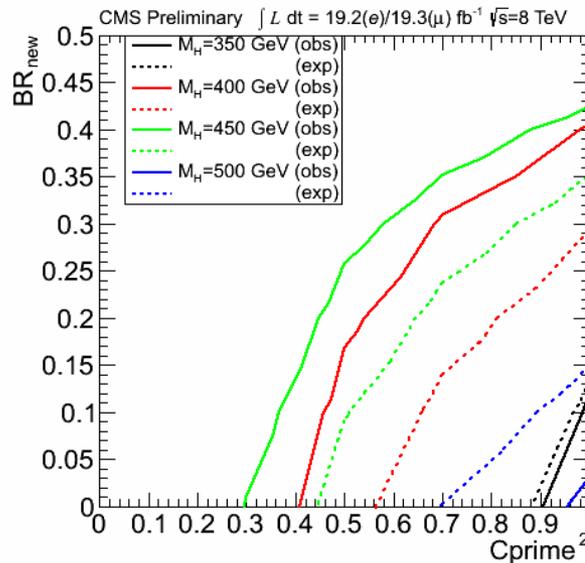
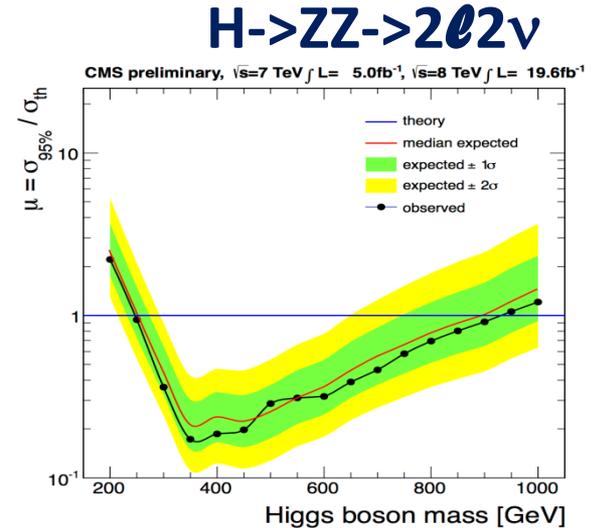
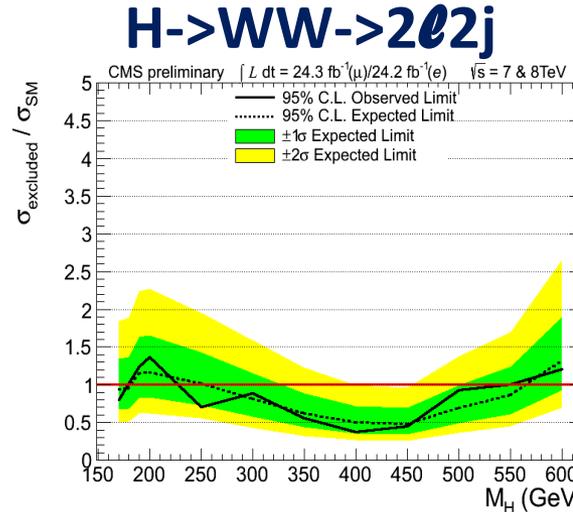
$$\Gamma_H = c' / [1 - BR_{H,new}] \Gamma_{H,SM}$$

$$BR_{H,i} = [1 - BR_{H,new}] BR_{H,SM,i}$$

$$c^2 + c'^2 = 1$$

$$c'^2 = 1 - \mu_h$$

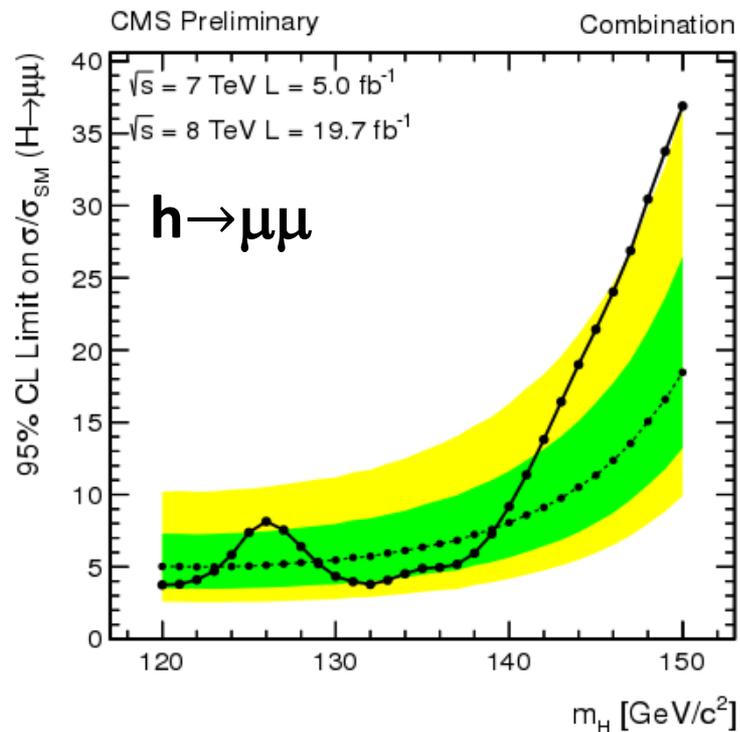
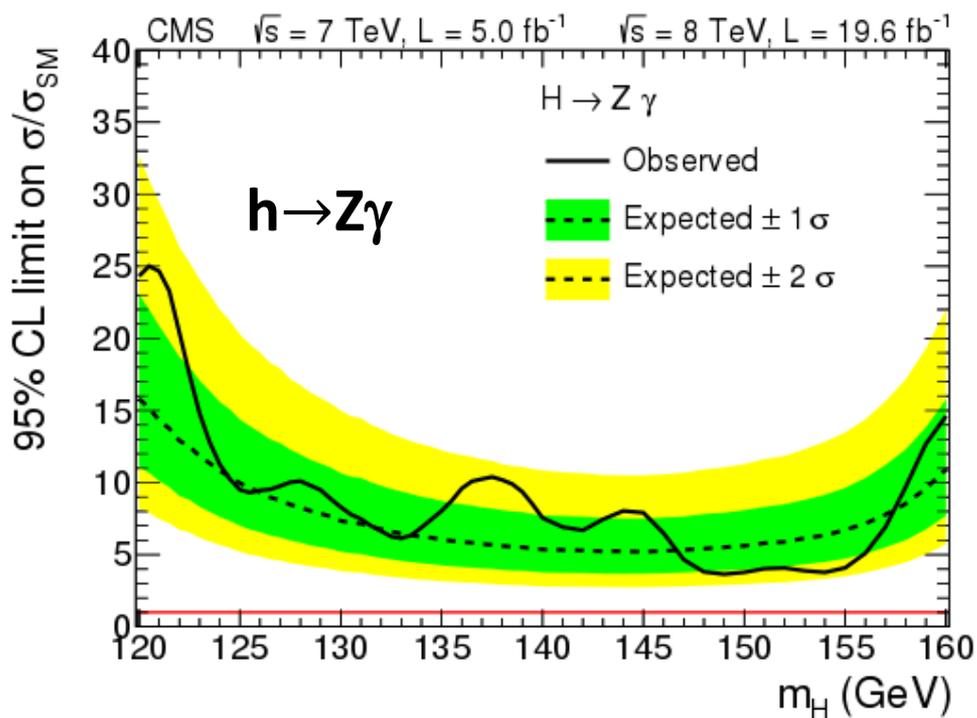
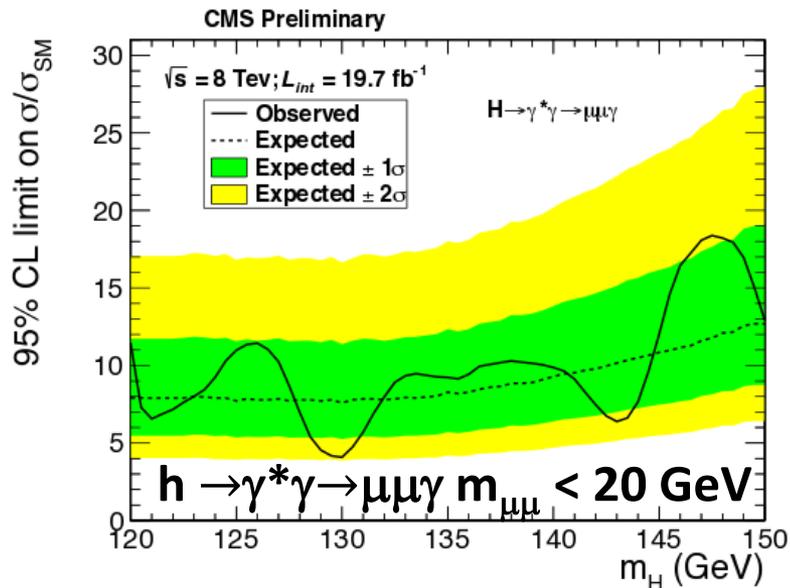
$$\mu_H = c'^2 [1 - BR_{H,new}]$$





Rare SM h decay

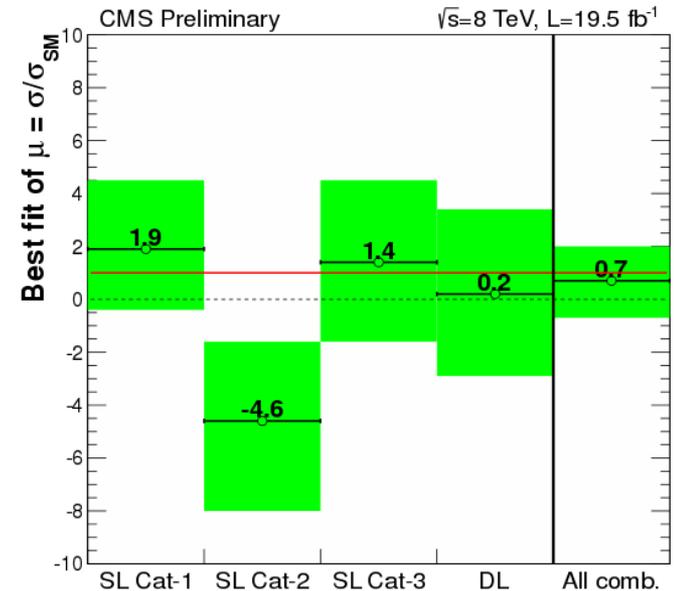
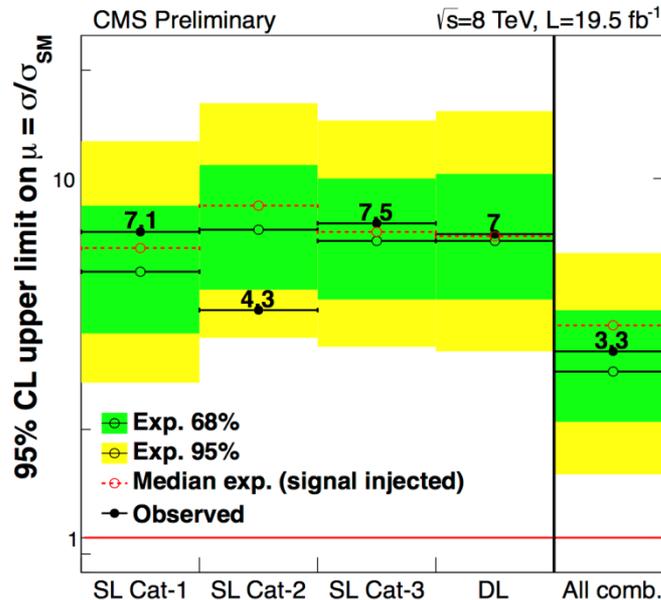
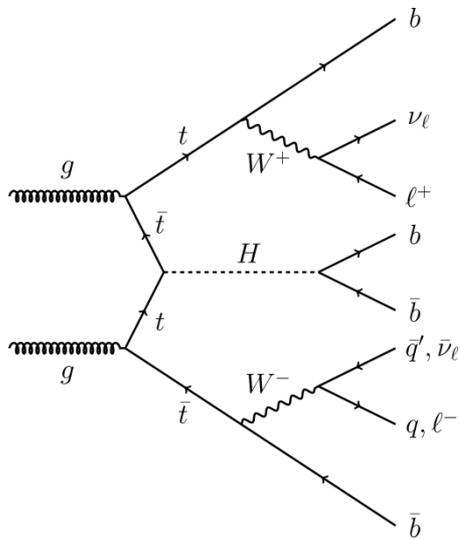
Decay mode	limit σ / σ_{SM}
$h \rightarrow \gamma^* \gamma \rightarrow \mu\mu\gamma$	~ 10
$h \rightarrow Z\gamma$	~ 10
$h \rightarrow \mu\mu$	7.4



Resent result on tth, h->bb:

limit 3.3 σ_{SM} at 95 % C.L.

- Not rare, but difficult and very important channel
 - supposed to be the hot subject with 13 TeV data
- $1\ell + \geq 5$ jets or $2\ell + \geq 4$ jets with ≥ 2 b-tagged jets
- ME method based on LO MadGraph
 - must use NLO generation in Run2 analyses !



Conclusions

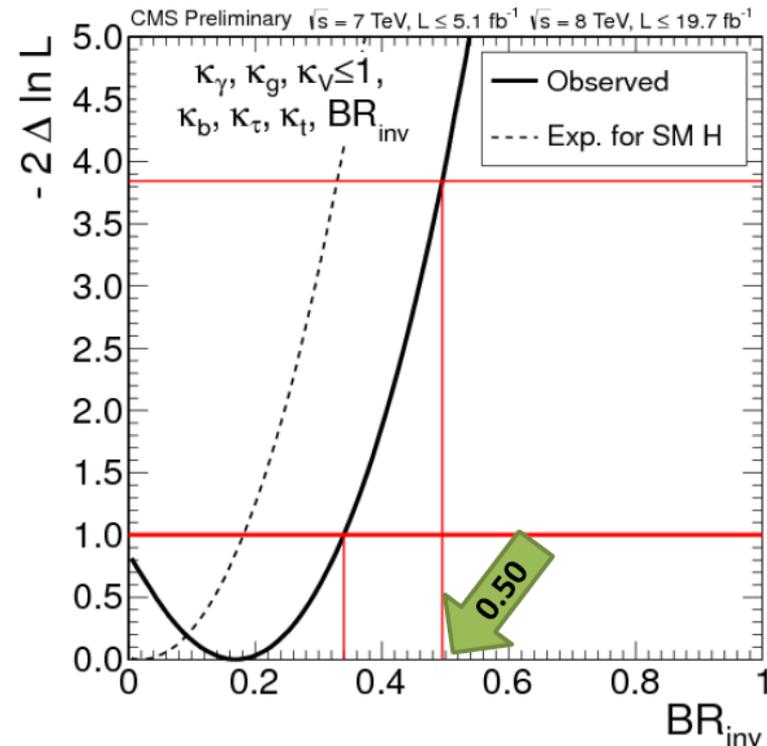
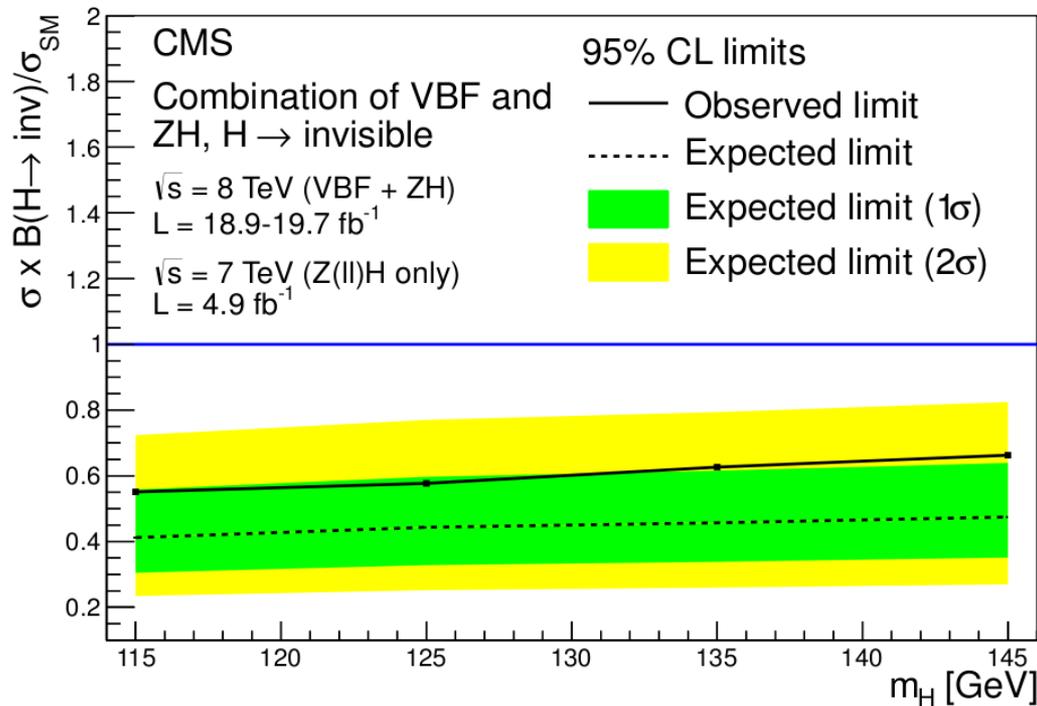
- **Very reach physics program for BSM Higgs boson searches and $h(125)$ boson measurements at LHC**
- **We expect to have another discovery in the Higgs sector during LHC or HL-LHC operation**

BACK UP

References to CMS results shown

- $h \rightarrow ZZ \rightarrow 4l$, arXiv:1312.5353 Phys. Rev. D 89 (092007) , width: PAS-14-002
- Tensor structure from $H \rightarrow WW$ and combined with $H \rightarrow ZZ$; HIG-14-012, HIG-14-014
- $h \rightarrow \gamma\gamma$, (p)
- $h \rightarrow WW$, arXiv:1312.1129, JHEP 01 (2014) 096
- Combination for h decays to fermions, HIG-13-033, arXiv:1401.6527
 - $h \rightarrow \tau\tau$, arXiv:1401.5041, JHEP 05 (2014) 104
 - Vh , $h \rightarrow bb$, arXiv:1310.3687, Phys. Rev. D 89, 012003 (2014)
- $t\bar{t}h$ combination, HIG-13-029 (to be paper), $t\bar{t}H$ with ME HIG-14-010
- $pp \rightarrow q\bar{t}h$, $h \rightarrow \gamma\gamma$ HIG-14-001, $h \rightarrow bb$ HIG-14-015
- $h \rightarrow Z\gamma/\gamma^*\gamma \rightarrow \ell\ell\gamma$, $h \rightarrow \mu\mu$, ee , arXiv:1307.5515, HIG-14-003, HIG-13-007
- $h \rightarrow$ invisible, arXiv:1404.1344
- $h \rightarrow \mu\tau$, HIG-14-005
- $\phi \rightarrow \tau\tau$, $\mu\mu$ HIG-13-021, HIG-13-024
- $H \rightarrow hh \rightarrow \gamma\gamma bb$, HIG-13-032
- $H \rightarrow hh$, $A \rightarrow Zh$ with multi-leptons and di- γ s, HIG-13-025
- Heavy $X \rightarrow \gamma\gamma$, HIG-14-006
- $H^+ \rightarrow \tau\nu$ with τ_h , no CADI line yet
- $H^+ \rightarrow tb$, $\tau\nu$ (with τ_θ), HIG-13-026
- $H^+ \rightarrow cs$, HIG-13-035
- $H \rightarrow WW, ZZ$, heavy Higgs combination HIG-13-031,
 - heavy Higgs HIG-14-007 ($ZZ \rightarrow 2l2q$), 14-008 ($WW \rightarrow l\nu j$) , HIG-13-027(jj)/008 (J) ($WW \rightarrow l\nu jj$), HIG-13-014 ($H-ZZ \rightarrow 2\ell 2\nu$)
- Couplings from combination, HIG-14-009

Constraining $BR(h \rightarrow \text{inv})$ using direct $h \rightarrow \text{inv}$ searches AND “visible” modes

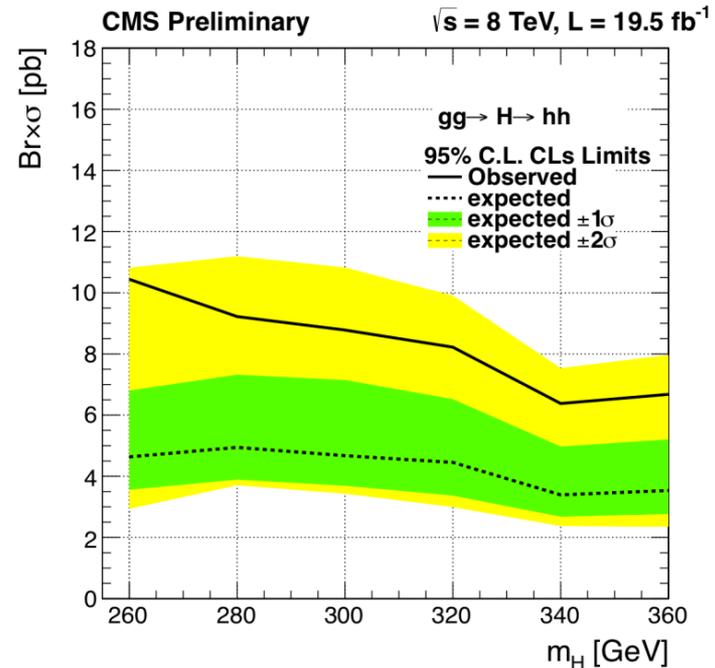
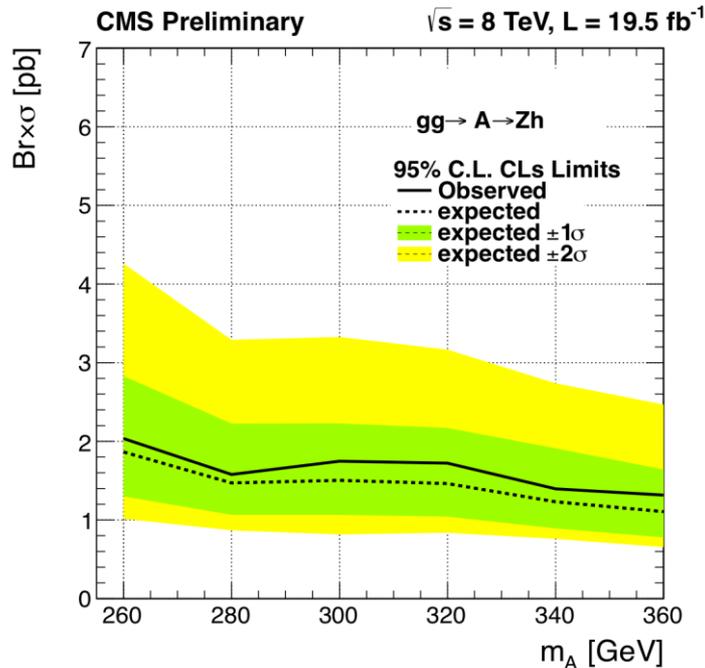


$BR_{\text{invisible}} < 0.50$ at 95% CL

(“invisible” - any decay mode that goes into “non-interacting” particles)

Low $\tan\beta$: $H \rightarrow hh$ and $A \rightarrow Zh$ with multi-lepton and di-photon final states

- Limits on $\sigma \times \text{Br}$ for $gg \rightarrow H \rightarrow hh$ ($gg \rightarrow A \rightarrow Zh$) are given assuming SM BRs for h and no contribution from $gg \rightarrow A \rightarrow Zh$ ($gg \rightarrow H \rightarrow hh$)
- In hMSSM (A. Djouadi et.al. arXiv:1307.5205):
 - $\sigma(gg \rightarrow A) \text{BR}(A \rightarrow Zh) = 1.7 \text{ pb}$ for $m_A = 300 \text{ GeV}$, $\tan\beta = 2.0$ - close to observed limit !
 - $\sigma(gg \rightarrow H) \text{BR}(H \rightarrow hh) = 3.9 \text{ pb}$ for $m_A = 300 \text{ GeV}$, $\tan\beta = 1.0$ – lower than observed limit



Background estimation in $\Phi \rightarrow \tau\tau$

$Z/\gamma^* \rightarrow \tau\tau$:

- Embedding: in $Z \rightarrow \mu\mu$, replace μ by sim. τ decay
- Normalized from $Z \rightarrow \mu\mu$ events

$t\bar{t}$:

- From simulation
- Normalization from sideband

QCD:

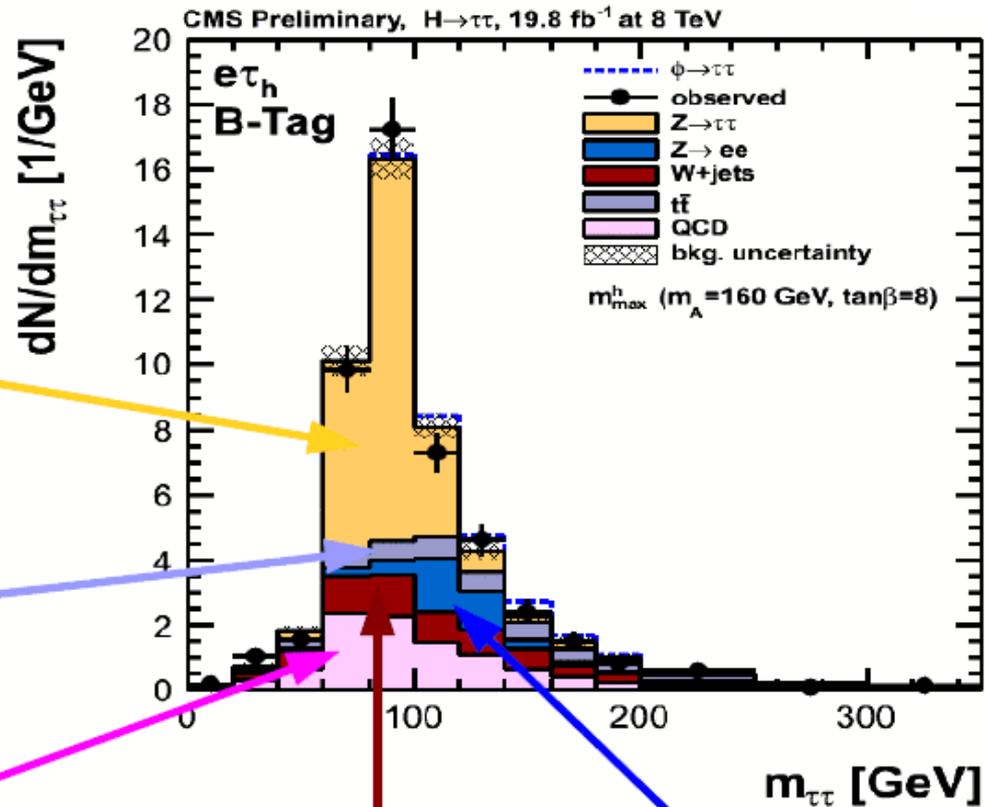
- Normalization & shape taken from SS/OS or fakerate

Di-boson/W+jets:

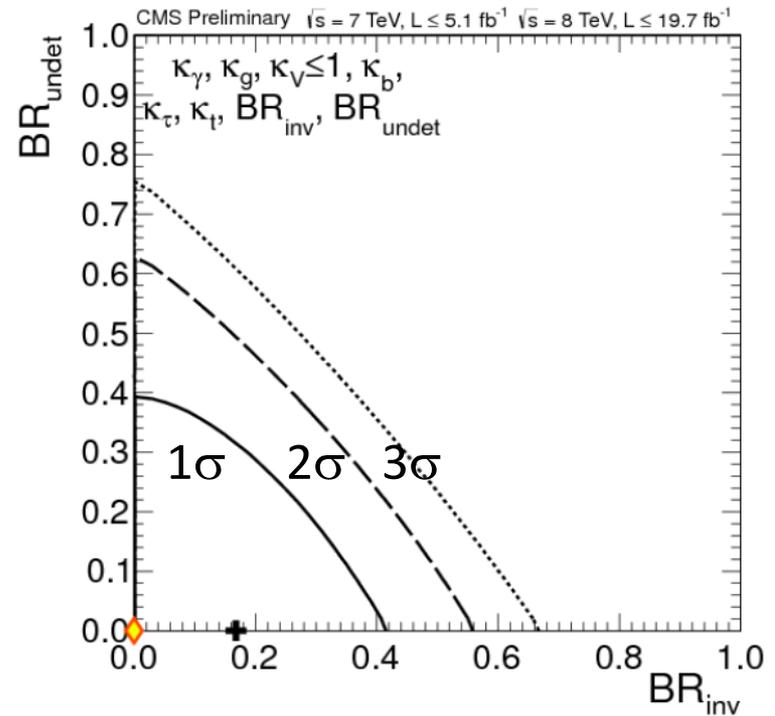
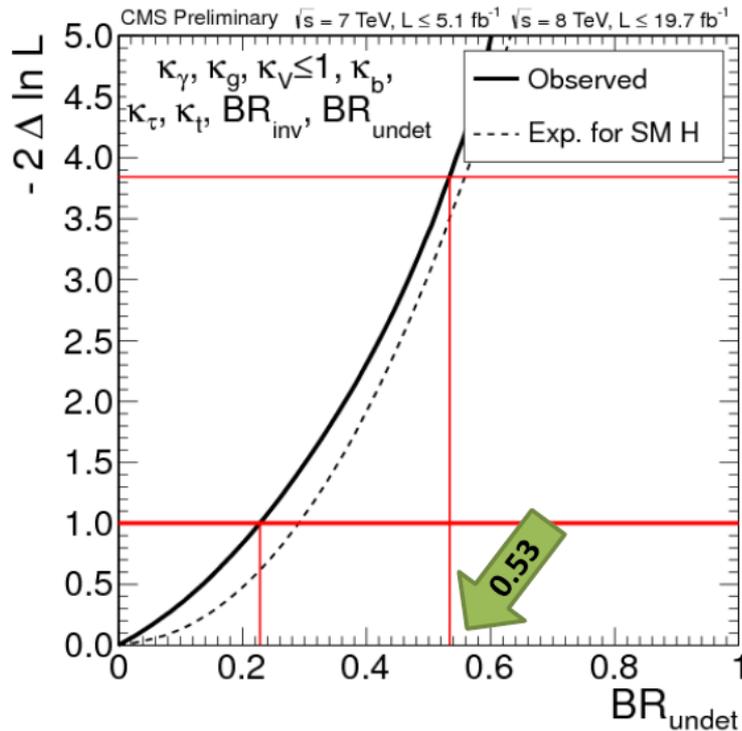
- From simulation or data
- Normalization from sideband

$Z/\gamma^* \rightarrow ee (\mu\mu)$:

- From simulation or data
- Corrected for jet $\rightarrow \tau$, $e/\mu \rightarrow \tau$ fakerate



Constraining “invisible” and “undetectable” modes using direct $h \rightarrow \text{inv}$ searches AND “visible” modes



$BR_{\text{undet}} < 0.53$ at 95% CL

(“undetectable” - any mode, SM or BSM, that can’t be seen,
e.g. due. to bkg or other reasons)