

# Relation between neutrino sector and proton decay in flipped SU(5) unification

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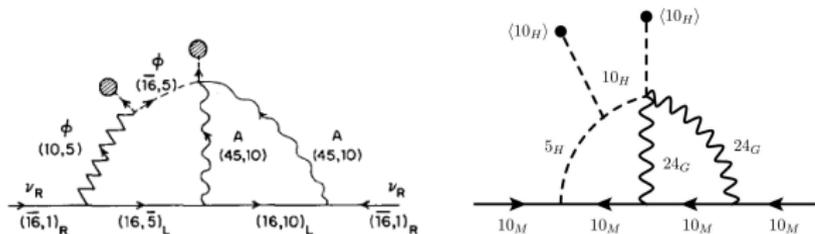
Institute of Particle and Nuclear Physics, FMP CU

Joint work with Michal Malinský (IPNP) and Carolina Arbeláez Rodríguez (University of Valencia)



# Outline

- 1 Witten's mechanism (=loop generation of RH neutrino masses) in SO(10) unification

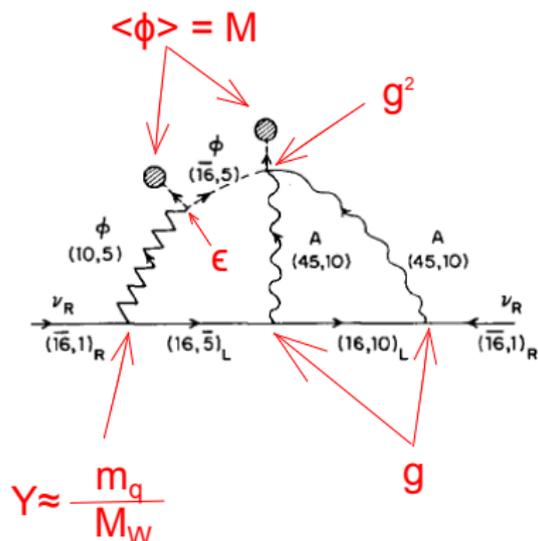


- 2 Witten's mechanism in non-SUSY flipped SU(5) unification
- 3 Interplay between neutrino sector and proton decay predictions
- 4 Witten's mechanism in SUSY: Qualitatively the same conclusions are valid if split SUSY considered!



# Witten's mechanism

- Edward Witten: Neutrino Masses in the Minimal  $O(10)$  Theory, Phys. Lett. B91, 81 (1980)



We can then estimate a right-handed neutrino mass

$$m_{\nu_R} = (m_q/M_W)\epsilon(\alpha/\pi)^2 M. \quad (3)$$

...

Thus, from (4),  $m_{\nu_L} = m_q^2/m_{\nu_R}$ , so, using our previous estimate for  $m_{\nu_R}$ , we find

$$m_{\nu_L} = m_q [\epsilon(\alpha/\pi)^2]^{-1} M_W/M. \quad (5)$$

With our previous numbers,  $M = 10^{15}$  GeV and  $\epsilon = 1/10$ , this means that for each generation  $m_{\nu_L} = 10^{-7} m_q$ .

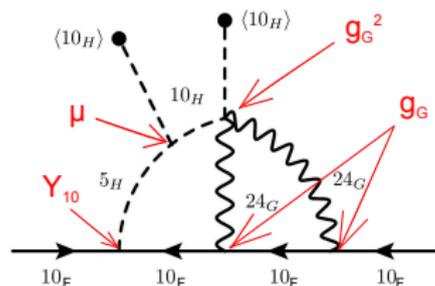
For the electron neutrino, this means  $m_{\nu_e} \approx 1$  eV, while for the muon neutrino  $m_{\nu_\mu} = 10^{-7} m_c \approx 100$  eV and for the  $\tau$  neutrino  $m_{\nu_\tau} = 10^{-7} m_t$  and might be 1–10 keV, depending on the  $t$  quark mass.

- Precise measurements of  $\sin \theta_W \Rightarrow$  gauge coupling unification for non-SUSY models only if  $M \sim 10^{13}$  GeV (intermediate scale)  $\Rightarrow m_{\nu_L}$  too big.



# Witten's mechanism in the flipped SU(5) unification

$$M_{\nu_R} = \left(\frac{1}{16\pi^2}\right)^2 g_G^4 Y_{10} \mu \frac{\langle 10_H \rangle^2}{M_G^2} \times \mathcal{O}(1)$$



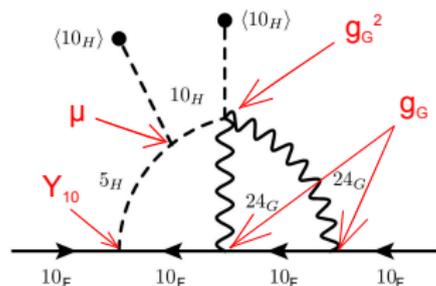
- $M_G \approx g_G \langle 10_H \rangle$  (VB mass),  $\mu \leq \lambda \langle 10_H \rangle$  (non-tachyonic spectrum).

Perturbativity of couplings  $\Rightarrow$

$$\max_{i,j \in \{1,2,3\}} |(M_{\nu_R})_{ij}| \leq \frac{\alpha_G}{16\pi^2} M_G \times \mathcal{O}(1)$$

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- Seesaw mechanism:

$$M_{\nu_L} = -M_{\nu}^D (M_{\nu_R})^{-1} (M_{\nu}^D)^T \Rightarrow$$

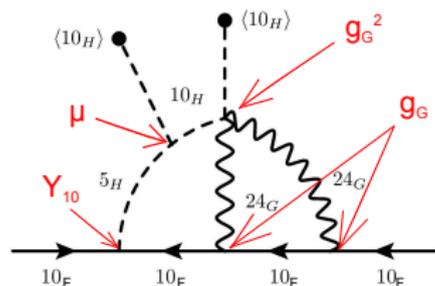
$$M_{\nu_R} = -D_u U_{\nu}^{\dagger} D_{\nu}^{-1} U_{\nu}^* D_u$$

( $D_{u/\nu}$  – diagonal mass matrices for up-type quarks/light neutrinos)

$U_{\nu}$  – unitary matrix such that  $M_{\nu_L} = U_{\nu}^T D_{\nu} U_{\nu}$ )

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$\Rightarrow$  **constraint on  $U_{\nu}$  (dependent on light neutrino masses!)**

# GUT = Theories of Proton Decay

- In the simplest flipped  $SU(5)$  scenarios (with down quark mass matrix symmetric):

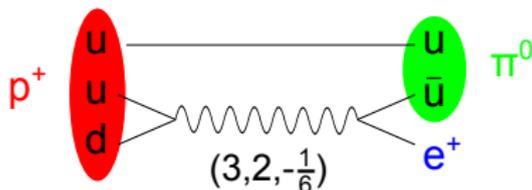
$$\Gamma(p \rightarrow K^+ \bar{\nu}) = 0$$

$$\Gamma(p \rightarrow \pi^+ \bar{\nu}) = C_1 \left( \frac{g_G}{M_G} \right)^4$$

$$\frac{\Gamma(p \rightarrow \pi^0 e^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} = \frac{1}{2} |(V_{CKM})_{11}|^2 |(V_{PMNS} U_\nu)_{\alpha 1}|^2$$

$$\frac{\Gamma(p \rightarrow \eta e^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} = \frac{C_2}{C_1} |(V_{CKM})_{11}|^2 |(V_{PMNS} U_\nu)_{\alpha 1}|^2$$

$$\frac{\Gamma(p \rightarrow K^0 e^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} = \frac{C_3}{C_1} |(V_{CKM})_{12}|^2 |(V_{PMNS} U_\nu)_{\alpha 1}|^2$$



- $C_i$  - calculable factors,  $U_\nu$  constrained in our model!
- Lower bound on proton lifetime determined experimentally: e.g.  $\tau(p \rightarrow \pi^0 + e^+) > 8.2 \times 10^{33} \text{ y}$  [Super-Kamiokande Collaboration (2012)]

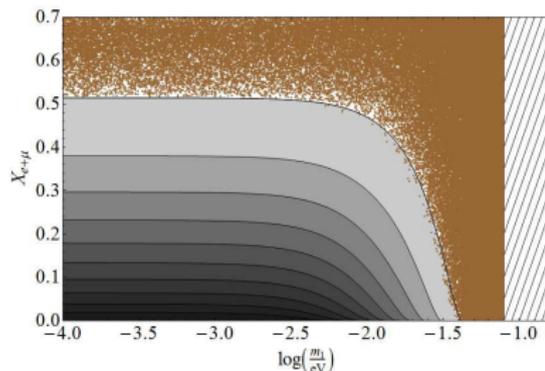
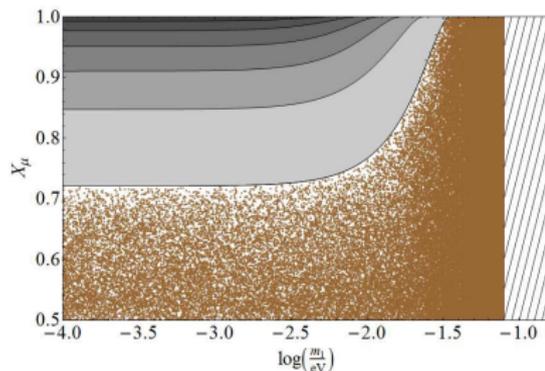
# Results

- $|(V_{PMNS} U_\nu)_{\alpha 1}|^2$  computed if constraints on  $U_\nu$  applied and the mass of the lightest neutrino  $m_1$  varied from  $10^{-4}$  eV to 0.08 eV.

⇒ Global upper/lower limits on

$$X_\mu \equiv \frac{\Gamma(p \rightarrow \pi^0 \mu^+)}{\frac{1}{2} \Gamma(p \rightarrow \pi^+ \bar{\nu}) |(V_{CKM})_{11}|^2}$$
$$X_{e+\mu} \equiv \frac{\Gamma(p \rightarrow \pi^0 e^+) + \Gamma(p \rightarrow \pi^0 \mu^+)}{\frac{1}{2} \Gamma(p \rightarrow \pi^+ \bar{\nu}) |(V_{CKM})_{11}|^2}$$

- Brown dots = points consistent with constraints on  $U_\nu$ ,  
full lines = analytically computed bounds of the regions with consistent points
- The model parameterized by the multiplicative constant in the expression for the loop-generated RH neutrino mass ⇒ different curves.
- Absolute scale:  $\Gamma^{-1}(p \rightarrow \pi^+ \bar{\nu}) \sim 10^{39} \text{y}$

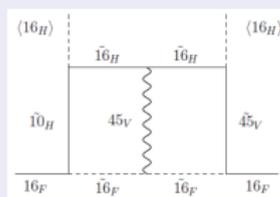
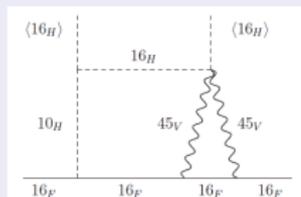


# Witten's mechanism in SUSY SO(10)

- ☺ Gauge unification works for SUSY SO(10) without intermediate scale  $\Rightarrow$  VEV entering the Witten's loop of order  $M_G \sim 10^{16}$  GeV  $\Rightarrow$  realistic neutrino masses
- ☹ Soft SUSY breaking @TeV  $\Rightarrow$  Witten's loop compensated by superpartner contributions (non-renormalization theorem)

## B. Bajc and G. Senjanović, Phys. Lett. B 610 (2005) 80

- Implementation of Witten's mechanism in maximally split SUSY SO(10) model
- Gauginos and higgsinos @ TeV  $\Rightarrow$  ☺ remains
- Sfermions @ GUT scale  $\Rightarrow$  ☹ solved
- Loop contributions to charged fermion masses  $\Rightarrow$  possible cure for wrong mass relations like  $M_d = M_l$  for down-type quark and charged lepton mass matrices

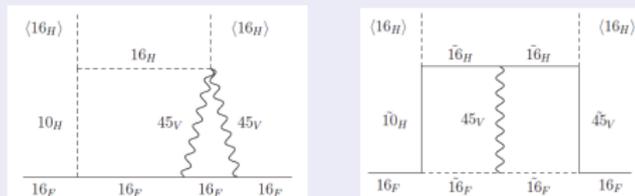


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**Qualitatively no change if our model considered in the split SUSY context!**

# Conclusion

- Viable model based on  $SU(5) \times U(1)$  gauge group (“flipped  $SU(5)$ ”) constructed
- To generate Majorana masses of RH neutrinos no new fields added – generated via **Witten’s mechanism**  $\Rightarrow$  “minimal” setting
- Predictions on potentially measurable quantities (all proton decay partial widths) calculable
- N.B.: Alternative option for generating RH neutrino masses in flipped  $SU(5)$  model: adding  $50_H$  scalar representation. No correlation between neutrino sector and proton decay in this case!
- Our results are valid qualitatively in the **split SUSY** context as well
- For further details see [C. A. Rodríguez, H. Kolečová, and M. Malinský, Phys.Rev. D89 \(2014\) 055003, arXiv:1309.6743 \[hep-ph\]](#)

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**Thank you for your attention!**