

# Lepton flavour violating lepton decays in the supersymmetric inverse seesaw

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Cédric Weiland

in collaboration with A. Abada, M. E. Krauss, W. Porod, F. Staub and A. Vicente

Instituto de Física Teórica, Universidad Autónoma de Madrid/CSIC, Spain

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# Neutrino masses and lepton flavour violation

- $P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta} \neq 0$  only if  $\Delta m_{kj}^2 = m_k^2 - m_j^2$  and  $U_\nu \neq \mathbb{1}$
- SM: no  $\nu$  mass term, lepton flavour is conserved  
⇒ need new Physics, e.g. seesaw mechanism
- Neutrino oscillations ⇒ neutral lepton flavour violation. Why not charged lepton flavour violation (cLFV) ?
- cLFV arises from higher order processes:  
negligible in the SM
- If observed:
  - Evidence of New Physics
  - Might probe the origin of lepton mixing
  - Might probe the origin of new physics

# cLFV

- Complementary to other New Physics searches

- High energy: LHC
- High intensity:
  - B factories: Rare decays, etc
  - Neutrino dedicated experiments:  $U_{PMNS}$  non-unitarity, NSI, etc
  - Other low energy experiments:  $(g - 2)_\mu$ , EDM, LUV, etc

- Can occur in many channels

- Radiative decays, e.g.  $\text{Br}(\mu \rightarrow e\gamma) < 5.7 \times 10^{-13}$  [MEG, 2013]
- 3-body lepton decays, e.g.  $\text{Br}(\tau \rightarrow 3\mu) < 2.1 \times 10^{-8}$  [Belle, 2010]
- $\mu - e$  conversion, e.g.  $\text{CR}(\mu \rightarrow e, \text{Au}) < 7 \times 10^{-13}$  [SINDRUM II, 2006]
- Meson decays, e.g.  $\text{Br}(B_d^0 \rightarrow e\mu) < 2.8 \times 10^{-9}$  [LHCb, 2013]
- Z decays, e.g.  $\text{Br}(Z^0 \rightarrow e\mu) < 1.7 \times 10^{-6}$  [OPAL, 1995]
- Higgs decays, e.g.  $\text{Br}(H \rightarrow \tau\mu) < 1.57\%$  [CMS-PAS-HIG-14-005]

# Why supersymmetry?

- The SM doesn't only lack neutrino masses, e.g. the hierarchy problem
- Extended frameworks to address SM issues:
  - Strongly coupled theories (e.g. Technicolor, Composite Higgs)
  - Extra-dimensions (e.g. Randall-Sundrum, Large extra dimension)
  - Extending the SM field content/gauge group (e.g. 2HDM, Little Higgs, GUT)
  - Supersymmetric extensions (e.g. **MSSM**)
- Advantages of SUSY
  - Most general extension of the Poincaré algebra
  - Gauge coupling unification
  - Dark matter candidate
  - Graviton naturally appears in local supersymmetry

# Supersymmetric seesaw models

- No  $\nu_R$  in the MSSM  $\Rightarrow$  Massless neutrinos  
→ Implement a seesaw mechanism
- Seesaw mechanism: Consider new fields at this scale ( $\sim M_R$ ) and Majorana mass terms  $\Rightarrow$  Generate  $m_\nu$  in a **renormalizable** way
- **Unique** dimension 5 operator for all seesaw mechanisms  
→ Violates lepton number L  $\Rightarrow$  **Majorana neutrinos**

$$\delta\mathcal{L}^{d=5} = \frac{1}{2} c_{ij} \frac{\bar{L}_i \tilde{H} \tilde{H}^T L_j^C}{\Lambda} + \text{h.c.}$$

- To distinguish the several seesaw mechanisms, either
  - Directly produce the heavy states (LHC, future collider)
  - Look for dimension  $\geq 6$  operators effects  $\rightarrow$  **cLFV**

# The inverse seesaw mechanism

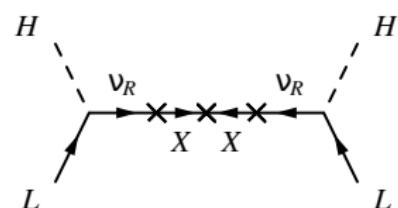
- Inverse seesaw  $\Rightarrow$  Consider fermionic gauge singlets  $\nu_{Ri}$  ( $L = +1$ ) and  $X_i$  ( $L = -1$ ) [Mohapatra and Valle, 1986]

$$\mathcal{L}_{\text{inverse}} = -Y_\nu^{ij} \overline{L}_i \tilde{H} \nu_{Rj} - M_R^{ij} \overline{\nu_{Ri}^C} X_j - \frac{1}{2} \mu_X^{ij} \overline{X_i^C} X_j + \text{h.c.}$$

with  $m_D = Y_\nu v$ ,  $M^\nu = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & 0 & M_R \\ 0 & M_R^T & \mu_X \end{pmatrix}$

$$m_\nu \approx \frac{m_D^2 \mu_X}{m_D^2 + M_R^2}$$

$$m_{N_1, N_2} \approx \mp \sqrt{m_D^2 + M_R^2} + \frac{M_R^2 \mu_X}{2(m_D^2 + M_R^2)}$$



2 scales:  $\mu_X$  and  $M_R$

# The supersymmetric inverse seesaw model

- MSSM extended by singlet chiral superfields  $\hat{N}_i$  and  $\hat{X}_i$  with  $L = -1$  and  $L = +1$
- Superpotential:

$$\begin{aligned} \mathcal{W} = & Y_d \hat{Q} \hat{H}_d \hat{D} + Y_u \hat{Q} \hat{H}_u \hat{U} + Y_e \hat{L} \hat{H}_d \hat{E} - \mu \hat{H}_d \hat{H}_u \\ & + Y_\nu \hat{L} \hat{H}_u \hat{N} + M_R \hat{N} \hat{X} + \frac{1}{2} \mu_X \hat{X} \hat{X} \end{aligned}$$

- New couplings, e.g.  $A_{Y_\nu} Y_\nu \tilde{L} \tilde{N} H_u + \text{h.c.}$
- Work with a flavour-blind mechanism for SUSY breaking,  
**CMSSM-like** boundary conditions
- Right-handed sneutrino mass:

$$M_{\tilde{N}}^2 = \textcolor{blue}{m_{\tilde{N}}^2} + \textcolor{blue}{M_R^2} + Y_\nu^\dagger Y_\nu v_u^2 \sim (1 \text{TeV})^2$$

⇒ Natural Yukawa couplings with a TeV new Physics scale

# cLFV in supersymmetric seesaw models

- Typically in SUSY, cLFV appears through RGE-induced slepton mixing  $(\Delta m_L^2)_{ij}$   
 [Borzumati and Masiero, 1986, Hisano et al., 1996, Hisano and Nomura, 1999]  
 $\Rightarrow (\Delta m_L^2)_{ij} \propto (Y_\nu^\dagger Y_\nu)_{ij} \ln \frac{M_{GUT}}{M_R}$
- Contribute to all cLFV observables  
 → Dominant in most of the SUSY seesaw models
- Type I seesaw ( $Y_\nu \sim 1$ ,  $M_R \sim 10^{14}$ GeV)  $\rightarrow (\Delta m_L^2)_{ij} \propto 5$
- Inverse seesaw ( $Y_\nu \sim 1$ ,  $M_R \sim 1$ TeV)  $\rightarrow (\Delta m_L^2)_{ij} \propto 30$   
 → one-loop  $\tilde{N}$ -mediated processes are no longer suppressed

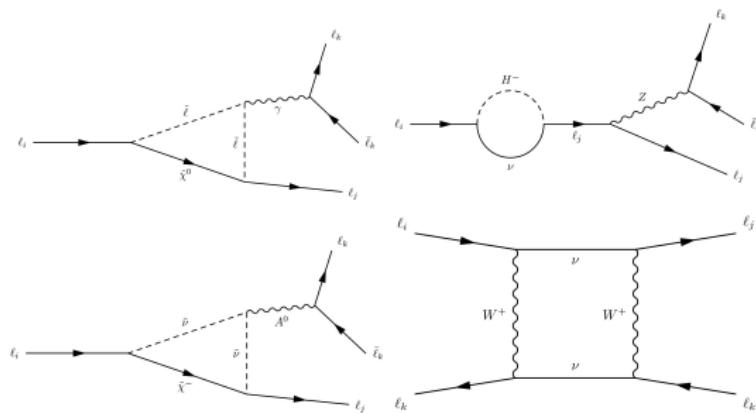
[Deppisch and Valle, 2005, Hirsch et al., 2010, Abada et al., 2012, Ilakovac et al., 2012, Krauss et al., 2013]

## Similar enhancement in non-SUSY contributions

[Ilakovac and Pilaftsis, 1995, Deppisch et al., 2006, Forero et al., 2011, Alonso et al., 2013, Dinh et al., 2012]

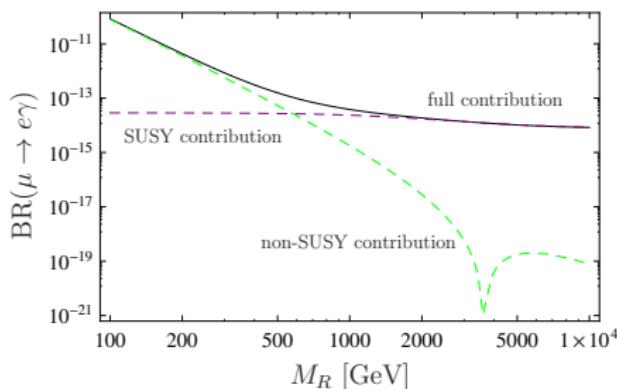
# Diagrams

- In the Feynman-'t Hooft gauge, including both SUSY and non-SUSY:  
More than 100 classes of diagrams

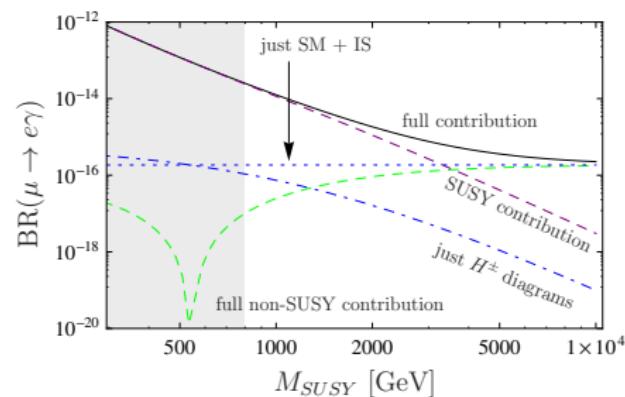


- $\gamma, Z, h_i, A_i$ -penguins and boxes
- Formulas computed using the FlavorKit interface
- Checked against the literature when possible
- Numerics done with SARAH/Spheno using 2 loops RGEs
- Enhancement from:  
 $\mathcal{O}(1) Y_\nu$  couplings  
-TeV scale  $\nu_R, \tilde{N}$

# Radiative cLFV decays



$$m_0 = M_{1/2} = 1 \text{ TeV}, A_0 = -1.5 \text{ TeV}$$

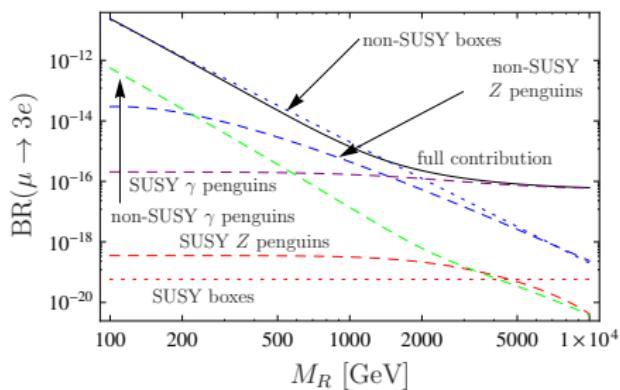


$$M_R = 2 \text{ TeV} \mathbb{1}, \\ M_{\text{SUSY}} = m_0 = M_{1/2} = -A_0$$

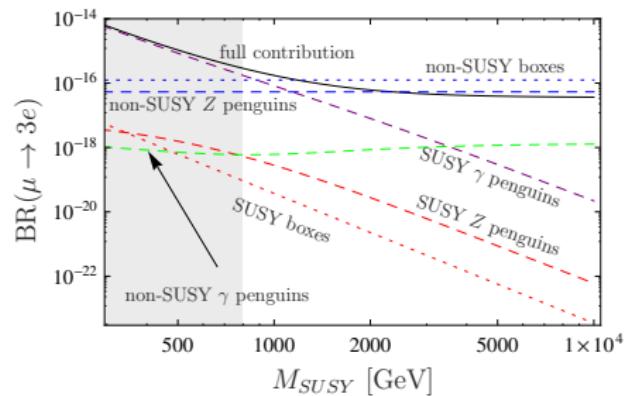
$$\tan \beta = 10, \text{ sign}(\mu) = +, \mu_X = 10^{-5} \text{ GeV} \mathbb{1}, B_{\mu_X} = 100 \mu_X, B_{M_R} = 100 M_R$$

- Reach the current upper limit:  $\text{Br}(\mu \rightarrow e\gamma) < 5.7 \times 10^{-13}$  [MEG, 2013]  
Expected sensitivity:  $6 \times 10^{-14}$  [MEG upgrade]
- Dominant contribution from the **lightest scale** ( $M_R$  or  $M_{\text{SUSY}}$ )

# 3-body cLFV decays



$$m_0 = M_{1/2} = 1 \text{ TeV}, A_0 = -1.5 \text{ TeV}$$

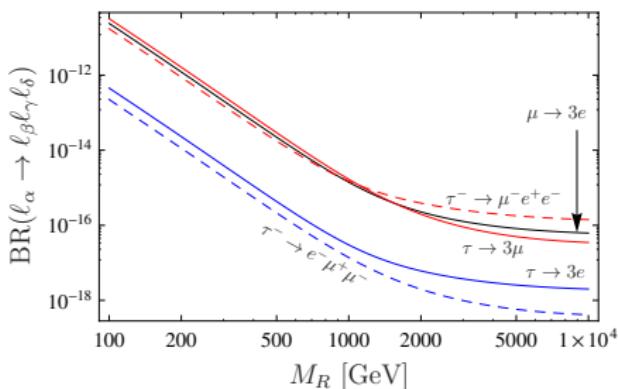


$$M_R = 2 \text{ TeV} \mathbb{1},$$

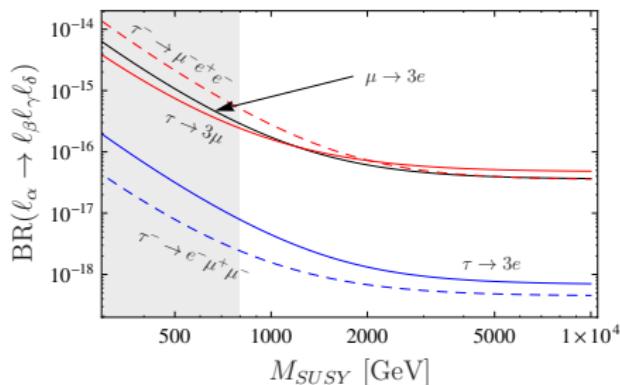
$$M_{\text{SUSY}} = m_0 = M_{1/2} = -A_0$$

- **Saturate current UL**:  $\text{Br}(\mu \rightarrow eee) < 1.0 \times 10^{-12}$  [SINDRUM, 1988]  
Expected sensitivity:  $10^{-15} - 10^{-16}$  [Mu3e proposal]
- Dominant non-SUSY contribution: **boxes and Z-penguins**
- Dominant SUSY contribution:  **$\gamma$ -penguins**
- Higgs-penguins sub-dominant, except at  $\tan \beta \geq 50$  ( $\tan^6 \beta$  enhanced)

# 3-body cLFV decays



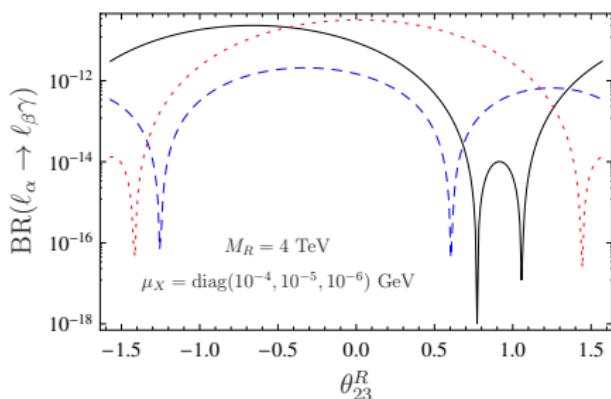
$$m_0 = M_{1/2} = 1 \text{ TeV}, A_0 = -1.5 \text{ TeV}$$



$$M_R = 2 \text{ TeV} \mathbb{1}, \\ M_{SUSY} = m_0 = M_{1/2} = -A_0$$

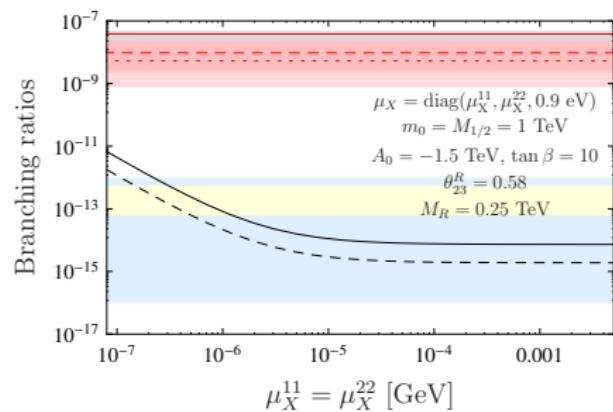
- Contributions with **similar behaviours** for both  $\mu$  and  $\tau$  decays
- $R = \mathbb{1}, \theta_{13} \ll \theta_{12}, \theta_{23}$   
→ Similar Br for  $\tau \rightarrow \mu, \mu \rightarrow e$  transitions, much smaller for  $\tau \rightarrow e$
- Strong suppression** of  $\text{Br}(\tau \rightarrow e\mu^+ e^-)$  and  $\text{Br}(\tau \rightarrow \mu e^+ \mu^-)$ :  
2 LFV vertices are needed

# Impact of active-sterile neutrino misalignment



$m_0 = M_{1/2} = 1 \text{ TeV}, A_0 = -1.5 \text{ TeV}$

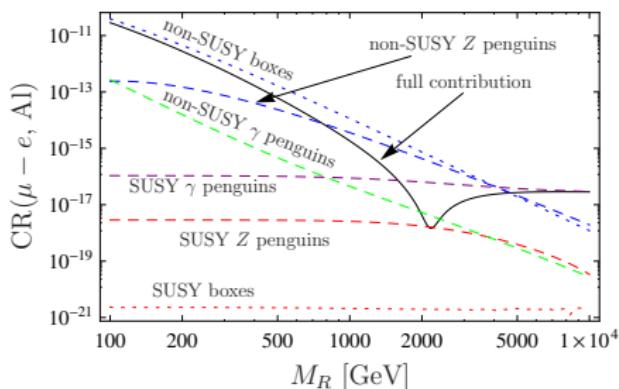
$\mu \rightarrow e\gamma$   $\tau \rightarrow e\gamma$   $\tau \rightarrow \mu\gamma$



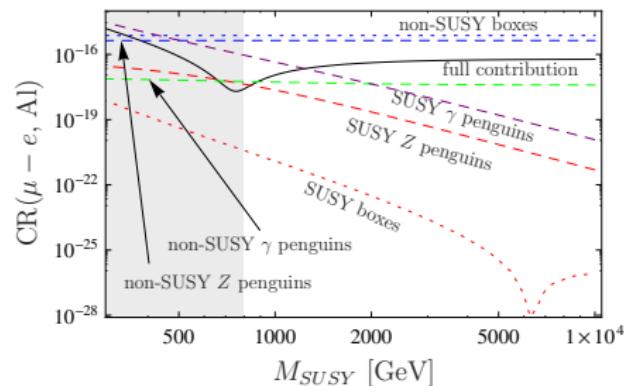
$\mu \rightarrow e\gamma$  (full),  $\mu \rightarrow 3e$  (dashed),  
 $\tau \rightarrow \mu\gamma$  (full),  $\tau \rightarrow 3\mu$  (dashed),  
 $\tau^- \rightarrow \mu^- e^+ e^-$  (dotted)

- Shaded areas: Expected sensitivities of future experiments
- $R \neq \mathbb{1}$  impacts relative size of Br
- Large enhancement of cLFV  $\tau$  decays

# Neutrinoless $\mu - e$ conversion



$$m_0 = M_{1/2} = 1 \text{ TeV}, A_0 = -1.5 \text{ TeV}$$

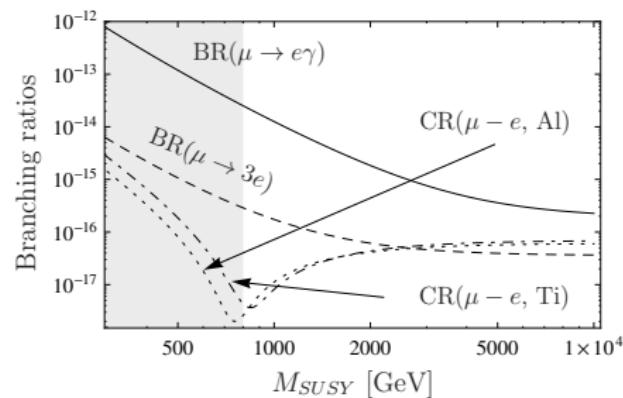
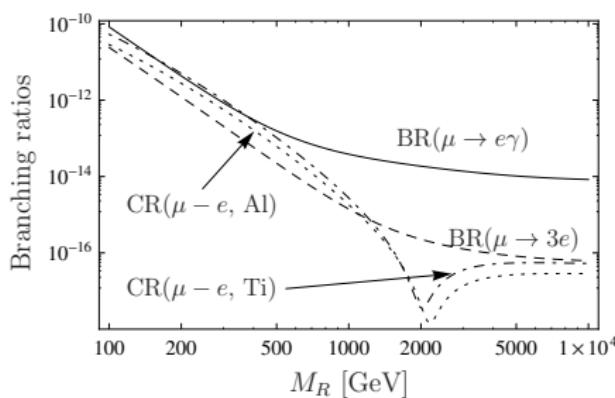


$$M_R = 2 \text{ TeV} \mathbb{1},$$

$$M_{SUSY} = m_0 = M_{1/2} = -A_0$$

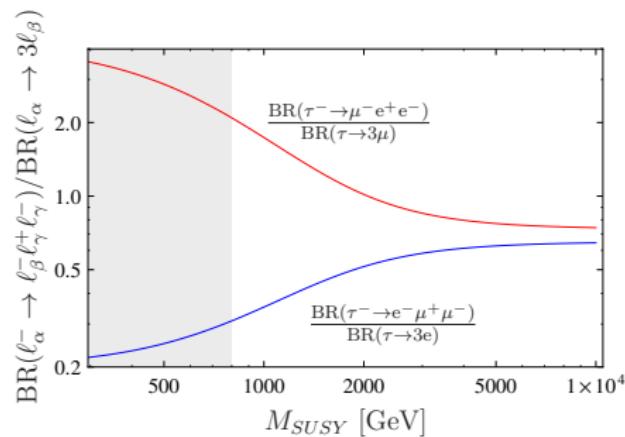
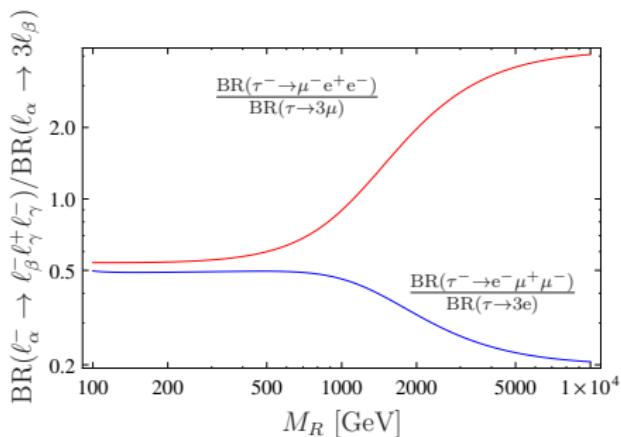
- **Saturate current UL**:  $CR(\mu - e, Au) < 7.0 \times 10^{-13}$  [SINDRUM II, 2006]  
Expected sensitivity:  $10^{-14}$  [DeeMe],  $10^{-17} - 10^{-18}$  [Mu2e, COMET/PRISM]
- Dips: partial cancellation between up quark and down quark contributions
- Otherwise similar to  $\mu \rightarrow eee$

# Comparison of cLFV decays



- $\mu \rightarrow e\gamma$ : largest Br and the lowest current UL ( $5.7 \times 10^{-13}$ )  
→ Most constraining observable today
- $\mu \rightarrow 3e$  conversion: best mid-term sensitivity ( $\sim 10^{-15}$ )  
→ Should be the most constraining by 2016.
- $\mu - e$  conversion: best long-term sensitivity (down to  $10^{-18}$ )  
→ Should be the most constraining around 2020.

# Finding the dominant contribution



- LFV  $\tau$  decays: factor 100 sensitivity improvement in Belle II
- Ratios: sensitive to the dominant contribution (SUSY or non-SUSY)

# Conclusions

- First complete calculation with both SUSY and non-SUSY contributions
- At low  $M_R$  / high  $M_{SUSY}$ : dominant contributions from non-SUSY boxes and Z-penguins
- At low  $M_{SUSY}$  / high  $M_R$ : dominant contributions from SUSY  $\gamma$ -penguins
- All observables can already be used to constrain the parameter space
- Most promising observable:
  - short-term:  $\mu \rightarrow e\gamma$
  - mid-term:  $\mu \rightarrow 3e$
  - long-term:  $\mu - e$  conversion
- Use ratios of  $\tau$  decays to find the dominant contribution

# Backup slides

# Modified Casas-Ibarra parameters and neutrino input

- Casas-Ibarra parametrization adapted to the inverse seesaw:

$$Y_\nu = \frac{\sqrt{2}}{v_u} V^\dagger D_{\sqrt{X}} R D_{\sqrt{m_\nu}} U_{\text{PMNS}}^\dagger$$

- Input parameters:  $M_R = 2 \text{ TeV}$ ,  $\mu_X = 10^{-5} \text{ GeV}$ ,  $m_{\nu_1} = 10^{-4} \text{ eV}$ ,  $R$  matrix
- Neutrino oscillation best-fit values [Forero et al., 2014]:

$$\Delta m_{21}^2 = 7.60 \cdot 10^{-5} \text{ eV}^2, \quad m_{31}^2 = 2.48 \cdot 10^{-3} \text{ eV}^2, \\ \sin^2 \theta_{12} = 0.323, \quad \sin^2 \theta_{23} = 0.467, \quad \sin^2 \theta_{13} = 0.0234$$

